

Axion in Reach of ALPS II? ■

Andreas Ringwald

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- The axion window in photon coupling vs. mass parameter space is much wider than previously thought:



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Redefining the Axion Window

Luca Di Luzio,^{1,*} Federico Mescia,^{2,†} and Enrico Nardi^{3,‡}

¹*Institute for Particle Physics Phenomenology, Department of Physics, Durham University, Durham DH1 3LE, United Kingdom*

²*Departamento de Física Quàntica i Astrofísica, Institut de Ciències del Cosmos (ICCUB),*

Universitat de Barcelona, Martí Franquès 1, E08028 Barcelona, Spain

³*INFN, Laboratori Nazionali di Frascati, C.P. 13, 100044 Frascati, Italy*

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A major goal of axion searches is to reach inside the parameter space region of realistic axion models. Currently, the boundaries of this region depend on somewhat arbitrary criteria, and it would be desirable to specify them in terms of precise phenomenological requirements. We consider hadronic axion models and classify the representations R_Q of the new heavy quarks Q . By requiring that (i) the Q 's are sufficiently short lived to avoid issues with long-lived strongly interacting relics, (ii) no Landau poles are induced below the Planck scale; 15 cases are selected which define a phenomenologically preferred axion window bounded by a maximum (minimum) value of the axion-photon coupling about 2 times (4 times) larger than is commonly assumed. Allowing for more than one R_Q , larger couplings, as well as complete axion-photon decoupling, become possible.

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The photo-philic QCD axion

Marco Farina,^a Duccio Pappadopulo,^b Fabrizio Rompineve,^{b,c} and Andrea Tesi^d

^a*New High Energy Theory Center, Department of Physics, Rutgers University, 136 Frelinghuysen Road, Piscataway, NJ 08854, U.S.A.*

^b*Center for Cosmology and Particle Physics, Department of Physics, New York University, New York, NY 10003, U.S.A.*

^c*Institute for Theoretical Physics, University of Heidelberg, Philosophenweg 19, 69120 Heidelberg, Germany*

^d *Enrico Fermi Institute, Department of Physics, University of Chicago, 5620 S. Ellis Ave, Chicago, IL 60637, U.S.A.*

E-mail: farina.phys@gmail.com, duccio.pappadopulo@gmail.com,
f.rompineve@thphys.uni-heidelberg.de, atesi@uchicago.edu

ABSTRACT: We propose a framework in which the QCD axion has an exponentially large coupling to photons, relying on the “clockwork” mechanism. We discuss the impact of present and future axion experiments on the parameter space of the model. In addition to the axion, the model predicts a large number of pseudoscalars which can be light and observable at the LHC. In the most favorable scenario, axion Dark Matter will give a signal in multiple axion detection experiments and the pseudo-scalars will be discovered at the LHC, allowing us to determine most of the parameters of the model.

KEYWORDS: Beyond Standard Model, CP violation

ARXIV EPRINT: [1611.09855](https://arxiv.org/abs/1611.09855)



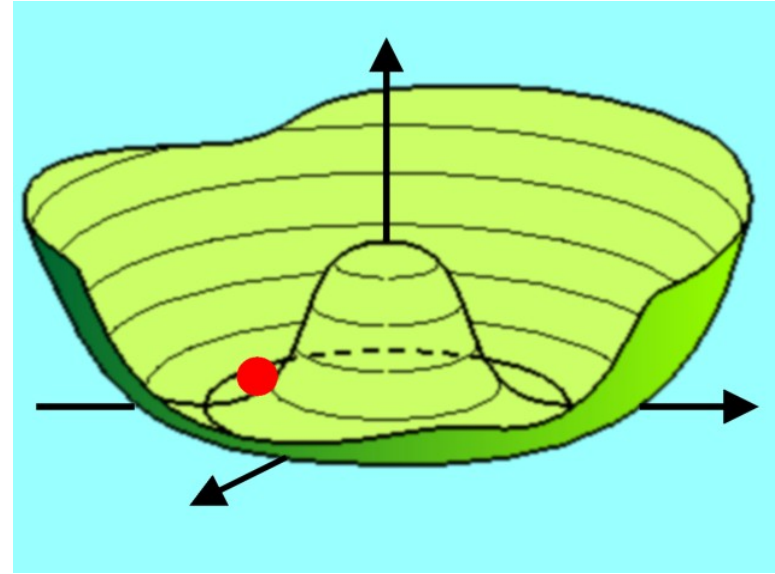
Axionic solution of strong CP problem

➤ A singlet complex scalar field σ featuring a global $U(1)_{PQ}$ symmetry is added to SM

➤ Symmetry is broken by vev $\langle \sigma \rangle = v_{PQ}/\sqrt{2}$

$$\sigma(x) = \frac{1}{2} (v_{PQ} + \rho(x)) e^{iA(x)/v_{PQ}}$$

- Excitation of modulus: $m_\rho \sim v_{PQ}$
- Excitation of angle: NGB $m_A \ll v_{PQ}$



[Raffelt]

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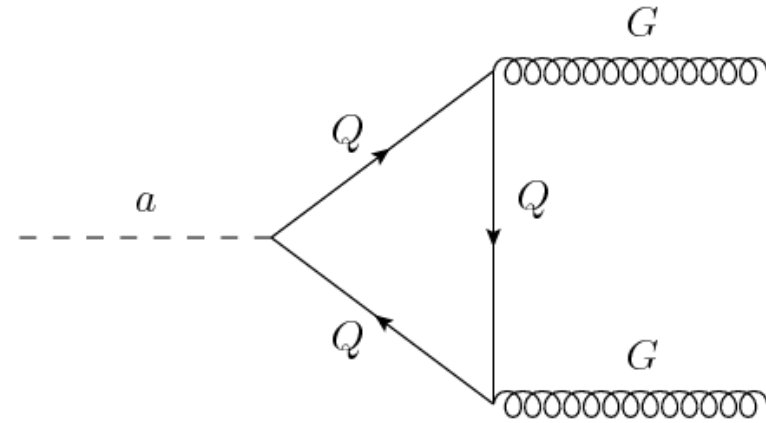
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> Quarks (SM or extra) carry PQ charges such that $U(1)_{\text{PQ}}$ is anomalously broken due to a gluonic triangle anomaly

$$\partial_\mu J_{U(1)_{\text{PQ}}}^\mu = -\frac{\alpha_s}{8\pi} N G_{\mu\nu}^a \tilde{G}^{a\mu\nu} - \frac{\alpha}{8\pi} E F_{\mu\nu} \tilde{F}^{\mu\nu}$$



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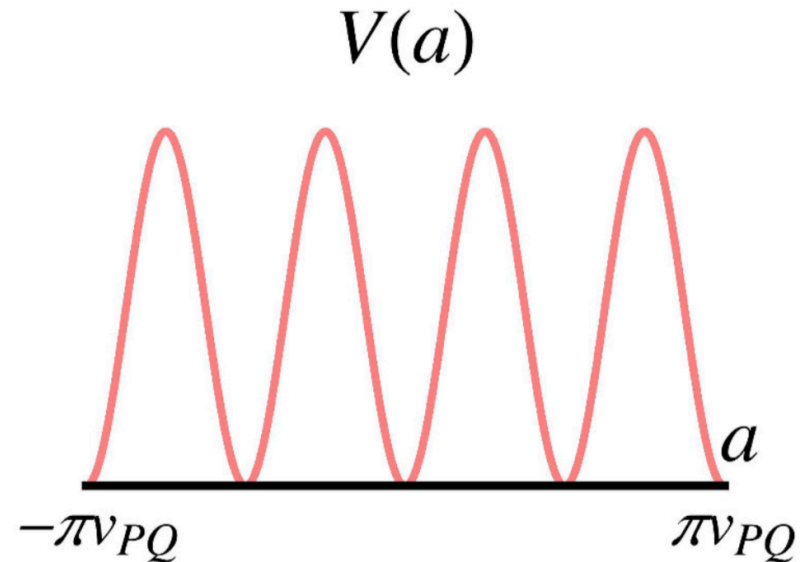
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> No strong CP problem, since axion field acts as x-dependent theta parameter

$$\mathcal{L} \supset -\frac{\alpha_s}{8\pi} \frac{A(x)}{f_A} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} - \frac{\alpha}{8\pi} \frac{E}{N} \frac{A(x)}{f_A} F_{\mu\nu} \tilde{F}^{\mu\nu}; \quad f_A = v_{PQ}/N$$

QCD dynamics: $\langle A(x) \rangle = 0; \quad m_A = 57.0(7) \left(\frac{10^{11} \text{GeV}}{f_A} \right) \mu\text{eV},$



[Peccei,Quinn 77; Weinberg 78; Wilczek 78]



How wide is the axion window?

- Coupling to photons most important for experimental searches

$$\mathcal{L} \supset -\frac{\alpha}{8\pi} \frac{E}{N} \frac{A(x)}{f_A} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

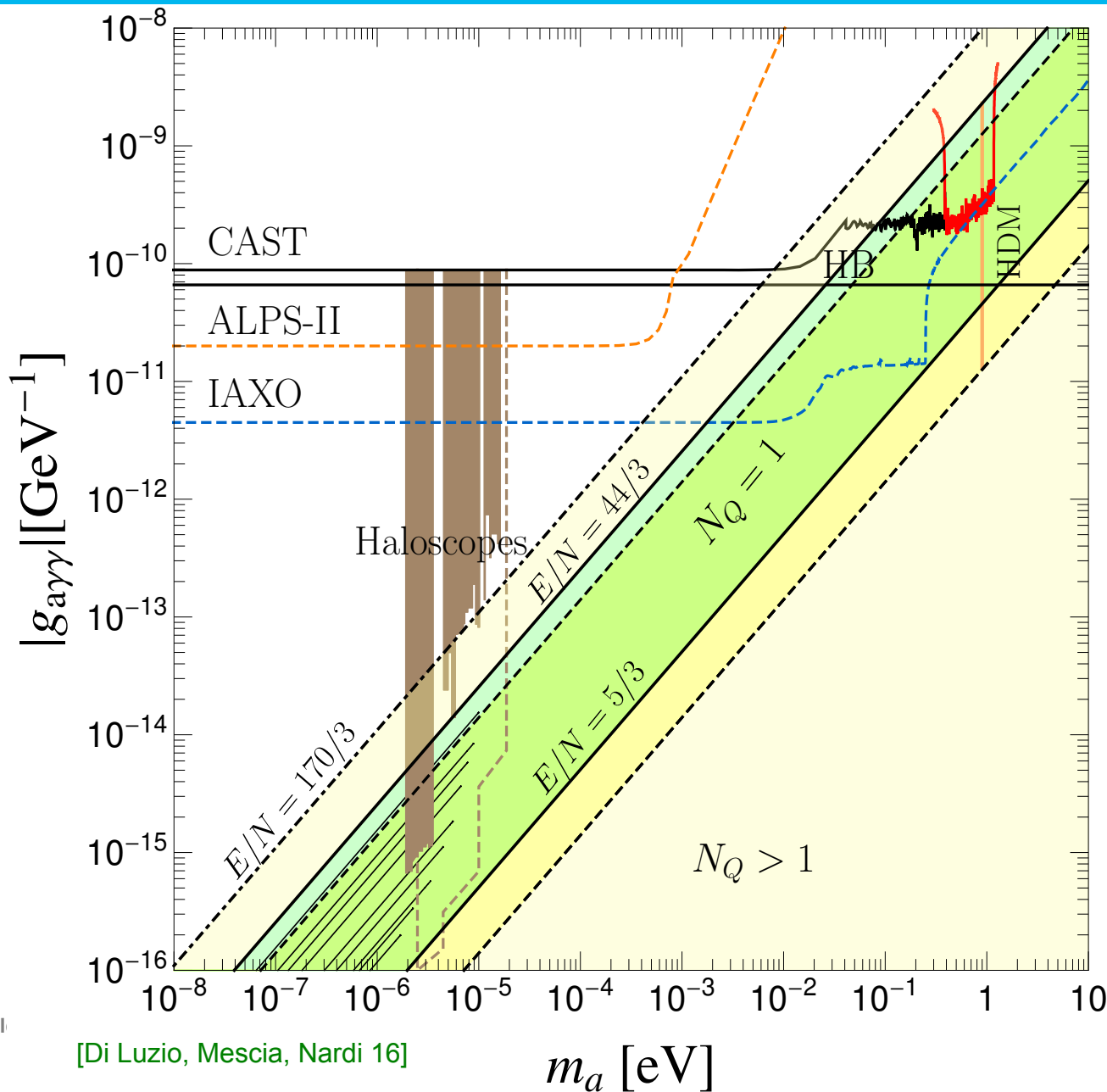
- How wide is the axion window, i.e. what is realistic range of E/N ?
- Consider KSVZ-type axion models: Anomaly induced by heavy fermions Q in the representation $R_Q = (\mathcal{C}_Q, \mathcal{I}_Q, \mathcal{Y}_Q)$ under $SU(3)_C \times SU(2)_I \times U(1)_Y$
- Fifteen cases survive phenomenological requirements
 - Q sufficiently short lived to avoid issues with long-lived strongly interacting relics,
 - no Landau poles induced below Planck scale

R_Q	O_{Qq}	$\Lambda_{LP}^{R_Q} [\text{GeV}]$	E/N	N_{DW}
$R_1: (3, 1, -\frac{1}{3})$	$\bar{Q}_L d_R$	$9.3 \cdot 10^{38} (g_1)$	2/3	1
$R_2: (3, 1, +\frac{2}{3})$	$\bar{Q}_L u_R$	$5.4 \cdot 10^{34} (g_1)$	8/3	1
$R_3: (3, 2, +\frac{1}{6})$	$\bar{Q}_R q_L$	$6.5 \cdot 10^{39} (g_1)$	5/3	2
$R_4: (3, 2, -\frac{5}{6})$	$\bar{Q}_L d_R H^\dagger$	$4.3 \cdot 10^{27} (g_1)$	17/3	2
$R_5: (3, 2, +\frac{7}{6})$	$\bar{Q}_L u_R H$	$5.6 \cdot 10^{22} (g_1)$	29/3	2
$R_6: (3, 3, -\frac{1}{3})$	$\bar{Q}_R q_L H^\dagger$	$5.1 \cdot 10^{30} (g_2)$	14/3	3
$R_7: (3, 3, +\frac{2}{3})$	$\bar{Q}_R q_L H$	$6.6 \cdot 10^{27} (g_2)$	20/3	3
$R_8: (3, 3, -\frac{4}{3})$	$\bar{Q}_L d_R H^{\dagger 2}$	$3.5 \cdot 10^{18} (g_1)$	44/3	3
$R_9: (\bar{6}, 1, -\frac{1}{3})$	$\bar{Q}_L \sigma d_R \cdot G$	$2.3 \cdot 10^{37} (g_1)$	4/15	5
$R_{10}: (\bar{6}, 1, +\frac{2}{3})$	$\bar{Q}_L \sigma u_R \cdot G$	$5.1 \cdot 10^{30} (g_1)$	16/15	5
$R_{11}: (\bar{6}, 2, +\frac{1}{6})$	$\bar{Q}_R \sigma q_L \cdot G$	$7.3 \cdot 10^{38} (g_1)$	2/3	10
$R_{12}: (8, 1, -1)$	$\bar{Q}_L \sigma e_R \cdot G$	$7.6 \cdot 10^{22} (g_1)$	8/3	6
$R_{13}: (8, 2, -\frac{1}{2})$	$\bar{Q}_R \sigma l_L \cdot G$	$6.7 \cdot 10^{27} (g_1)$	4/3	12
$R_{14}: (15, 1, -\frac{1}{3})$	$\bar{Q}_L \sigma d_R \cdot G$	$8.3 \cdot 10^{21} (g_3)$	1/6	20
$R_{15}: (15, 1, +\frac{2}{3})$	$\bar{Q}_L \sigma u_R \cdot G$	$7.6 \cdot 10^{21} (g_3)$	2/3	20



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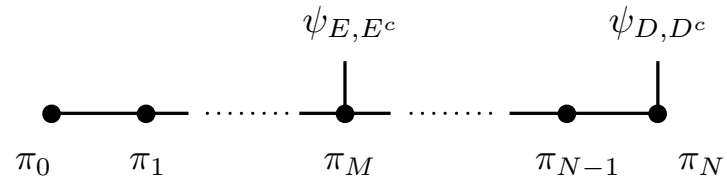
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- ALPS II discovery reach nearly touches extended axion window



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- > Even wider window in clockwork axion models involving additional N pseudo NGBs

$$\mathcal{L} \supset \frac{1}{2} \sum_{j=0}^N \partial_\mu \pi_j \partial^\mu \pi_j - \frac{m^2}{2} \sum_{j=0}^{N-1} (\pi_j - q \pi_{j+1})^2 - \frac{\alpha_s}{8\pi} \frac{\pi_N}{f} G_{\mu\nu} \tilde{G}^{\mu\nu} - \frac{\alpha_s}{8\pi} \frac{\pi_M}{f} F_{\mu\nu} \tilde{F}^{\mu\nu}$$



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> Particle spectrum: one massless NGB A and N pseudo NGBs with mass $\propto m$

> Couplings of axion:

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▪ Photon coupling exponentially enhanced compared to gluon coupling



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[Giudice, McCullough 16]

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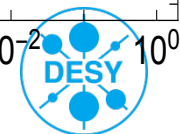
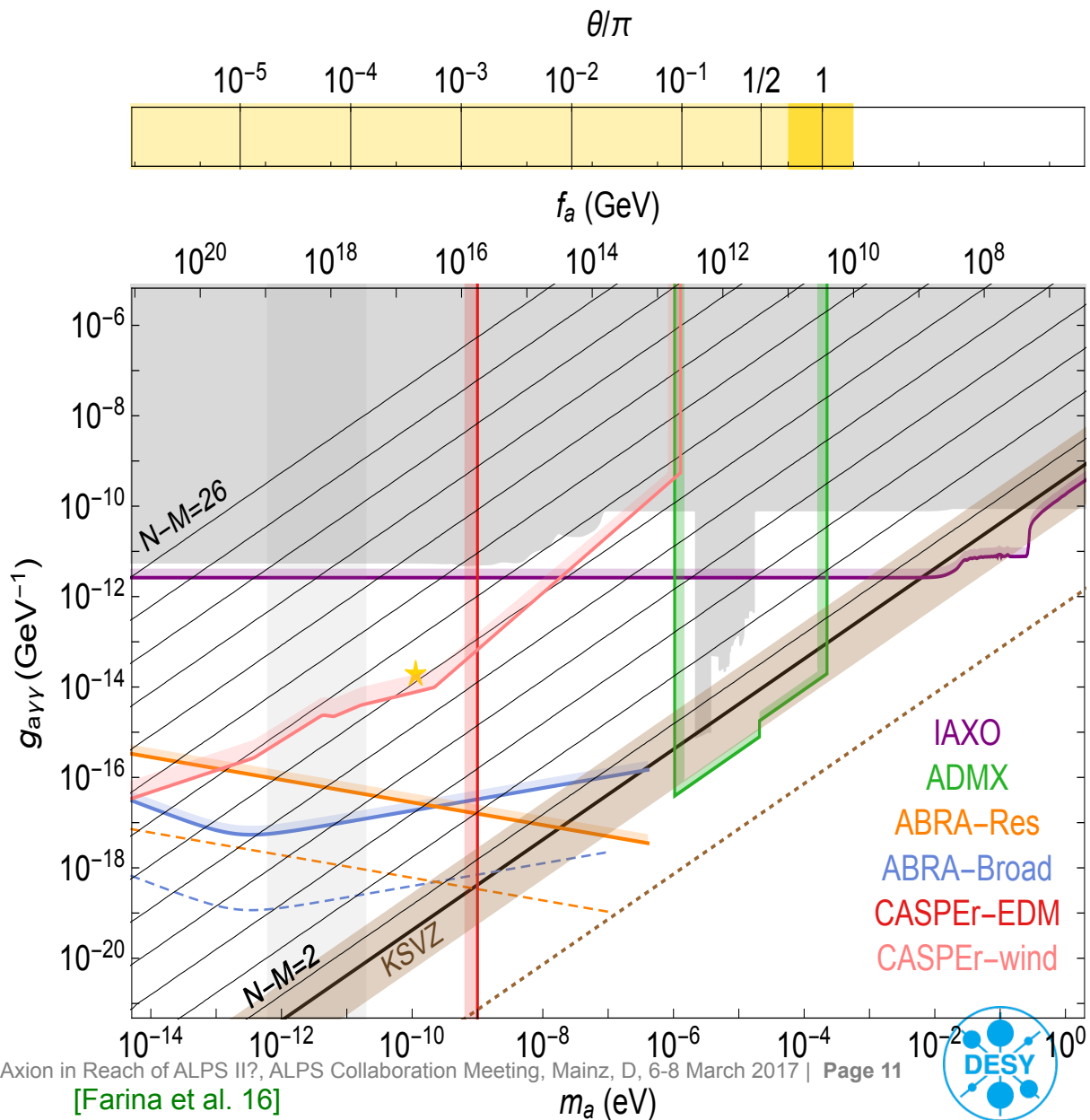
- Photon coupling exponentially enhanced compared to gluon coupling
- Boosted decay constant:

$$f_A \equiv \frac{q^N}{\mathcal{N}_0} f; \quad \mathcal{N}_0 \equiv \sqrt{\frac{q^2 - 1}{q^2 - q^{-2N}}}$$



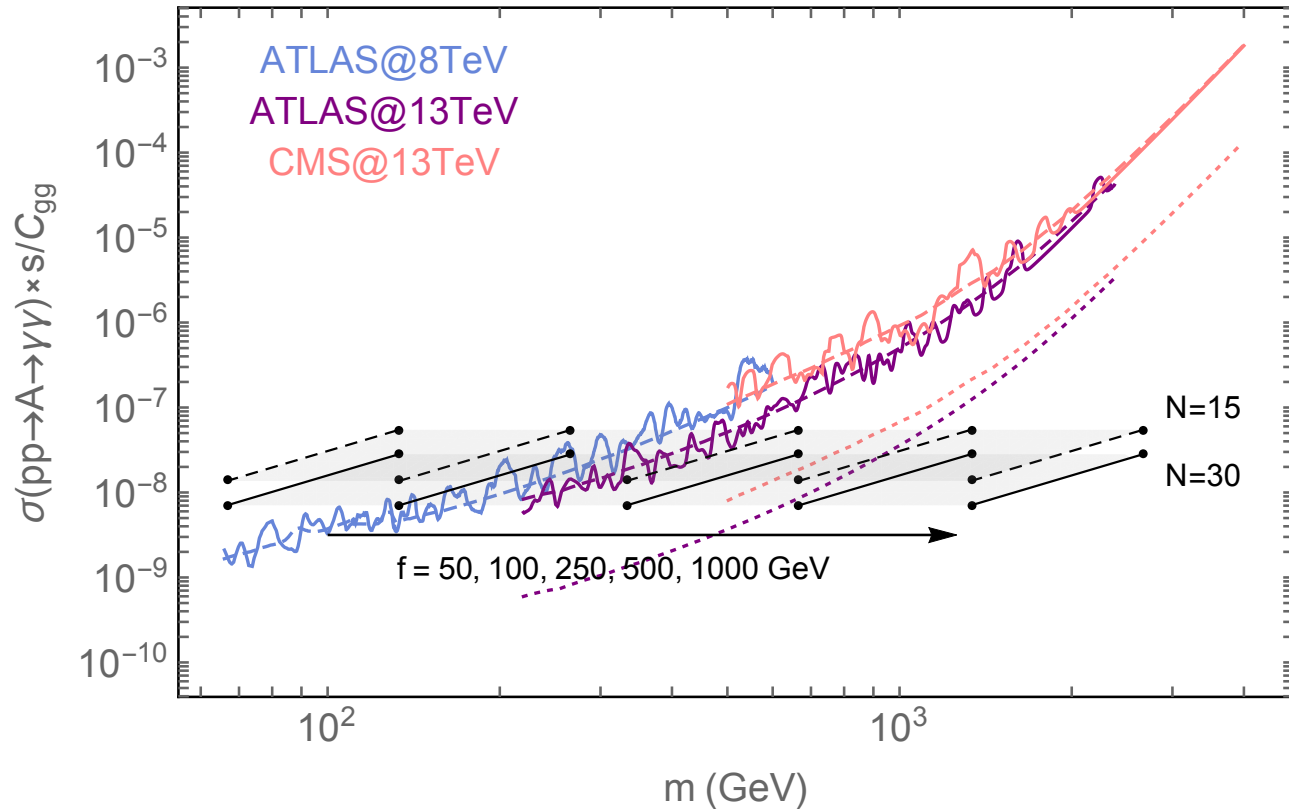
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 - Additional ALPs may be searched for at LHC



[Farina et al. 16]

