

BERGISCHE UNIVERSITÄT WUPPERTAL



The Unavoidable: Monte Carlos for the LHC Helac-Phegas



Malgorzata Worek

Bergische Universität Wuppertal

Helmholtz Alliance

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Outline of the Talk



- Introduction & Motivation
- State-of-the-Art for Multi-leg Processes
- Main Obstacles and Computational Complexity
- Leading Order level MC Programs: HELAC-PHEGAS
- Main Features & Numerical Results
- Matching to Parton Shower
- Summary & Outlook

Cross Sections @ LHC

proton - (anti)proton cross sections



• Total inelastic pp cross section

 $\sigma_{tot} \simeq 100 \, mb$

• Event rate at low/high luminosity

$$R = \frac{Nevt}{sec} = \sigma \cdot L \simeq 10^8 / s \ (10^9 / s)$$

• Interesting events are rare !

 $\sigma(pp \rightarrow W) \simeq 150 \, nb \simeq 10^{-5} \sigma_{tot}$

• We need luminosity to be high to maximise number of events

Final States @ LHC



Process	Events/s	Events/year
Jet E _T > 100 GeV	10 ³	10 ¹⁰
Jet E _T > 1 TeV	1.5x10 ⁻²	1.5 x10 ⁵
W -> lv	20	2x10 ⁸
bb	5x10 ⁵	5x10 ¹²
tt	1	10 ⁷
WW -> lv lv	6x10 ⁻³	6x10⁴

September 2009 ???

• Low luminosity:

 $L = 10^{33} \, cm^{-2} \, s^{-1} \Rightarrow 10 \, fb^{-1} / year$

- High luminosity
 - $L = 10^{34} \, cm^{-2} \, s^{-1} \Rightarrow 100 \, fb^{-1} / year$

Provide description of these events, features of new physics processes (rates, distributions, details of the final states, overall multiplicities, etc.)

MC Generators

Monte Carlo Generators

• Physics program of two main LHC experiments ATLAS and CMS:

- Discovery of Higgs boson(s)
- Search for signal of new physics beyond the SM
- Remember: **Present day signals = Tomorrow backgrounds**
- Signal events dug out from a bulk of background events
- Backgrounds due to SM processes
- Mostly **QCD** processes, accompanied by additional electroweak bosons
- Final state high number of jets or identified particles
- Reliable predictions for multi-particle final states needed !
- All this can be described by Monte Carlo generators

Benchmarking against real data turns MC simulation into powerful tool !

ttH event @ LHC

- Hard interaction big red blob
- Decay of top quarks and Higgs small red blobs
- QCD bremsstrahlung ISR & FSR
- Multiple interactions
- Hadronisation of final state partons
- Hadrons decays
- QED bremsstrahlung Photons radiations
- Event generators rely on factorisation of such event into different well-defined phases corresponding to different kinematic regimes
- Hard process calculated in fixed order perturbation theory in α_s
- QCD evaluation described by parton shower perturbation theory beyond fixed order - LL
- Hadronisation Λ_{QCD} phenomenological models with parameters to be fitted to data
- Underlying events phenomenological models beyond factorisation



Borrowed from SHERPA people

Factorisation Theorem

• Cross section for hard scattering process initiated by two hadrons

$$\sigma^{had.} = \sum_{ij} \int dx_i dx_j f_{i/h_1}(x_1, \mu_F^2) f_{j/h_2}(x_2, \mu_F^2) \hat{\sigma}^{part.}(x_1 p_1, x_2 p_2, \alpha_s(\mu_R^2), Q^2/\mu^2)$$

- Long distance (hadronic) physics factored out and absorbed into PDF
- Short distance (partonic) physics $Q \gg M_{had}$
- Factorisation is not exact but corrections are $O(M_{had}/Q)$
- Asymptotic freedom $\rightarrow \alpha_s$ small at high Q
- $\hat{\sigma}^{part}$ can be computed as a perturbation series in α_s
- More terms included in perturbative expansion weaker dependence on $\mu_F = \mu_R = \mu$
- Cross section to all orders independent of scales

$$\frac{\partial \sigma}{\partial \mu_F} = \frac{\partial \sigma}{\partial \mu_R} = 0$$

Process Calculation



• For a given process $gg \rightarrow u\overline{u}s\overline{s}gg$

• With some cuts $p_T = \sqrt{p_x^2 + p_y^2}$, $\eta = -\ln \tan(\theta/2)$, $\Delta R = \sqrt{\Delta \phi^2 + \Delta \eta^2}$

- We ask to compute cross section $\hat{\sigma}^{part.}$ @ LHC
- Steps
- Find all Feynman diagrams
- Compute them to get an amplitude
- Sum over all colour and helicity configurations
- Square the amplitude
- Integrate over the phase space

Biggest Obstacles



- Complexity of calculation based on Feynman diagrams ~ n!
- Flavours of partons never detected (b-tagging)
- For given jet configuration very many contributing subprocesses
- Neither the colour nor the spin of any parton is observed
- Amplitude with p quark and q gluons $(2 \times 3)^p (2 \times 8)^q$ contributions
- Amplitude peaks in complicated ways inside the momentum phase space
- Straightforward integration is impractical

Automated calculation of the ME based on recursive equations

MC summation over helicity and colour as well as over flavour

Search for efficient mappings - importance and stratified sampling

Number of FD



N=8	N=9	N=10	N=11	N=12	N=13
15495	231280	4016775	79603720	1773172275	43864374400

QCD with 3 (identical) fermion pairs

N=8 N=9 N=10 N=11	N=12	N=13
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roughly grows like



Tools @ LHC



- Standard Model and beyond tools @ tree level
- Parton-level tools which are completely self-contained and automated and provide amplitudes and integrators on their own

AlpGen, AMEGIC++/Sherpa, Helac-Phegas, MadGraph/MadEvent, O'Mega/WHIZARD, ...

 General purpose Monte Carlo programs (parton shower, hadronisation, multiple interactions, hadrons decays, ...)

Herwig, Herwig++, Pythia 6.4, Pythia 8, Sherpa, ...

Generators are not perfect ! Shop around and compare several approaches before drawing conclusions. Blind usage of a generator is not encouraged !

T. Sjöstrand



Features of a MC generator





HELAC-PHEGAS

http://helac-phegas.web.cern.ch/helac-phegas/

A. Cafarella, C. G. Papadopoulos, M. Worek e-Print: arXiv:0710.2427 [hep-ph]





Extensibility



- Standalone version with build-in CTEQ6L1
- Interface to the LHAPDF





 $p_{T~i} > 60~GeV, ~~\theta_{ij} > 30^o ~~|\eta_i| < 3$

# jets	3	4	5	6	7	8
$\sigma(nb)$	91.41	6.54	0.458	$2.97 imes 10^{-2}$	$2.21 imes 10^{-3}$	$2.12 imes10^{-4}$
% Gluon	45.7	39.2	35.7	35.1	33.8	26.6







• Automatic multi-channel phase-space mapping

- Self-adapting procedure to reshape the generated phase space density
- For example: per mille level tt + 0,1,2 jets (6,7,8 final states) with full off-shell and finite width effects





- Straightforward inclusion of new physics effects
- New models: for example MSSM
- New couplings: for example effective Hγγ and Hgg couplings





- **3**ⁿ complexity due to the Dyson-Schwinger recursive algorithm
- Monte Carlo summation over:
 - helicity configurations (completed)
 - color configurations (completed)
 - subprocesses
- Trivial parallelization over clusters (LSF at CERN)





- Arbitrary cuts, distributions and scale choices
- Configuration via scripts
- Compilers: Lahey Fujitsu, Intel Fortran, GNU gfortran/g95
- Multiprecision numerics



Dyson-Schwinger Recursion

- Alternative to Feynman Diagrams representation
- Express n-point Green's functions in terms of 1-, 2-,... (n-1)-point functions
- Diagrammatic representation, e.g. QED-like theory
- Interaction of a spinor field to a gauge boson

$$\boldsymbol{b}^{\mu}(\boldsymbol{P}) = \sum_{n=1}^{n} \delta(\boldsymbol{P} = \boldsymbol{P}_{i}) \boldsymbol{b}^{\nu}(\boldsymbol{P}_{i}) + \sum_{\boldsymbol{P} = \boldsymbol{P}_{i} + \boldsymbol{P}_{2}} (\boldsymbol{i}\boldsymbol{g}) \Pi^{\mu}_{\nu} \overline{\psi}(\boldsymbol{P}_{2}) \boldsymbol{\gamma}^{\nu} \psi(\boldsymbol{P}_{1}) \epsilon(\boldsymbol{P}_{1}, \boldsymbol{P}_{2})$$

$$b_{\mu}(P) = \longrightarrow \qquad \psi(P) = -$$
 $\bar{\psi}(P) = -$

- Sub-amplitude with off-shell boson of momentum P
- Blobs denote sub-amplitude with the same structure

Dyson-Schwinger Recursion

• Fermion with momentum P

• Antifermion with momentum P

Building Amplitude



- Off-shell fields building blocks of any process
- Used iteratively, at each step two (three) momenta are combined
- Initial conditions for the external particles:

$$b^{\mu}(p_i) = \epsilon^{\mu}_{\lambda}(p_i), \lambda = \pm 1, 0$$

$$\psi(p_i) = \begin{cases} u_{\lambda}(p_i) & \text{if } p_i^0 \ge 0\\ v_{\lambda}(-p_i) & \text{if } p_i^0 \le 0 \end{cases}$$

$$\bar{\psi}(p_i) = \begin{cases} \bar{u}_{\lambda}(p_i) & \text{if } p_i^0 \ge 0\\ \bar{v}_{\lambda}(-p_i) & \text{if } p_i^0 \le 0 \end{cases}$$

• Amplitude can be calculated by any of the following relations:

$$\mathcal{A}(p_1, \dots, p_n) = \begin{cases} b_0^{\mu}(P_i)b_{\mu}(p_i)\\ \bar{\psi}_0(P_i)\psi(p_i)\\ \bar{\psi}(p_i)\psi_0(P_i) \end{cases}$$



- DS equations both for: full and colour ordered amplitudes
- Colour ordered ordinary approach SU(N) algebra
- Quarks and gluons treated differently

$$M(\{p_i\}_1^n; \{\varepsilon_i\}_1^n; \{a_i\}_1^n) = \sum_{I \in P(2,...,n)} Tr(t^{a_{\sigma_i(1)}}t^{a_{\sigma_i(2)}}...t^{a_{\sigma_i(n)}}) A_I(\{p_i\}_1^n; \{\varepsilon_i\}_1^n)$$

$$\sum_{[\varepsilon_i]_1^n; [a_i]_1^n} |M(\{p_i\}_1^n; \{\varepsilon_i\}_1^n; \{a_i\}_1^n)|^2 = \sum_{\varepsilon} \sum_{IJ} A_I C_{IJ} A_J^*$$

$$C_{IJ} = \sum_{IJ} Tr(t^{a_1} t^{a_{\sigma_i(2)}} \dots t^{a_{\sigma_i(n)}}) Tr(t^{a_1} t^{a_{\sigma_i(2)}} \dots t^{a_{\sigma_i(n)}})^*$$

- Colour ordered U(N) type colour algebra, gluon = qq pair
- Each colour amplitude proportional to D_{I}
- σ_{I} I-th permutation of the set 1,2,...,n

$$M(\{p_i\}_1^n; \{\varepsilon_i\}_1^n; \{a_i\}_1^n) = \sum_{I \in P(2,...,n)} D_I A_I(\{p_i\}_1^n; \{\varepsilon_i\}_1^n)$$

$$D_I = \delta_{1\sigma_I(1)} \delta_{2\sigma_I(2)} \dots \delta_{n\sigma_I(n)}$$

$$C_{IJ} = \sum_{IJ} D_I D_J = N_C^{\alpha}, \quad \alpha = \langle \sigma_{1,} \sigma_2 \rangle$$

• Exact colour treatment, efficient for low colour charge

P. Draggiotis, R. Kleiss, C. G. Papadopoulos Phys. Lett. B439, 157 (1998)

• Simplification for gluon fields

$$\boldsymbol{G}_{AB} \equiv \sum_{a=1}^{8} \boldsymbol{t}_{AB}^{a} \boldsymbol{G}^{a}$$

$$A, B = 1, 2, 3$$

- New objects traceless 3x3 matrices in colour space
- Diagonalization of the colour structure of 3-gluon vertex

$$f^{abc} t^{a}_{AB} t^{b}_{CD} t^{c}_{EF} = - \frac{i}{4} (\delta_{AD} \delta_{CF} \delta_{EB} - \delta_{AF} \delta_{CB} \delta_{ED})$$

• 3-gluon vertex in a new representation

- Shows the colour flow in the real physical process
- Gluon represented by $q\bar{q}$ states in colour space
- Colour remains unchanged on an interrupted colour line

- Quarks and antiquarks already in this representation
- Additionally representation independent identity is used

$$\sum_{a=1}^{8} t_{ij}^{a} t_{kl}^{a} = \frac{1}{2} (\delta_{il} \delta_{kj} - \frac{1}{3} \delta_{ij} \delta_{kl}) \qquad i, j, k, l = 1, 2, 3$$

• Recursion equations modified to reflect the new colour structure

Next step - make the computation of the colour part of an amplitude more efficient !

LCA

- Leading Colour Approximation
- In the limit $N_c \rightarrow \infty$ only diagonal terms survive in colour matrix
- Interference between different colour flows vanish in this limit

Process	$\sigma_{\rm \tiny LCA} \pm \varepsilon ~({\rm nb})$	ε (%)	$\sigma_{\rm lca}/\sigma_{\rm exact}$
$\begin{array}{c} gg \rightarrow 2g \\ gg \rightarrow 3g \\ gg \rightarrow 4g \\ gg \rightarrow 5g \\ gg \rightarrow 6g \end{array}$	$\begin{array}{c} (0.46060 \pm 0.00308) \times 10^{4} \\ (0.15040 \pm 0.00159) \times 10^{3} \\ (0.12613 \pm 0.00187) \times 10^{2} \\ (0.09806 \pm 0.00196) \times 10^{1} \\ (0.69370 \pm 0.01736) \times 10^{-1} \end{array}$	$0.7 \\ 1.0 \\ 1.5 \\ 2.0 \\ 2.5$	1.0 1.0 1.1 1.0 0.9
$\begin{array}{l} gg \rightarrow u \bar{u} \\ gg \rightarrow g u \bar{u} \\ gg \rightarrow 2 g u \bar{u} \\ gg \rightarrow 3 g u \bar{u} \\ gg \rightarrow 4 g u \bar{u} \end{array}$	$\begin{array}{c} (0.44170 \pm 0.00239) \times 10^2 \\ (0.40539 \pm 0.00456) \times 10^1 \\ (0.56107 \pm 0.01008) \times 10^0 \\ (0.60952 \pm 0.01771) \times 10^{-1} \\ (0.56299 \pm 0.01552) \times 10^{-2} \end{array}$	$0.5 \\ 1.1 \\ 1.8 \\ 2.9 \\ 2.7$	1.2 1.1 1.1 1.2 1.1
$\begin{array}{l} gg \rightarrow c \overline{c} c \overline{c} \\ gg \rightarrow g c \overline{c} c \overline{c} \\ gg \rightarrow 2 g c \overline{c} c \overline{c} \\ gg \rightarrow 2 g c \overline{c} c \overline{c} \end{array}$	$\begin{array}{c} (0.31180 \pm \ 0.00612) \times 10^{-2} \\ (0.73387 \pm \ 0.02022) \times 10^{-3} \\ (0.11462 \pm \ 0.00377) \times 10^{-3} \end{array}$	2.0 2.7 3.3	$1.2 \\ 1.2 \\ 1.2$

$$\sum_{a,\varepsilon} |\mathcal{M}(\{p_i\}_1^n, \{\varepsilon_i\}_1^n, \{a_i\}_1^n)|^2 = g^{2n-4} N_c^{n-2} (N_c^2 - 1) \sum_{\varepsilon} \sum_I |\mathcal{A}_I|^2.$$

MC over Colour

- Incoherent sum over colour performed via MC over "real" colour (completed, available in the next version)
- All possible configurations $N_{CC}^{ALL} = N_{C}^{q+q}$
- MC one particular colour-anticolour configuration randomly selected
- Number of colours and anticolours of each type the same
- $|\mathbf{M}|^2$ multiplied by the number of non zero colour configurations \mathbf{N}_{cc}

Process	$N_{\rm cc}^{\rm all}$	N_{cc}	$\rm N_{cc}/\rm N_{cc}^{all}$	N_{cc}^{F} (%)
$\begin{array}{l} gg \rightarrow u \bar{u} \\ gg \rightarrow g u \bar{u} \\ gg \rightarrow 2g u \bar{u} \\ gg \rightarrow 3g u \bar{u} \\ gg \rightarrow 4g u \bar{u} \\ gg \rightarrow 5g u \bar{u} \\ gg \rightarrow 6g u \bar{u} \end{array}$	729 6561 59049 531441 4782969 43046721 387420489	93 639 4653 35169 272835 2157759 17319837	$\begin{array}{c} 0.1276 \\ 0.0974 \\ 0.0788 \\ 0.0662 \\ 0.0570 \\ 0.0501 \\ 0.0447 \end{array}$	93.5 91.6 92.6 94.6 96.4 97.8 98.6
$\begin{array}{l} gg \rightarrow c \bar{c} c \bar{c} \\ gg \rightarrow g c \bar{c} c \bar{c} \\ gg \rightarrow 2g c \bar{c} c \bar{c} \\ gg \rightarrow 3g c \bar{c} c \bar{c} \\ gg \rightarrow 4g c \bar{c} c \bar{c} \end{array}$	6561 59049 531441 4782969 43046721	639 4653 35169 272835 2157759	$\begin{array}{c} 0.0974 \\ 0.0788 \\ 0.0662 \\ 0.0570 \\ 0.0501 \end{array}$	99.1 98.8 99.0 99.3 99.6

Process	$N_{\rm cc}^{\rm All}$	N_{cc}	$\rm N_{cc}/\rm N_{cc}^{All}$	N_{cc}^{F} (%)
$\begin{array}{c} gg \rightarrow 2g \\ gg \rightarrow 3g \\ gg \rightarrow 4g \\ gg \rightarrow 5g \\ gg \rightarrow 6g \\ gg \rightarrow 7g \\ gg \rightarrow 8g \end{array}$	6561 59049 531441 4782969 43046721 387420489 3486784401	639 4653 35169 272835 2157759 17319837 140668065	$\begin{array}{c} 0.0974 \\ 0.0788 \\ 0.0662 \\ 0.0570 \\ 0.0501 \\ 0.0447 \\ 0.0403 \end{array}$	$59.1 \\68.4 \\77.4 \\85.0 \\90.4 \\94.0 \\96.4$

C. G. Papadopoulos, M. Worek Eur. Phys. J. C50 (2007) 843

MC over Colour

• Total cross sections for processes with gluons

	1 0				LHC
Process	$\sigma_{\rm EXACT} \pm \varepsilon \ ({\rm nb})$	ε (%)	$\sigma_{\rm \scriptscriptstyle MC}\pm\varepsilon~({\rm nb})$	ε (%)	
$\begin{array}{c} gg \rightarrow 2g \\ gg \rightarrow 3g \\ gg \rightarrow 4g \\ gg \rightarrow 5g \\ gg \rightarrow 6g \end{array}$	$\begin{array}{l} (0.46572 \pm 0.00258) \times 10^4 \\ (0.15040 \pm 0.00159) \times 10^3 \\ (0.11873 \pm 0.00224) \times 10^2 \\ (0.10082 \pm 0.00198) \times 10^1 \\ (0.74717 \pm 0.01490) \times 10^{-1} \end{array}$	0.5 1.0 1.9 1.9 2.0	$\begin{array}{l} (0.46849 \pm 0.00308) \times 10^{4} \\ (0.15127 \pm 0.00110) \times 10^{3} \\ (0.12116 \pm 0.00134) \times 10^{2} \\ (0.09719 \pm 0.00142) \times 10^{1} \\ (0.76652 \pm 0.01862) \times 10^{-1} \end{array}$	$0.6 \\ 0.7 \\ 1.1 \\ 1.5 \\ 2.4$	

• Comparison of the computational time

Process	$\sigma_{\rm MC} \pm \varepsilon \ ({\rm nb})$	ε (%)
$\begin{array}{c} gg \rightarrow 7g \\ gg \rightarrow 8g \\ gg \rightarrow 9g \end{array}$	$\begin{array}{r} (0.53185 \pm 0.01149) \times 10^{-2} \\ (0.33330 \pm 0.00804) \times 10^{-3} \\ (0.13875 \pm 0.00430) \times 10^{-4} \end{array}$	2.1 2.4 3.1

Process	$t_{\rm EXACT}^{\rm CF}$	$t_{\rm MC}$	$t_{\rm exact}/t_{\rm mc}$
$\begin{array}{c} gg \rightarrow 2g \\ gg \rightarrow 3g \\ gg \rightarrow 4g \\ gg \rightarrow 5g \\ gg \rightarrow 6g \end{array}$	$\begin{array}{c} 0.315 \times 10^{0} \\ 0.329 \times 10^{1} \\ 0.383 \times 10^{2} \\ 0.517 \times 10^{3} \\ 0.987 \times 10^{4} \end{array}$	$\begin{array}{c} 0.554 \times 10^{0} \\ 0.143 \times 10^{1} \\ 0.372 \times 10^{1} \\ 0.105 \times 10^{2} \\ 0.362 \times 10^{2} \end{array}$	$\begin{array}{c} 0.57 \\ 2.30 \\ 10.29 \\ 49.24 \\ 272.65 \end{array}$

C. G. Papadopoulos, M. Worek Eur. Phys. J. C50 (2007) 843

Helicity Treatment

• Summation over helicity configurations of the external partons

 \rightarrow MC integration over a phase variable

• Polarization vector for gluons:

 $\epsilon^{\mu}_{\phi}(\mathbf{p}) = \mathbf{e}^{i\phi} \epsilon^{\mu}(\mathbf{p}, +) + \mathbf{e}^{-i\phi} \epsilon^{\mu}(\mathbf{p}, -)$

• For incoming quarks e.g.:

$$\boldsymbol{u}_{\phi}(\boldsymbol{p}) = \boldsymbol{e}^{i\phi} \boldsymbol{u}_{+}(\boldsymbol{p}) + \boldsymbol{e}^{-i\phi} \boldsymbol{u}_{-}(\boldsymbol{p})$$

• ϕ - random number $\phi \in (0, 2\pi)$

• By integrating over phase we can get the correct sum over polarisations

Monte Carlo Integration

- Use random numbers for integration \vec{x}_i
- Evaluate, estimate of the integral *I*(*f*)

$$\langle I(f) \rangle = \frac{1}{N} \sum_{i=1}^{N} f(\vec{x}_i)$$

• Quality of estimate given by error estimator – Variance

$$\langle E(f) \rangle^2 = \frac{1}{N-1} [\langle I^2(f) \rangle - \langle I(f) \rangle^2]$$

- Minimize *E(f)*
- Problem: Large fluctuation in integrand f
- Solution: Smart sampling methods !

Importance Sampling

- MC carried out using sets of random points picked from arbitrary distribution
- Uniform distributions give very poor estimates of high-dimensional integrals
- Sampling points from a given probability function !
- Choosing points from a distribution g(x) which concentrates the points where the function f(x) being integrated is large
 improve convergence behaviour
- Function g(x) is chosen to be a reasonable approximation to f(x)
- When f(x)/g(x) is ~ constant E(f/g) is small

Stratified Sampling

• Decompose integral in M sub-integrals

$$\langle I(f) \rangle = \sum_{j=1}^{M} \langle I_j(f) \rangle$$

 $\langle E(f) \rangle^2 = \sum_{j=1}^{M} \langle E_j(f) \rangle^2$

- Overall variance smallest if equally distributed !
- Sample where the fluctuations are
- Divide interval in bins
- Adjust bin-size or weight per bin such that variance identical in all bins

Multichannel Sampling

Method for parton level event generators ! Translate Feynman diagrams into channels ! s- and t- channel propagators as building blocks

• Use sum of functions

 $g(\vec{x}) = \sum_{j=1}^{N} \alpha_{j} g_{j}(\vec{x})$

- Conditions on weights like stratified sampling
- Combination of importance sampling and stratified sampling
- Select $g_j(\vec{x})$ with probability $\alpha_j \rightarrow \vec{x}_j$
- Calculate total $g(\vec{x}_j)$ and partial $g_j(\vec{x}_j)$ weight
- Add f(x_j)/g(x_j) to total result and f(x_j)/g_j(x_j) to partial (channel-) results
 After N sampling steps update apriori weights

Unweighting & Hit or Miss

• Take over-estimator g(x)

 $g(x) > f(x) \qquad \forall x \in [x_{\min}, x_{\max}]$

- Select *x* according to *g*
- Accept or reject with f(x)/g(x)
- Obvious guaranteed g(x):

 $g(x) = Max\{f(x)\}$

• Compare actual $f(\vec{x})$ with maximal value during sampling \rightarrow unweighted events

Performance

Combining ME & PS

How to simulate hard processes with additional hard radiation

Matrix element:

- Exact at some given order in α_s all interferences are included
- High energetic and well separated partons
- Soft/Collinear regions are not adequately described luck of multiple unresolved gluon emission
- Difficult to match to hadronisation models

Parton Showers:

- Include logarithmically enhanced soft and collinear contributions of parton emissions
- Able to connect both hard and fragmentation scales
- Not enough high energetic gluons are emitted that have large angle from the shower initiator

Clearly two descriptions complement each other !

Combining ME & PS

Goal

- All jet emissions correct at tree level + LL
- Soft emission correctly resumed in PS

Problem

• Double Counting - jet can appear both from relatively hard emission during shower evolution and from inclusion of higher order ME

Solution

• Matching algorithm

Recipe

- Separate jet production/evolution by $Q_{iet} k_T$ algorithm
- Produce jets according to LO MEs
- Reweight with Sudakov form factor + running α_s weight
- Veto jet production in parton shower
- Process Independent Implementation !

Combining ME & PS

A few algorithms along these lines:

S. Catani, F. Krauss, R. Khun, B. R. Webber JHEP 0111 (2001) 063 L. Lonnblad JHEP 0205 (2002) 046 F. Krauss JHEP 0208 (2002) 015 M. L. Mangano, M. Moretti, R. Pittau Nucl. Phys. B. 632 (2002) 343 M. L. Mangano, M. Moretti, F. Piccinini, M. Treccani JHEP 0701 (2007) 013

- CKKW-L for e+e- dependence on the resolution parameter is shifted beyond NLL accuracy, proposal to extend procedure to hadronic collisions but no proof of NLL accuracy
- MLM alternative proposal, LL accuracy

Differ mainly:

- Jet definition used for the ME evaluation
- Way the ME rejection weights are constructed
- Details concerning starting conditions of jet vetoing inside PS

Have similar systematics:

- Residual dependence on the phase space separation cut Q_{iet}
- Variations with the number of ME legs
- Dependencies on the internal jet algorithm

 $pp \rightarrow W + jets @ LHC$

AlpGen – angular-ordered PS in HERWIG with MLM matching

Ariadne – matrix elements MadGraph, p_T ordered dipole PS with CKKW-L, PYTHIA

HELAC – mass-ordered PS in PYTHIA with MLM matching

MadEvent – mass-ordered PS in PYTHIA with MLM matching

SHERPA – mass-ordered PS with CKKW matching, PYTHIA

J. Alwall et al. Eur. Phys. J. C53 (2008) 473

- di spectra scale in parton level event where i jets are clustered into i-1 jets using k_T algorithm
- ΔR separations
- Curves normalized to unit area

J. Alwall et al. Eur. Phys. J. C53 (2008) 473

HELAC-PHEGAS

Systematics @ LHC

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Combining Tools

- **VBFNLO** Warped Higgsless Kaluza-Klein model of narrow spin1 resonances
- HELAC-PHEGAS Most prominent background processes
- Full off-shell and finite width effects for final states with two tagging jets and four leptons
- Double forward jet-tagging techniques
- Dedicated cuts on the observable jets and charged leptons
- Substantial sensitivity to strong interactions in EWSB sector

C. Englert, B. Jager, M. Worek, D. Zeppenfeld arXiv:0810.4861 [hep-ph]

How to use Helac-Phegas

- Edit once for all the file **myenv**
 - **FC** = Fortran compiler
 - **FORTRAN_LIBRARIES** = path to your Fortran libraries
 - **LHAPDFLIB** = path to Les Houches Accord PDF's libraries
 - **LHAPDFSETS** = path to PDF sets
- User interface
 - **run.sh** shell script that reads input files compiles and runs HELAC
 - **default.inp** default values (can not be modified)
 - **user.inp** select process, collider, energy, modify default values
 - getqcdscale.h define QCD scale to be used by PDF and α_s

. run.sh user.inp./run.sh user.inp myenv-xxx

Code Structure

Main Files:

- main_mc.f \rightarrow main file
- intpar.f \rightarrow integer arithmetic
- master_new.f \rightarrow master file for DS solution
- pan1.f \rightarrow non-dressed vertices and amplitude calculation
- pan2.f \rightarrow dressed vertices
- physics.f \rightarrow all couplings

Main common blocks:

- common_int.h ۵
- common_helc.h ۵
- common_phegas_mom.h
- \rightarrow common/helac_int/n,io(20),ifl(20)
 - \rightarrow common/helac_helc/ipol(20),icol(20,2)
- \rightarrow common/phegas_mom/pmom(20,5)

From LO to NLO

- If one is content with Born level diagrams possible to go to quite high orders
- 8-10 partons in the final state
- Well separated to avoid phase space regions where divergences become troublesome
- To improve accuracy of prediction higher order calculations are needed
- Benefits of higher order calculations are well know
 - Less sensitivity to unphysical input scales
 - First predictive normalisation of observables at NLO
 - Confidence that cross sections are under control
- There are many programs currently available for predictions
 - Author-controlled non-public
 - Single process or class of processes
 - more generic programs MC@NLO, MCFM, NLOJETS++, VBFNLO,...

Summary & Outlook

- **HELAC–PHEGAS:** Framework for high energy phenomenology
 - Standard Model fully included tested
 - Ready to be used for LHC, TeVatron, ILC
 - High colour charge processes (MC summation over color)
 - Multijet production to be available very soon

• HELAC NEW FEATURES:

- MSSM development phase (so far Higgs sector only)
- Effective Hyy and Hgg (tested phase)
- Anomalous couplings, new resonances, ...

• HELAC GENERAL PLANS:

- Contribute to **ATLAS** and **CMS** generator groups in all stages (interfacing, validation, tuning, installation, configuration, user help, **physical analysis**...)
- Make **HELAC-PHEGAS** an option for the **LHC** !

Summary & Outlook

- Broad range of fully automatic LO Monte Carlo programs
- 8-10 partons in the final state
- LO + LL description for all jet emissions
- Performing NLO calculations on a case-by-case basis is not a way for the future

Automated approaches combining algebraic and numerical recipes appears both promising and feasible !

• First results have already been presented

Lecture by C. Papadopoulos

W+3jet
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