### (Towards) CLS simulations at physical pion mass

Daniel Mohler

Zeuthen, April 11th, 2017





#### Representing the CLS effort With plots from Hubert Simma, Jakob Simeth

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## Outline

Introduction - The CLS 2+1 flavor ensembles

- The CLS 2+1 flavor ensembles Key features
- Landscape of CLS ensembles
- Towards the physical point
  - Autocorrelation times towards the physical point
  - Statistical uncertainty towards physical light-quark masses
- Simulations and challenges: "X200"
  - Thermalization strategy
  - Current status

#### Conclusions and Outlook

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#### Coordinated Lattice Simulations - Members

- Berlin (NIC/DESY-Zeuthen/HU Berlin)
- CERN
- Mainz
- Madrid
- Münster
- Odense/ CP3-Origins
- Regensburg
- Rome (Roma I, Roma II)
- Wuppertal

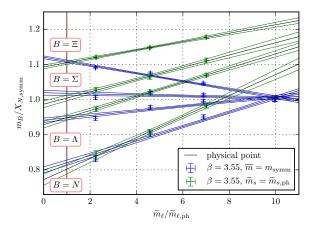
## The CLS 2+1 flavor ensembles – Key features

- Open boundary conditions to avoid topological freezing for  $a \rightarrow 0$
- Twisted mass reweighting (for the light quarks)
- Simulation along trajectory with fixed Tr(M)
- Additional simulations along trajectories with fixed strange quark mass  $m_s = \text{const.}$  and with  $m_s = m_l$
- Flexible simulations with OpenQCD

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http://luscher.web.cern.ch/luscher/openQCD/
```

- Nested hierarchical integrators
- Hasenbusch-style mass preconditioning with an arbitrary number of pseudofermion pairs
- Rational approximation (+ reweighting) for the strange quark
- Deflation acceleration and chronological solver
- A number of solvers

## Baryon masses: trajectory with fixed Tr(M) vs. $m_s = \text{const.}$



plot from Bali et al. RQCD, arXiv:1702.01035

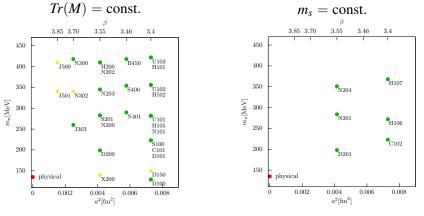
• Example of octet baryon masses at  $a \approx 0.064$  fm (from RQCD)

• Illustrates typical behavior

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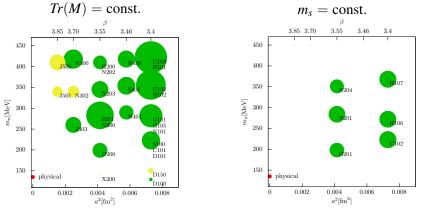
## CLS 2+1 flavor ensembles: Overview



plots by Jakob Simeth, RQCD

- Letters in the name denote the aspect ratio T/L; First digit encodes  $\beta$
- Ensembles at 5 lattice spacings and with a range of  $M_{\pi} \leq 420 \text{MeV}$
- Ensembles to control (or exploit) finite volume effects

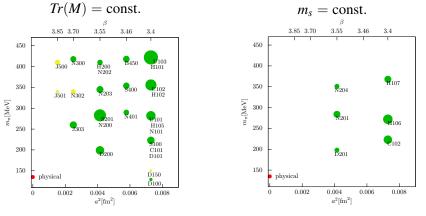
## CLS 2+1 flavor ensembles: Statistics



plots by Jakob Simeth, RQCD

- > 4000 MDU for many ensembles Typically save 1 configuration every 4 MDU
- target statistics chosen considering largest  $\tau_{int}$  (YM action density)

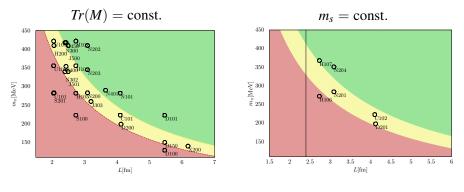
## CLS 2+1 flavor ensembles: Statistics



plots by Jakob Simeth, RQCD

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#### CLS 2+1 flavor ensembles: Volumes used



plots by Jakob Simeth, RQCD

- red:  $m_{\pi}L \leq 4$ ; yellow:  $4 \leq m_{\pi}L \leq 5$ ; green  $5 \leq m_{\pi}L$
- Most ensembles with  $m_{\pi}L \geq 4$
- Some smaller volumes to check finite size effects

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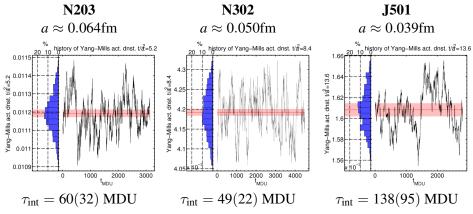
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#### Autocorrelation towards the continuum limit

Action density at  $t_0$  as defined by  $t^2 \langle E \rangle = 0.3$ 



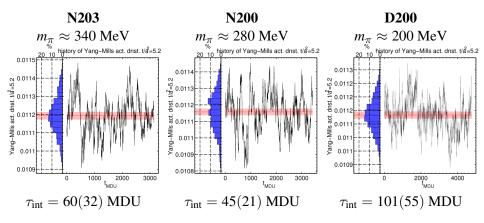
- Autocorrelation time is expected to increase significantly
- Uncertainty is still sizable

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#### Autocorrelation towards physical quark masses

Action density at  $t_0$  as defined by  $t^2 \langle E \rangle = 0.3$ 



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#### Noise/Signal at light quark masses - Introduction

• For the nucleon we have (argument by Parisi, Lepage)

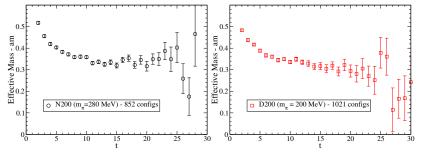
$$N\sigma_{N,\mathbf{p}=0}^{2} = \left\langle C_{N}(\mathbf{p}=0,t;m)^{2} \right\rangle - \left\langle C_{N}(\mathbf{p}=0,t;m) \right\rangle^{2}$$
$$\propto Z_{3\pi} e^{-3m_{\pi}t} + Z_{N}^{2} e^{-2m_{N}t}$$

• The noise to signal ratio therefore degrades exponentially

$$rac{\sigma_N(t)}{\langle C_N(t)
angle}\simeq rac{1}{\sqrt{N}}\mathrm{e}^{\left(m_N-rac{3}{2}m_\pi
ight)t}$$

• Similar argument for Nuclei, heavy mesons, etc.

## Nucleon effective masses

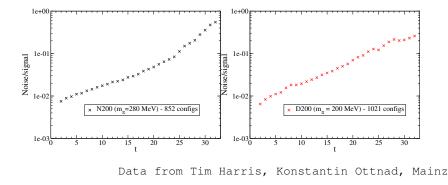


Data from Tim Harris, Konstantin Ottnad, Mainz

#### Setup

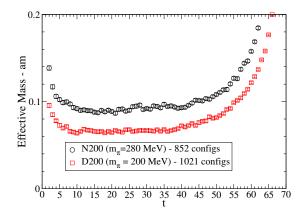
- All-mode-averaging (AMA)
  - 12 ( $n_c \times n_D$ ) exact inversions and 16  $\times$  12 sloppy inversions
- Results from sources in a single timeslice
- Effective mass from the local-smeared correlator

## Nucleon noise/signal



- Slope in (most of) plateau region does not reach asymptotic value (given by  $m_N \frac{3}{2}m_\pi$ )
- Suggests that in practice noise/signal scaling is not as severe
- Exponential growth qualitatively observed

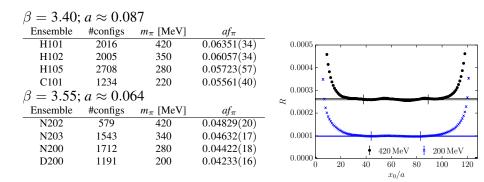
## Pion effective masses



- Strong effects from open boundary visible
- These are well understood
- There are plenty of usable timeslices
- Note: Thermal effects (with periodic bc) can also be a nuisance

#### Pion decay constant

Results from Bruno, Korzec, Schaefer, arXiv:1608.08900



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Simulations and challenges: "X200"

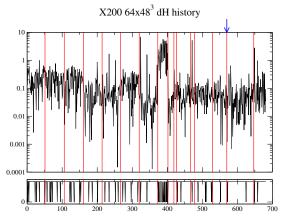
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#### "X200" –Description and thermalization strategy

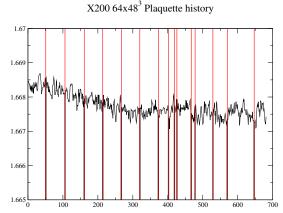
- Physical ud/s run at  $\beta = 3.55$  ( $a \approx 0.064$  fm)
- To keep  $m_{\pi}L \ge 4$ :  $L^3 \times T = 96^3 \times 192$
- Thermalization strategy:
  - Start from an SU(3) run with 3 light quarks and periodic boundary conditions
  - Perform a number of runs to thermalize this small volume  $(L^3 \times T = 48^3 \times 64)$
  - **③** Triple the time extent  $48^3 \times 64 \rightarrow 48^3 \times 192$
  - **(**) Double the spatial extent  $48^3 \times 192 \rightarrow 96^3 \times 192$
- At this fairly coarse lattice spacing periodic boundary conditions are chosen

#### Small volume run - dH history



Acceptance history for the thermalization of a small volume physical quark mass run. The vertical red lines indicate changes of run parameters. The arrow indicates where the target quark masses have been reached.

#### Small volume run - Plaquette history



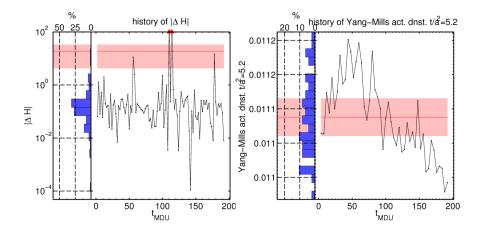
Plaquette history for the thermalization of a small volume physical quark mass run.

• Overall no significant difficulty in (partially) thermalizing the small volume

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# Intermediate run of size $L^3 \times T = 48^3 \times 192$



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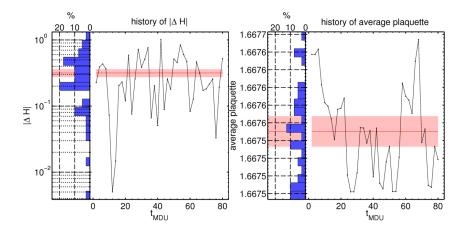
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#### Challenges

- Runs with  $L^3 \times T = 48^3 \times 64$  proceeded smoothly
- Runs at intermediate volume needed various minor parameter adjustments (more frequent updates of the deflation subspace)
- Runs with  $L^3 \times T = 96^3 \times 192$ 
  - Run only stable with large deflation blocksize  $6^4 \rightarrow 6 \times 4 \times 8^2 \rightarrow 8 \times 4 \times 8^2$
  - Large deflation blocksize was needed in order to maintain a manageable size of the little Dirac operator
    - $\rightarrow$  Iteration counts higher than desirable/ deflation not as efficient
  - Indicates that a multigrid setup with 3 levels might be preferable for this lattice volume (but not obvious that it would pay off)
  - Even larger lattices would likely need further algorithmic improvements

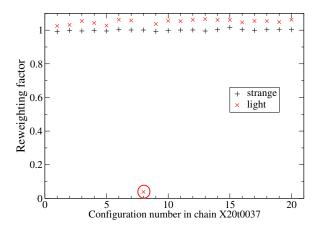
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## JUQUEEN $\rightarrow$ Cluster MOGON II (JGU Mainz)



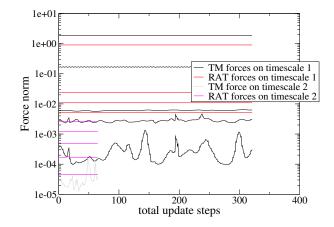
- First stable run not yet fully thermalized
- Run uses local partition of size  $24 \times 8^3$  and 692 nodes/ 13824 cores
- Made possible by early usage time on Mogon II

## JUQUEEN $\rightarrow$ Cluster MOGON II (JGU Mainz)



• Reweighting factors for first stable run; mostly small fluctuations

## Further improvements: A look at the force norms

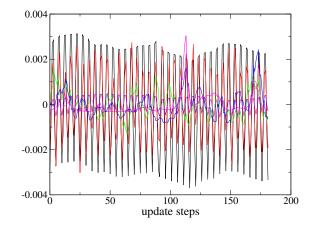


- integrators: lvl0 and lvl1: 4th order Omelyan integrator; lvl2: 2nd order Omelyan integrator
- lvl 2 forces are updated less often
- Lead to a slight adjustment in Hasenbusch masses

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## After some initial changes: A look at the norm fluctuations



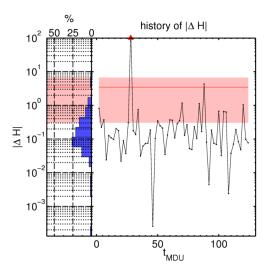
- after various tests (rearrangements of forces, further Hasenbusch masses, etc.)
- It seems that observed dH is largely driven by the force fluctuations

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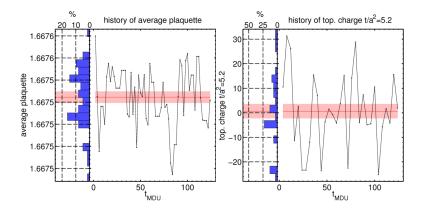
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#### Plots from the current run: dH

Acceptance: 0.790(53) Shown: 124 MDU, 31 configurations Completed (as of today): 188 MDU, 47 configurations



## Plots: average Plaquette & topological charge



- Proper analysis will need a much longer chain
- Looks very promising

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- Large library of CLS 2+1 flavor ensembles
- Many physics studies started (and a number close to publication)
- Stable run at (very close to) physical  $m_l, m_s$  with  $a \approx 0.064$  fm
- Not enough statistics for a detailed study of autocorrelation, pion masses, nucleon masses, decay constants, ...
- might already profit from a lvl3 multigrid setup at the current lattice volume



## Thank you!

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