

# Hypothesis Testing and Confidence Intervals

## Terascale Statistics School

Roger Barlow

7th March 2017



- Hypothesis Testing: signal or background?
- Hypothesis testing: Is there a discovery?
- $\chi^2$  and  $p$ -values
- The need for  $5\sigma$
- Confidence Intervals
- Credible intervals

# Making Choices ('Hypothesis Testing')

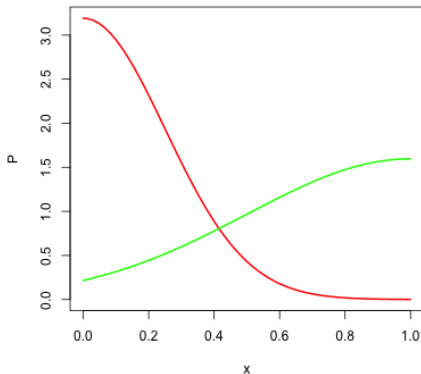
Classification problem: Separation  
into two (or more) classes: e.g.  
Signal and Background events

Hypothesis is 'This is a signal(red)  
not background(green)'

Indicator variable  $x$  - may be  
multidimensional and/or output from  
ML system

Behaviour known - from simulations  
or from control samples

Put a cut somewhere and accept all that pass Accept that some signal  
events will be lost (efficiency) and some background events accepted  
(purity)



# Type I and Type II errors

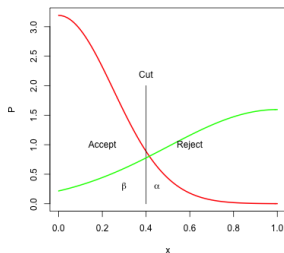
Two sorts of errors possible:

Type I: reject a signal event. Lowers efficiency

Type II: accept a background event. Lowers purity, increases contamination

$\alpha$  is the probability of a type I error  $\int S(x)dx$  over the rejection region. Called 'Significance'

$\beta$  is the probability of a type II error  $\int B(x)dx$  over acceptance region.  $1 - \beta$  called 'Power'

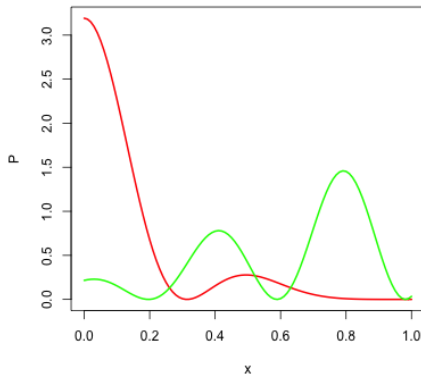


# Neyman-Pearson Lemma

How to choose the best acceptance/rejection region even in cases where  $S(x)$  and  $B(x)$  have lots more structure.

$\alpha$  is imposed. You want to minimise  $\beta$ .

Add regions of  $x$  with the highest  $S(x)/B(x)$  to the acceptance region until  $\alpha$  is achieved.  $\beta$  is then minimised, as replacing any  $\Delta x$  with an equivalent region from the rejection region would bring in more  $B(x)$



# Composite Hypothesis

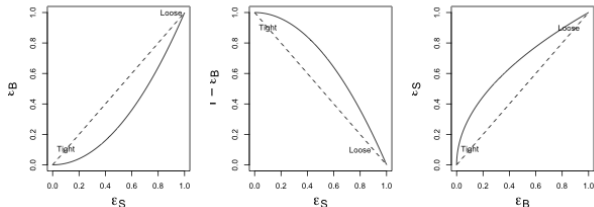
Suppose you have several different sources of background - and their relative sizes are not known so you can't just lump them together

Or a hypothesis containing an adjustable parameter, or parameters

Doesn't affect  $\alpha$

Need to quote worst (i.e. largest)  $\beta$ . Then can say background contamination probability  $\leq \beta$

# ROC plots



Varying cut value varies  $\alpha$ , or  $1 - E_S$  and  $\beta$ , or  $E_B$

Plot  $E_S$  against  $E_B$ , or  $1 - E_S$  against  $E_B$ , or...

'ROC plots' : also called efficiency plots, efficiency/contamination plots, and (wrongly!) efficiency/purity plots Show how good your separation technique is by how far they bend from diagonal. ROC stands for Receiver Operating Characteristic (jargon but sounds good)

# The Null Hypothesis $H_0$

You want to test a theory. Do an experiment. Data may show the theory is wrong - but how do you show it's right using frequentist tools?

'The data are compatible with the theory being true' says nothing.

Have to test 'The theory are not compatible with the theory being false'

Construct **Null Hypothesis**  $H_0$  which is that your theory is false.

Examples: There are no dark matter events in your sample

Web advertising does not increase sales

There are no bumps in the mass plot

New Treatment has no effect on cure rate

.....

'Every experiment may be said to exist only to have the chance to disprove the null hypothesis' - Fisher



# Type I and Type II errors revisited

Type I: reject a signal event  $\rightarrow$  Probability of rejecting the  $H_0$  when it is true.  
'False positive' probability

Description of  $\alpha$  as 'significance' makes sense. Set  $\alpha = 0.05$  and get a result in the reject region. Hooray! Null hypothesis rejected at significance level of 0.05.

Also quote as **Confidence level**  $1 - \alpha$ . Here 95%.

If  $H_0$  passes the test I say, with 95% confidence 'There is no effect'. Statement is either true or false, but is a member of an ensemble of statements, 95% true.

Legal point: I say with 95% confidence that a statement is true if it belongs to an ensemble of which *at least* 95% of statements are true. Extra two words mean (i) a statement true with 96% confidence is also true with 95% confidence and (ii) composite hypotheses are covered by taking the worst case.

Type II: accept a background event  $\rightarrow$  fail to reject  $H_0$  given the chance

# $\chi^2$ and Goodness of Fit

Tool often used in Hypothesis testing

Data  $\{x_i, y_i\}$  and  $H_0$  is that  $y = f(x)$

$$\text{Define } \chi^2 = \sum_{i=1}^N \left( \frac{y_i - f(x_i)}{\sigma_i} \right)^2$$

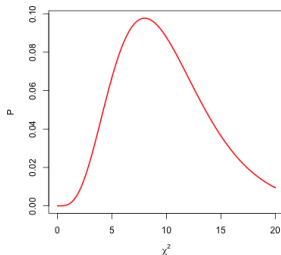
Expectation  $\chi^2 = N$ . (broadly obvious, and can be proved)

Figure shows  $N = 20$

$$P(\chi^2; N) = \frac{2^{-N/2}}{\Gamma(N/2)} \chi^{N-2} e^{-\chi^2/2}$$

Comes from integrating over multiple Gaussians.

Use  $\chi^2$  as indicator variable for  $H_0$  (i.e. does the curve  $y = f(x)$  go through the data points  $\{x_i, y_i \pm \sigma_i\}$ )



## More about $\chi^2$

You will also see  $\chi^2 = \sum_{i=1}^N \frac{(y_i - f(x_i))^2}{y_i}$ . This is a simplified version and applies only to histograms

If  $f(x)$  actually  $f(x; a)$  and  $a$  adjusted to minimise  $\chi^2$ , replace  $N$  by  $N_D$ : number of data points minus number of adjusted parameters. (Nice property of  $\chi^2$ : assumes fitting acts to reduce dimensionality of  $\chi$  space.)

If  $\chi^2$  is large, reject  $H_0$ . Calculate  $p = \int_{\chi^2}^{\infty} P(\chi'^2; N) d\chi'^2$  - probability of getting a  $\chi^2$  value this large or worse.

If  $\chi^2$  is small - be suspicious and check your error calculations.

$\chi^2$  is a really nice tool: minimising it gives estimates (and their errors) *and* its value gives goodness-of-fit. Contrast likelihood, where the actual value is meaningless.

$p$  values are not just for  $\chi^2$ : - probability of getting a value this extreme or worse for any measure of agreement.

## The Difference between $p$ and $\alpha$

In one sense, none. Both are  $\int_x^\infty P(x') dx'$

In another sense, plenty.  $\alpha$  is computed *before* you see the data and is a property of the test.  $p$  is a property of the data.

# Wilks' Theorem

Getting a lot of use recently

Start with variant of  $\chi^2$  test. Suppose 2 models  $H_0$  and  $H_1$ , where  $H_1$  is similar to  $H_0$  but has more parameters (e.g.  $H_0$  is straight line fit,  $H_1$  is quadratic)

The difference  $\Delta\chi^2$  between the two is also  $\chi^2$  distributed, with  $N_D$  the number of new parameters.

Use to answer question 'Do the extra parameters really help?' Sort of decoupled from total  $\chi^2$  values

Wilks extends this to likelihoods:  $-2\Delta\ln L$  has a  $\chi^2$  distribution. (Hence name 'Likelihood Ratio test')

Proviso (1) applies in large  $n$  limit and (2) Models must be nested: for particular parameter value  $H_1$  reverts to  $H_0$

# Five sigma

$p$ -values often expressed in terms of Gaussian  $\sigma$ .

E.g. probability of Gaussian exceeding  $3\sigma$  is 0.13%. (1-tailed). Call a  $p$ -value of 0.13% or below '3 sigma'

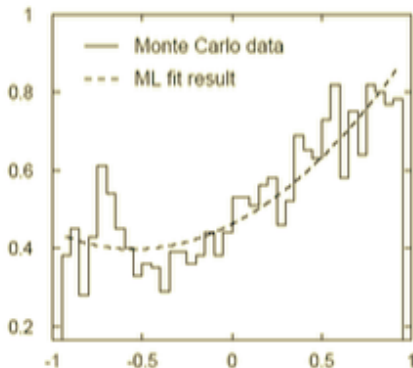
Do these calculations in R with `1-pnorm(Nsigma)` and `qnorm(1-p)`

For a 'discovery' claim we need 5 sigma. (Probability below  $3 \times 10^{-7}$ )

Harsh? Yes, but justified by history

# Look Elsewhere effect

Multiply number of particle physicists by number of plots they can draw per day times days in year - many thousand. There will be 3 sigma effects!



If you histogram random numbers, you see lots of apparently significant bumps

# Blind Analysis



Cuts must be performed without looking at the final sample

Otherwise the analyst will tweak cut values to enhance the signal size, or the effect they are looking for.

The story of the split  $A_2$  meson is a cautionary tale

We are better today but not immune - as the  $Z(750)$  shows

Statistics texts warn against publication bias - if 20 studies are done, one will (probably) report a 5% significance effect. OK, except it may be the only one that gets published.

We probably err in the reverse direction. If a measurement agrees with the Standard Model, we publish. If it doesn't we check and wait for more data.



# Confidence Intervals

Given Gaussian measurement (say)  $M = 123 \pm 4.5 \text{ GeV}$ .

We know this was sampled from a distribution  $\propto \exp^{-\frac{1}{2}(\frac{M-\mu}{4.5})^2}$

But  $\mu$  is unknown and we want to say something about it

Argument:

$M$  will be within 4.5 GeV of  $\mu$  in 68% of all cases

$\mu$  will be within 4.5 GeV of  $M$  in 68% of all cases

If I say ' $\mu$  lies between 118.5 and 127.5 GeV' I have a 68% chance of being correct. (The statement is a member of an ensemble of statements of which 68% are true)

118.5 to 125.5 GeV is a 68% confidence interval

*Frequentists manage to say something useful about uncertainty on measurement without violating their probability definition!*

The choice (68%, 95%, whatever) is up to the author. Tradeoff between precision and confidence.

The ends of the interval do not have to be symmetric, provided integral is chosen value. Other choices include lower limit, upper limit, shortest interval and central (probability of both tails equal)

# Gaussians and Likelihood

Measurement may not come labeled as  $\mu \pm \sigma$

For true  $\mu$ , log likelihood is  $-\frac{1}{2} \left( \frac{M-\mu}{\sigma} \right)^2 + \text{constant term}$

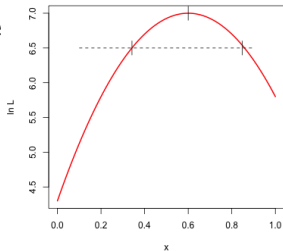
This is a parabola.  $\Delta \ln L = -\frac{1}{2}$  at  $x = \pm\sigma$

$\frac{1}{\sigma^2}$  is the second derivative.

We do not have access to the true  $\ln L$  but we do have access to  $\ln L$  at  $\hat{\mu}$  and the second derivative is (hopefully) not that different

So estimate true  $\sigma$  from our  $\frac{d^2 \ln L}{d\mu^2}$  and estimate that from points where  $\Delta \ln L = -\frac{1}{2}$

This defines a 68% confidence interval.



## $\chi^2$ again

$\Delta \ln L = -\frac{1}{2}$  translates to  $\Delta \chi^2 = +1$

In 1-D, probability of  $\chi^2 < 1$  is 68%

In 2-D, probability of  $\chi^2 < 1$  is 39% (in R, `pchisq(1.0,2)`)

If you want to find a region for which the chance is 68.3%, need  $\Delta \chi^2 = 2.30$  (`qchisq(0.683,2)`)

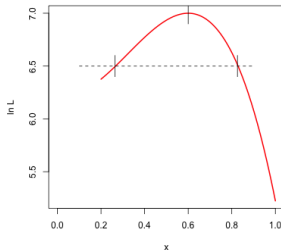
For general n-D, and other probability values, use `pchisq` and `qchisq`

# Asymmetric Statistical Errors

Maybe your  $\ln L$  is not a nice parabola

$\Delta \ln L$  points are then not symmetric.

Argument for quoting them is: could transform to some  $\mu'$  for which  $\ln L$  was parabolic. Extract symmetric errors. Then transform back.



OK **but what later?** How should you use combination of errors when errors are asymmetric?

Standard procedure is to combine positive and negative errors separately  
This is **WRONG** as it violates the Central Limit Theorem

For a better method, see RB 'Asymmetric Statistical Errors'  
[arXiv:physics/0406120](https://arxiv.org/abs/physics/0406120)

# Confidence bands

Constructed Horizontally, read vertically

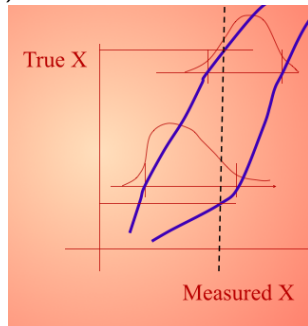
Gaussian measures are easy to handle because of their symmetry. What about more general  $P(M; \mu)$  giving probability that a true  $\mu$  will result in a measured  $M$

Consider  $M, \mu$  plane. For each  $\mu$ , select the range of  $M$  according to value and choice (i.e. 68% or 95%, central or shortest...)

This defines the confidence belt.

Then take your measurement of  $M$ . We say with 68% (or whatever) confidence that  $(M, \mu)$  lies inside the belt. So we read off the limits on  $\mu$

Note that the upper/lower limits in construction give the lower/upper limits in readout



# Poisson confidence intervals

Extra complication as  $x$  – axis is discrete.  $\int_{M \in \text{region}} P(M, \mu) dM$  becomes  $\sum_{r \in \text{region}} P(r, \mu)$

May not be possible to choose region to satisfy probability exactly: inclusion may bring it over the value - but that's OK due to the 'at least'

Most Poisson studies are interested in setting one-sided *upper limits*. Numbers are small. (If they were large we could use the Gaussian approximation).

So if you see  $N$  you want  $\sum_{r=0}^N e^{-\mu} \frac{\mu^r}{r!} = 0.05$  (or whatever)

If the true value is  $\mu$ , or higher, the probability of only seeing  $N$ , or fewer, is only 0.05, or less.

Magic number: for  $N = 0$ , 95% confidence,  $\mu = 3.00$ . If you see zero events, you know with 95% confidence that the true value is at most 3.0

# The big problem

What do you do if you see 0 events - and your expected background is 3.1?

See  $CL_S$ , Feldman-Cousins and other topics in later lectures



# Bayesian credible intervals

Everything is easy.

Take prior. Multiply by likelihood, and normalise to get your posterior.

Find regions for which  $\int_{\mu \in \text{region}} P(\mu | \text{data}) d\mu = C$  where  $C$  is the desired probability content and the region may be symmetric, central, whatever

Poisson Upper limit (i.e. credible interval  $[0, \mu]$ ) with uniform prior

$$P(\mu | r) = \frac{e^{-\mu} \mu^r / r!}{\int e^{-\mu'} \mu'^r / r! d\mu'}$$

Integral can be done by parts and you end up with  $\sum_{r=0}^N e^{-\mu} \frac{\mu^r}{r!} = 0.05$  – same as Frequentists!

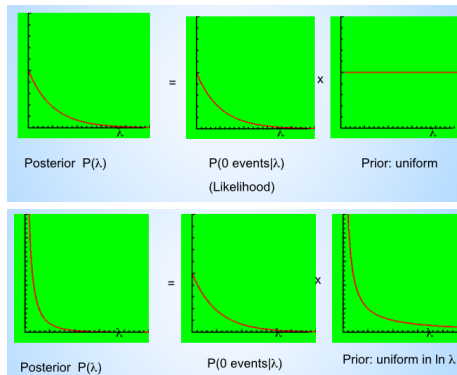
*This is just a coincidence - doesn't hold for upper limits. But has saved a lot of arguments.*

# Sensitivity to Priors

Consider case of zero observed events

95% upper limit is 3.00

95% upper limit much smaller than 3.00



# Coverage

Another use for Toy Monte Carlos

Check out that your confidence interval procedure makes sense by

- 1) Choose some true value  $\mu$
- 2) Generate lots of random pseudo-experimental results
- 3) Construct the confidence interval for each
- 4) Count how many times the interval encloses the true value

Ideally the result of (4) should be the same as you specified when you chose the constructor.

If more ('over-coverage') this is OK, thanks to the 'at least' clause, but inefficient. If less ('under-coverage') then something is wrong,

Note: coverage is a function of  $\mu$ , not a single number. Typical sawtooth plots.

Bayesians do not care about coverage - but check it anyway

# Tau decay puzzle

Here is just one page of PDG limits on 'forbidden' tau decays, all quoted at 90% Confidence

10% of these results are allowed to be wrong.

Frankly, I don't think any of them are.

What's going on?

A frequentist can invoke the 'at least' clause in the definition of confidence

But what about the Bayesians?

$e^- \gamma$	LF	< 3.3	$\times 10^{-8}$ CL=90%
$\mu^- \gamma$	LF	< 4.4	$\times 10^{-8}$ CL=90%
$e^- \pi^0$	LF	< 8.0	$\times 10^{-8}$ CL=90%
$\mu^- \pi^0$	LF	< 1.1	$\times 10^{-7}$ CL=90%
$e^- K_S^0$	LF	< 2.6	$\times 10^{-8}$ CL=90%
$\mu^- K_S^0$	LF	< 2.3	$\times 10^{-8}$ CL=90%
$e^- \eta$	LF	< 9.2	$\times 10^{-8}$ CL=90%
$\mu^- \eta$	LF	< 6.5	$\times 10^{-8}$ CL=90%
$e^- \rho^0$	LF	< 1.8	$\times 10^{-8}$ CL=90%
$\mu^- \rho^0$	LF	< 1.2	$\times 10^{-8}$ CL=90%
$e^- \omega$	LF	< 4.8	$\times 10^{-8}$ CL=90%
$\mu^- \omega$	LF	< 4.7	$\times 10^{-8}$ CL=90%
$e^- K^*(892)^0$	LF	< 3.2	$\times 10^{-8}$ CL=90%
$\mu^- K^*(892)^0$	LF	< 5.9	$\times 10^{-8}$ CL=90%
$e^- \bar{K}^*(892)^0$	LF	< 3.4	$\times 10^{-8}$ CL=90%
$\mu^- \bar{K}^*(892)^0$	LF	< 7.0	$\times 10^{-8}$ CL=90%
$e^- \eta'(958)$	LF	< 1.6	$\times 10^{-7}$ CL=90%
$\mu^- \eta'(958)$	LF	< 1.3	$\times 10^{-7}$ CL=90%
$e^- f_0(980) \rightarrow e^- \pi^+ \pi^-$	LF	< 3.2	$\times 10^{-8}$ CL=90%
$\mu^- f_0(980) \rightarrow \mu^- \pi^+ \pi^-$	LF	< 3.4	$\times 10^{-8}$ CL=90%
$e^- \phi$	LF	< 3.1	$\times 10^{-8}$ CL=90%
$\mu^- \phi$	LF	< 8.4	$\times 10^{-8}$ CL=90%
$e^- e^+ e^-$	LF	< 2.7	$\times 10^{-8}$ CL=90%
$e^- \mu^+ \mu^-$	LF	< 2.7	$\times 10^{-8}$ CL=90%
$e^+ \mu^- \mu^-$	LF	< 1.7	$\times 10^{-8}$ CL=90%