

# Algebraic properties of the monopole formula

Marcus Sperling

Fakultät für Physik, Universität Wien

September 28, 2017



universität  
wien

FWF

Der Wissenschaftsfonds.

Based on arXiv:1605.00010 & 1611.07030 with A.Hanany

# Outline

- ① SUSY vacua and Hilbert series
- ② Monopole formula
- ③ Algebraic properties
- ④ Conclusions and outlook

# Outline

- 1 SUSY vacua and Hilbert series
- 2 Monopole formula
- 3 Algebraic properties
- 4 Conclusions and outlook

# Moduli space of supersymmetric vacua

**Ex:**  $4d \mathcal{N} = 1$  theory

- **gauge group**  $G$ : vector multiplets  $V_a$
- **matter content**: representation  $\mathcal{R}$  for chiral multiplets  $X_i$
- superpotential  $W$ :  $G$ -invariant polynomial in  $X_i$

# Moduli space of supersymmetric vacua

**Ex:**  $4d \mathcal{N} = 1$  theory

- **gauge group**  $G$ : vector multiplets  $V_a$
- **matter content**: representation  $\mathcal{R}$  for chiral multiplets  $X_i$
- superpotential  $W$ :  $G$ -invariant polynomial in  $X_i$
- **scalar potential**:  $V(X, X^\dagger) = \sum_i |F_i|^2 + \frac{g^2}{2} \sum_a (D_a)^2$  with  $F_i := \frac{\partial W}{\partial X_i}$  and  $D_a := \sum_{i,j} X_i^\dagger (T^a)_j^i X^j$

# Moduli space of supersymmetric vacua

**Ex:**  $4d \mathcal{N} = 1$  theory

- **gauge group**  $G$ : vector multiplets  $V_a$
- **matter content**: representation  $\mathcal{R}$  for chiral multiplets  $X_i$
- superpotential  $W$ :  $G$ -invariant polynomial in  $X_i$
- **scalar potential**:  $V(X, X^\dagger) = \sum_i |F_i|^2 + \frac{g^2}{2} \sum_a (D_a)^2$  with  $F_i := \frac{\partial W}{\partial X_i}$  and  $D_a := \sum_{i,j} X_i^\dagger (T^a)_j^i X^j$
- **Moduli space** of SUSY vacua [Luty, Taylor '95]

$$\mathcal{M} := \{(X, X^\dagger) \mid V(X, X^\dagger) = 0\} / G$$

# Moduli space of supersymmetric vacua

**Ex:**  $4d \mathcal{N} = 1$  theory

- **gauge group**  $G$ : vector multiplets  $V_a$
- **matter content**: representation  $\mathcal{R}$  for chiral multiplets  $X_i$
- superpotential  $W$ :  $G$ -invariant polynomial in  $X_i$
- **scalar potential**:  $V(X, X^\dagger) = \sum_i |F_i|^2 + \frac{g^2}{2} \sum_a (D_a)^2$  with  $F_i := \frac{\partial W}{\partial X_i}$  and  $D_a := \sum_{i,j} X_i^\dagger (T^a)_j^i X_j$
- **Moduli space** of SUSY vacua [Luty, Taylor '95]

$$\mathcal{M} = \{ (X, X^\dagger) \mid F_i = 0 \ \forall i, \ D_a = 0 \ \forall a \} / G$$

# Moduli space of supersymmetric vacua

**Ex:**  $4d \mathcal{N} = 1$  theory

- **gauge group**  $G$ : vector multiplets  $V_a$
- **matter content**: representation  $\mathcal{R}$  for chiral multiplets  $X_i$
- superpotential  $W$ :  $G$ -invariant polynomial in  $X_i$
- **scalar potential**:  $V(X, X^\dagger) = \sum_i |F_i|^2 + \frac{g^2}{2} \sum_a (D_a)^2$  with  $F_i := \frac{\partial W}{\partial X_i}$  and  $D_a := \sum_{i,j} X_i^\dagger (T^a)_j^i X^j$
- **Moduli space** of SUSY vacua [Luty, Taylor '95]

$$\mathcal{M} \cong \{X \mid F_i = 0 \ \forall i\} / G^{\mathbb{C}}$$

$\implies \mathcal{M}$  is Kähler



# Chiral ring and Hilbert series

- **gauge-inv. scalar chiral operator**  $\mathcal{O}$ , i.e.  $\bar{Q}_{\dot{\alpha}}\mathcal{O}(x) = 0$   
 $\longrightarrow$  VEV  $\langle \mathcal{O} \rangle =$  holomorphic fct. on  $\mathcal{M}$

# Chiral ring and Hilbert series

- **gauge-inv. scalar chiral operator**  $\mathcal{O}$ , i.e.  $\bar{Q}_{\dot{\alpha}}\mathcal{O}(x) = 0$   
 $\longrightarrow$  VEV  $\langle \mathcal{O} \rangle =$  holomorphic fct. on  $\mathcal{M}$
- **chiral ring**  $\mathfrak{R}$  spanned by  $[\mathcal{O}]_{\sim}$  subject to **relations**  
 $F_i = \partial_{X_i} W(X)|_{\mathfrak{R}} = 0$  [Lerche, Vafa, Warner '89]

# Chiral ring and Hilbert series

- **gauge-inv. scalar chiral operator**  $\mathcal{O}$ , i.e.  $\bar{Q}_{\dot{\alpha}}\mathcal{O}(x) = 0$   
→ VEV  $\langle \mathcal{O} \rangle = \text{holomorphic fct. on } \mathcal{M}$
- **chiral ring**  $\mathfrak{R}$  spanned by  $[\mathcal{O}]_{\sim}$  subject to **relations**  
 $F_i = \partial_{X_i} W(X)|_{\mathfrak{R}} = 0$  [Lerche, Vafa, Warner '89]
- **physics assumption:** chiral ring = coordinate ring of  $\mathcal{M}$   
→ Hilbert series [Pouliot '99; Benvenuti, Feng, Hanany, He '06]

# Chiral ring and Hilbert series

- **gauge-inv. scalar chiral operator**  $\mathcal{O}$ , i.e.  $\bar{Q}_{\dot{\alpha}}\mathcal{O}(x) = 0$   
→ VEV  $\langle \mathcal{O} \rangle =$  holomorphic fct. on  $\mathcal{M}$
- **chiral ring**  $\mathfrak{R}$  spanned by  $[\mathcal{O}]_{\sim}$  subject to **relations**  
 $F_i = \partial_{X_i} W(X)|_{\mathfrak{R}} = 0$  [Lerche, Vafa, Warner '89]
- **physics assumption:** chiral ring = coordinate ring of  $\mathcal{M}$   
→ Hilbert series [Pouliot '99; Benvenuti, Feng, Hanany, He '06]

$$\mathfrak{R} = \bigoplus_r \mathcal{H}_r, \quad \mathcal{H}_r = \{\mathcal{O}_i \mid \bar{Q}\mathcal{O}_i = 0, R\mathcal{O}_i = r\mathcal{O}_i\}$$

$$\text{R-charge: } [R, \bar{Q}] = \bar{Q}, \quad RW = 2W$$

# Chiral ring and Hilbert series

- **gauge-inv. scalar chiral operator**  $\mathcal{O}$ , i.e.  $\bar{Q}_{\dot{\alpha}}\mathcal{O}(x) = 0$   
→ VEV  $\langle \mathcal{O} \rangle =$  holomorphic fct. on  $\mathcal{M}$
- **chiral ring**  $\mathfrak{R}$  spanned by  $[\mathcal{O}]_{\sim}$  subject to **relations**  
 $F_i = \partial_{X_i} W(X)|_{\mathfrak{R}} = 0$  [Lerche, Vafa, Warner '89]
- **physics assumption:** chiral ring = coordinate ring of  $\mathcal{M}$   
→ Hilbert series [Pouliot '99; Benvenuti, Feng, Hanany, He '06]

$$\mathfrak{R} = \bigoplus_r \mathcal{H}_r, \quad \mathcal{H}_r = \{\mathcal{O}_i \mid \bar{Q}\mathcal{O}_i = 0, R\mathcal{O}_i = r\mathcal{O}_i\}$$

$$\textbf{R-charge:} \quad [R, \bar{Q}] = \bar{Q}, \quad RW = 2W$$

$$H_{\mathfrak{R}}(t) := \sum_r \dim(\mathcal{H}_r) t^r = \frac{Q(t)}{\prod_i (1 - t^{d_i})}$$

# Outline

- ① SUSY vacua and Hilbert series
- ② Monopole formula
- ③ Algebraic properties
- ④ Conclusions and outlook

# bare monopole operators

**for  $3d \mathcal{N} = 2$  theories:** new class of chiral operators

→ 't Hooft monopole operators [['t Hooft '78](#); [Borokhov, Kapustin, Wu '02](#)]

# bare monopole operators

for  $3d \mathcal{N} = 2$  theories: new class of chiral operators

→ 't Hooft monopole operators [['t Hooft '78](#); [Borokhov, Kapustin, Wu '02](#)]

**bare  $\mathcal{N} = 2$  monopole operator  $V_m$  :**

$$A_{\pm} \sim \frac{m}{2} (\pm 1 - \cos(\theta)) d\phi$$
$$\sigma \sim \frac{m}{2r}$$



# bare monopole operators

for  $3d \mathcal{N} = 2$  theories: new class of chiral operators

→ 't Hooft monopole operators [’t Hooft ’78; Borokhov, Kapustin, Wu ’02]

bare  $\mathcal{N} = 2$  monopole operator  $V_m$  :

$$A_{\pm} \sim \frac{m}{2} (\pm 1 - \cos(\theta)) d\phi$$
$$\sigma \sim \frac{m}{2r}$$

Dirac quantisation condition [Englert & Windey]

$$e^{2\pi i m} = \mathbb{1}_G \quad \Leftrightarrow \quad m \in \Gamma_{\hat{G}}$$

$\Gamma_{\hat{G}}$  = weight lattice of GNO dual group [Goddard, Nuyts, Olive ’77]

# bare monopole operators

for  $3d \mathcal{N} = 2$  theories: new class of chiral operators

→ 't Hooft monopole operators [’t Hooft ’78; Borokhov, Kapustin, Wu ’02]

bare  $\mathcal{N} = 2$  monopole operator  $V_m$  :

$$A_{\pm} \sim \frac{m}{2} (\pm 1 - \cos(\theta)) d\phi$$
$$\sigma \sim \frac{m}{2r}$$

Dirac quantisation condition [Englert & Windey]

$$e^{2\pi i m} = \mathbb{1}_G \quad \Leftrightarrow \quad m \in \Gamma_{\hat{G}}$$

$\Gamma_{\hat{G}}$  = weight lattice of GNO dual group [Goddard, Nuyts, Olive ’77]

- ▶  $\mathcal{N} = 2$  v-plet  $(A, \sigma) \rightarrow$  chiral multiplet  $V_m$  for each  $m$
- ▶ for each  $m \in \Gamma_{\hat{G}}$  exists a **single**  $\mathcal{N} = 2$  bare monopole  $V_m$

# Coulomb branch of $3d \mathcal{N} = 4$

**Specialise to  $\mathcal{N} = 4$**

$$G : (\mathcal{N}=4 \text{ vector}) = (\mathcal{N}=2 \text{ vector}) \oplus (\mathcal{N}=2 \text{ chiral } \Phi \in \text{adj})$$

$$\mathcal{R} : (\mathcal{N}=4 \text{ hyper}) = (\mathcal{N}=2 \text{ chiral } X \in \mathcal{R}) \oplus (\mathcal{N}=2 \text{ chiral } X \in \bar{\mathcal{R}})$$

$$\text{with } W = \bar{X}\Phi X \rightarrow \text{locally } \mathcal{M} = \mathcal{M}_C + \mathcal{M}_H$$

# Coulomb branch of $3d \mathcal{N} = 4$

**Specialise to  $\mathcal{N} = 4$**

$$G : (\mathcal{N}=4 \text{ vector}) = (\mathcal{N}=2 \text{ vector}) \oplus (\mathcal{N}=2 \text{ chiral } \Phi \in \text{adj})$$

$$\mathcal{R} : (\mathcal{N}=4 \text{ hyper}) = (\mathcal{N}=2 \text{ chiral } X \in \mathcal{R}) \oplus (\mathcal{N}=2 \text{ chiral } X \in \bar{\mathcal{R}})$$

$$\text{with } W = \bar{X}\Phi X \rightarrow \text{locally } \mathcal{M} = \mathcal{M}_C + \mathcal{M}_H$$

**Coulomb branch  $\mathcal{M}_C$ :**

Dressing of bare monopole  $V_m$  by **residual gauge theory  $T_m$**

- ▶ residual gauge group  $H_m = \text{Stab}_G(m)$
- ▶ residual matter fields  $\Phi_\alpha$  for  $\alpha(m) = 0$

# Coulomb branch of $3d \mathcal{N} = 4$

**Monopole formula** counts dressed  $\mathcal{N} = 2$  monopole operators

$$H(t) = \sum_{m \in \Gamma_{\hat{G}}/\mathcal{W}} t^{\Delta(m)} P(t, m)$$

[Cremonesi, Hanany, Zaffaroni '13]

# Coulomb branch of $3d \mathcal{N} = 4$

**Monopole formula** counts dressed  $\mathcal{N} = 2$  monopole operators

$$H(t) = \sum_{m \in \Gamma_{\hat{G}}/\mathcal{W}} t^{\Delta(m)} P(t, m)$$

[Cremonesi, Hanany, Zaffaroni '13]

- **R-charge** of bare monopole  $V_m$

$$\Delta(m) = \frac{1}{2} \sum_{\rho \in \mathcal{R}} |\rho(m)| - \sum_{\alpha \in \Phi_+} |\alpha(m)|$$

[Borokhov, Kapustin, Wu '02; Gaiotto, Witten '08; ...]

# Coulomb branch of $3d \mathcal{N} = 4$

**Monopole formula** counts dressed  $\mathcal{N} = 2$  monopole operators

$$H(t) = \sum_{m \in \Gamma_{\hat{G}}/\mathcal{W}} t^{\Delta(m)} P(t, m)$$

[Cremonesi, Hanany, Zafferoni '13]

- ▶ **R-charge** of bare monopole  $V_m$

$$\Delta(m) = \frac{1}{2} \sum_{\rho \in \mathcal{R}} |\rho(m)| - \sum_{\alpha \in \Phi_+} |\alpha(m)|$$

[Borokhov, Kapustin, Wu '02; Gaiotto, Witten '08; ...]

- ▶ **Dressing** by residual gauge theory  $T_m$

$$P(t, m) = \frac{1}{\prod_{i=1}^r (1 - t^{d_i(m)})} = H_{T_m}(t)$$

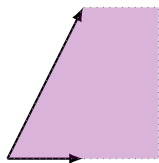
# Outline

- ① SUSY vacua and Hilbert series
- ② Monopole formula
- ③ Algebraic properties
- ④ Conclusions and outlook

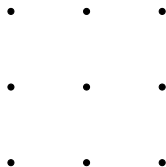


# Introduction of fan and Hilbert bases

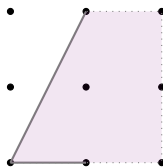
Summation range  $S = \Gamma_{\hat{G}}/\mathcal{W}$



$\cap$



$=$



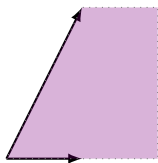
Weyl chamber  $\hat{G}$   
 $=$  **polyhed. cone**  $\sigma$

weight lattice  
 $\Gamma_{\hat{G}}$

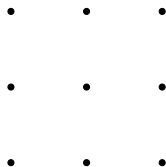
**monoid**  
 $S = \sigma \cap \Gamma_{\hat{G}}$

# Introduction of fan and Hilbert bases

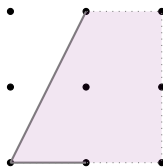
Summation range  $S = \Gamma_{\hat{G}}/\mathcal{W}$



$\cap$



$=$



Weyl chamber  $\hat{G}$   
 $=$  **polyhed. cone**  $\sigma$

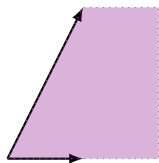
weight lattice  
 $\Gamma_{\hat{G}}$

**monoid**  
 $S = \sigma \cap \Gamma_{\hat{G}}$

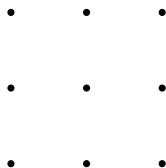
- ▶ cone  $\sigma$  generated by its edges
- ▶ monoid  $S$  generated by **Hilbert basis**

# Introduction of fan and Hilbert bases

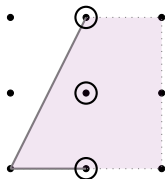
Summation range  $S = \Gamma_{\hat{G}}/\mathcal{W}$



$\cap$



$=$



Weyl chamber  $\hat{G}$   
 $=$  **polyhed. cone**  $\sigma$

weight lattice  
 $\Gamma_{\hat{G}}$

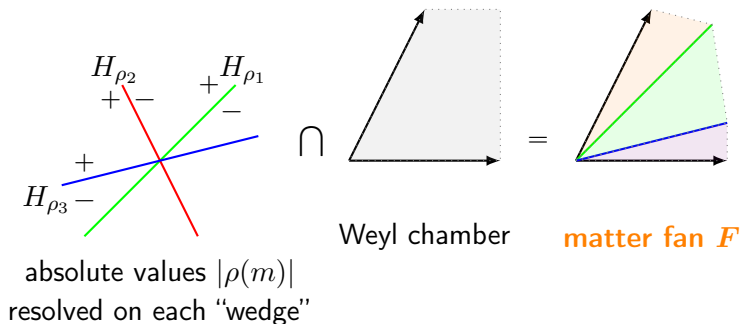
**monoid**  
 $S = \sigma \cap \Gamma_{\hat{G}}$

- ▶ cone  $\sigma$  generated by its edges
- ▶ monoid  $S$  generated by **Hilbert basis**

# Introduction of fan and Hilbert bases

**Problem:** absolute values  $|\rho(m)|$  need be resolved

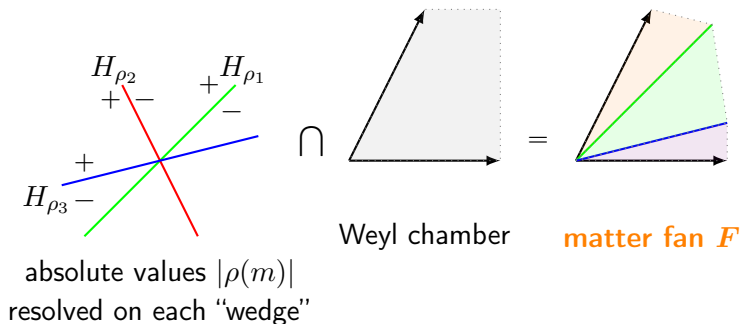
→ turn into **guiding principle!**



# Introduction of fan and Hilbert bases

**Problem:** absolute values  $|\rho(m)|$  need be resolved

→ turn into **guiding principle!**



restrict to weight lattice → **monoids** generated by **Hilbert bases**

# Algebraic properties

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

# Algebraic properties

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

- sufficient set of chiral ring generators

# Algebraic properties

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

- sufficient set of chiral ring generators
  - ▶ **Casimir invariants** of  $G \leftrightarrow P(t, 0)$



# Algebraic properties

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

- sufficient set of chiral ring generators
  - ▶ **Casimir invariants** of  $G \leftrightarrow P(t, 0)$
  - ▶ **bare monopole operators** via Hilbert bases  $\leftrightarrow H_{\omega(S_{\tau})}(t)$   
→ identify all charges

# Algebraic properties

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

- sufficient set of chiral ring generators
  - ▶ **Casimir invariants** of  $G \leftrightarrow P(t, 0)$
  - ▶ **bare monopole operators** via Hilbert bases  $\leftrightarrow H_{\omega(S_{\tau})}(t)$   
→ identify all charges
  - ▶ **dressed monopole operators**  $\leftrightarrow H_{M_{\tau}^{\text{Dress}}}(t)$   
→ identify their finite number and all charges

$$H(t) = P(t, 0) \sum_{\tau \in F} H_{M_{\tau}^{\text{Dress}}}(t) H_{\omega(S_{\tau})}(t)$$

- sufficient set of chiral ring generators
  - ▶ **Casimir invariants** of  $G \leftrightarrow P(t, 0)$
  - ▶ **bare monopole operators** via Hilbert bases  $\leftrightarrow H_{\omega(S_{\tau})}(t)$   
→ identify all charges
  - ▶ **dressed monopole operators**  $\leftrightarrow H_{M_{\tau}^{\text{Dress}}}(t)$   
→ identify their finite number and all charges
- pole structure
  - ▶ order of pole of  $H(t)$  at  $t = 1$  is  $\text{rk}(G) = \dim(\mathcal{M}_C)$
  - ▶ order of pole of  $H(t)$  at  $t \rightarrow \infty$  is  $\text{rk}(G) = \dim(\mathcal{M}_C)$

# Outline

- ① SUSY vacua and Hilbert series
- ② Monopole formula
- ③ Algebraic properties
- ④ Conclusions and outlook

## Geometric structure unveiled

- rewrite monopole formula
  - ▶ convenient for **implementation via** Macaulay2
  - ▶ **prove algebraic properties** of HS
  - ▶ appearing object mathematically well-studied
- identify sufficient set of **chiral ring generators**

# Conclusions and Outlook

## Geometric structure unveiled

- rewrite monopole formula
  - ▶ convenient for **implementation via** Macaulay2
  - ▶ **prove algebraic properties** of HS
  - ▶ appearing object mathematically well-studied
- identify sufficient set of **chiral ring generators**

## Future directions

- Hilbert series with background charges  $\longrightarrow$  resolution of  $\mathcal{M}_C$
- implications for  $\mathcal{N} = 2$  monopole formula

[Cremonesi '15; Cremonesi, Mekareeya, Zafferoni '16]