Constraints on the SM from the Weak Gravity Conjecture



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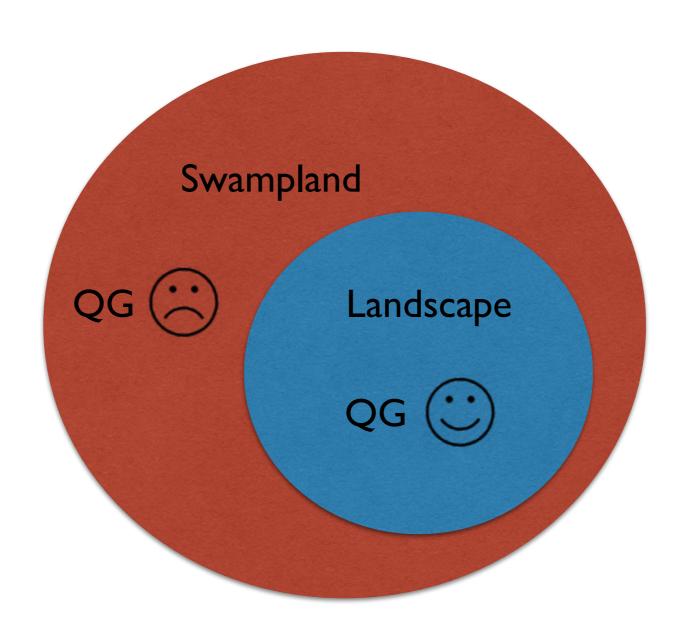


Ibanez, Martin-Lozano, IV [arXiv:1706.05392 [hep-th]]

Ibanez, Martin-Lozano, IV [arXiv: 1707.05811 [hep-th]]

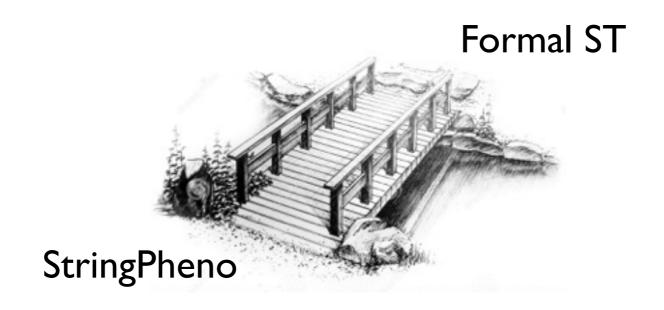
DESY Theory Workshop, 2017

What are the constraints that an effective theory must satisfy to be embedded in quantum gravity?



Quantum Gravity Conjectures

Motivated many times by observing recurrent features of the string landscape and "model building failures"



They can have significant implications in low energy physics!

Weak Gravity Conjecture

Weak Gravity Conjecture: [Arkani-Hamed et al.'06]

Given an abelian p-form gauge field, there must exist an electrically charged state with $T \leq Q$

Sharpened WGC:) [Ooguri-Vafa'16]

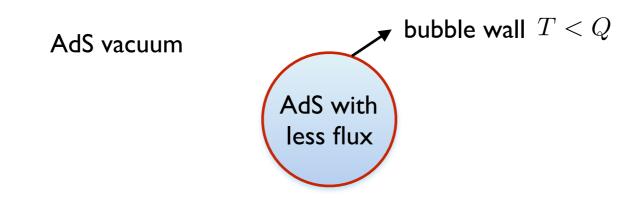
Bound is saturated only for a BPS state in a SUSY theory

Geometry supported by fluxes



Brane charged under the flux with $T \leq Q$

[Maldacena et al.'99]

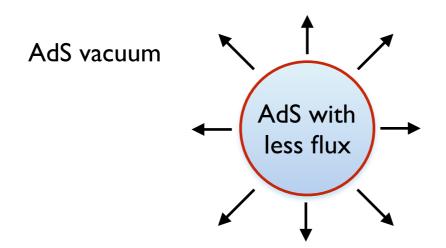


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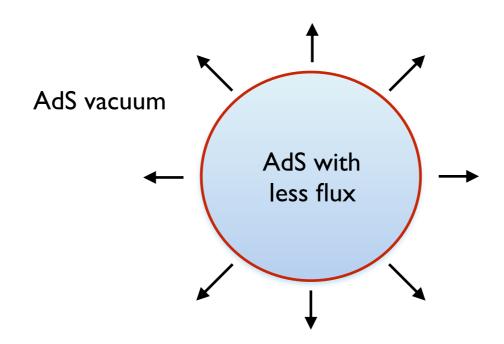


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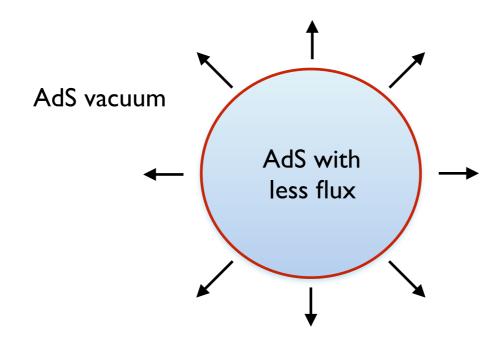
(non-susy)
Geometry supported

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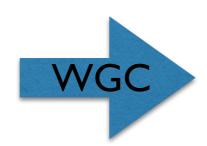
Brane charged under the flux with T < Q

[Maldacena et al.'99]



(non-susy)

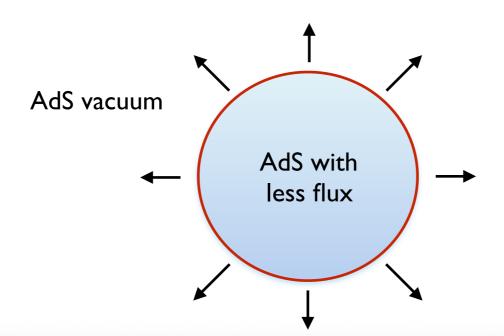
Geometry supported by fluxes



Brane charged under the flux with T < Q

[Maldacena et al.'99]

 $\red{!}$ In AdS, a brane with T < Q describes an instability



Non-susy AdS vacua supported by fluxes are at best metastable

[Ooguri-Vafa'16]

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Same conjecture in [Freivogel-Kleban'16]

[Ooguri-Vafa'16]

Non-susy AdS vacua are at best metastable

Same conjecture in [Freivogel-Kleban'16]

Non-susy stable AdS vacua cannot be embedded in quantum gravity!

Implications for:

- Holography
- String landscape
- Low energy physics?

Standard Model + Gravity on S^1 : [Arkani-Hamed et al.'07] (also [Arnold-Fornal-Wise'10])

$$V(R)\simeq rac{2\pi r^3\Lambda_4}{R^2}$$
 + Casimir energy tree-level one-loop corrections

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Depending on the light mass spectra and the cosmological constant, we can get AdS, Minkowski or dS vacua in 3d

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But AdS vacua are not consistent with quantum gravity!



Assumption: Background independence

If our 4d SM is consistent with QG



Compactifications of SM should also be consistent

We should not get stable non-susy AdS vacua from compactifying the SM!

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Absence of these vacua Constraints on SM (light espectra)

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- Absence of these vacua
 Constraints on SM (light espectra)
- There is some hidden instability

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 - ▶ Instability appearing upon compactification (periodic b.c. → no bubbles of nothing)

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Instability already in 4 dimensions —— Transfered to 3d

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(periodic b.c. → no bubbles of nothing)

Instability already in 4 dimensions —— Transfered to 3d

Assumption: Background independence

A 4d bubble instability also yields a 3d instability if

$$R_b < l_{AdS_3}$$

We

Therefore, the 3d vacuum will be stable if:

(periodic b.c. → no bubbles of not

Instability already in 4 dimensions

Transfered to 3d

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Instability already in 4 dimensions —— Transfered to 3d unless $l_{AdS_3} < R_{\text{bubble}} < l_{dS_4}$

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If our 4d SM is consistent with QG



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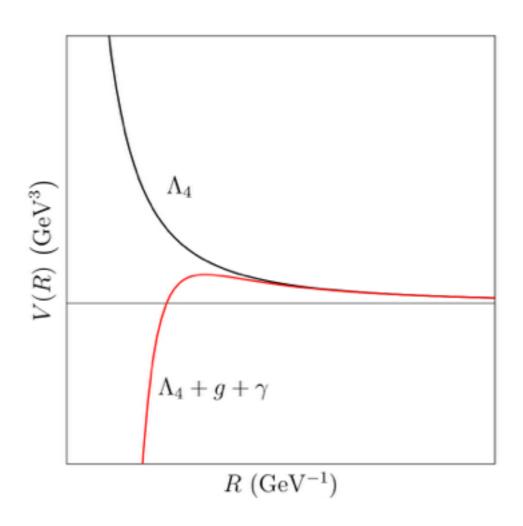
Standard Model + Gravity on S^1 :

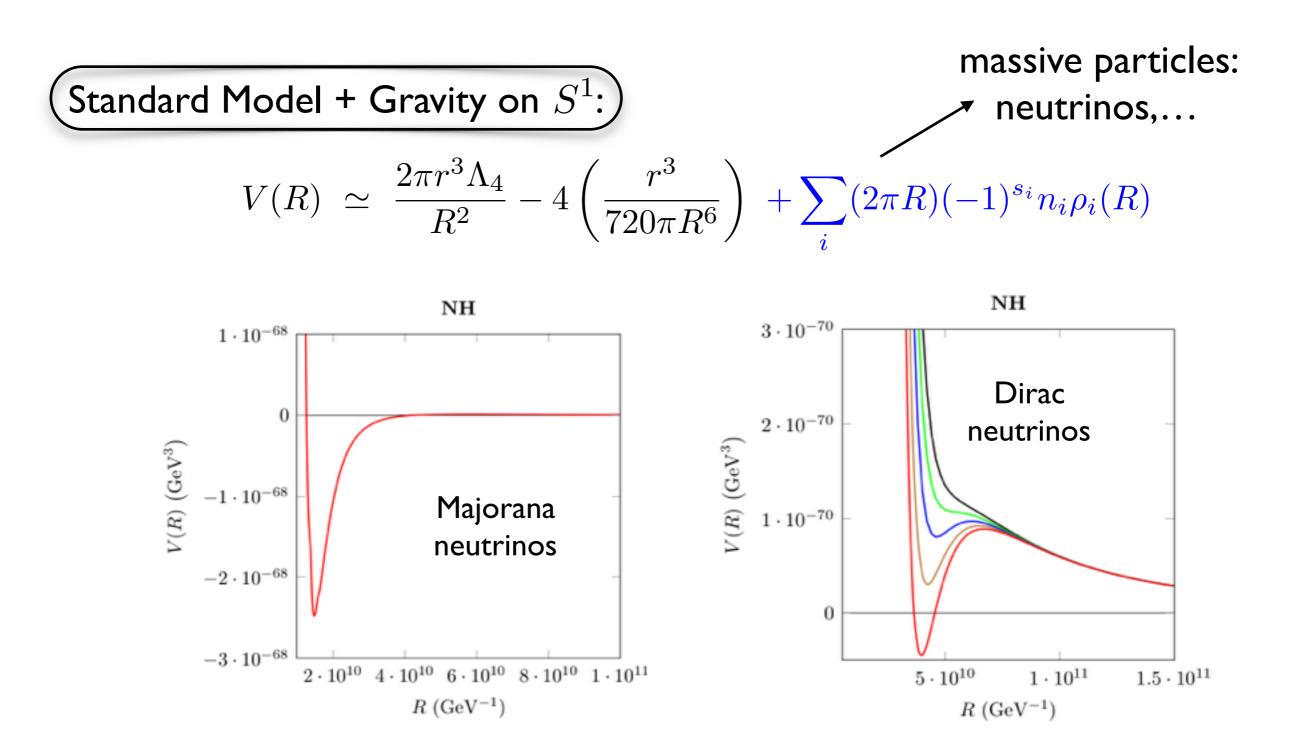
$$V(R) \simeq rac{2\pi r^3 \Lambda_4}{R^2}$$
 + Casimir energy



$$V(R) \simeq \frac{2\pi r^3 \Lambda_4}{R^2} - 4\left(\frac{r^3}{720\pi R^6}\right)$$

massless particles: graviton, foton





The more massive the neutrinos, the deeper the AdS vacuum

Constraints on neutrino masses

Majorana:

There is an AdS vacuum for any value of m_{ν}

Majorana neutrinos ruled out!

Dirac:

	NH	IH
No vacuum	$m_{\nu_1} < 6.7 \text{ meV}$	$m_{\nu_3} < 2.1 \text{ meV}$
dS_3 vacuum	$6.7 \text{ meV} < m_{\nu_1} < 7.7 \text{ meV}$	$2.1 \text{ meV} < m_{\nu_3} < 2.56 \text{ meV}$
AdS_3 vacuum	$m_{\nu_1} > 7.7 \; {\rm meV}$	$m_{\nu_3} > 2.56 \text{ meV}$

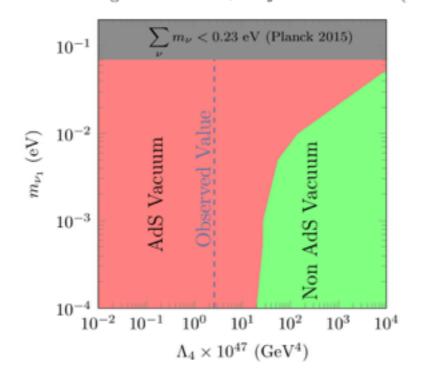
Absence of AdS vacuum requires

$$m_{\nu_1} < 7.7 \text{ meV (NH)}$$

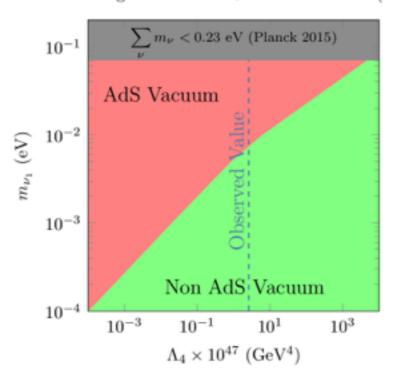
$$m_{\nu_1} < 2.1 \text{ meV (IH)}$$

Lower bound on the cosmological constant

Cosmological Constant + Majorana Neutrinos (NH)



Cosmological Constant + Dirac Neutrinos (NH)



The bound for Λ_4 scales as $m_{
u}^4$

(as observed experimentally)

$$\Lambda_4 \ge \frac{a(n_f)30(\Sigma m_i^2)^2 - b(n_f, m_i)\Sigma m_i^4}{384\pi^2}$$

with
$$a(n_f) = 0.184(0.009) \\ b(n_f, m_i) = 5.72(0.29)$$
 for Majorana (Dirac)

First argument (not based on cosmology) to have $\Lambda_4 \neq 0$

Adding BSM physics

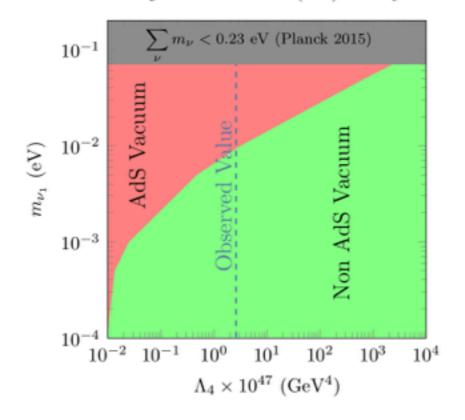
Light fermions

Positive Casimir contribution — helps to avoid AdS vacuum

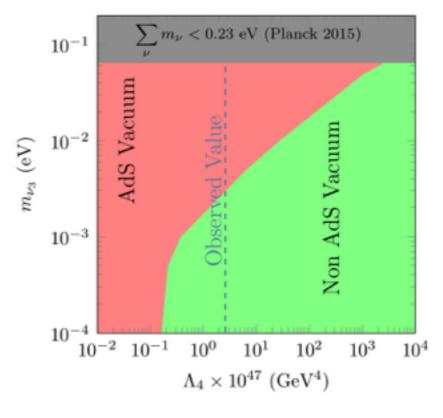
Majorana neutrinos are consistent if adding $m_\chi \lesssim 2 \,\,\mathrm{meV}$

example. For $m_\chi = 0.1~{\rm meV}$:

C.C. + Majorana Neutrinos (NH) + Weyl fermion



C.C. + Majorana Neutrinos (IH) + Weyl fermion

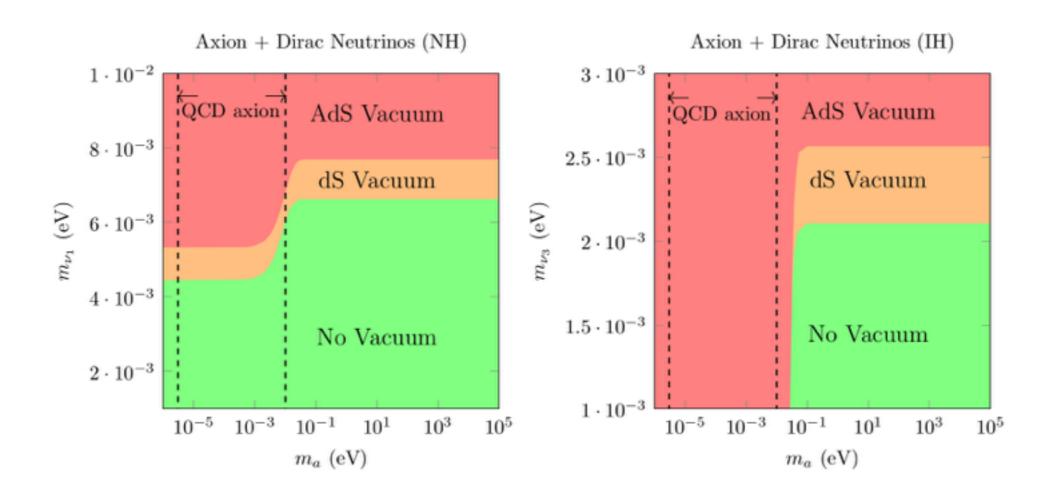


Adding BSM physics

Axions

1 axion: negative contribution — bounds get stronger

Multiple axions: can destabilise AdS vacuum



Bounds on the SM + light BSM physics

Model	Majorana (NI)	Majorana (IH)	Dirac (NH)	Dirac (IH)
SM (3D)	no	no	$m_{\nu_1} \le 7.7 \times 10^{-3}$	$m_{\nu_3} \le 2.56 \times 10^{-3}$
SM(2D)	no	no	$m_{\nu_1} \le 4.12 \times 10^{-3}$	$m_{\nu_3} \le 1.0 \times 10^{-3}$
SM+Weyl(3D)	$m_{\nu_1} \le 0.9 \times 10^{-2}$	$m_{\nu_3} \le 3 \times 10^{-3}$	$m_{\nu_1} \le 1.5 \times 10^{-2}$	$m_{\nu_3} \le 1.2 \times 10^{-2}$
	$m_f \le 1.2 \times 10^{-2}$	$m_f \le 4 \times 10^{-3}$		
SM+Weyl(2D)	$m_{\nu_1} \le 0.5 \times 10^{-2}$	$m_{\nu_3} \le 1 \times 10^{-3}$	$m_{\nu_1} \le 0.9 \times 10^{-2}$	$m_{\nu_3} \le 0.7 \times 10^{-2}$
	$m_f \le 0.4 \times 10^{-2}$	$m_f \le 2 \times 10^{-3}$		
SM+Dirac(3D)	$m_f \le 2 \times 10^{-2}$	$m_f \le 1 \times 10^{-2}$	yes	yes
SM+Dirac(2D)	$m_f \le 0.9 \times 10^{-2}$	$m_f \le 0.9 \times 10^{-2}$	yes	yes
$SM+1 \operatorname{axion}(3D)$	no	no	$m_{\nu_1} \le 7.7 \times 10^{-3}$	$m_{\nu_3} \le 2.5 \times 10^{-3}$
				$m_a \ge 5 \times 10^{-2}$
$SM+1 \operatorname{axion}(2D)$	no	no	$m_{\nu_1} \le 4.0 \times 10^{-3}$	$m_{\nu_3} \le 1 \times 10^{-3}$
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$\geq 2(10)$ axions	yes	yes	yes	yes

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$\geq 2(10)$ axions (yes	yes	yes	yes

Majorana neutrinos are consistent if adding:

- A Weyl (or Dirac) fermion $m_f \leq 10 \,\,\mathrm{meV}$
- Multiple axions

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$\geq 2(10)$ axions	yes	yes	yes	yes

Compactifications of SM on T_2 — qualitatively similar, but a bit stronger

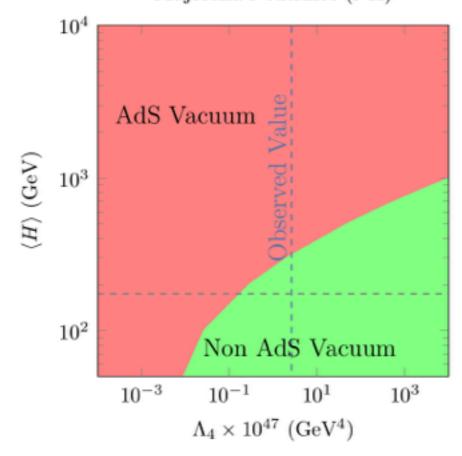
(see also [Hamada-Shiu'17])

Upper bound on the EW scale

Majorana case

$$\langle H \rangle \lesssim \frac{\sqrt{2}}{Y_{\nu_1}} \sqrt{M \Lambda^{1/4}}$$

Majorana Neutrinos (NH)

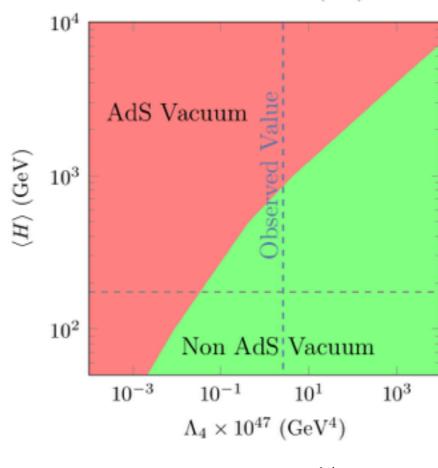


$$M = 10^{10} \text{ GeV}, Y = 10^{-3}$$

Dirac case

$$\langle H \rangle \lesssim 1.6 \frac{\Lambda^{1/4}}{Y_{\nu_1}}$$

Dirac Neutrinos (NH)



$$Y = 10^{-14}$$

Conclusions

- Consistency with quantum gravity implies constraints on low energy physics:
 - Lower bound on the cosmological const. of order the neutrino masses
 - Upper bound on the EW scale in terms of the cosmological const.
- Assumptions taken:
 - Validity of the Ooguri-Vafa Conjecture
 - Non-perturbative stability of 3D SM vacua
- New approach to fine-tuning or hierarchy problems?

 UV/IR mixing? (see also [Luest-Palti'17])

Thank you!

back-up slides

Casimir energy

Potential energy in 3d:

$$V(R) = \frac{2\pi r^3 \Lambda_4}{R^2} + \sum_{i} (2\pi R) \frac{r^3}{R^3} (-1)^{s_i} n_i \rho_i(R)$$

Casimir energy density:

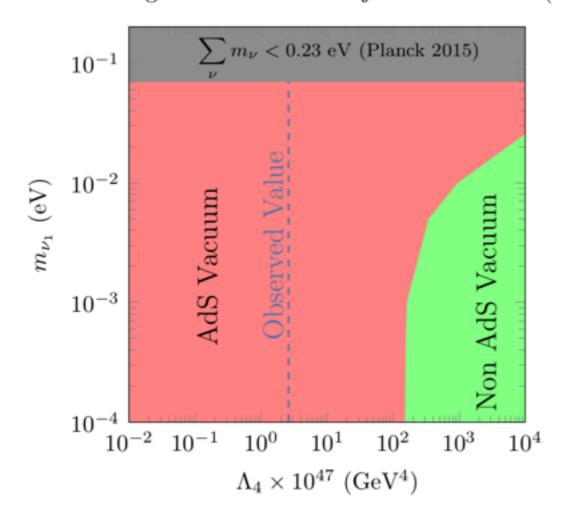
$$\rho(R) = \mp \sum_{n=1}^{\infty} \frac{2m^4}{(2\pi)^2} \frac{K_2(2\pi Rmn)}{(2\pi Rmn)^2}$$

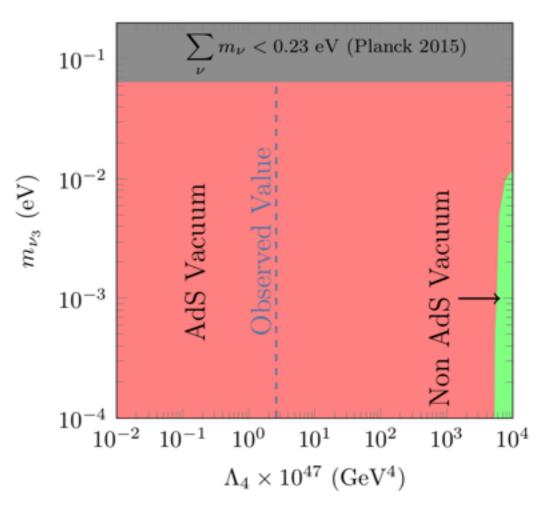
For small mR:

$$\rho(R) = \mp \left[\frac{\pi^2}{90(2\pi R)^4} - \frac{\pi^2}{6(2\pi R)^4} (mR)^2 + \frac{\pi^2}{48(2\pi R)^4} (mR)^4 + \mathcal{O}(mR)^6 \right]$$

Cosmological Constant + Majorana Neutrinos (NH)

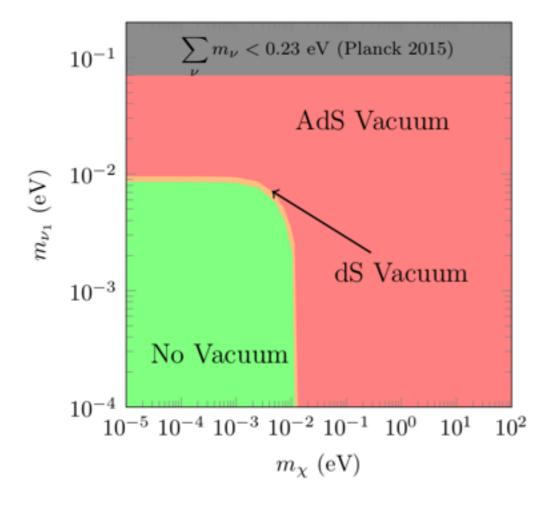






Adding light fermions





Weyl Fermion + Dirac Neutrinos (NH)

