Chirally rotated Schrödinger Functional: first checks at tree-level of PT

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- Lattice regularisation: $V = L^3 \times T$, spacing a
 - inverse lattice spacing 1/a: ultraviolet cutoff
 - observables on the lattice:

 $\langle \mathcal{O} \rangle_{latt} = \langle \mathcal{O} \rangle_{cont} +$ cutoff effects

- remove cutoff \Leftrightarrow continuum limit: $a \rightarrow 0$
- reach continuum limit as fast as possible:

$$\langle \mathcal{O} \rangle_{\text{latt}} = \langle \mathcal{O} \rangle_{\text{cont}} + O(a^2)$$

better than

$$\langle \mathcal{O} \rangle_{latt} = \langle \mathcal{O} \rangle_{cont} + O(\textbf{O})$$

• Big effort to find a lattice regulator with leading $O(a^2)$ effects

- add counterterms to remove the O(a)
- ► get leading O(a²) automatically

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• Our lattice action: Wilson twisted mass fermions at maximal twist

► automatic O(a) improvement (R. Frezzotti and G.C. Rossi, hep-lat/0306014)

$$\langle \mathcal{O} \rangle_{\text{latt}} = \langle \mathcal{O} \rangle_{\text{cont}} + O(o^2)$$

- Many observables need to be renormalised
- Want renormalisation scheme which:
 - ▶ is non-perturbative
 - ► is mass independent: massless renormalisation
 - keeps automatic O(a) improvement
- Final goal: twisted mass lattice QCD simulations for $N_f = 2 + 1 + 1$ (dynamical up, down, strange and charm)

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- Schrödinger Functional schemes: (M. Lüscher et al., hep-lat/9207009)
 - finite space-time volume: $V = L^3 \times T$
 - * can use $\frac{1}{L}$ as renormalisation scale
 - finite size techniques: non-perturbative running with 1/L
 - boundary conditions:
 - ★ periodic in spatial directions
 - * Dirichlet in time direction: (S. Sint, hep-lat/9312079), (M. Lüscher, hep-lat/0603029)

$$P_{\pm}\psi(x)|_{x_0=0} = 0$$
 $P_{\pm}\psi(x)|_{x_0=T} = 0$
 $P_{\pm} = \frac{1}{2} (1 \pm \gamma_0)$

 \implies non-zero bound in the spectrum of the Dirac operator (1/27)

• Lattice regularisation and Schrödinger Functional b.c:

- non-zero bound in eigenvalue spectrum of the Dirac operator
 - \implies allows lattice simulations at the chiral point
 - \Rightarrow non-perturbative and massless renormalisation of QCD
- General problem of SF schemes (with lattice regularisation)
 O(a) effects from the boundaries
 - \implies add boundary counterterms

Standard Schrödinger Functional b.c break chiral symmetry

Wilson twisted mass fermions and standard SF b.c:

Incompatible with bulk automatic O(a) improvement

• Let's try with a chiral rotation of the standard SF formulation

(S. Sint, hep-lat/0511034)

 \implies

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Ohiral rotation of the the quark fields (with maximal twist):

$$\psi(x) \to e^{i\frac{\alpha}{2}\gamma_5\tau^3}\psi(x) \qquad \bar{\psi}(x) \to \bar{\psi}(x) e^{i\frac{\alpha}{2}\gamma_5\tau^3} \quad ; \quad \alpha = \pi/2$$

Chirally rotated SF boundary conditions

(S. Sint, hep-lat/0511034), (S. Sint, talk at Lattice 2008)

$$\begin{split} & \mathcal{Q}_{+}\psi(x)|_{x_{0}=0}=0 \qquad \mathcal{Q}_{-}\psi(x)|_{x_{0}=T}=0 \\ & \mathcal{Q}_{\pm}=\frac{1}{2}\left(\mathbbm{1}\pm i\gamma_{0}\gamma_{5}\tau^{3}\right) \end{split}$$

- in the continuum it is a change of basis: same theory as SF
- $\gamma_5 \tau^1$ -symmetry of the b.c: automatic O(a) improvement

- Fermion lattice action of a theory with boundaries:
 - want (massless) Wilson fermions in the bulk
 - need to define the lattice action near the desired time boundaries
 - use orbifold techniques (Y. Taniguchi, hep-lat/0412024)

• Get: desired boundary conditions at tree-level up to O(a) effects

(S. Sint, private comunication)

$$Q_{+}(1-\frac{1}{2}\sigma\partial_{0}^{*})\psi(x)|_{x_{0}=0}=0$$
 $Q_{-}(1+\frac{1}{2}\sigma\partial_{0})\psi(x)|_{x_{0}=T}=0$

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Relevant (d=3) boundary operator

- $\gamma_5 \tau^1$ -odd: breaking of flavour and parity symmetries
- ▶ bulk O(a) effects \Leftrightarrow loose bulk automatic O(a) improvement
- add d = 3 finite boundary counterterm with coefficient Z_f

$$\delta S = (\mathcal{Z}_f - 1) \sigma^3 \sum_{\vec{x}} \left(\bar{\psi} \psi |_{x_0 = 0} + \bar{\psi} \psi |_{x_0 = T} \right)$$

• need to tune \mathcal{Z}_f to restore the symmetries

 \implies bulk automatic O(a) improvement

- non-perturbative tuning
 - ★ massless scheme: $m_0 \rightarrow m_c$
 - ★ bulk improvement: $Z_f \rightarrow Z_f^*$

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- Irrelevant (d=4) boundary operator
 - boundary O(a) effects
 - add d = 4 boundary counterterm with coefficient d_s

$$\delta S = a(d_s - 1)a^3 \sum_{\vec{x}} \left(\bar{\psi} \gamma_k D_k \psi |_{x_0 = 0} + \bar{\psi} \gamma_k D_k \psi |_{x_0 = T} \right)$$

- tune d_s to cancel the O(a) boundary effects
- perturbative tuning

* first need to find the tree-level value: $d_s^0 (d_s = d_s^0 + d_s^1 O(g_0^2))$

• Theory with χ SF b.c. is automatic O(a) improved at tree-level if:

- $\sim m_0 = 0$
- ► $Z_f = 1$
- $d_s^0 = \frac{1}{2}$

• Example:



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Definitions:

$$g_A^{ab}(x_0)_{\pm} = -\langle A_0^a(x)\mathcal{Q}_{\pm}^b \rangle \qquad g_P^{ab}(x_0)_{\pm} = -\langle P^a(x)\mathcal{Q}_{\pm}^b \rangle$$

$$\mathcal{Q}_{\pm}^{a} = a^{b} \sum_{\vec{y},\vec{z}} \bar{\zeta}(\vec{y}) \gamma_{5} \frac{1}{2} \tau^{a} \mathcal{Q}_{\pm} \zeta(\vec{z}) e^{i\vec{p}(\vec{y}-\vec{z})}$$

$$\zeta(\vec{x}) = \psi(x)|_{x_0=a} \qquad \bar{\zeta}(\vec{x}) = \bar{\psi}(x)|_{x_0=a}$$

• Set:
$$d_s = 1/2$$

$$u_s = 1/2$$

• Tuning of m_0 :

$$m_{PCAC,-} \equiv \frac{\partial_0 g_A(T/2)_-}{2g_P(T/2)_-} = 0$$

• Tuning of Z_f : $g_{A}(T/2)_{-}=0$







m_{PCAC} vs. m₀



- The use of proper tuning conditions and the precise non-perturbative determination of m_c and Z^{*}_f are needed for
 - computing bulk automatic O(a)-improved observables
 from our lattice twisted mass simulations
 - universality: obtain the correct theory in the continuum limit

proper setup for non-perturbative renormalisation needed for many physical quantities: moments of PDFs, form factors, ...

• At tree-level the theory is automatic O(a) improved if

 $har m_0 = 0$

 \Longrightarrow

- ► $\mathcal{Z}_f = 1$
- $d_s = 1/2$

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- Short-time goal:
 - first quenched studies
 - lattice perturbation theory
- Real goal:

extend this to dynamical simulations with $N_f = 2 + 1 + 1$

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