

On the gauge dependence of Quantum Electrodynamics

How to go beyond the Feynman gauge

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The Landau-Khalatnikov-Fradkin transform

$$G'(x - y) = G(x - y) \exp [ie(M(x - y) - M(0))]$$

- ▶ $M(x - y)$: describes the propagator of a scalar field & characterizes the gauge dependence
- ▶ gauge dependent terms factorize
- ▶ gauge dependent terms exponentiate



Renormalized self-energy of the electron

$$-\Sigma \left(\alpha, \xi, L = \ln \left[-p^2/\mu^2 \right] \right)_{\widetilde{\text{MOM}}}$$

=

$$\begin{aligned} & -\frac{1}{3!} \xi^3 L^3 & -\frac{3}{2} \xi L^2 & -\frac{3}{2} L & \left(\frac{\alpha}{4\pi} \right)^3 \\ & & +\frac{1}{2!} \xi^2 L^2 & +\frac{3}{2} L & \left(\frac{\alpha}{4\pi} \right)^2 \\ & & & -\frac{1}{1!} \xi L & \left(\frac{\alpha}{4\pi} \right)^1 \end{aligned}$$

Renormalized self-energy of the electron

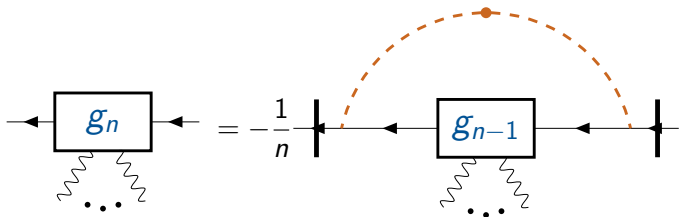
$$\begin{aligned}
 & -\Sigma\left(\alpha, \xi, L = \ln[-p^2/\mu^2]\right)_{\text{MOM}} \stackrel{?}{=} -\Sigma(\alpha, 0, L)_{\text{MOM}} \exp\left[-\xi L \left(\frac{\alpha}{4\pi}\right)\right] \\
 & = \\
 & + \frac{1}{4!} \xi^4 L^4 + \frac{3}{2} \frac{\xi^2}{2!} L^3 + \left(\frac{3}{2} \xi\right) L^2 \quad \left(\frac{\alpha}{4\pi}\right)^4 \\
 & \quad - \frac{1}{3!} \xi^3 L^3 \quad - \frac{3}{2} \xi L^2 \quad - \frac{3}{2} L \quad \left(\frac{\alpha}{4\pi}\right)^3 \\
 & \quad \quad + \frac{1}{2!} \xi^2 L^2 \quad + \frac{3}{2} L \quad \left(\frac{\alpha}{4\pi}\right)^2 \\
 & \quad \quad \quad - \frac{1}{1!} \xi L \quad \left(\frac{\alpha}{4\pi}\right)^1
 \end{aligned}$$

Renormalized self-energy of the electron

$$\begin{aligned}
 & -\Sigma\left(\alpha, \xi, L = \ln[-p^2/\mu^2]\right)_{\text{MOM}} \stackrel{?}{=} -\Sigma(\alpha, 0, L)_{\text{MOM}} \exp\left[-\xi L \left(\frac{\alpha}{4\pi}\right)\right] \\
 & = \\
 & + \frac{1}{4!} \xi^4 L^4 + \frac{3}{2} \frac{\xi^2}{2!} L^3 + \left(\frac{3}{2} \xi + \frac{9}{8}\right) L^2 + \left(\frac{1027}{8} + 400\zeta(3) - 640\zeta(5)\right) L \left(\frac{\alpha}{4\pi}\right)^4 \\
 & \quad - \frac{1}{3!} \xi^3 L^3 \quad - \frac{3}{2} \xi L^2 \quad - \frac{3}{2} L \left(\frac{\alpha}{4\pi}\right)^3 \\
 & \quad \quad + \frac{1}{2!} \xi^2 L^2 \quad + \frac{3}{2} L \left(\frac{\alpha}{4\pi}\right)^2 \\
 & \quad \quad \quad - \frac{1}{1!} \xi L \left(\frac{\alpha}{4\pi}\right)^1
 \end{aligned}$$

Characterization of the ξ dependence

$$G_c(\varepsilon, \alpha, \xi) = g_0(\varepsilon, \alpha) + g_1(\varepsilon, \alpha)\xi + g_2(\varepsilon, \alpha)\xi^2 + \dots$$



Beyond the Feynman gauge: an algorithm

- ❶ derive ε -expansion of the connected GF in the Feynman gauge

$$g_0(\varepsilon, \alpha) = c_1(\varepsilon)\alpha + c_2(\varepsilon)\alpha^2 + c_3(\varepsilon)\alpha^3 + \dots$$

- ❷ derive $F(\varepsilon, a)$ of the auxiliary graph



- ❸ $\xi^0 \Rightarrow \xi^1$ terms:

$$g_1(\varepsilon, \alpha) = F(\varepsilon, 1)\alpha + F(\varepsilon, 1 + \varepsilon)c_1(\varepsilon)\alpha^2 + F(\varepsilon, 1 + 2\varepsilon)c_2(\varepsilon)\alpha^3 + \dots$$

- ❹ for ξ^2, ξ^3, \dots terms: iterate step 3)

Higher dimensional QED

$$\mathcal{L}_{d=6} = -\frac{1}{4}(\partial_\mu F_{\nu\rho})(\partial^\mu F^{\nu\rho}) - \frac{1}{2\xi}(\partial_\mu \partial^\nu A_\nu)(\partial^\mu \partial^\rho A_\rho) + i\bar{\psi}\not{D}\psi$$

$$\begin{aligned}\mathcal{L}_{d=8} = & -\frac{1}{4}(\partial_\mu \partial_\nu F_{\rho\sigma})(\partial^\mu \partial^\nu F^{\rho\sigma}) - \frac{1}{2\xi}(\partial_\mu \partial_\nu \partial^\rho A_\rho)(\partial^\mu \partial^\nu \partial^\sigma A_\sigma) + i\bar{\psi}\not{D}\psi \\ & + \frac{g_2^2}{32}(F_{\mu\nu}F^{\mu\nu})(F_{\rho\sigma}F^{\rho\sigma}) + \frac{g_3^2}{8}F_{\mu\nu}F^{\nu\rho}F_{\rho\sigma}F^{\sigma\mu}\end{aligned}$$

(Gracey '16)

- ▶ photon propagator decomposes as usual
- ▶ photon self interactions are transversal
- ▶ ξ characterization applies across dimensions



The renormalized electron propagator

$$\frac{\partial}{\partial \xi} S_0(\alpha, \xi, L) = ie^2 \int \frac{d^D p}{(2\pi)^D} \frac{1}{[(q+p)^2]^2} \frac{1}{\not{q}} S_0(\alpha, \xi, \ln -q^2/\mu^2) \frac{1}{\not{q}}$$

- ▶ **IR divergence:** subtract counterterm which has same IR behaviour and respects $\widetilde{\text{MOM}}$ condition

- ▶ $\frac{\partial}{\partial \xi} S(\alpha, \xi, L) = L \frac{\alpha}{4\pi} S(\alpha, \xi, L)$

Proof of folklore statement

- ▶ $S(\alpha, \xi, L) = S(\alpha, 0, L) \exp \left[\xi L \frac{\alpha}{4\pi} \right]$

- ▶ $\gamma(\alpha, \xi) - \gamma(\alpha, 0) = \xi \frac{\alpha}{4\pi}$

The gauge dependence-infrared puzzle

gauge parameter

$$\text{QED} \quad \exp\left(\xi L \frac{\alpha}{4\pi}\right)$$

The gauge dependence-infrared puzzle

gauge parameter IR divergences

QED $\exp\left(\xi L \frac{\alpha}{4\pi}\right)$ $\exp\left(\text{diagram}\right)$

(Korthals Altes, de Rafael '76)

The gauge dependence-infrared puzzle

gauge parameter IR divergences

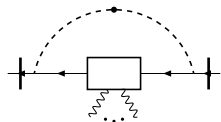
QED $\exp\left(\xi L \frac{\alpha}{4\pi}\right)$ $\exp\left(\text{triangle diagram}\right)$

(Korthals Altes, de Rafael '76)

QCD ? $\exp\left(\sum \text{webs}\right)$

Conclusion

- 1 simple characterization of ξ dependence



- 2 method applies to higher dimensional field theories

- 3 anomalous dimension of the electron depends on ξ at α^1 only

$$\gamma(\alpha, \xi) - \gamma(\alpha, 0) = \xi \frac{\alpha}{4\pi}$$

- 4 connects to IR divergences