

# Axion dark matter minicluster

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# QCD Axion

review talk by A. Ringwald

- zero-T mass fixed by  $f_{\text{PQ}}$

$$m(T = 0) \approx 5.7 \times 10^{-5} \text{ eV} \frac{10^{11} \text{ GeV}}{f_{\text{PQ}}}$$

- periodic potential

$$V(\theta, T) = m^2(T)(1 - \cos \theta)$$

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- below I will adopt the harmonic approximation

$$V(\theta, T) \approx \frac{1}{2} m^2(T) \theta^2$$

# Axion minicluster

Hogan, Rees, 1988; Kolb, Tkachev, 93, 94, 95

- PQ symmetry breaking after inflation: Axion field takes random values in different Hubble volumes
- O(1) density fluctuations when Axion mass switches on
- expect gravitationally bound objects with size  $\sim$  Hubble volume @ QCD PT

$$M \sim \frac{4\pi}{3} d_H^3(T_{\text{osc}}) \bar{\rho}(T_{\text{osc}}) \quad d_H \sim 1/H$$

$$M \sim 10^{-12} M_{\odot} (f_{\text{PQ}}/10^{11} \text{ GeV})^2$$

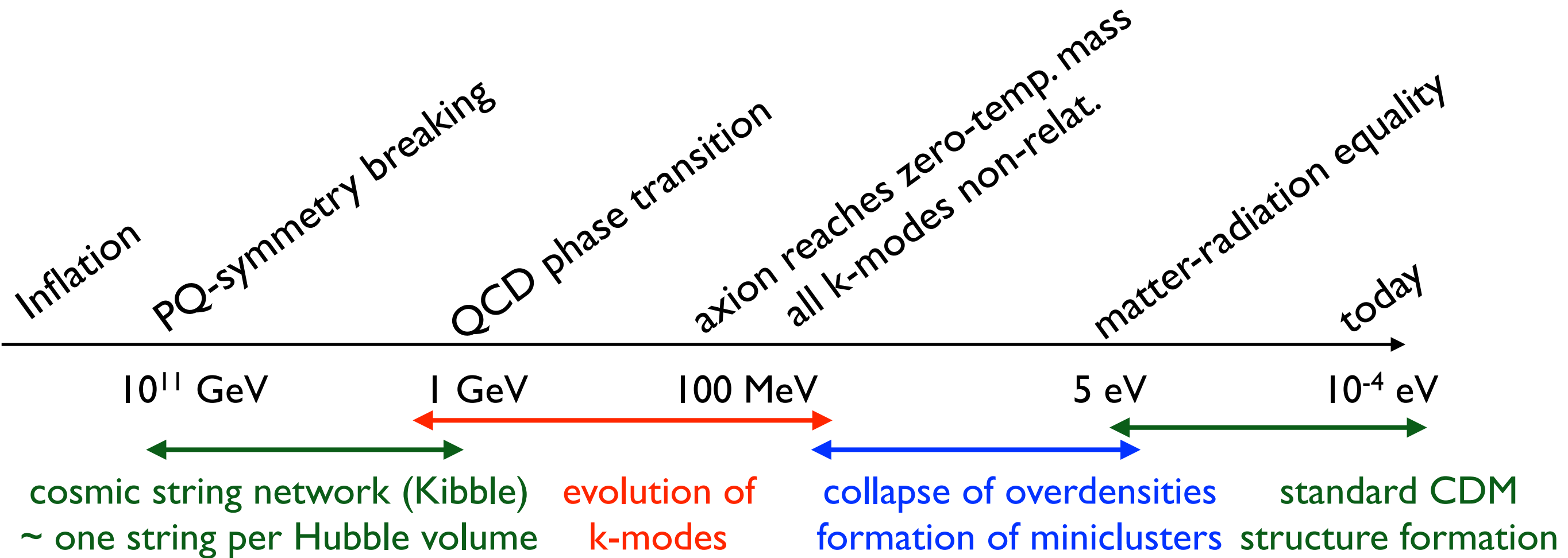
# Axion minicluster

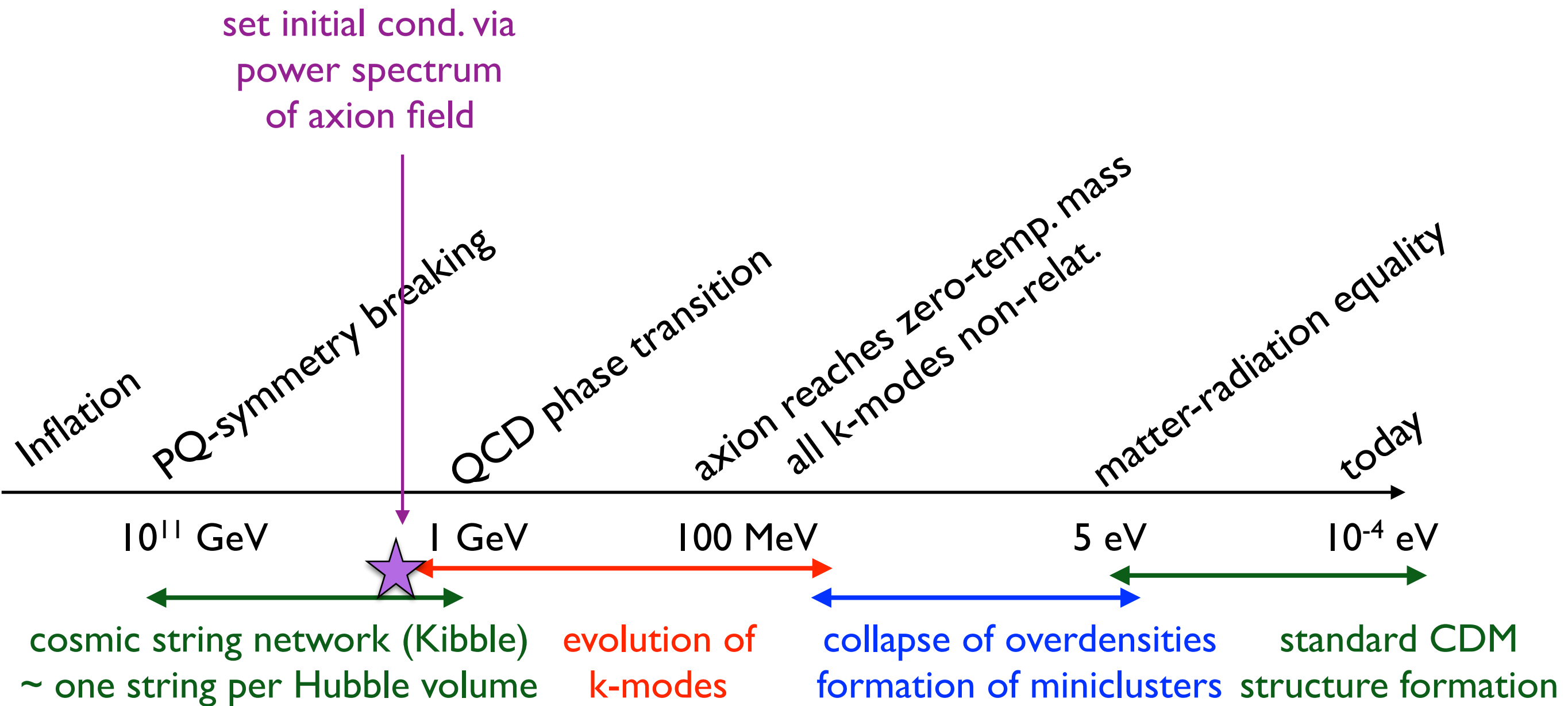
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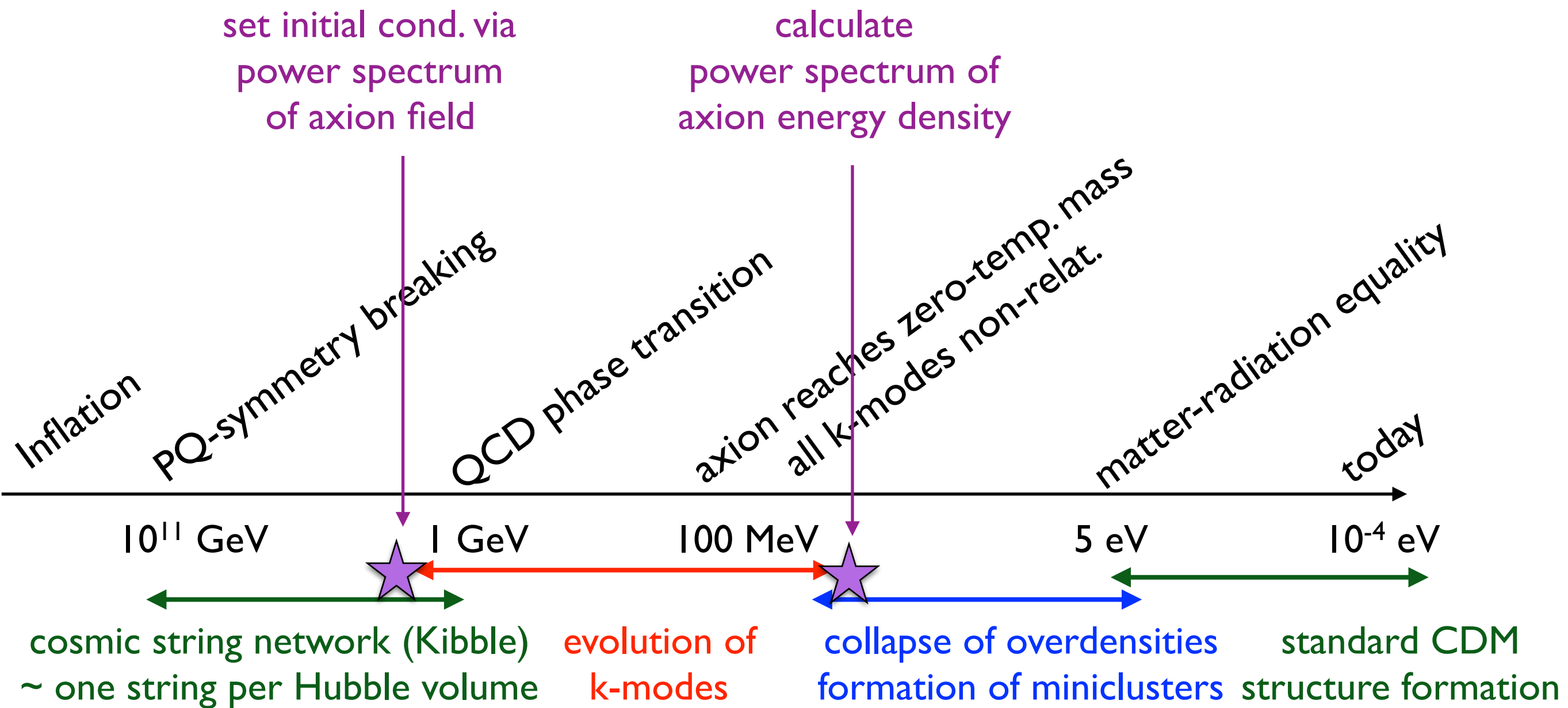
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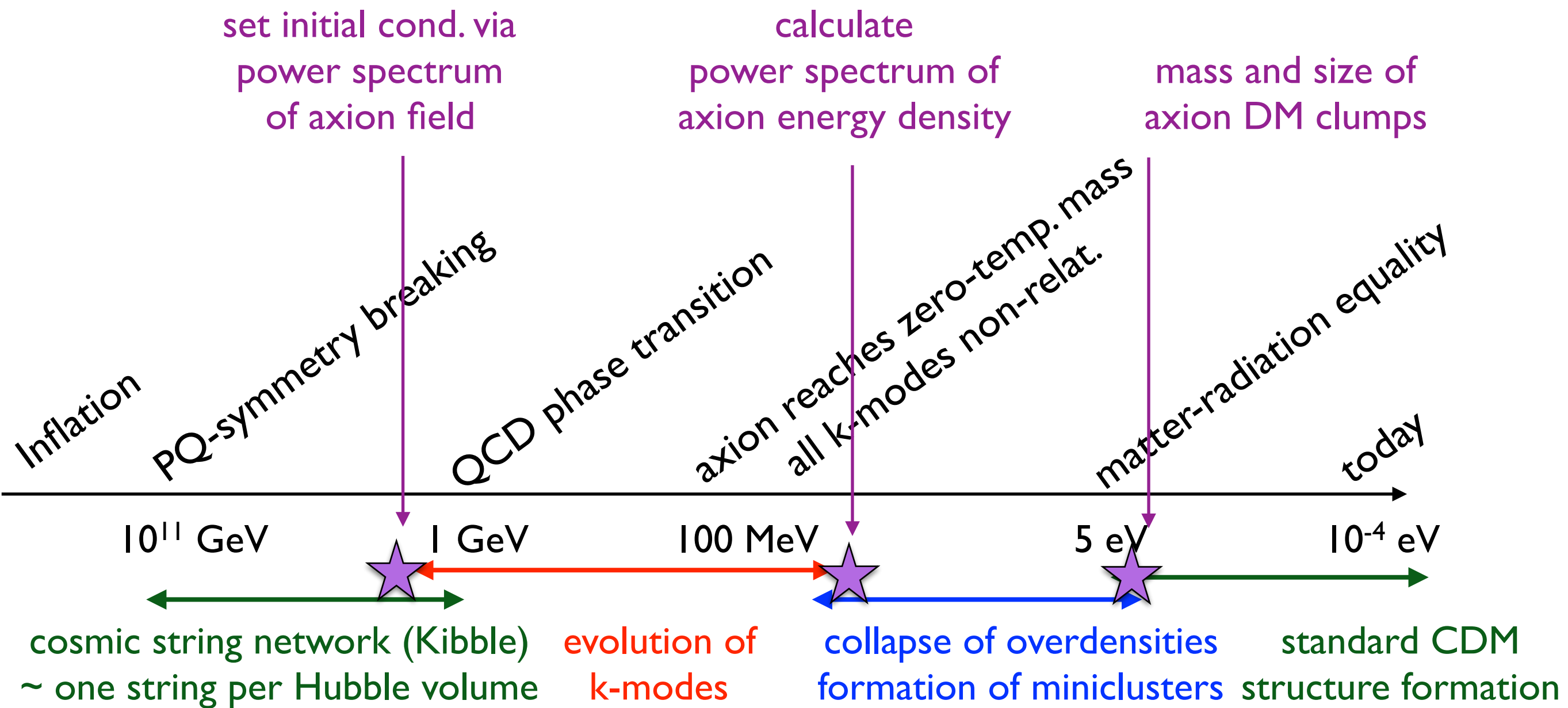
$$M \sim 10^{-12} M_{\odot} (f_{\text{PQ}}/10^{11} \text{ GeV})^2 \quad \rightarrow \text{see later}$$



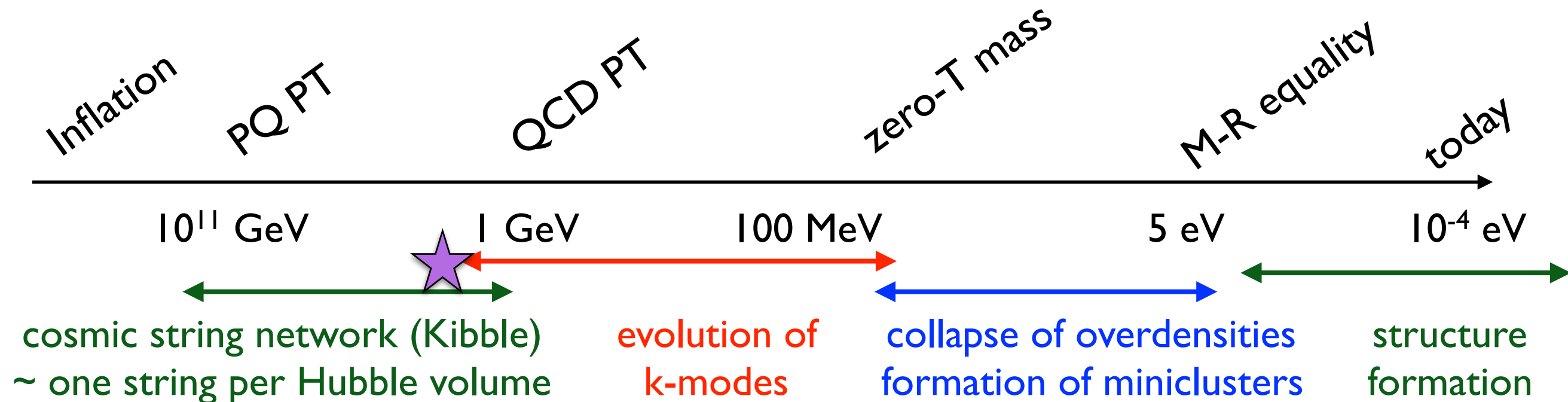








# Initial condition



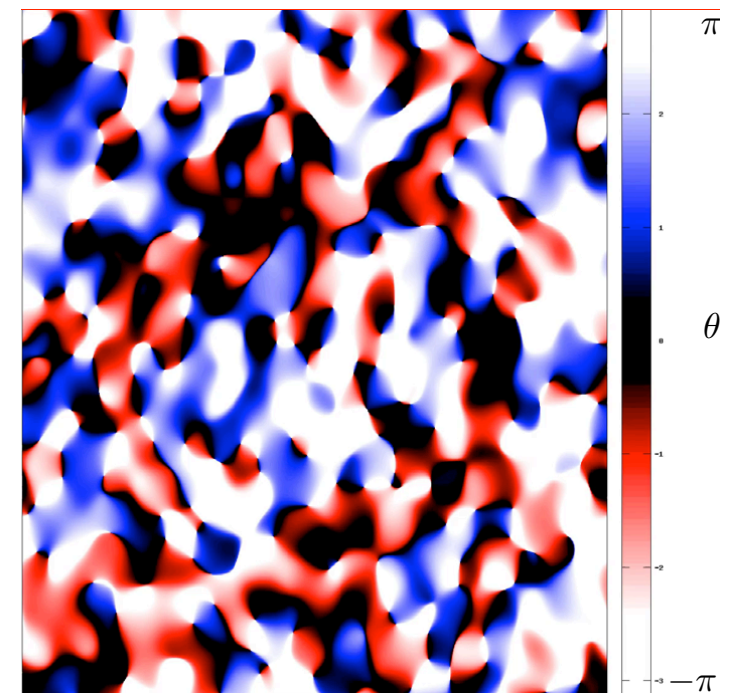
# Initial condition

- axion field smooth on scales  $<$  horizon  
uncorrelated on scales  $>$  horizon
- assume power spectrum for axion field  
w Gaussian cut-off

$$\langle \theta_k \theta_{k'}^* \rangle = (2\pi)^3 \delta^3(\vec{k} - \vec{k}') P_\theta(k)$$

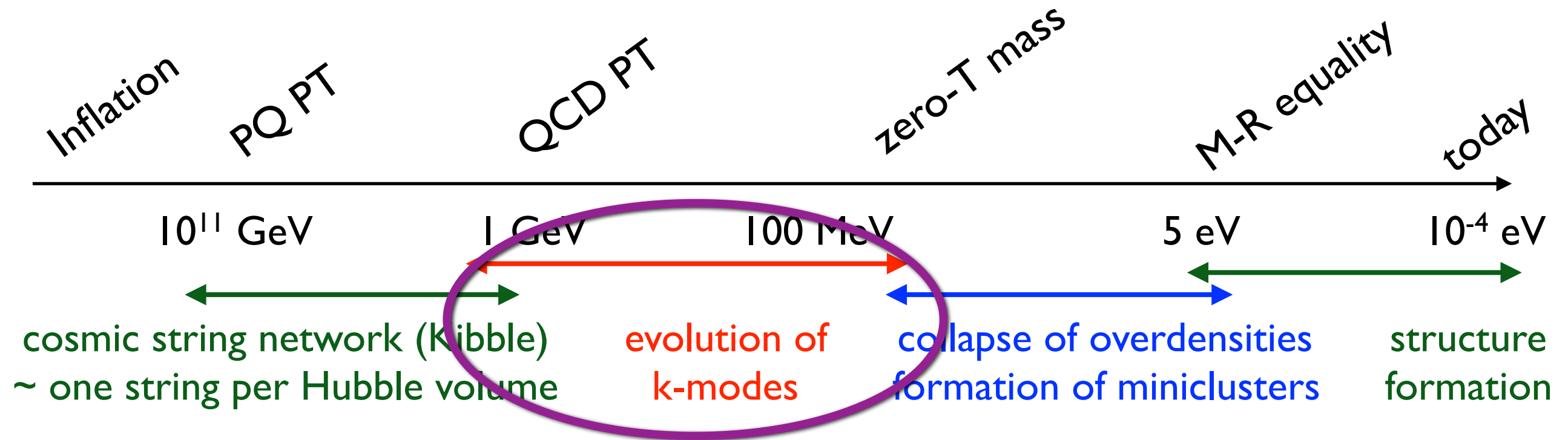
$$P_\theta^G(k) = \frac{8\pi^4}{3\sqrt{\pi}K^3} \exp\left(-\frac{k^2}{K^2}\right)$$

- normalization: fixed by flat distribution  $\langle \theta(\vec{x})^2 \rangle = \pi^2/3$
- cut-off: comoving horizon wave-number  $K = a_i H_i$



J. Redondo

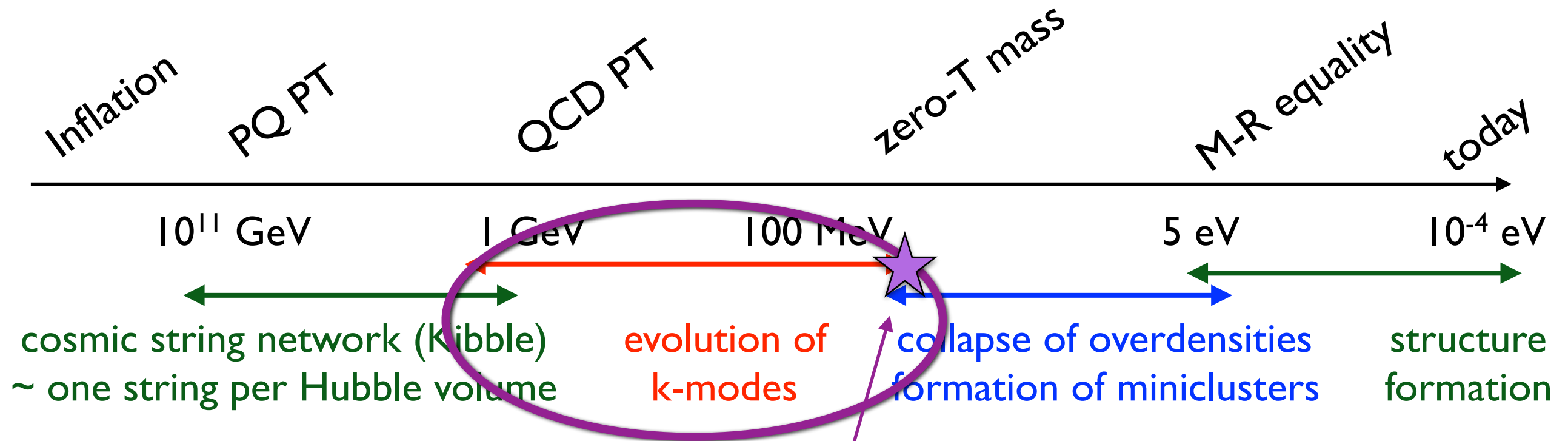
# Axion field evolution



- harmonic approximation of axion potential
- equation of motion including gradient terms:

$$\ddot{\theta}_k + 3H(T)\dot{\theta}_k + \omega_k^2\theta_k = 0, \quad \omega_k^2 \equiv \frac{k^2}{a^2} + m(T)^2$$

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- solve EoM to calculate axion energy density and density power spectrum

# Axion energy density

$$\rho(\vec{x}) = \frac{f_{\text{PQ}}^2}{2} \left[ \dot{\theta}^2 - \frac{1}{a^2} (\vec{\nabla} \theta)^2 + m^2(T) \theta^2 \right]$$

$$\theta_k(a) = \theta_k f_k(a)$$

- go to Fourier space for field

$$\rho(\vec{x}) = \frac{1}{(2\pi)^6} \frac{f_{\text{PQ}}^2}{2} \int d^3k d^3k' \theta_k \theta_{k'}^* F(k, k') e^{-i\vec{x}(\vec{k}-\vec{k}')} e^{-i\vec{x}(\vec{k}+\vec{k}')}$$

$$F(k, k') = \dot{f}_k \dot{f}_{k'} + \left( \frac{\vec{k} \cdot \vec{k}'}{a^2} + m^2(T) \right) f_k f_{k'}$$

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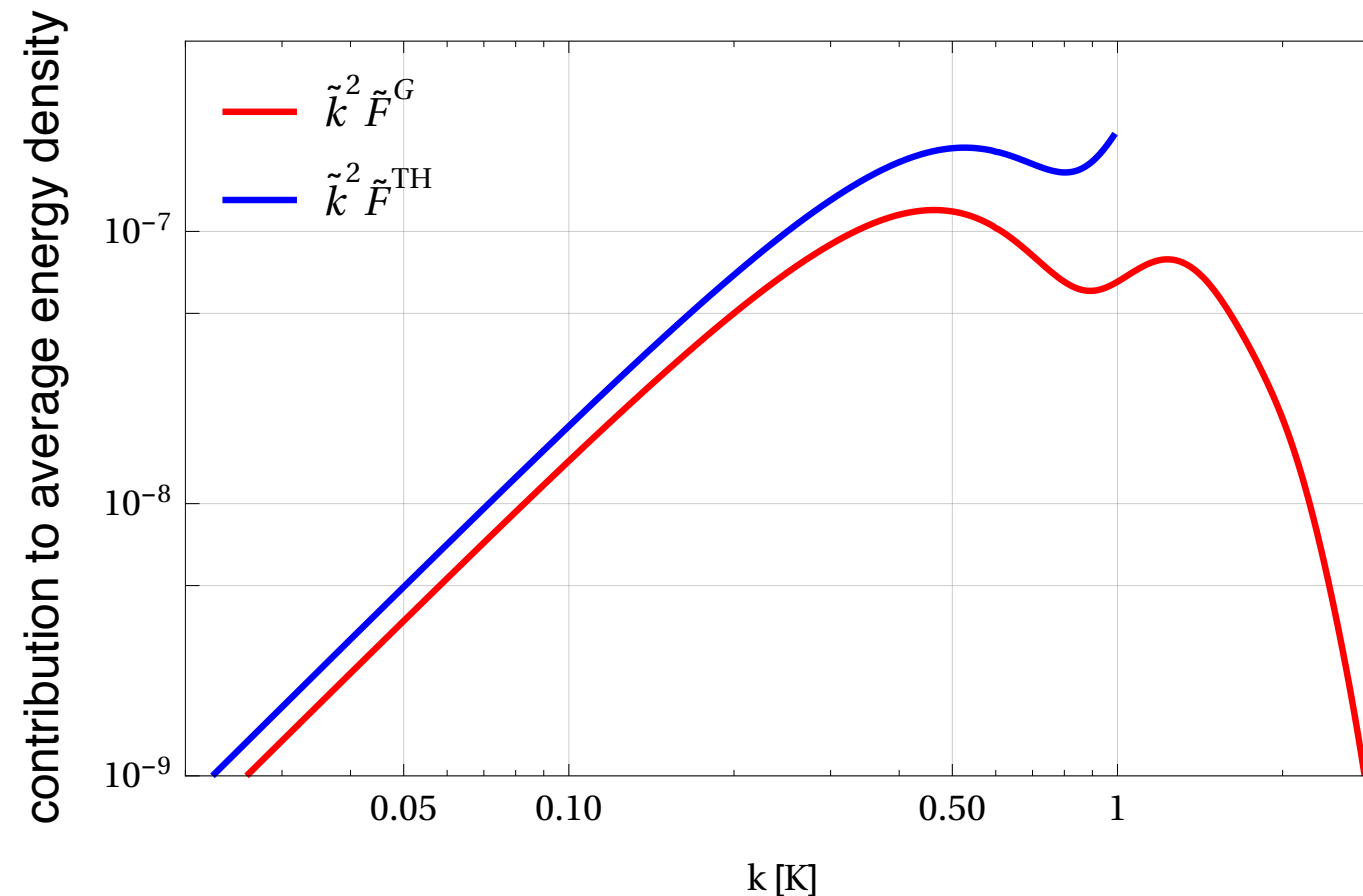
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- average energy density (without string contribution):

$$\bar{\rho} \equiv \langle \rho(\vec{x}) \rangle = \frac{1}{2\pi^2} \frac{f_{\text{PQ}}^2}{2} \int_0^\infty dk k^2 P_\theta(k) F(k, k)$$

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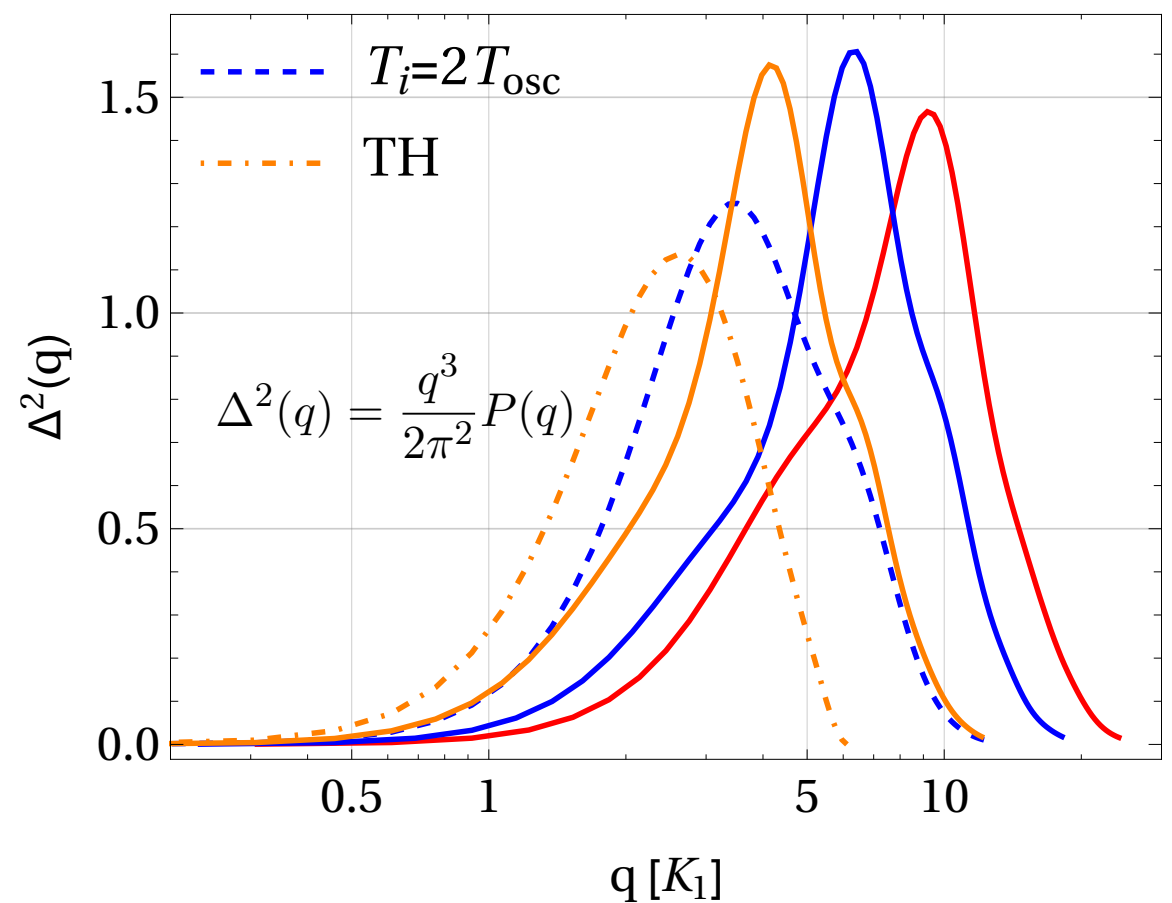
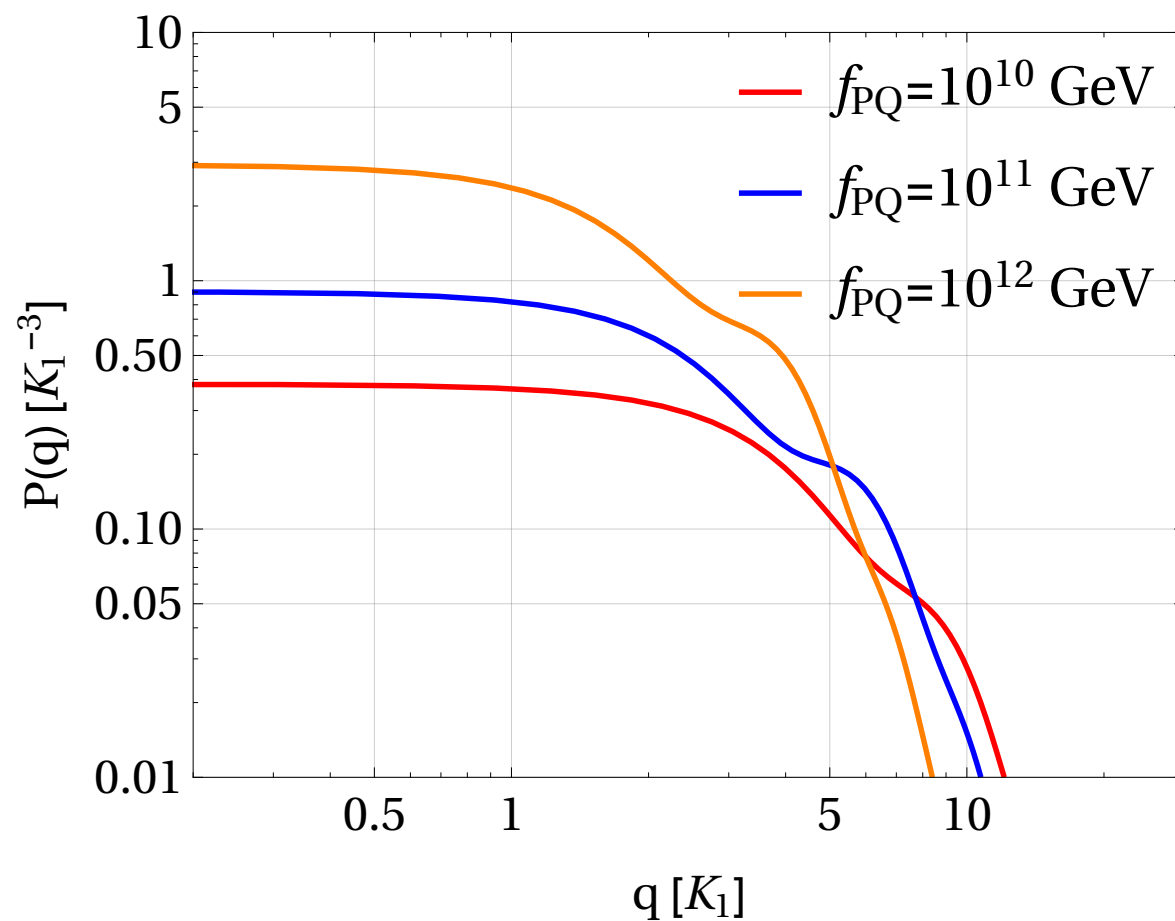


# Density power spectrum

$$P(q) = \frac{1}{V} \frac{\langle |\rho_q|^2 \rangle}{\bar{\rho}^2} = 2(2\pi)^3 \frac{\int d^3k P_\theta(|\vec{k}|) P_\theta(|\vec{k} - \vec{q}|) F(k, k - q)^2}{\left[ \int d^3k P_\theta(k) F(k, k) \right]^2}$$

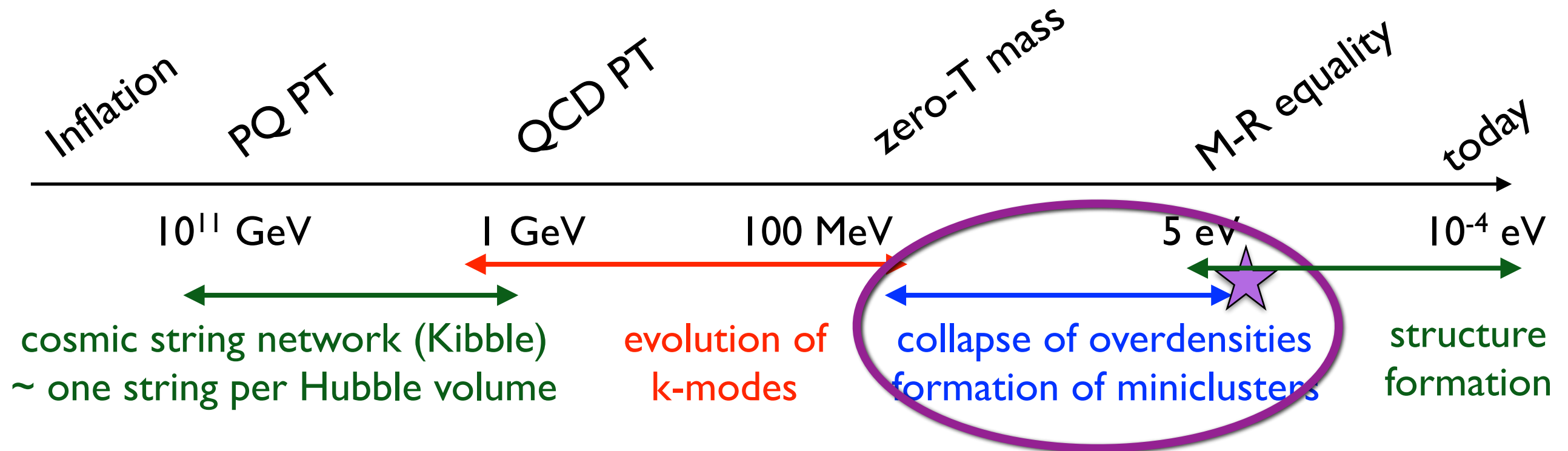
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- density fluctuations of order one
- charact. size a few times smaller than horizon @  $T_{osc}$

# Collapse of over-densities



- spherical collapse model for non-linear density fluctuations during radiation domination  
[Kolb, Tkachev, 94]
- modified Press-Schechter ansatz to calculate double-differential minicluster distribution in mass and size

# Double-differential mass function

- variance of the smoothed density contrast

$$\sigma_R^2 \equiv \langle \delta_R(\vec{x})^2 \rangle = \frac{1}{2\pi^2} \int_0^\infty dk k^2 P(k) \left| \tilde{W}_R(k) \right|^2$$

- Gaussian distribution for the smoothed contrast

$$f_{\text{sm}}(\delta; R) = \frac{1}{\sqrt{2\pi}\sigma_R} \exp\left(-\frac{\delta^2}{2\sigma_R^2}\right).$$

- derive distribution in  $\delta$  and  $R$ :

$$f(\delta, R) = -\frac{1}{\sigma_0} \frac{d\sigma_R}{dR} \frac{\delta^2}{\sigma_R^2} f_{\text{sm}}(\delta; R)$$

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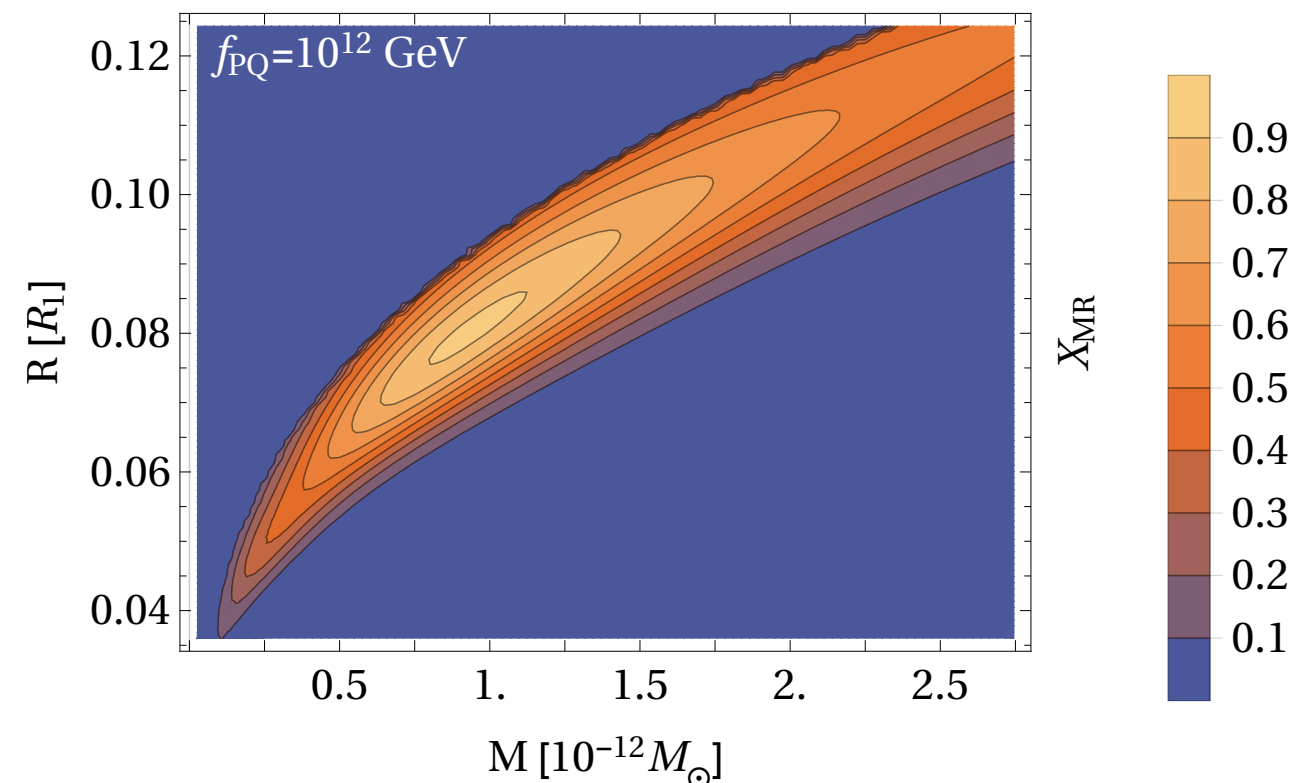
- double-differential distribution in  $M$  and  $R$  using:  $M = \frac{4\pi}{3} \bar{\rho} (1 + \delta) r^3$

$$\frac{dn}{dM dR} = \frac{3}{2\pi M R^3} f(\delta, R) \Theta[\delta - \delta_c(x)]$$

# Dimensionless double-differential mass function

$f_{\text{PQ}} [\text{GeV}]$	$M_{\text{peak}} [M_{\odot}]$	$M \text{ range } [M_{\odot}]$	physical sizes at turn around	
			$r_{\text{ta}}^{\text{peak}} [\text{km}]$	$r_{\text{ta}} \text{ range } [\text{km}]$
$10^{10}$	$4 \times 10^{-16}$	$[2 \times 10^{-17}, 1 \times 10^{-14}]$	$4 \times 10^4$	$[2 \times 10^4, 2 \times 10^5]$
$10^{11}$	$2 \times 10^{-14}$	$[5 \times 10^{-16}, 3 \times 10^{-13}]$	$2 \times 10^5$	$[4 \times 10^4, 7 \times 10^5]$
$10^{12}$	$8 \times 10^{-13}$	$[6 \times 10^{-14}, 2 \times 10^{-11}]$	$2 \times 10^6$	$[7 \times 10^5, 7 \times 10^6]$

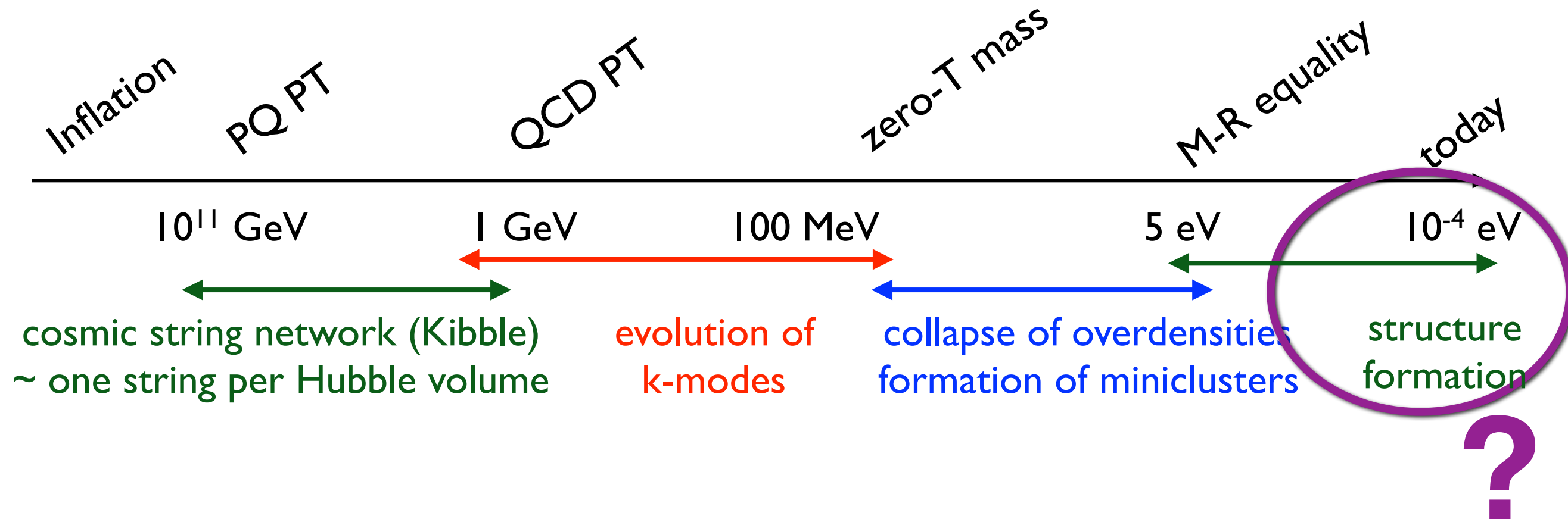
- MC masses span 3 orders, sizes span 1 order of magn.
  - peak-masses 2 orders of mag. smaller than naive estimates
- $$M \sim 10^{-12} M_{\odot} (f_{\text{PQ}} / 10^{11} \text{ GeV})^2$$
- (typical fluctuations smaller than horizon at  $T_{\text{osc}}$ )



# Outlook - approximations/assumptions

- initial power spectrum: should follow from evolution of string network (Kibble mechanism)
- harmonic approximation: anharmonic effects may lead to spikes in axion density [Kolb, Tkachev, 93]
- contribution from string/domain wall decays: likely to introduce additional energy density & fluctuations

# Axion DM today - in our galaxy?



- Do minicluster survive non-linear structure formation?
- Do they collapse to dense Axion-stars?  
Are Axion-stars stable?



# Axion DM today - in our galaxy?

- if a large fraction of the DM energy density is in MC-sized bound objects, the probability to meet one is very low (bad news for direct axion detection)
- depending on structure formation history, potentially interesting lensing signatures:
  - femto-lensing [Kolb, Tkachev, 95](#)
  - micro-lensing [Fairbairn, Marsh et al, 17](#)

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**Thank you!**

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