

From the trees to the forest

the search for the origin of neutrino mass

Michael A. Schmidt

14 September 2017

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THE UNIVERSITY OF
SYDNEY



CoEPP
ARC Centre of Excellence for
Particle Physics at the Terascale

Origin of neutrino mass: Neutrino mass generation mechanisms

- Dirac vs. Majorana neutrinos

⇒ Majorana mass generated by Weinberg operator



$$\mathcal{L}_\nu = \frac{1}{2} \frac{\kappa}{\Lambda} LHLH + \text{h.c.}$$

- **Effective operator** $LHLH$ suppressed by $\Lambda \gg \langle H \rangle \simeq 100\text{GeV} \gg m_\nu$
- Can be generated via seesaw mechanisms, minimal UV completions

Minkowski; Yanagida; Glashow; Gell-Mann, Ramond, Slansky; Mohapatra, Senjanovic.

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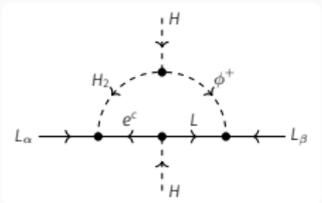


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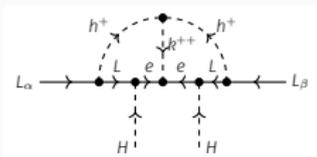
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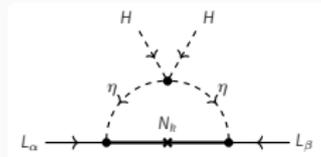
- **However not the only possibility to generate neutrino mass**



Zee (1980)



Zee (1986); Babu (1988) [Cheng, Li (1980)]



Ma hep-ph/0601225

- among many other possibilities

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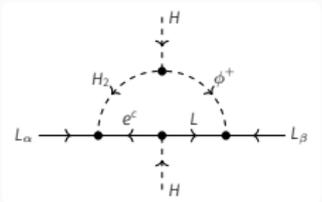


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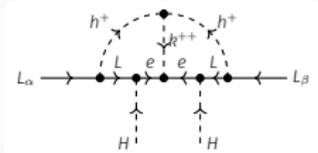
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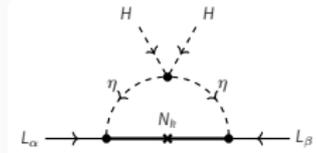
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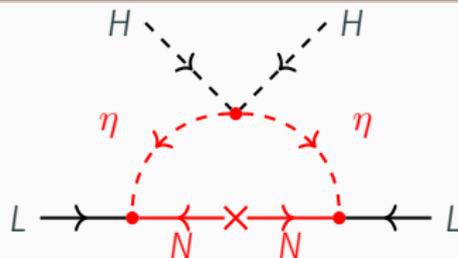
Ma hep-ph/0601225

- among many other possibilities

Why should I be interested in anything beyond the seesaw mechanisms?

Connection to other physics: Dark matter (DM)

- Scotogenic model Ma hep-ph/0601225
- Discrete symmetry $N \rightarrow -N, \eta \rightarrow -\eta$ to forbid tree-level contribution
- Lightest new particle is a DM candidate



$$N \sim (1, 1, 0) \quad \eta \sim (1, 2, \frac{1}{2})$$

- **Scalar DM:**

very similar to same as inert doublet model

Lopez-Honorez, Nezri, Oliver, Tytgat hep-ph/0612275

- **Fermionic WIMP DM:**

close connection with neutrino mass due to bounds from lepton flavour violation

recent study: Vicente, Yaguna 1412.2545; Lindner, Platscher, Yaguna, Merle 1608.00577

- **Fermionic FIMP DM:**

one neutrino (almost) massless. Otherwise DM phenomenology mostly decoupled from neutrino physics

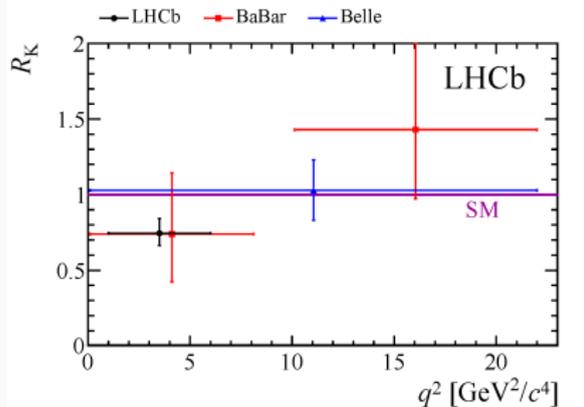
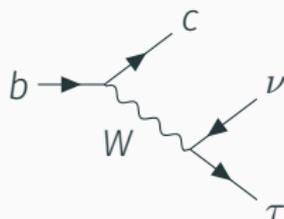
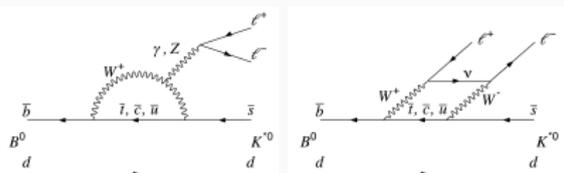
Molinaro, Yaguna, Zapata 1405.1259

Connection to other physics: B physics anomalies (1)

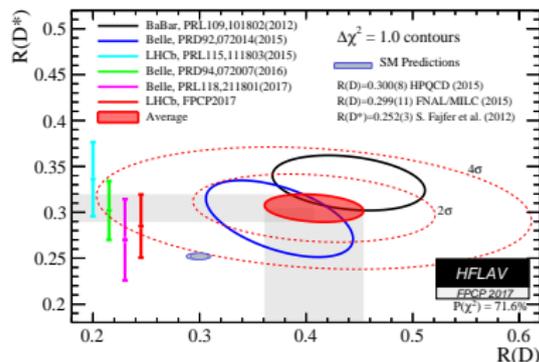
Hints for violations of LFU in $R_{K^{(*)}}$ and $R_{D^{(*)}}$

$$R_{K^{(*)}} = \frac{\Gamma(\bar{B} \rightarrow \bar{K}^{(*)} \mu^+ \mu^-)}{\Gamma(\bar{B} \rightarrow \bar{K}^{(*)} e^+ e^-)}$$

$$R_{D^{(*)}} = \frac{\Gamma(\bar{B} \rightarrow D^{(*)} \tau \bar{\nu})}{\Gamma(\bar{B} \rightarrow D^{(*)} l \bar{\nu})}$$



1406.6482

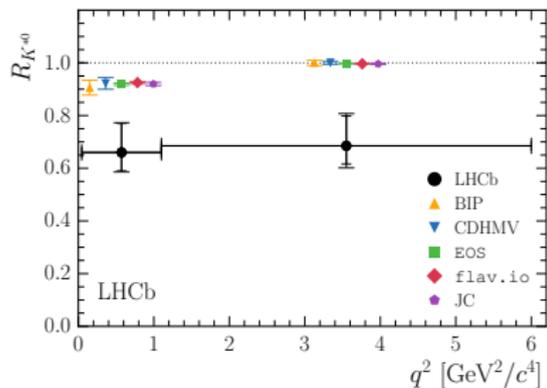
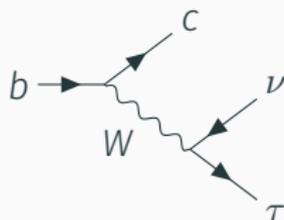
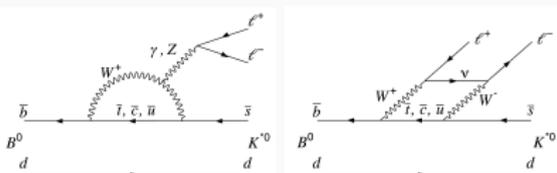


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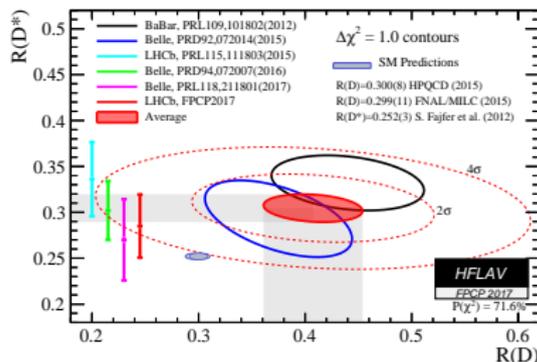
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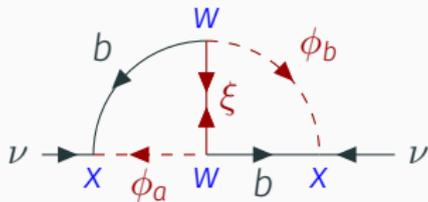
1705.05902



Connection to other physics: B-physics anomalies (3)

based on dimension-9 operator $\mathcal{O}_{11} = LLQd^cQd^c$ Angel, Cai, Rodd, MS, Volkas 1308.0463

Two LQs $\phi \sim (3, 1, -1/3)$ and Majorana fermion $\xi \sim (8, 1, 0)$

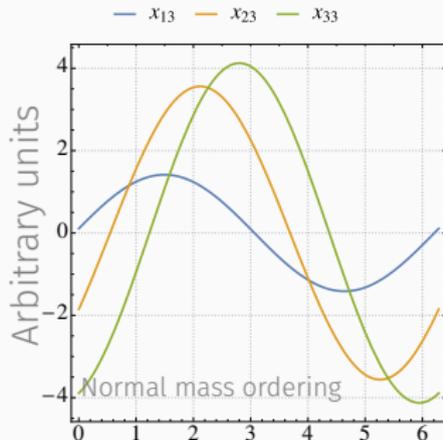


$$x = \begin{pmatrix} 0 & 0 & x_{13} \\ 0 & x_{22} & x_{23} \\ 0 & x_{32} & x_{33} \end{pmatrix}$$

Minimal scenario: only necessary to consider non-negligible w_{3a} (scale factor)

Important points:

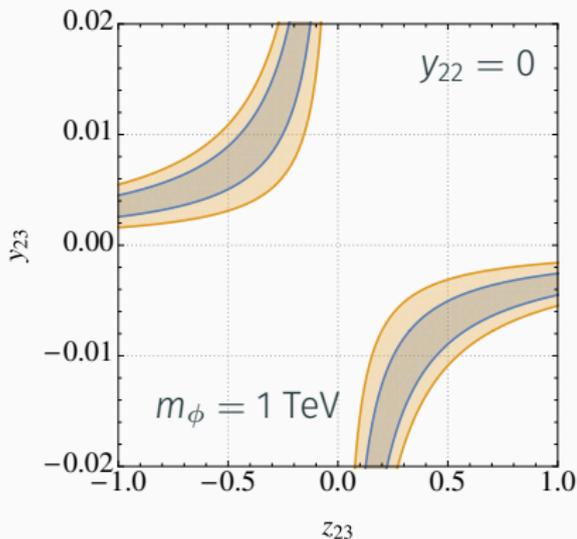
- x_{13} cannot be turned off *ad libitum*
 $\Rightarrow \mu N \rightarrow eN$ serious constraint
- inconsistent with hierarchy $|x_{23}| \gg |x_{33}|$
 needed for $R_{K^{(*)}}$ and $\tau \rightarrow \mu$ constraints



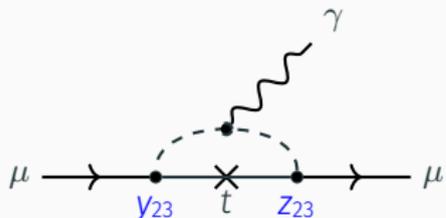
Cai, Gargalionis, MS, Volkas 1704.05849

Connections to other physics

- Anomalous magnetic moment $(g - 2)_\mu$
- New scalars can induce strong electroweak phase transition
- New bosons help with stability of electroweak vacuum
- baryogenesis, leptogenesis
- ...



$$\mathcal{L} \supset -z_{ij} e_L^i u_L^j \phi^\dagger + y_{ij} e_R^i u_R^j \phi + \text{h.c.}$$

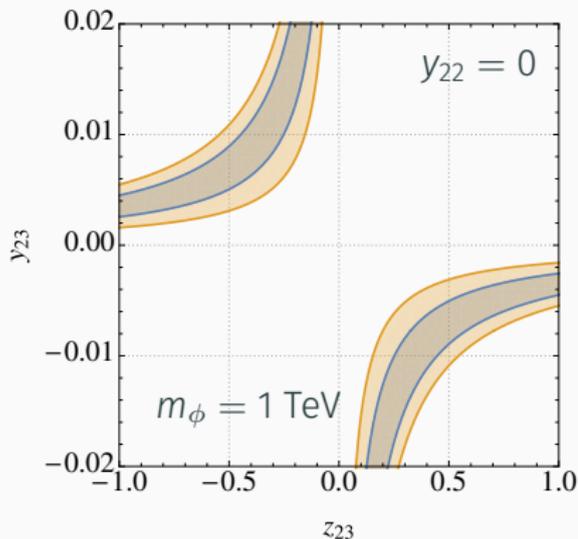


$$\text{Re}(y_{23}z_{23}) \approx \frac{-0.0039 \frac{m_\phi}{\text{TeV}}}{1 + 1.06 \ln \frac{m_\phi}{\text{TeV}}}$$

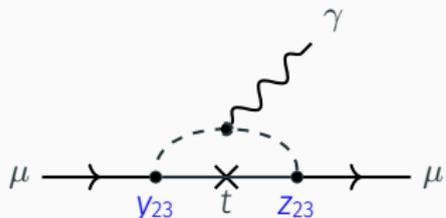
1 σ and 2 σ contours

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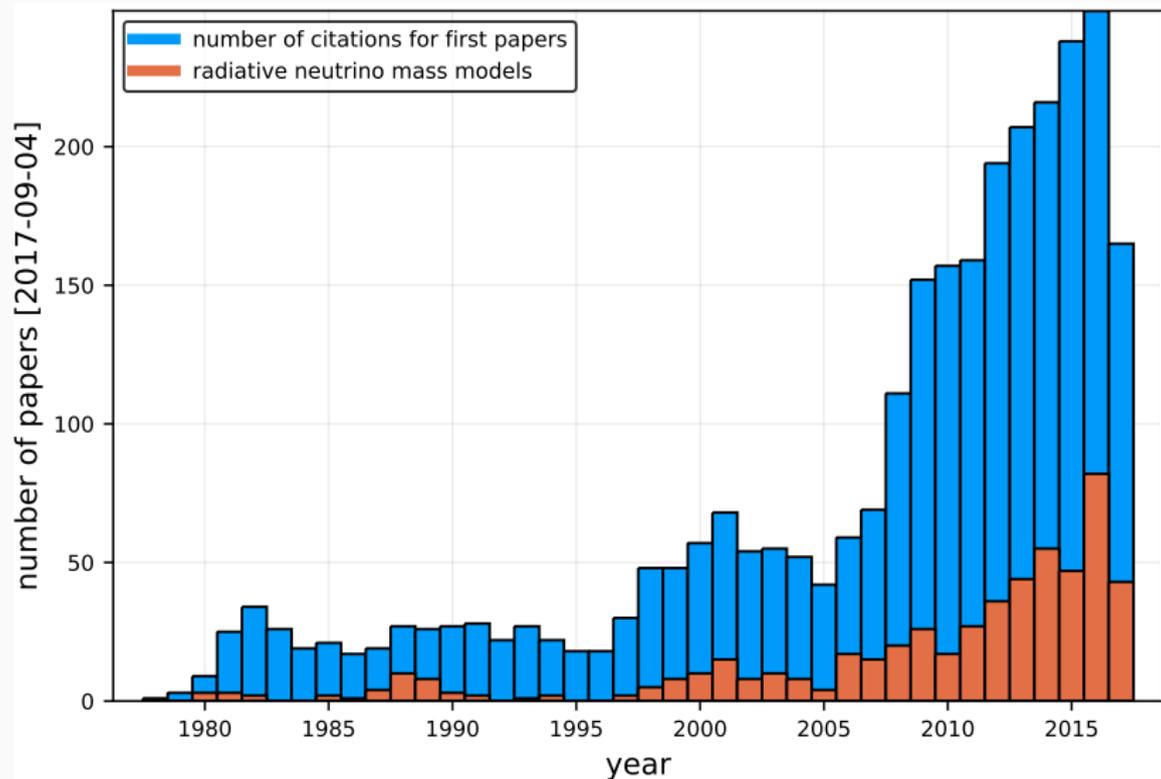


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1 σ and 2 σ contours

Interesting phenomenology testable in current/future experiments

Papers on radiative neutrino mass generation





Radiative transmission of lepton flavor hierarchies

Adisorn Adulpravitchai^a, Manfred Lindner^a, Alexander Merle^a, Rabindra N. Mohapatra^{b,*}

^a Max-Planck-Institut für Kernphysik
^b Maryland Center for Fundamental Physics

PHYSICAL REVIEW D **80**, 055031 (2009)

Confronting flavor symmetries and extended scalar sectors with lepton flavor violation bounds

Adisorn Adulpravitchai,^{*} Manfred Lindner,[†] and Alexander Merle[‡]

PHYSICAL REVIEW D **87**, 033006 (2013)

Lepton flavor at the electroweak scale: A complete A_4 model

Martin Holthausen,^{1,*} Manfred Lindner,^{1,†} and Michael A. Schmidt^{2,‡}

PHYSICAL REVIEW D **89**, 013007 (2014)

Dark matter and $U(1)'$ symmetry for the right-handed neutrinos

Manfred Lindner,^{*} Daniel Schmidt,[†] and Atsushi Watanabe[‡]

PHYSICAL REVIEW D **94**, 115027 (2016)

Fermionic WIMPs and vacuum stability in the scotogenic model

Manfred Lindner,^{*} Moritz Platscher,[†] and Carlos E. Yaguna[‡]

Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany

Alexander Merle[§]

ARTICLE

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1. Introduction

One of the standard model (SM) [1]. A simple introduces three masses addition for the light



A systematic approach: generalized Weinberg operator

Consider operators of type

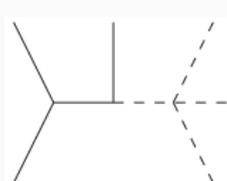
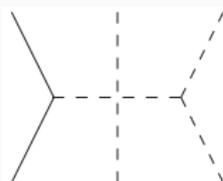
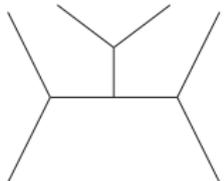
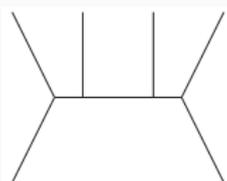
$$LLHH(H^\dagger H)^n$$

possibly with multiple Higgs fields

Bonnet, Hernandez, Ota, Winter 0907.3143

Construct all possible topologies:

- **tree-level topologies** Bonnet, Hernandez, Ota, Winter 0907.3143
- 1-loop topologies of Weinberg operator Bonnet, Hirsch, Ota, Winter 1204.5862
- 2-loop topologies of Weinberg operator Arisitizabal Sierra, Degee, Dorame, Hirsch 1411.7038
- 1-loop topologies of dimension-7 operator Cepedello, Hirsch, Helo 1705.01489



dimension-7 operator at tree-level



electroweak
triplet fermion
quadruplet scalar

Babu, Nandi, Tavartkiladze 0905.2710

The dashed lines always denote scalars and solid lines are either fermions or scalars.

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A systematic approach: $\Delta L = 2$ operators

- Black box theorem: Every $\Delta L = 2$ operator lead to neutrino mass

Schechter, Valle Phys. Rev. D25 (1982) 2951

dimension	5	7	9	11
field strings ¹	1	6	21	101
Lorentz structures ²	2	22	368	6632

¹no gauge fields, no Lorentz structure, no products of SM singlets (e.g. $LHLHH^\dagger H$)

²includes hermitean conjugates

- Consider all possible $\Delta L = 2$ operators Babu, Leung hep-ph/0106054; deGouvea, Jenkins 0708.1344
- UV completions Angel, Rodd, Volkas 1212.6111
- Also describes other $\Delta L = 2$ violating processes: neutrinoless double beta decay, LNV processes at a collider
- Indication of quantum numbers of new particles

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Other criteria: topology, complexity, flavour, common features, ...

See review for more details ...

Y. Cai, J. Herrero-Garcia, M.S. A. Vicente, R. Volkas [1706.08524]

From the trees to the forest: a review of radiative neutrino mass models

Yi Cai,^{a,b} Juan Herrero-García,^c Michael A. Schmidt,^d Avelino Vicente^e and Raymond R. Volkas^b

^aSchool of Physics, Sun Yat-sen University, Guangzhou, 510275, China

^bARC Centre of Excellence for Particle Physics at the Terascale, School of Physics, The University of Melbourne, VIC 3010, Australia

^cARC Centre of Excellence for Particle Physics at the Terascale, Department of Physics, The University of Adelaide, SA 5005, Australia

^dARC Centre of Excellence for Particle Physics at the Terascale, School of Physics, The University of Sydney, NSW 2006, Australia

^eInstituto de Física Corpuscular (CSIC-Universitat de València), Aptdo. 22085, E-46071 Valencia, Spain

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m.a.schmidt@sydney.edu.au, avelino.vicente@ific.uv.es,
r.volkas@unimelb.edu.au

Explanation for the lightness of neutrino masses is that neutrinos (typically Majorana) being generated radiatively at high energies and they are typically at the TeV scale. The suppression coming from the loop factor $1/16\pi^2$ is appealing. In particular, the leptonic flavor asymmetry is dependent on

Minimal UV completions of $\Delta L = 2$ dimension-7 operators

Y. Cai, J. Clarke, MS, R. Volkas 1410.0689

Any $\Delta L = 2$ operator induces Majorana mass term for neutrinos

Effective $\Delta L = 2$ operators of dimension 7

$$\mathcal{O}'_1 = LL\tilde{H}HHH$$

$$\mathcal{O}_2 = LLL\bar{e}H$$

$$\mathcal{O}_3 = LLQ\bar{d}H$$

$$\mathcal{O}_4 = LLQ^\dagger\bar{u}^\dagger H$$

$$\mathcal{O}_8 = L\bar{d}\bar{e}^\dagger\bar{u}^\dagger H$$

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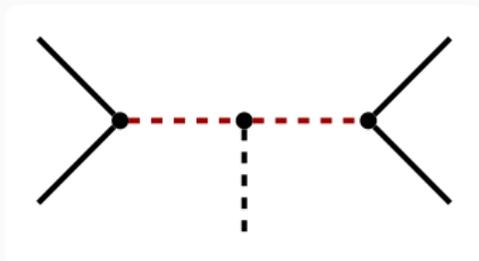
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Scalars: leptoquarks, singly charged scalars, EW doublets and quartets

Fermions: vector-like quarks/charged leptons mixing with third generation

Scalar	Scalar	Operator
$(1, 2, \frac{1}{2})$	$(1, 1, 1)$	$\mathcal{O}_{2,3,4}$
$(3, 2, \frac{1}{6})$	$(3, 1, -\frac{1}{3})$	$\mathcal{O}_{3,8}$
$(3, 2, \frac{1}{6})$	$(3, 3, -\frac{1}{3})$	\mathcal{O}_3

Leptoquarks $(3, 2, \frac{1}{6})$ and $(3, 1, -\frac{1}{3})$ used to explain R_K (and R_D)

Päs, Schumacher 1510.08757 Deppisch, Kulkarni, Päs, Schumacher 1603.07672

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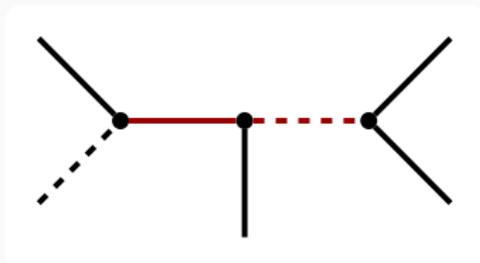
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Dirac fermion	Scalar	Operator
$(1, 2, -\frac{3}{6})$	$(1, 1, 1)$	\mathcal{O}_2
$(3, 2, -\frac{1}{6})$	$(1, 1, 1)$	\mathcal{O}_3
$(3, 1, \frac{2}{6})$	$(1, 1, 1)$	\mathcal{O}_3
$(3, 1, \frac{1}{6})$	$(3, 2, \frac{1}{6})$	\mathcal{O}_3
$(3, 2, -\frac{1}{6})$	$(3, 1, -\frac{1}{3})$	$\mathcal{O}_{3,8}$
$(3, 2, -\frac{1}{6})$	$(3, 3, -\frac{1}{3})$	\mathcal{O}_3
$(3, 3, \frac{2}{6})$	$(3, 2, \frac{1}{6})$	\mathcal{O}_3
$(3, 2, \frac{7}{6})$	$(1, 1, 1)$	\mathcal{O}_4
$(3, 1, -\frac{1}{3})$	$(1, 1, 1)$	\mathcal{O}_4
$(3, 2, \frac{7}{6})$	$(3, 2, \frac{1}{6})$	\mathcal{O}_8
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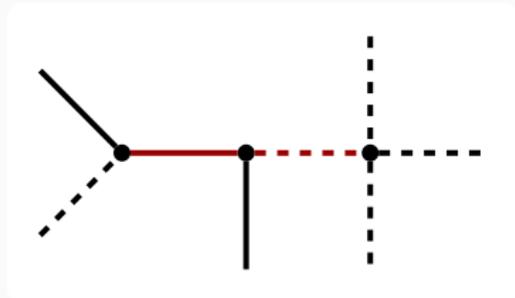
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Fermions: vector-like quarks/charged leptons mixing with third generation

Dirac fermion	Scalar	Operator
$(1, 3, -1)$	$(1, 4, \frac{3}{2})$	\mathcal{O}'_1

New particles at the LHC: (scalar) leptoquarks

Decay channels

$$\phi \rightarrow \ell q \quad \phi \rightarrow \nu q$$

Assuming 100% branching ratio

- first generation LQ (e): $m_{LQ} \gtrsim 1130$ GeV CMS-PAS-EXO-16-043
- second generation LQ (μ): $m_{LQ} \gtrsim 1165$ GeV CMS-PAS-EXO-16-007
- third generation LQ (τ): $m_{LQ} \gtrsim 900$ GeV CMS-PAS-EXO-16-023

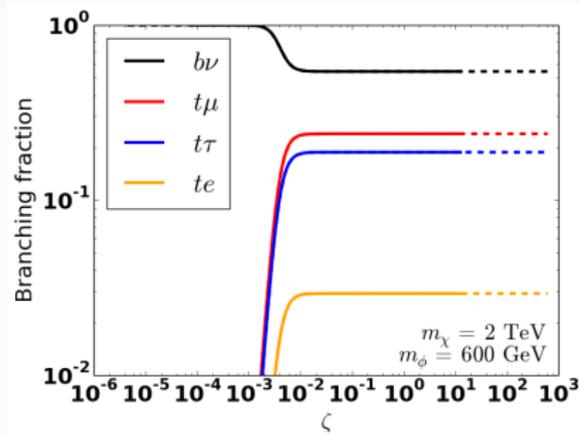
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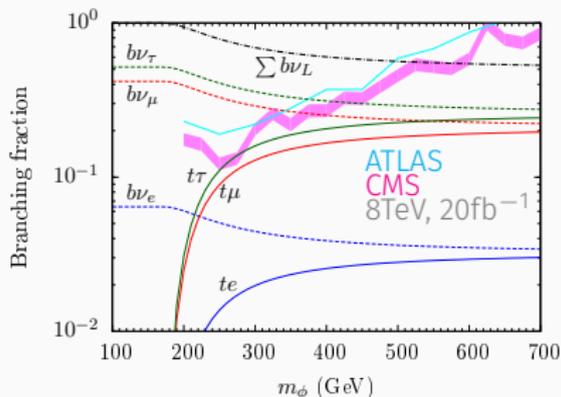
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Cai, Clarke, MS, Volkas 1410.0689



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1-loop model (O_{3b}): ζ controls relative size of Yukawa couplings (like Casas-Ibarra parameter)

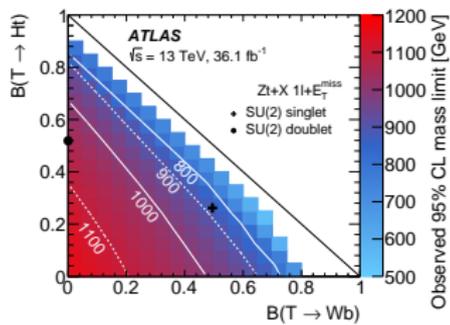
New particles at the LHC: vector-like quarks

B	T	(BY)	(XT)	(XTB)
$(3, 1, -\frac{1}{3})$	$(3, 1, \frac{2}{3})$	$(3, 2, -\frac{5}{6})$	$(3, 2, \frac{7}{6})$	$(3, 3, \frac{2}{3})$

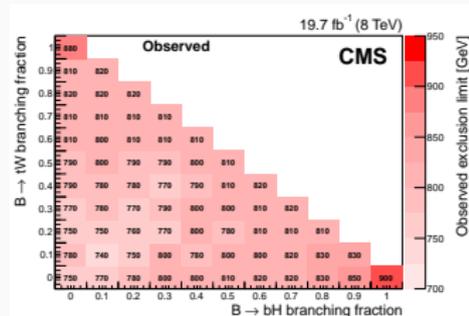
- Searched for at LHC looking using pair production
- Main decay channels to EW gauge bosons and Higgs

$$\begin{array}{lll}
 B: & B \rightarrow W^- t, & B \rightarrow Zb, & B \rightarrow Hb, \\
 T: & T \rightarrow W^+ b, & T \rightarrow Zt, & T \rightarrow Ht,
 \end{array}$$

$$\begin{array}{l}
 m_T \gtrsim 0.5 - 1.2 \text{ TeV} \\
 m_B \gtrsim 0.8 - 0.9 \text{ TeV}
 \end{array}$$



ATLAS 1705.10751



CMS 1507.07129

ATLAS 1504.04605 1409.5500 1505.04306 1705.10751 1606.03903 1509.04261 ATLAS-CONF-2016-032 CMS 1509.04177 1706.03408 1311.7667 1507.07129

New particles at the LHC: uncolored particles

Vector-like leptons

E	(NE)	(ED)	(NED)
$(1, 1, -1)$	$(1, 2, -\frac{1}{2})$	$(1, 2, -\frac{3}{2})$	$(1, 3, -1)$

currently no dedicated search at LHC

Uncolored scalars

- Charged scalars
 - doubly charged scalar: like-sign dilepton pairs (LNV)
 - singly charged scalar: similar to slepton pair production search
- Higher-dimensional EW multiplets (doublet, quadruplet, ...)
 - Drell-Yan production of charged scalar
 - possibly long lifetime, if small mass splitting
 - disappearing track signature

Conclusions

radiative neutrino mass interesting possibility

current experiments probe theory space

⇒ clear phenomenological signatures required

plethora of models

⇒ systematic approach needed

radiative neutrino mass interesting possibility

current experiments probe theory space

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plethora of models

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Thank you!