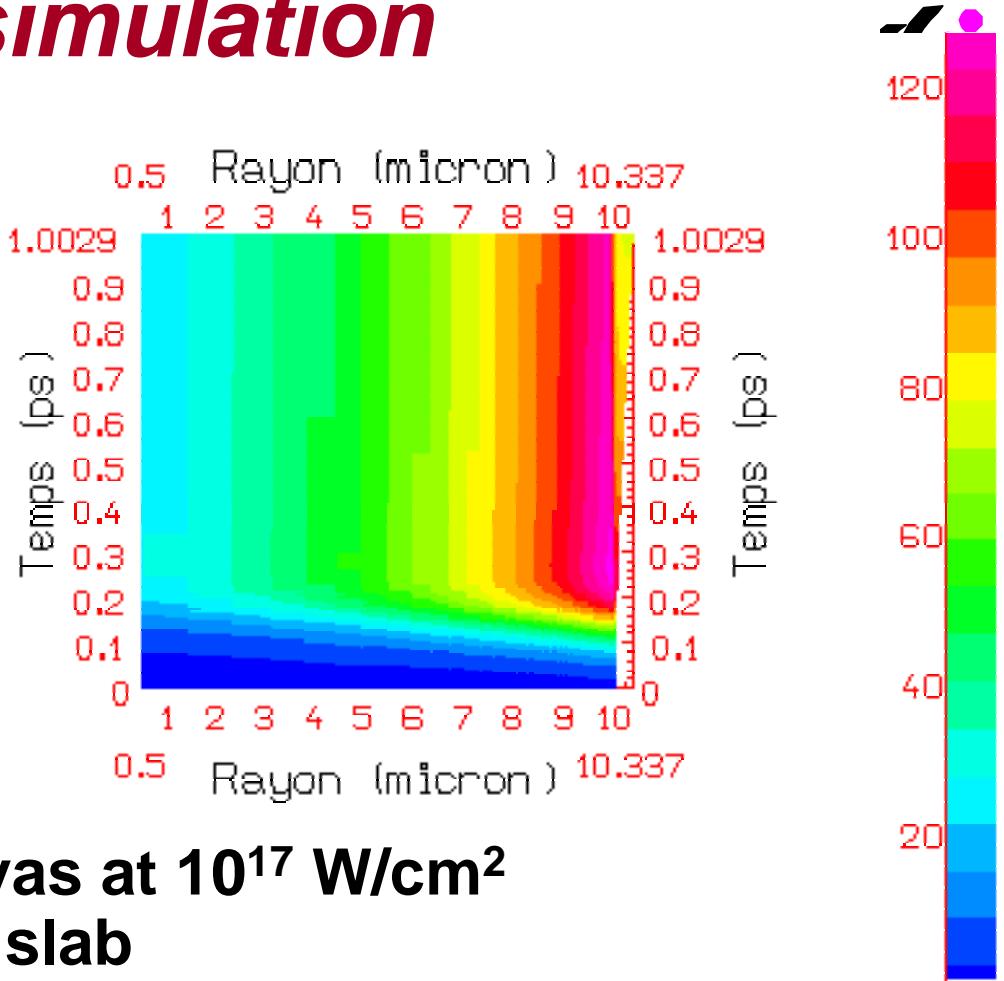
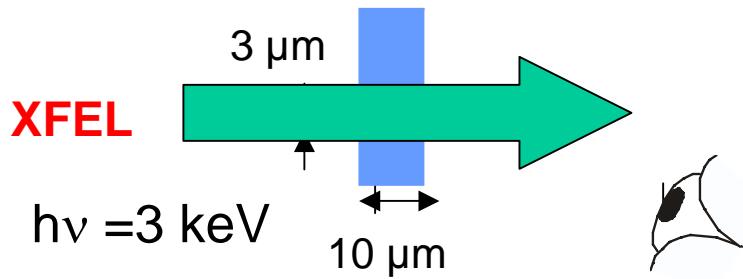


Frequency-domain interferometry for laser-produced plasma diagnostics

P. Audebert
Laboratoire LULI, Ecole Polytechnique

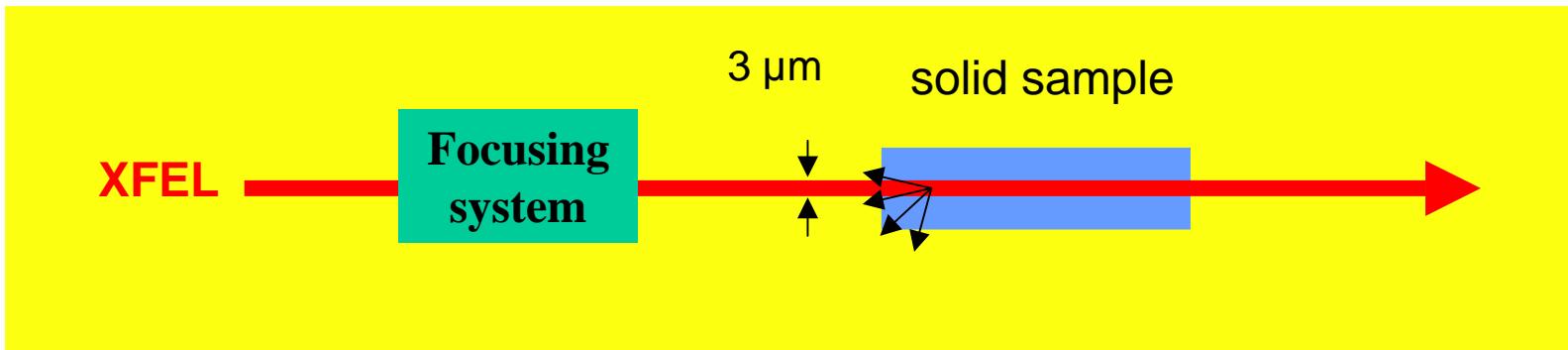
J.P. Geindre, C. Quoix, S. Rebibo, A. Mancic, J. Fuchs

XFEL Preliminary simulation



- ▶ Hydrodynamic 1D Code Chivas at 10^{17} W/cm^2
 $\hbar\nu = 3 \text{ keV}$, 100fs on 10μm Al slab
- ▶ Assuming we focus to a 3μm focal spot
- ▶ Target at solid density is heated up to 120 eV for more than a ps

Single beam experiment



Measurement

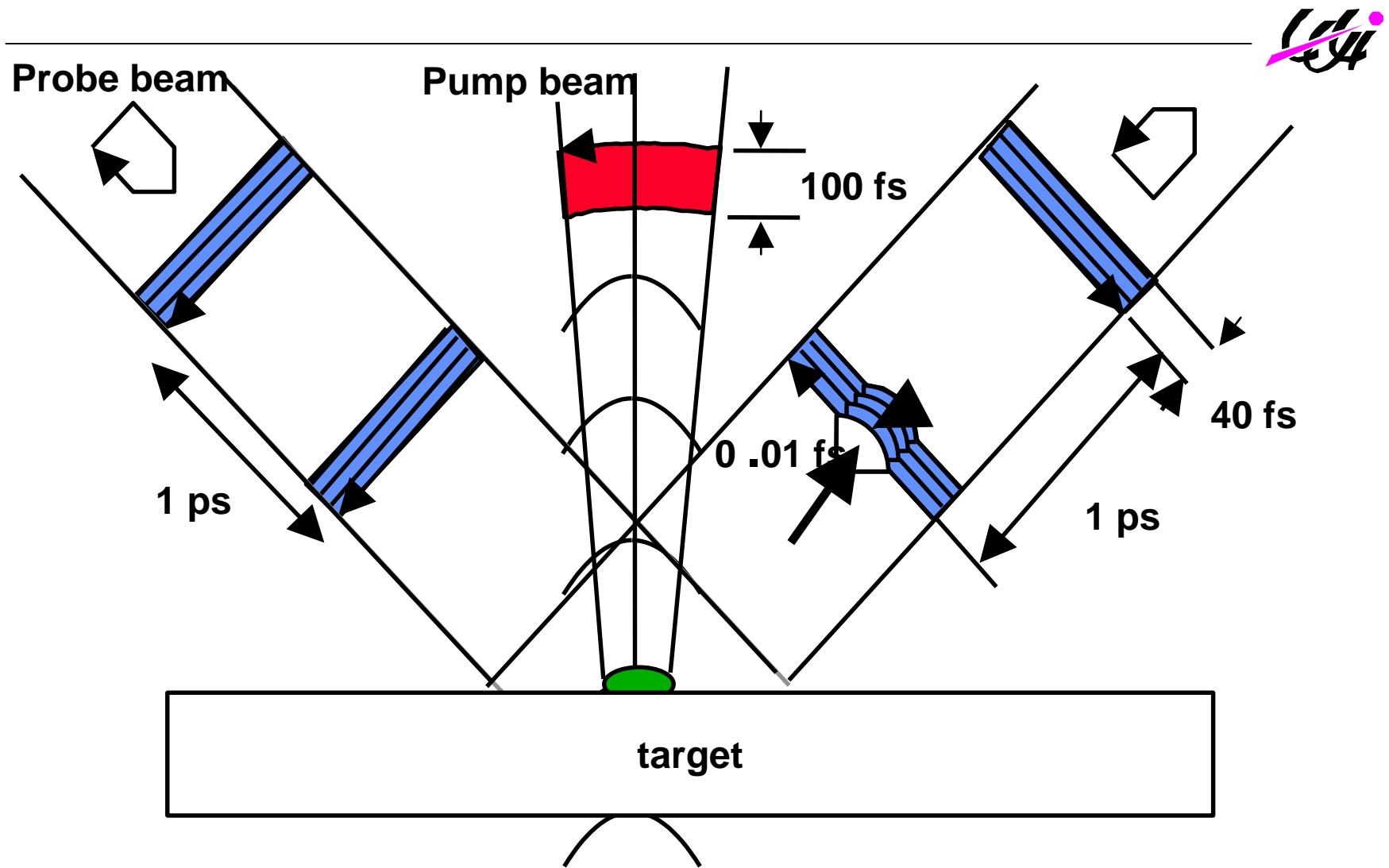
- X-FEL Scattering spectrum (Thomson)
- Transmitted energy (absorption)
- Plasma emission spectrum X-ray, XUV
time resolved (ps)
- Reflectivity measurement
- Pump probe experiment

Interferometry in Frequency Domain

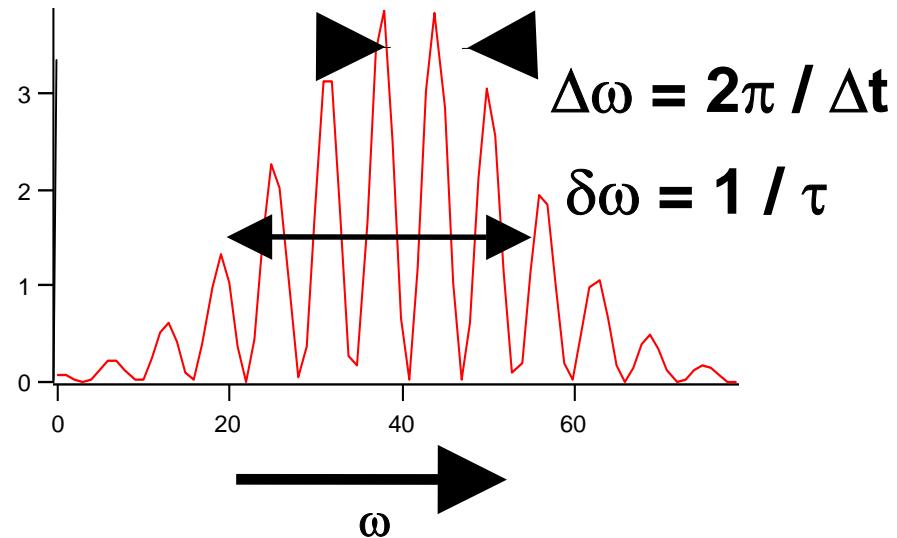
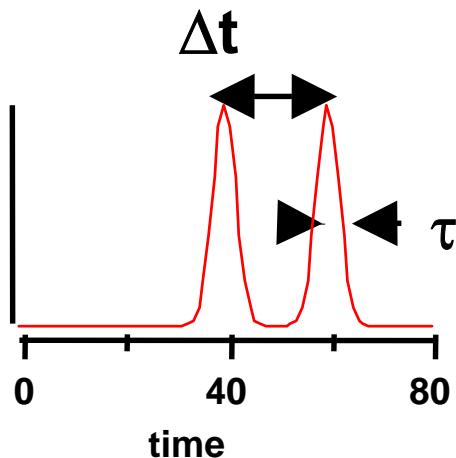


- Plasmas produced by high intensity laser are characterized by **very short lifetime, very steep gradients**, and small lateral dimensions.
- Different techniques as been used to have information of the plasma in time and space reflectometry, Schlieren, omboscopy, 2nd harmonic, spectral shifting
but record only change of amplitude of the probe.
- We have developed a technique :
the Interferometry in Fréquency Domain to access to both amplitude and phase of the probe.

Spectral interferometry



Phase-Space Interferometry



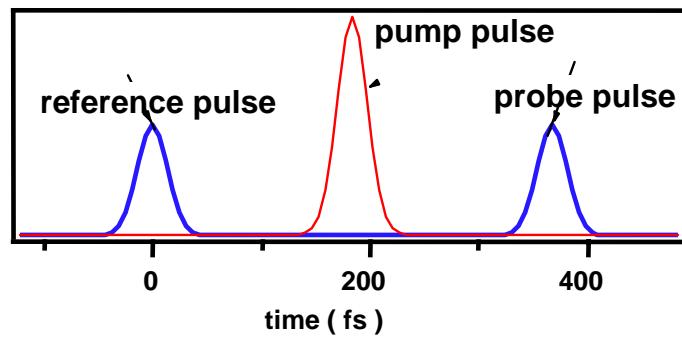
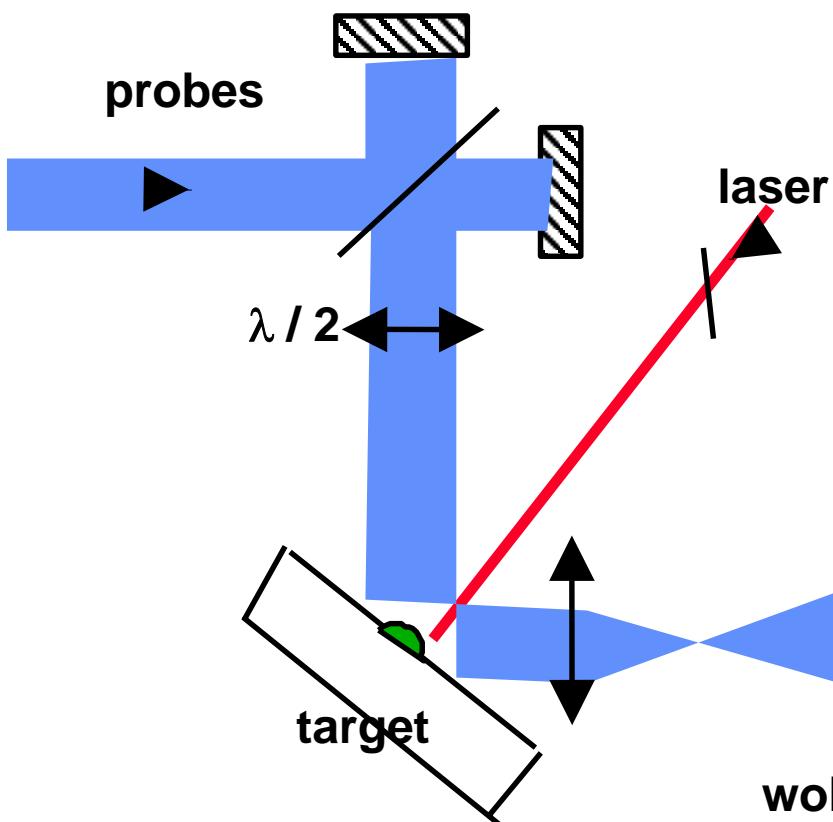
$$\left. \begin{aligned} E_1(t) &= E_0 \cdot e^{i\omega_0 t} \\ E_2(t) &= R \cdot e^{i\varphi} \cdot E_0(t) \cdot e^{i\omega_0 dt} \end{aligned} \right\}$$

The Spectrum is

$$S(\omega) = I_0(\omega) \cdot \{1 + R^2 + 2R \cdot \cos(\omega \cdot \Delta t + \varphi)\}$$

the interfringe separation is inversely proportional to Δt
where Δt is the time interval between the twin probe pulses

Experiment set up



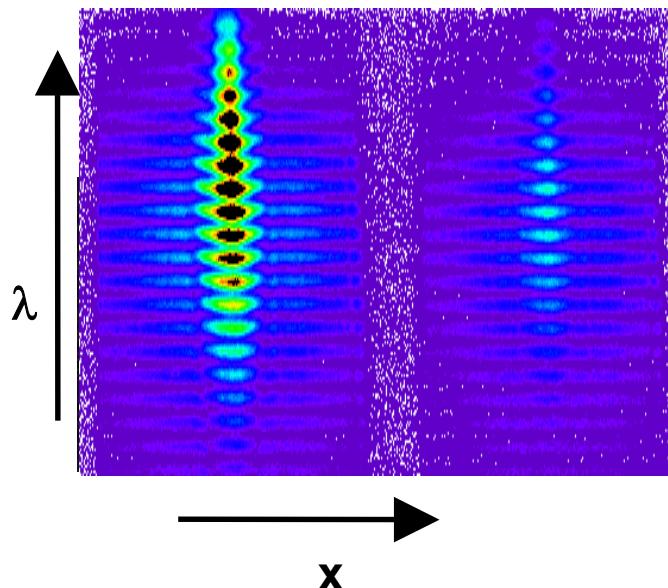
Focal spot



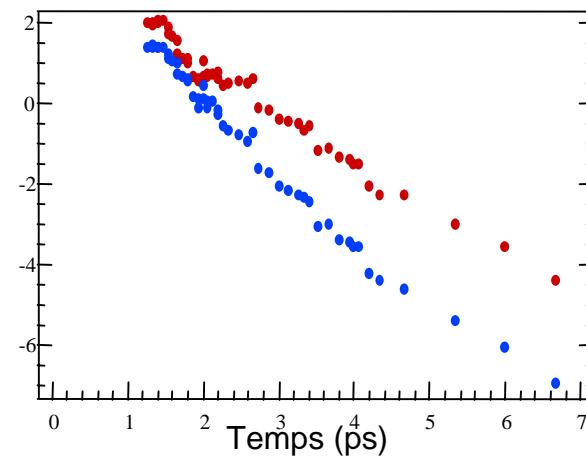
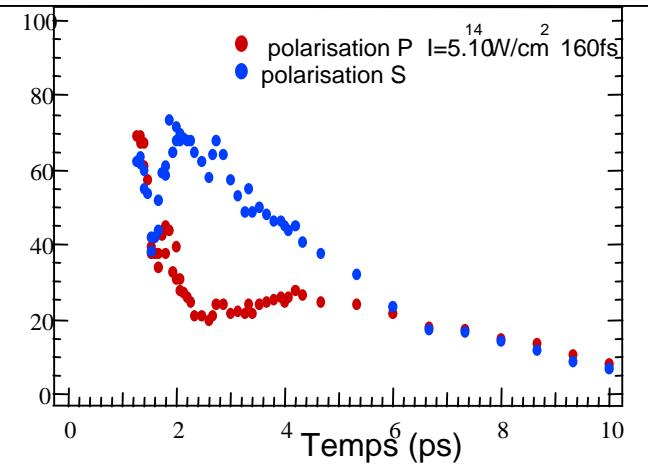
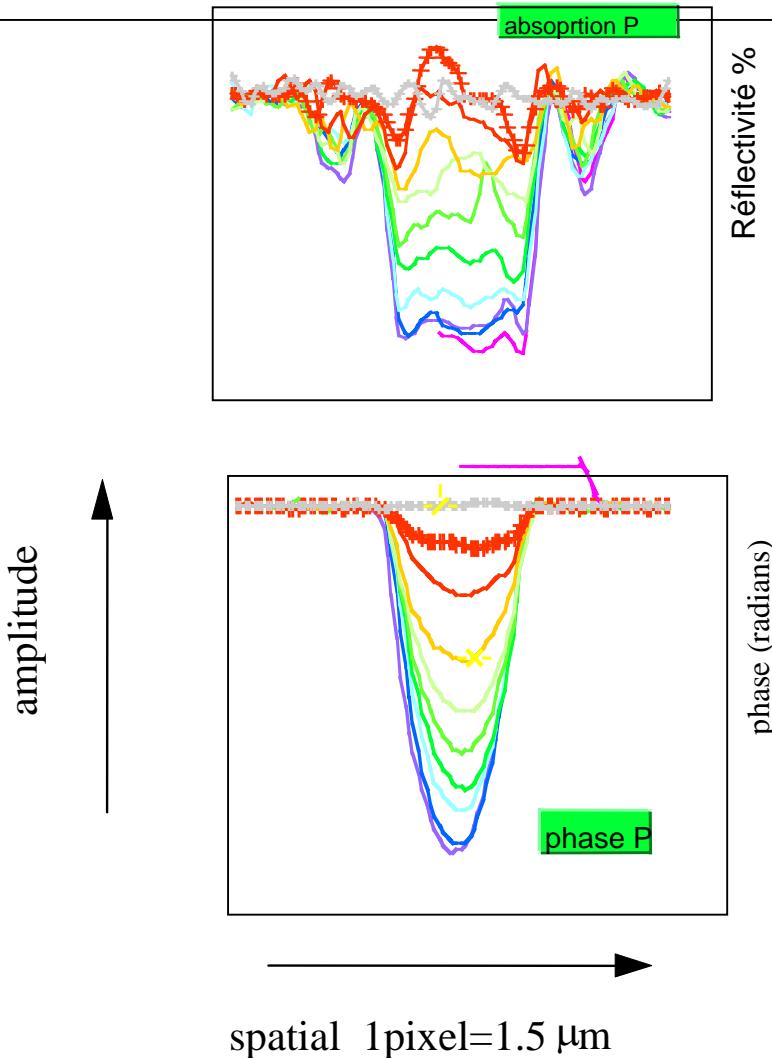
Polarization

P

S



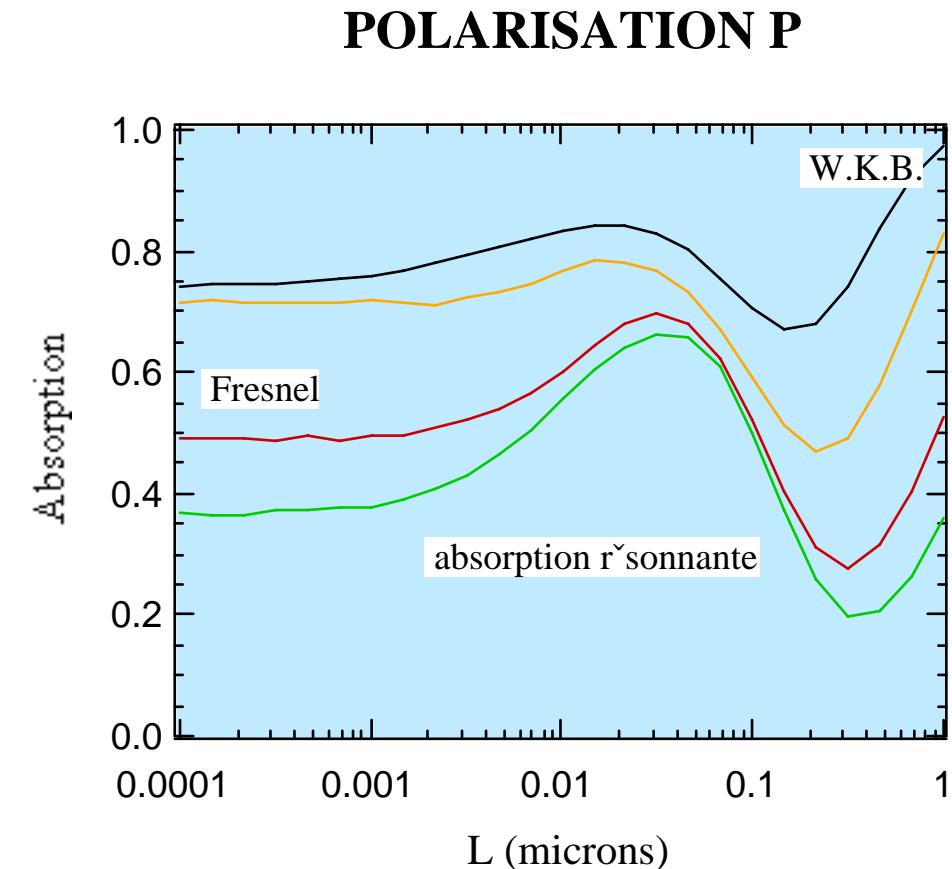
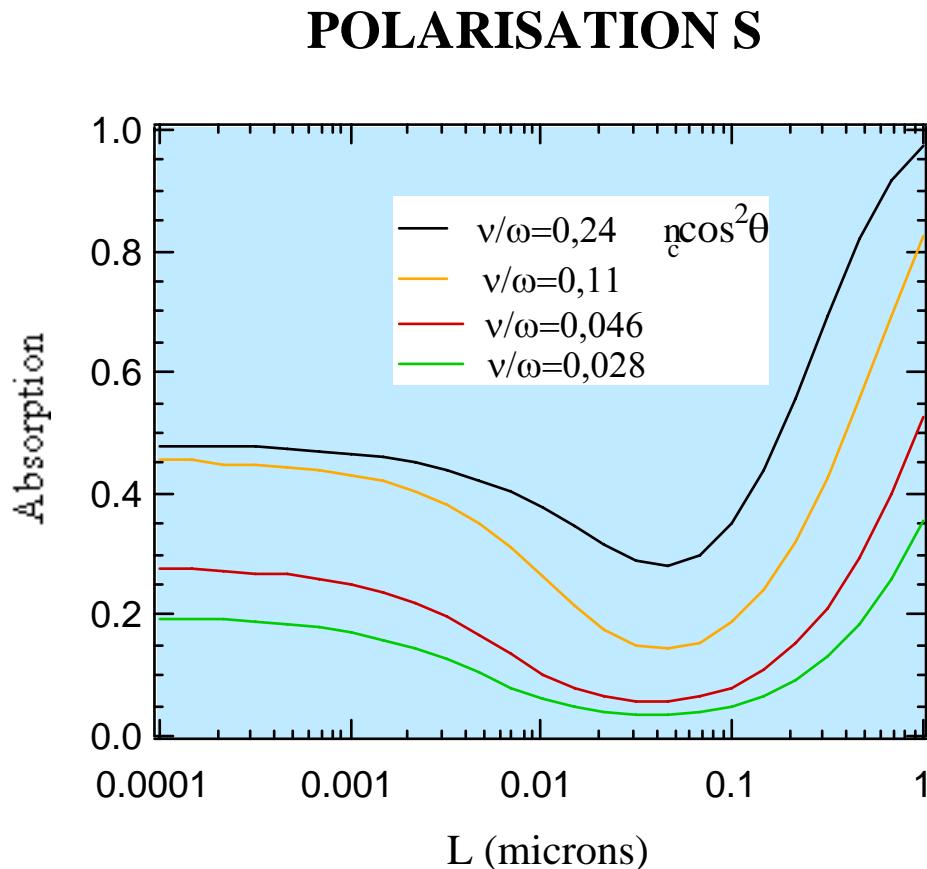
Experimentals RESULTS : time and space resolved for two polarisations



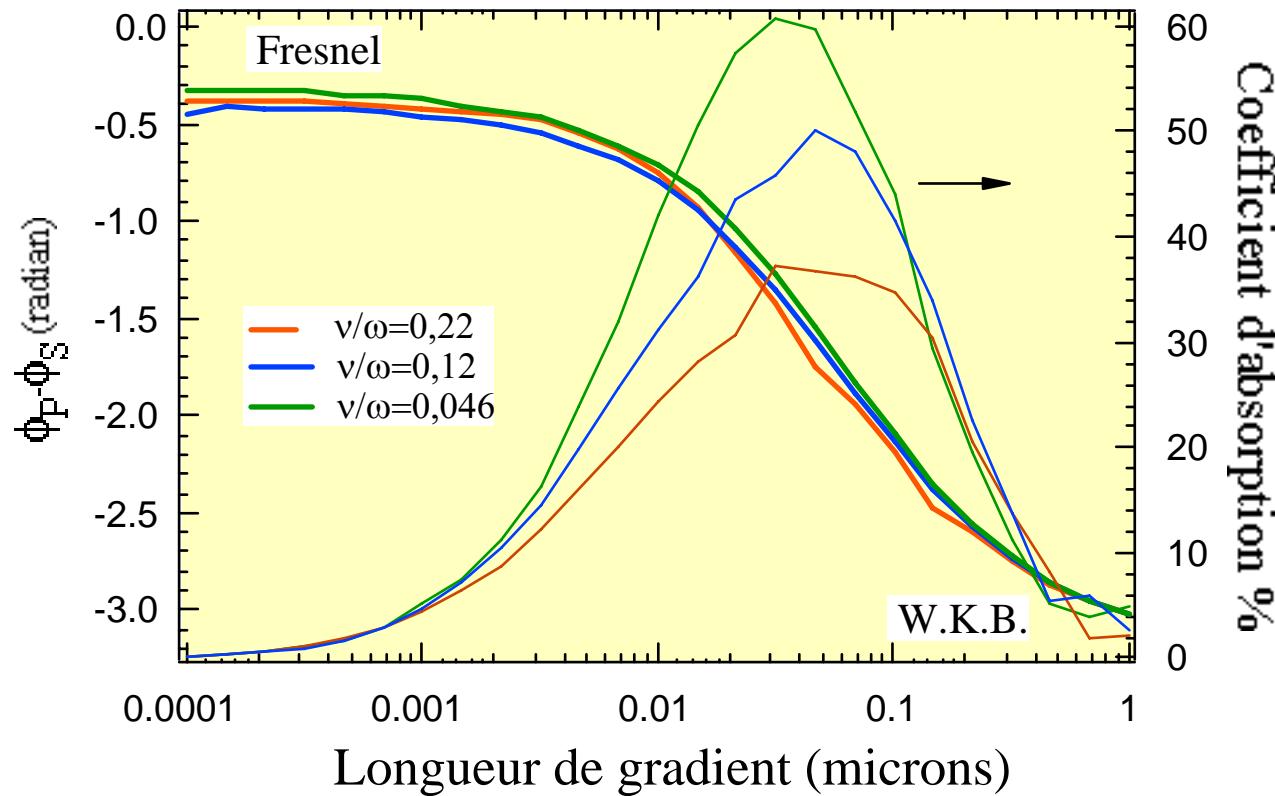
using an electromagnetic solver : Drude Model



absorption is function of: the electron density gradient L
- collision frequency ν



Gradient length measurement using polarisation phase difference



$\phi_P(t) - \phi_S(t) \Rightarrow$ density gradient L

measure nearly independant of the collision frequency

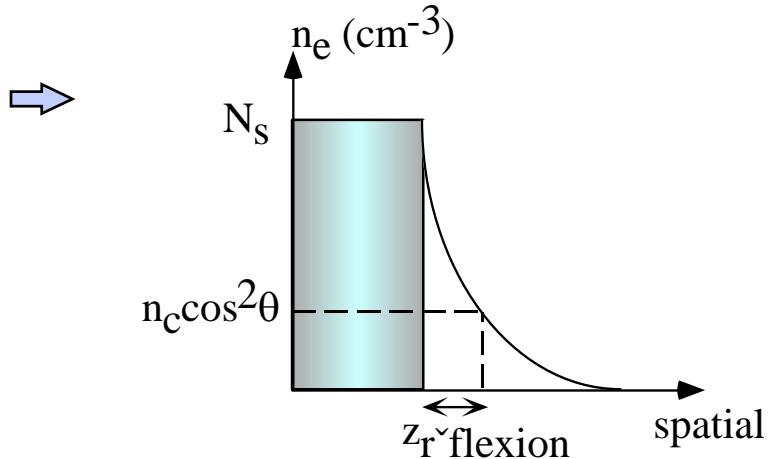
Characterization of the hydrodynamique expansion

$$\Phi_P(t) - \Phi_S(t)$$



$$L = \left(-\frac{1}{n_e} \frac{dn_e}{dz} \right)^{-1}$$

$$\Phi_S(t)$$



Density gradient scale length

Position of critical density
 $n_c \cos^2 \theta : z_{\text{réflexion}}$

S and P polarisation phase measurement as function of time
allow full characterization of the plasma

Spectral interferometry



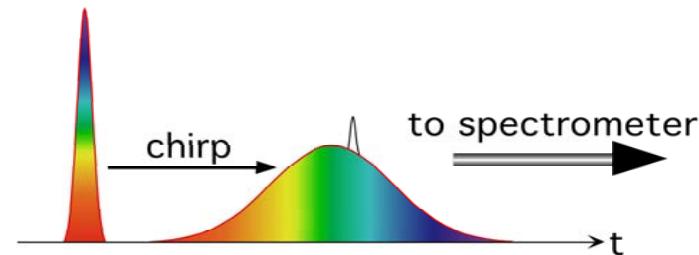
- ▶ Spectral interferometry, also called Fourier-Domain Interferometry (FDI), is a well known linear optical technique for measuring the phase and amplitude modification of a laser pulse.
- ▶ This technique is now widely used in ultrafast laser-matter interaction experiments.

- ▶ But this pump-probe technique give only one point in time for each shot

How to extend the twin pulses technique to single shot measurements?



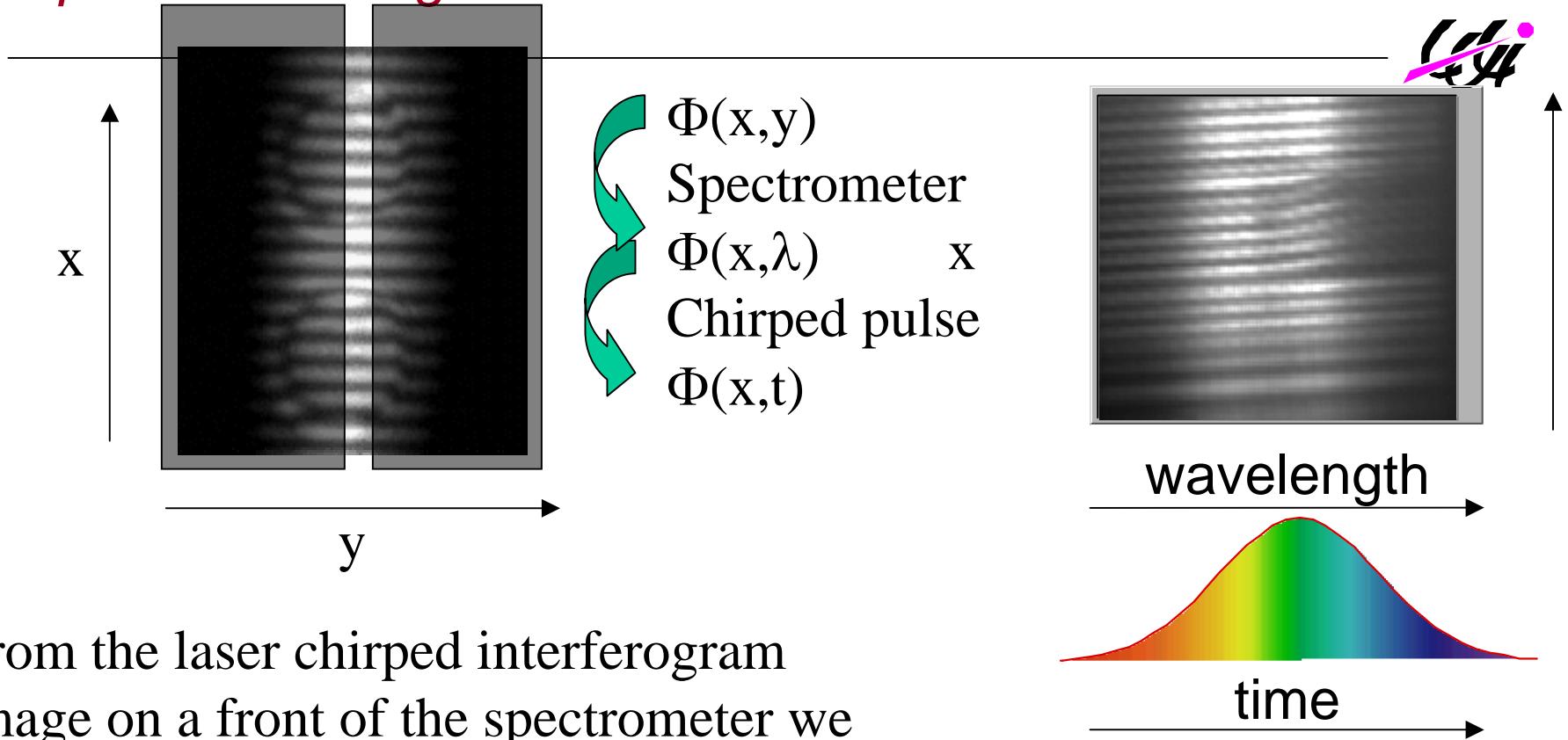
- The classical pump-probe method gives one point in time at each shot
- Single-shot measurements



FWHM Δt adapted to the temporal range

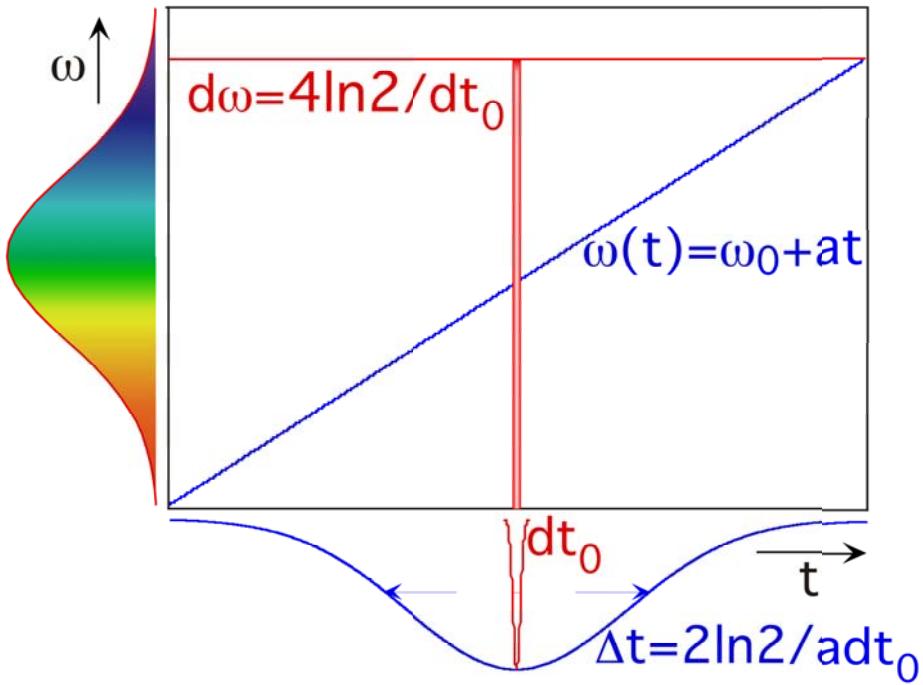
The phase and amplitude of the probe pulse are not independent any more

Single shot Fourier Domain Interferometry: spatial interferogram



From the laser chirped interferogram image on a front of the spectrometer we get information of the phase of the probe pulse reflected on the plasma as function of time and space

Direct method : Phase and reflectivity retrieval



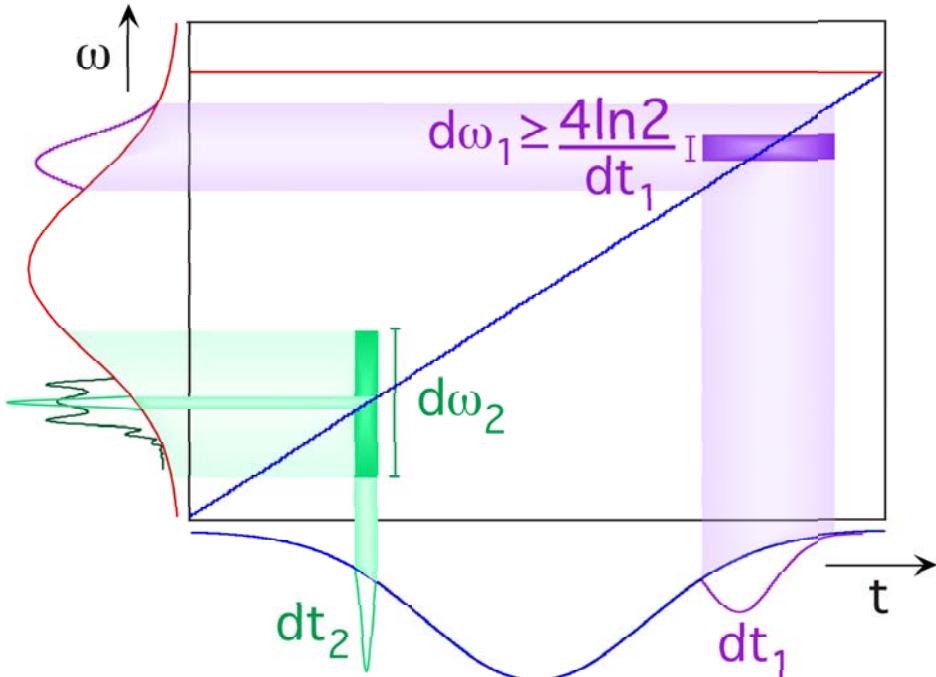
- ▶ We measure the **spectral phase and amplitude**
- ▶ **Reconstruction method when using chirped probe**
 - ❖ Time-frequency relation
- ▶ **Naïve way^{1,2}**
 - ❖ Spectral Frequency = Instantaneous Frequency

¹ A. Benazzi-Mounaix *et al.*, PRE **60**, n°3, 2488 (1999)

² C.Y. Chien *et al.*, Opt. Lett. **25**, n°8, 578 (2000)

Limits of the direct reconstruction

Need for a numerical method



► Localization criteria of the perturbation dt

- ❖ given by the Fourier limit and the time-frequency relation

$$dt / dt_0 \geq \sqrt{2} \cdot \sqrt{\Delta t / dt_0}$$



A shorter perturbation will alter the spectrum on a large zone
The phase and reflectivity temporal modulations
will mix up on the spectrum

simulations

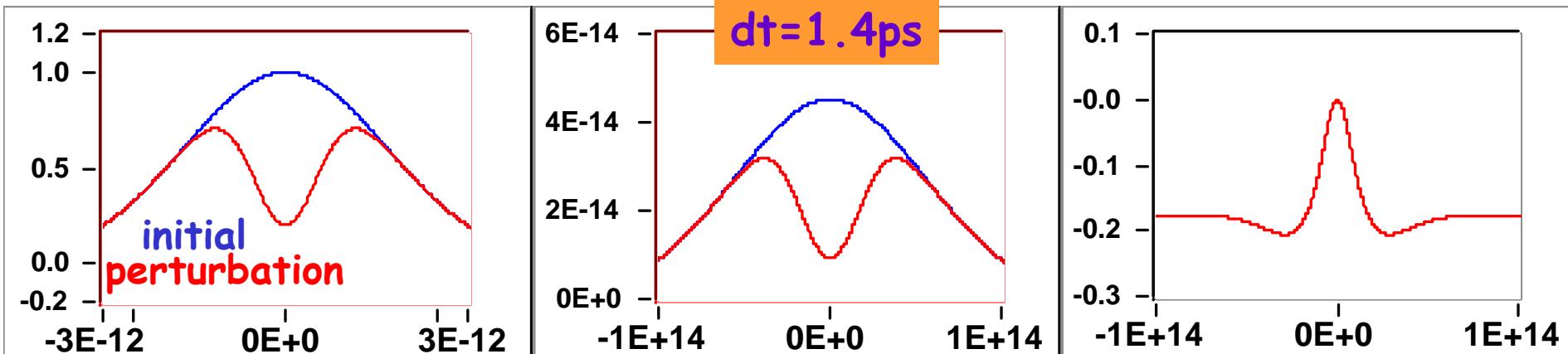
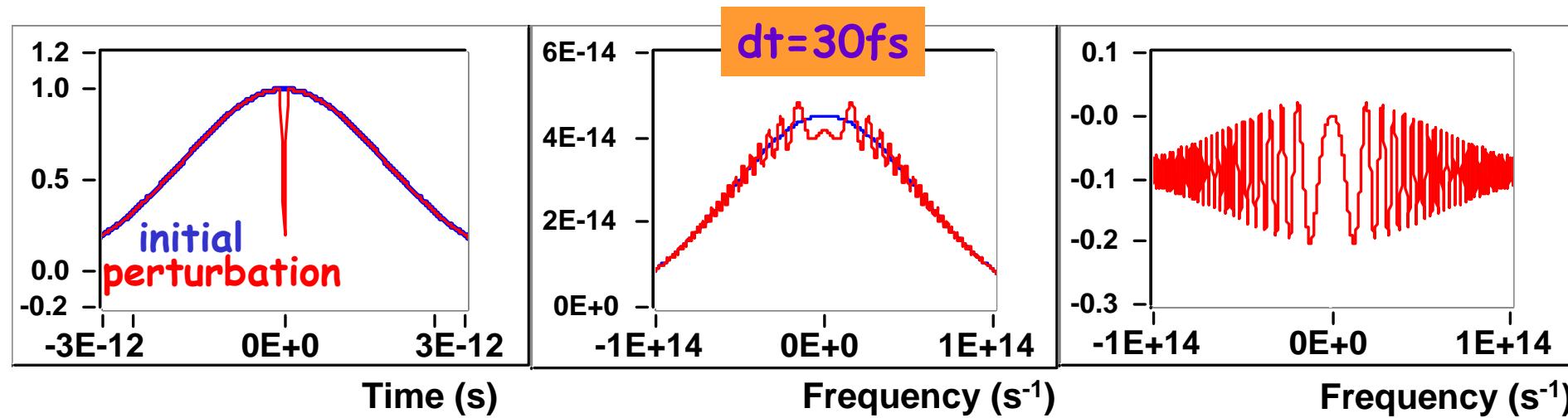
Amplitude Perturbation measure with a 30 fs probe chirp to 2.8ps



Time amplitude

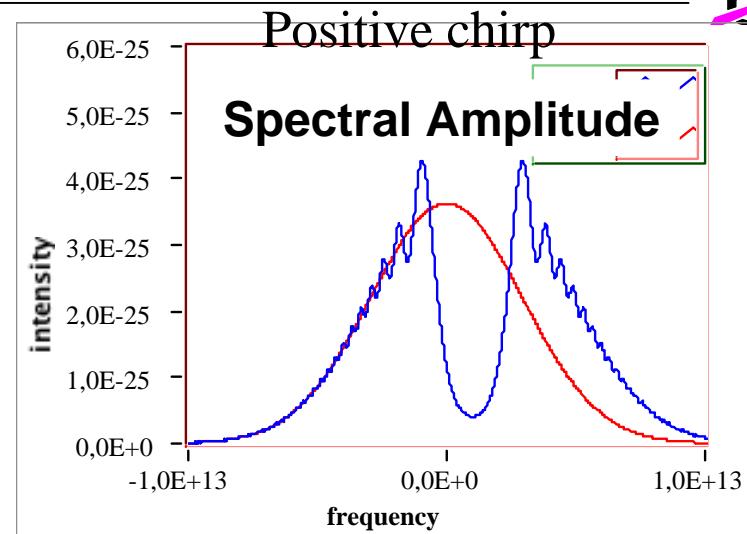
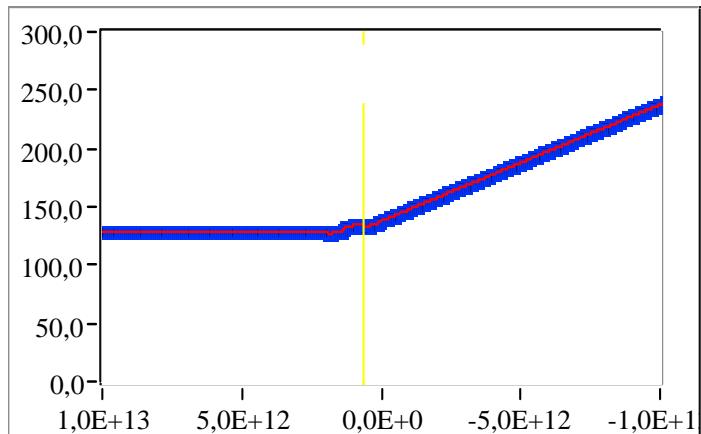
Spectral Amplitude

Phase spectrale

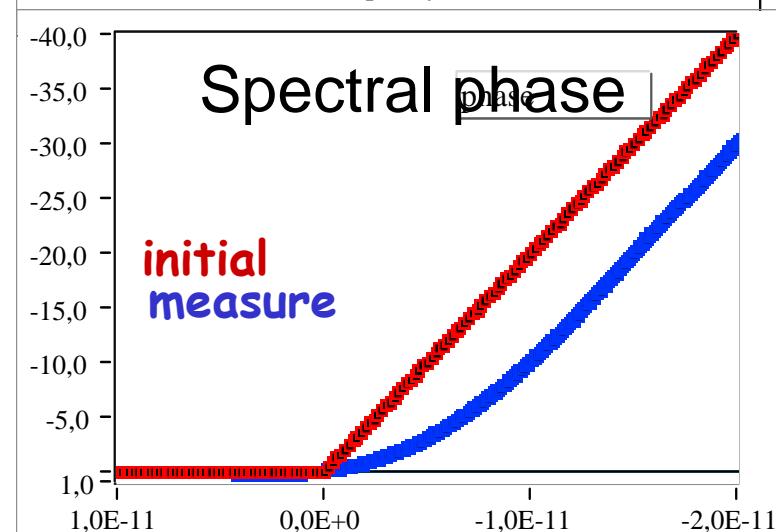


Simulations

Phase Perturbation measure with a 30 fs probe chirp to 2.8ps



When the phase perturbation is important there is a doppler effect which modify the measurement



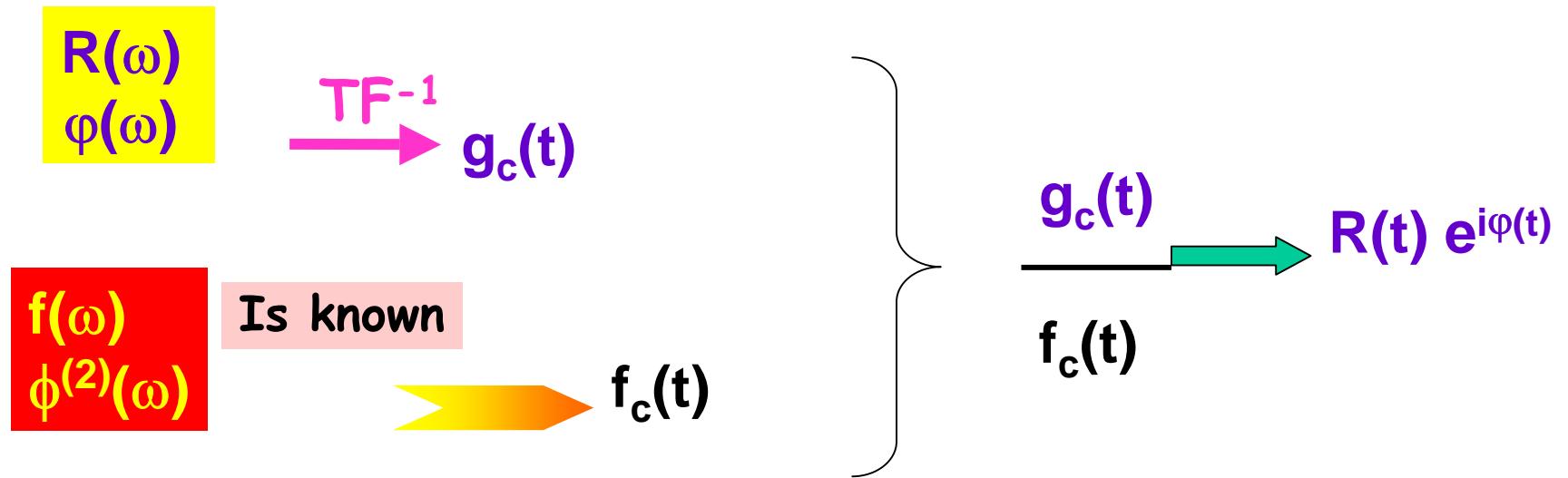
Amplitude and phase retrieval



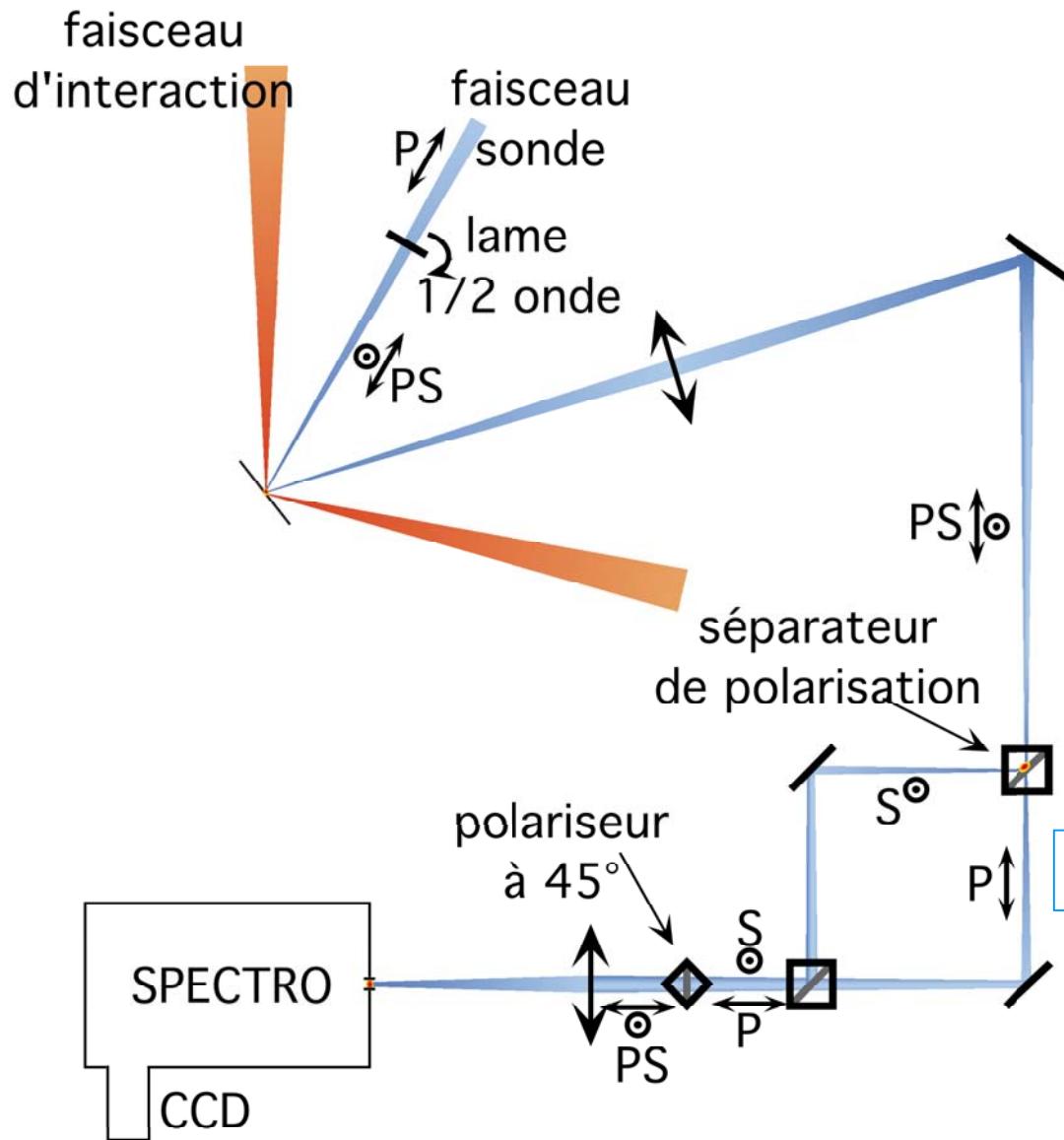
Initiale chirp probe $f_c(t) \Rightarrow f_c(\omega) = f(\omega) \cdot e^{i\phi^{(2)}\omega^2}$

Probe after the plasma $g_c(t) = f_c(t) \cdot R(t) e^{i\varphi(t)}$

- From the spectral interferometry we measure



Expérimental set up at LOA

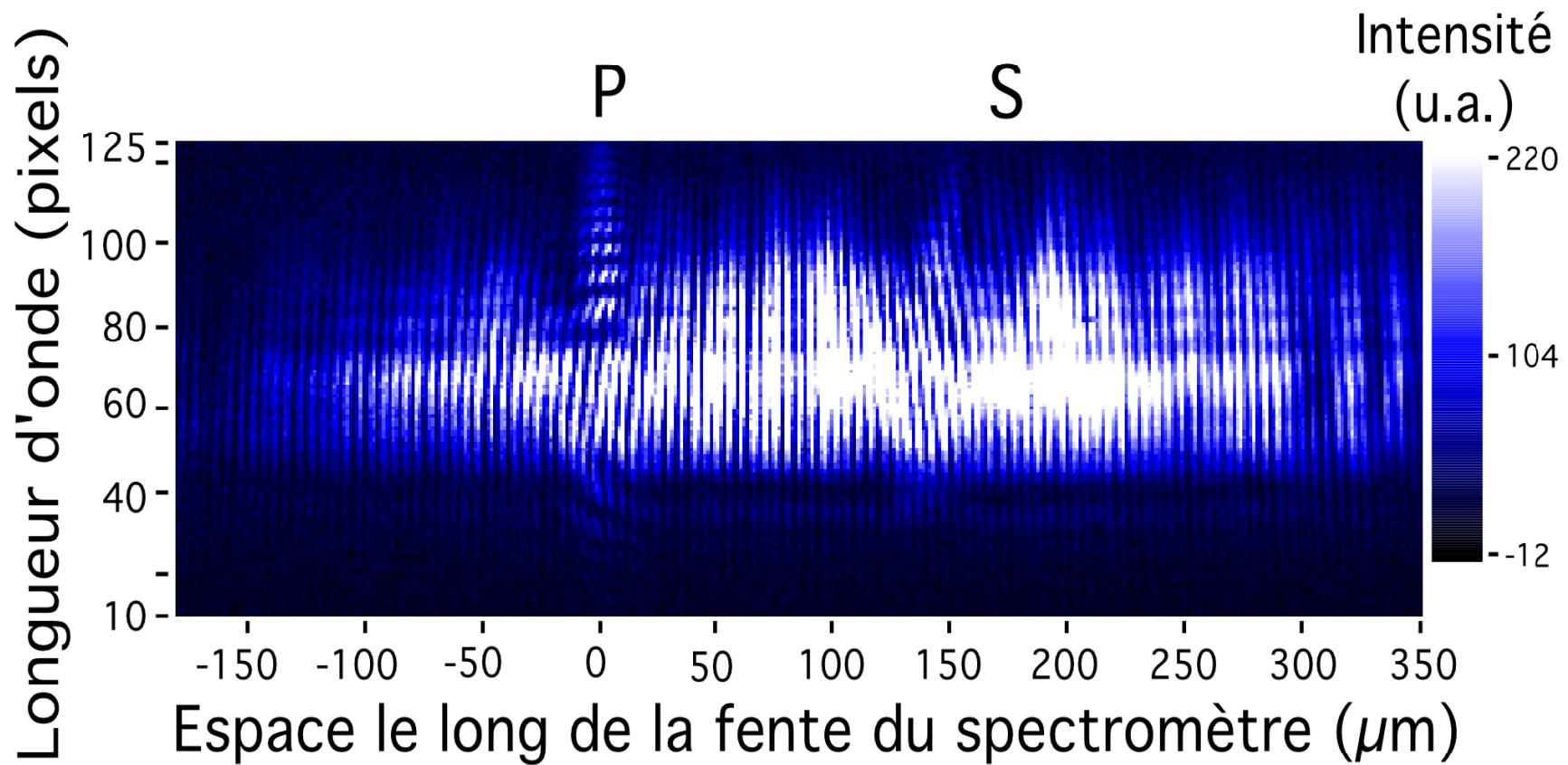


**Faisceau principal
51,5°
polarisation P**

**Faisceau sonde
22,5°
polarisé à 45°**

Mach-Zehnder

Interférogramme brut



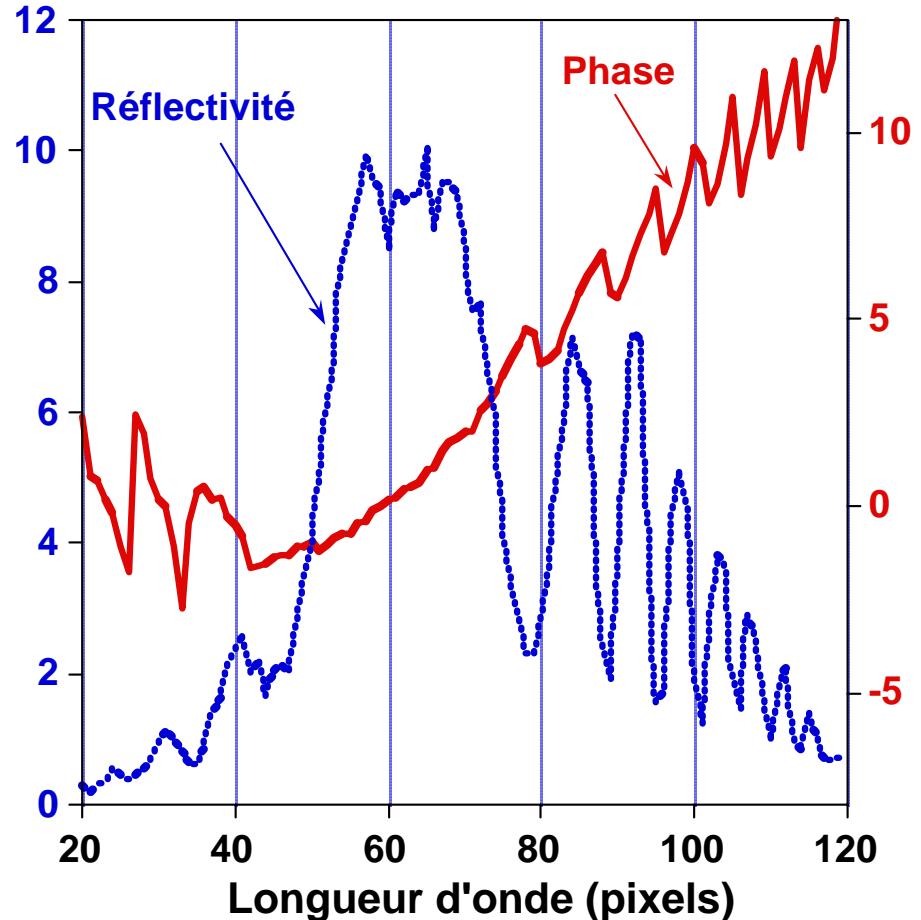
Le processus de reconstruction

Exemple

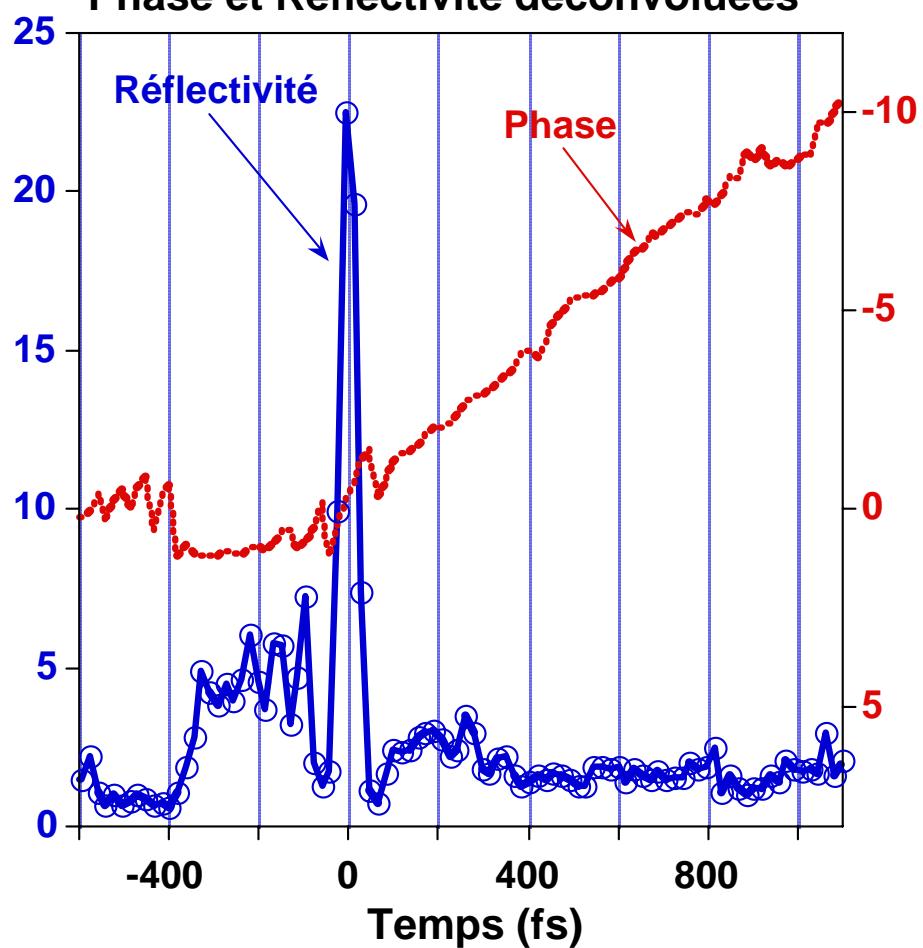
Centre de la tache focale, polarisation P



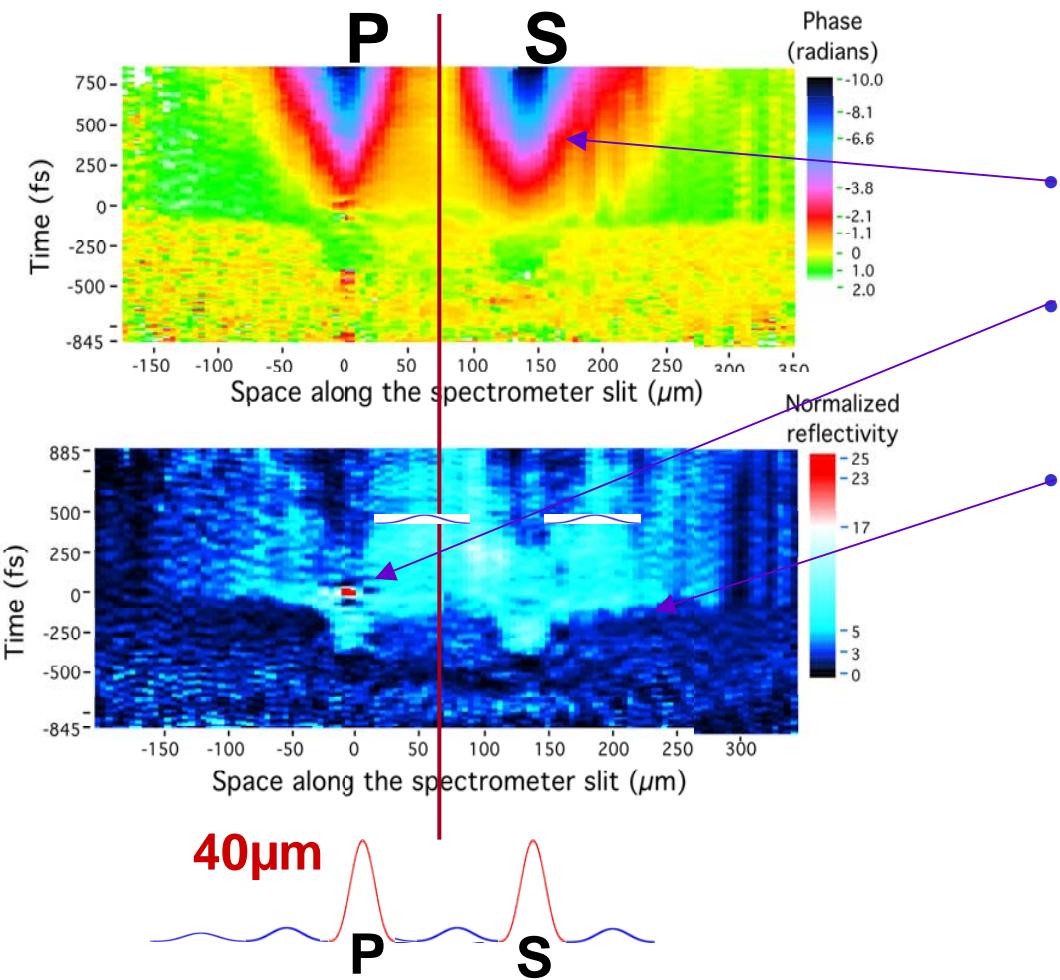
Phase et Réflectivité extraites



Phase et Réflectivité déconvolées



Phase and amplitude maps after deconvolution



Doppler phase

Hydrodynamic expansion

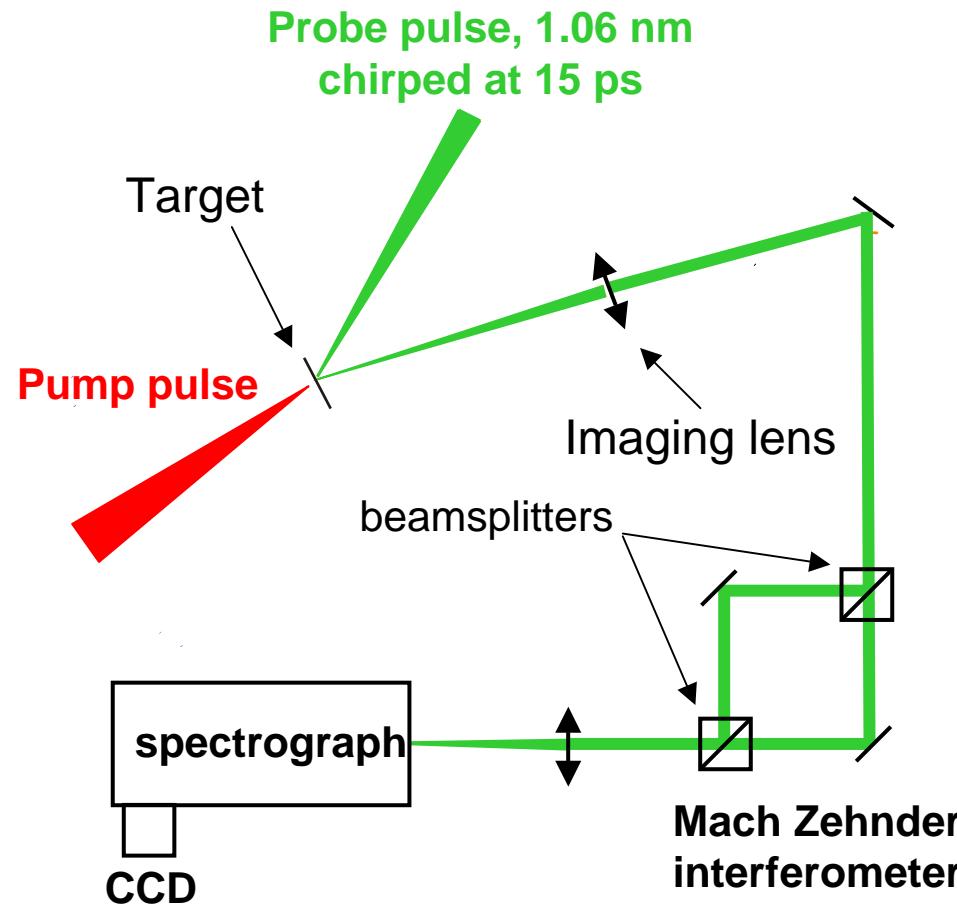
Scattering peak in P-polarization

Scattered light has the same polarization than the laser

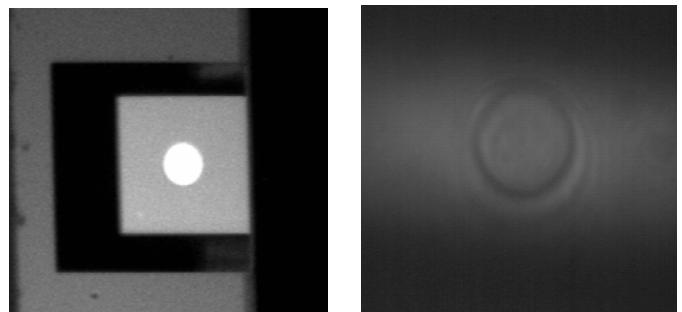
Phase jump of 1 radian outside of the focal spot, with sharp increase of the reflectivity

Single-shot Frequency Domain Interferometry

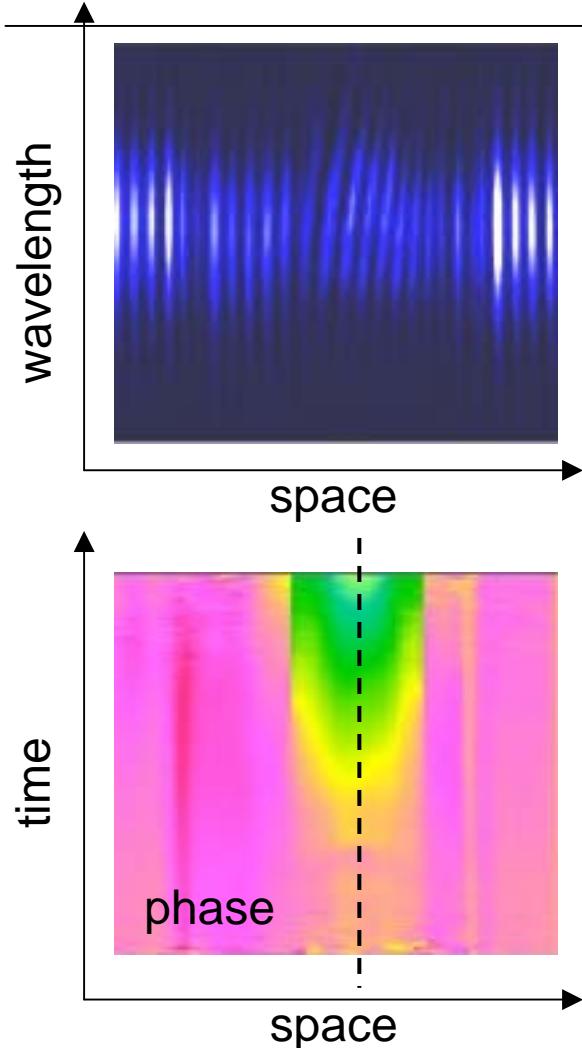
Of thin foil expansion heated by short pulse laser



500 Å Thin Aluminium
foil target

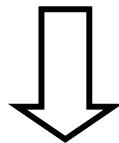


Single-shot FDI on 500 nm thin foil heated by 300fs laser at $3 \cdot 10^{15} \text{ W/cm}^2$

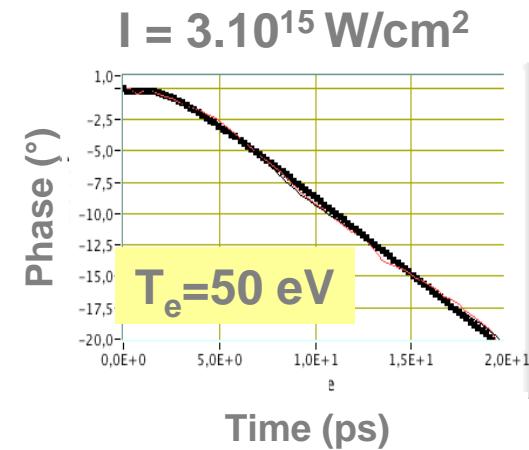


Simple model of hydrodynamic expansion

- total energy is constant
(adiabatic expansion)

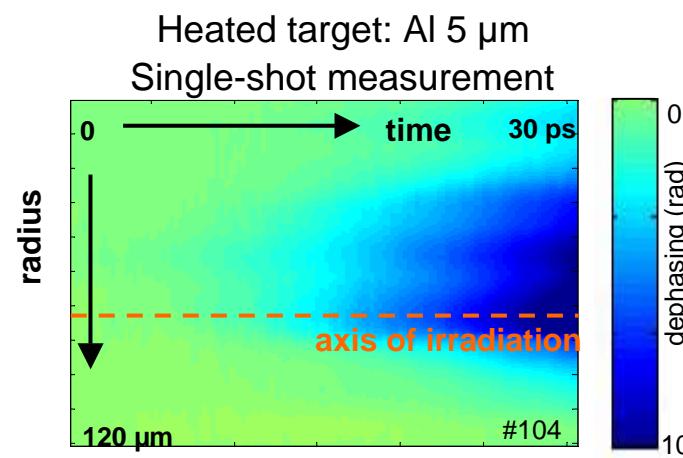
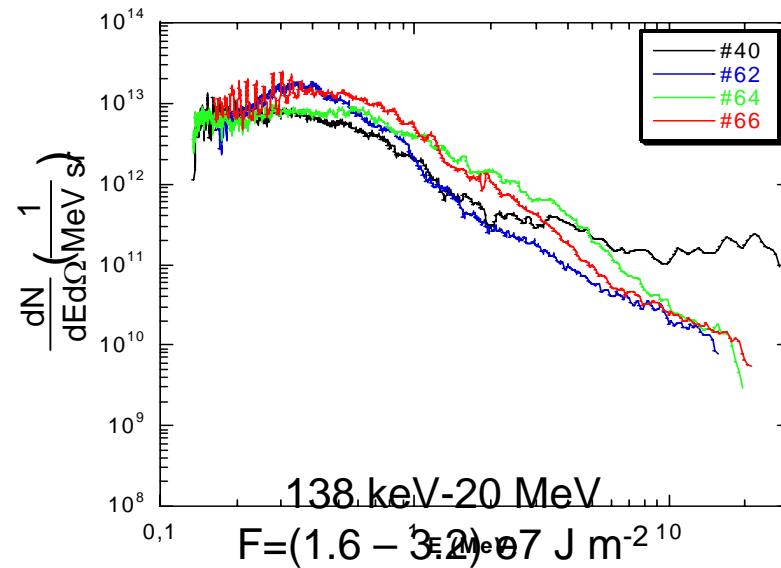
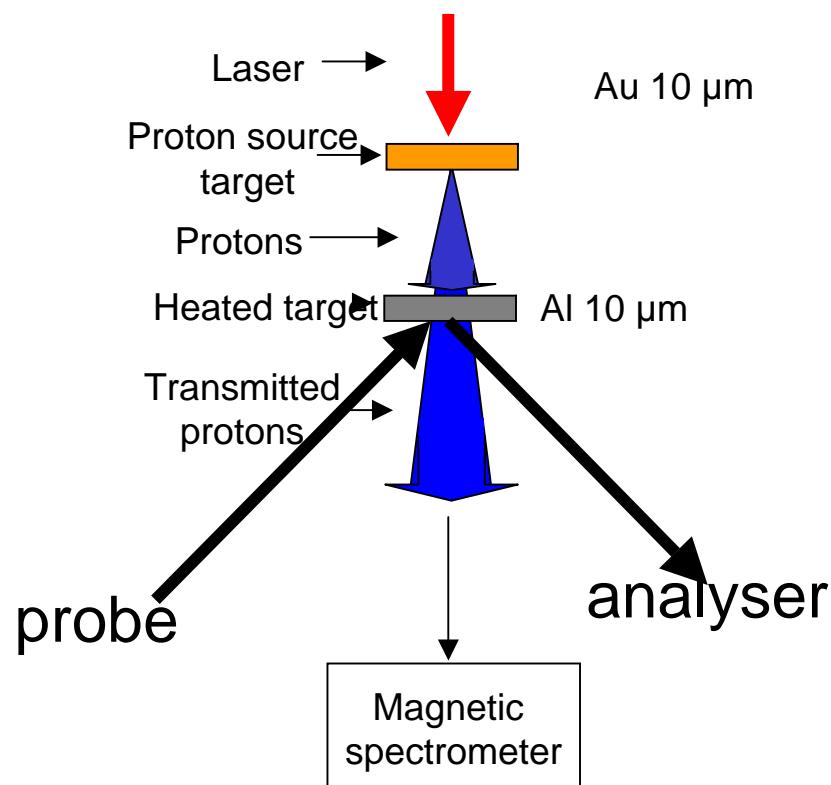


The phase variation
depends only on the
max temperature*



* Audebert et al., JQSRT 81, 19 (2003)

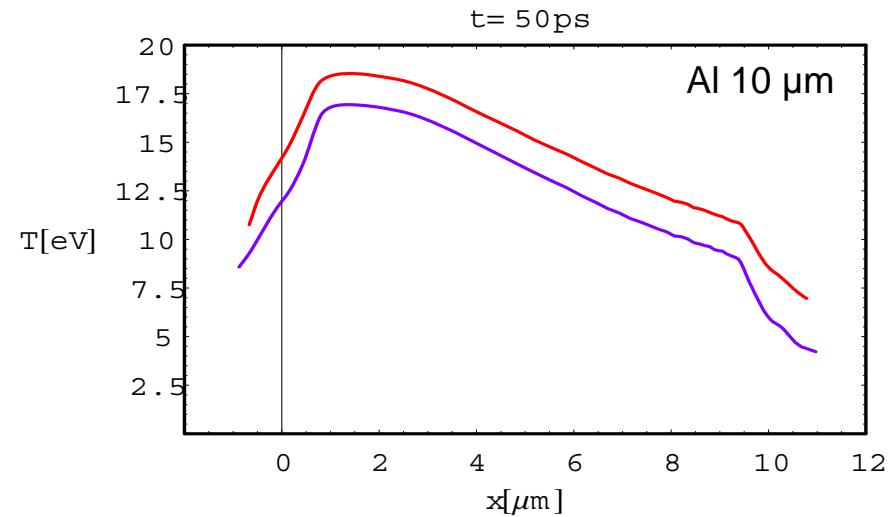
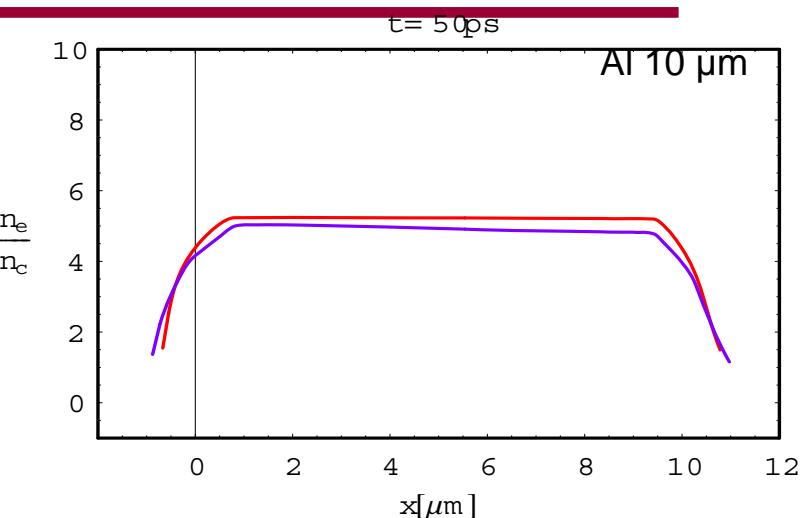
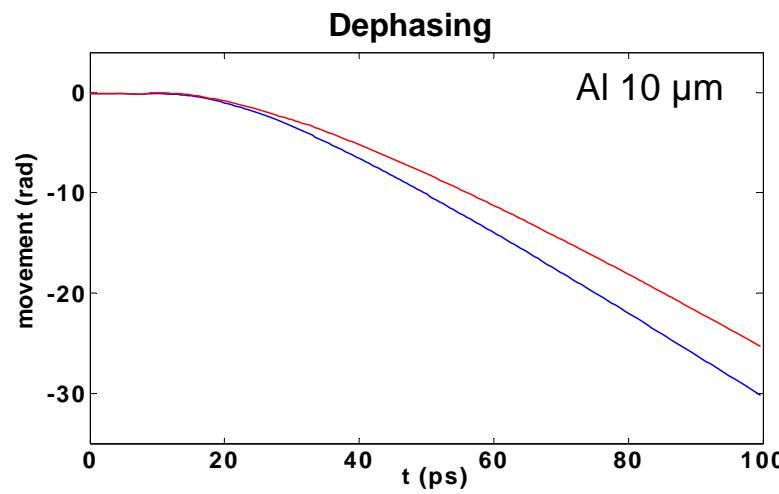
Target heated by laser accelerated protons



Simulation: with Different EOS on hydrocode



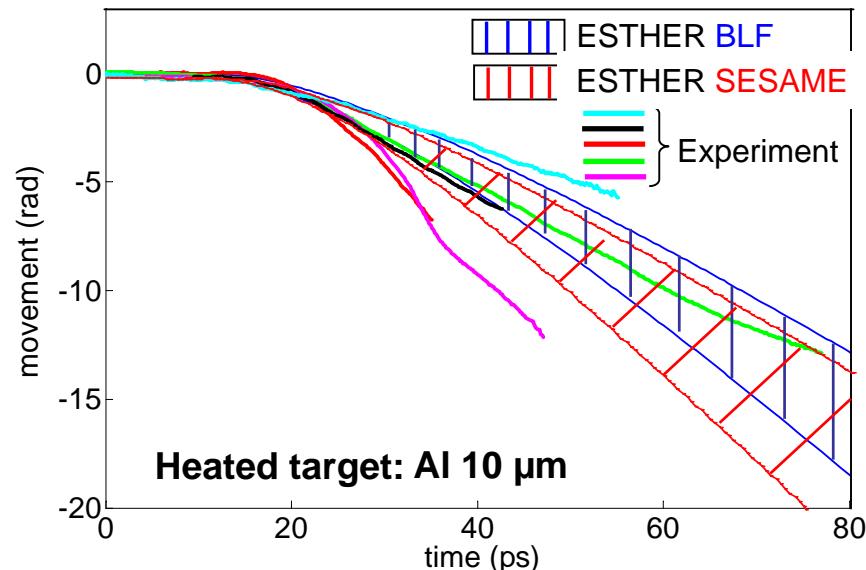
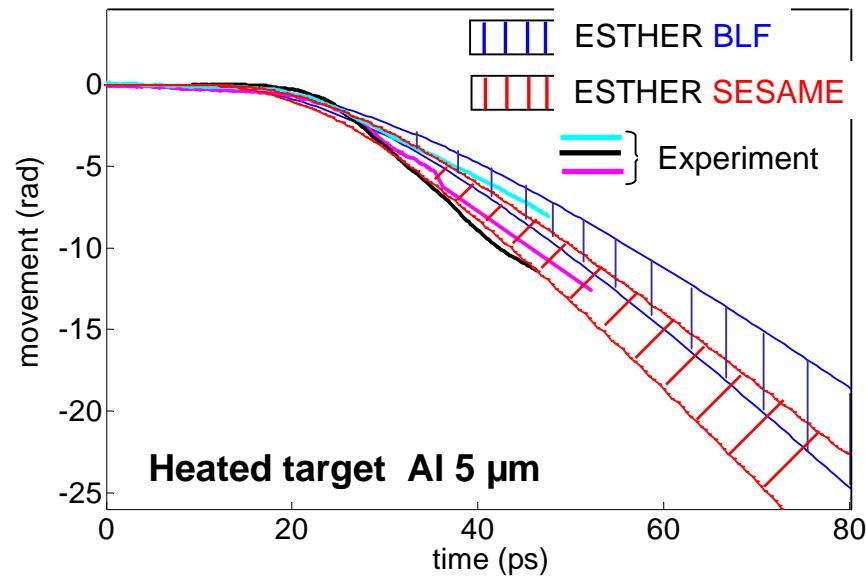
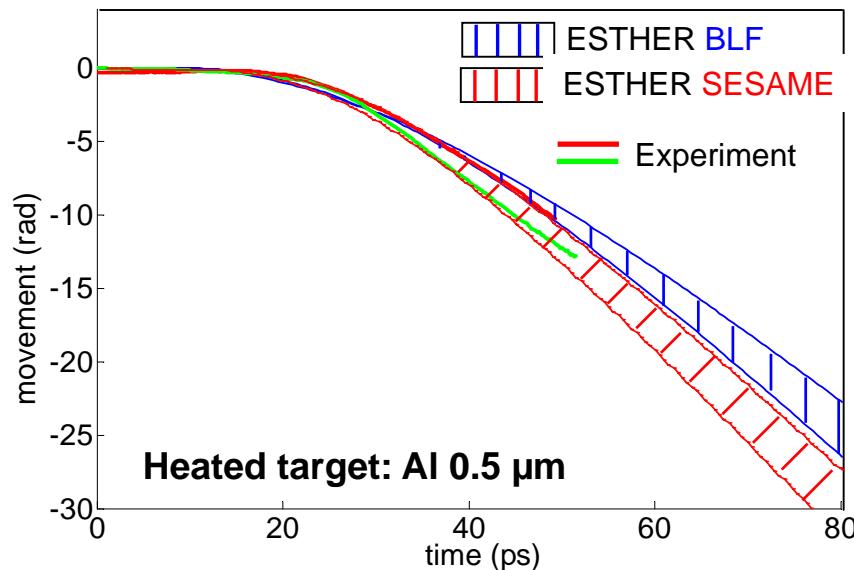
- **BLF** (Bushman-Lomonosov-Fortov)
 - more adapted for solids
 - take account of phase transitions
- **SESAME**
 - more adapted for plasma



Comparison: Experiment – ESTHER output



- A laser probe diagnostic is implemented in the code
- The interaction between the laser field and the target – Helmholtz wave equation
- Conductivity: Eberling & Palik – solid phase, Spitzer formula – plasma phase



Conclusions



- FDI is a powerful technique to measure as function of space and time ultrafast-evolving laser-produced plasmas.
- We have implemented using chirp pulse a design for single shot frequency-domain interferometry .
- This technique has been used to fully characterize different type of plasma of varying electron density gradient scale-length.