



Charged Composite Scalar Dark Matter

Maximilian Ruhdorfer
Technische Universität München



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based on arXiv:1707.07685 [hep-ph]

with R. Balkin, E. Salvioni and A. Weiler



Composite Higgs - A Reminder

- A light elementary scalar is **unnatural**
- Possible solution: Higgs is a **bound state** of a new strong sector

Description changes above confinement scale ($\sim \text{TeV}$)
 $\rightarrow$ Higgs mass is naturally "screened"

$$r_H \sim (\text{TeV})^{-1}$$

- Analogy to QCD Pions: Higgs arises as (approximate) **Goldstone boson** and is naturally light

$$\mathcal{G} \xrightarrow[H, \dots]{f} \mathcal{H}$$

$$\text{Minimal Model: } SO(5) \xrightarrow[H]{f} SO(4)$$

Agashe, Contino, Pomarol 2004



Composite Higgs - A Reminder

- Elementary sector couples **linearly** to strong sector

$$\begin{array}{c} \text{Elementary} \\ \text{fields} \end{array} \quad \begin{array}{c} W^\mu, t \end{array} \quad \sim \frac{gg_* f^2 W_\mu \rho^\mu}{\sim \epsilon \bar{t} \psi} \quad \begin{array}{c} \text{Strongly} \\ \text{interacting} \\ \text{sector} \end{array} \quad \begin{array}{c} \rho^\mu, \psi \end{array} \quad \begin{array}{l} \mathcal{G} \xrightarrow{f} \mathcal{H} \\ H, \chi, \dots \\ m_{\rho, \psi} \sim g_* f \end{array}$$

Elementary particles mix with strong sector resonances

$\Rightarrow$ SM particles are **partially composite** [Kaplan 1991](#)

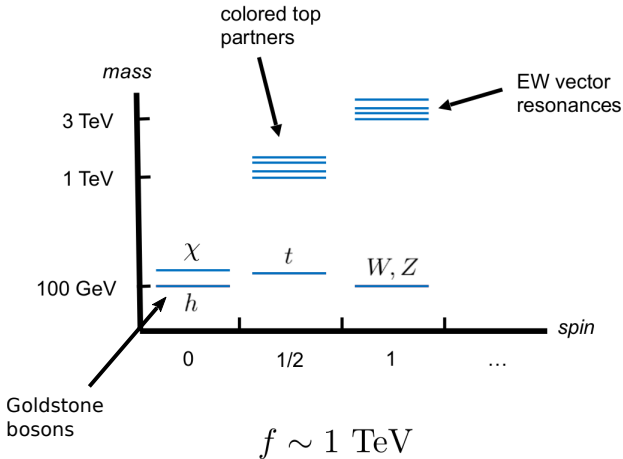
- Elementary couplings g, ϵ break global symmetry
 $\rightarrow$ Breaking generates Higgs potential at loop level

Strongest contribution from top + fermion resonances (**Top Partners**)

$$\mathcal{L} \sim \epsilon_q \bar{q}_L \mathcal{O}_q + \epsilon_t \bar{t}_R \mathcal{O}_t$$



Sketch of Particle Spectrum





Composite Scalar Dark Matter

- Bottom-up perspective: no reason for Higgs to go alone

$$\mathcal{G} \xrightarrow[H, \chi]{f} \mathcal{H} \quad \text{Additional pNGB } \chi \text{ as WIMP} \quad \text{Frigerio et al. 2012}$$

- χ is **naturally light** and **weakly coupled** at low energies

Dark matter stability?



Composite Scalar Dark Matter

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$$\mathcal{G} \xrightarrow[H, \chi]{f} \mathcal{H}$$

Additional pNGB χ as WIMP [Frigerio et al. 2012](#)

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Dark matter stability?

Extended symmetry structure can provide symmetry to stabilize the dark matter!

No additional symmetry or parity needed!



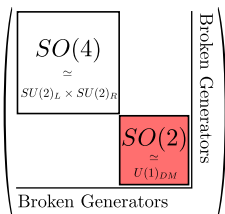
- Extend coset to $SO(7)/SO(6)$

$$SO(7) \xrightarrow{f} SO(6) \Rightarrow (H, \chi) \sim \mathbf{4}_0 + \mathbf{1}_{\pm 1}$$

Balkin, MR, Salvioni, Weiler,
1707.07685

- DM candidate is a **complex scalar** charged under **conserved** $U(1)_{\text{DM}}$

$$U(1)_{\text{DM}} \subset SO(6): \quad \chi \rightarrow e^{i\alpha} \chi$$



$$\Sigma \sim e^{i\pi^a X^a/f} \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}, \quad \vec{\pi} \sim \begin{pmatrix} h_1 \\ h_2 \\ h_3 \\ h_4 \\ \text{Im}(\chi) \\ \text{Re}(\chi) \end{pmatrix}$$



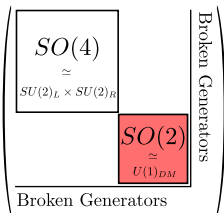
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Elementary Fermions in $\mathbf{7} \sim SO(7)$
respect $U(1)_{\text{DM}}$

BUT break Goldstone shift-symmetry



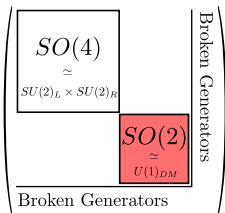
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$U(1)_{DM} \subset SO(6)$ automatically
preserved by UV-completion

$\Rightarrow$ UV-Safe Stabilization



Low-Energy Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \underbrace{\mathcal{L}_{nl\sigma m} + \mathcal{L}_t}_{\text{tree-level}} - \underbrace{V_{\text{eff}}}_{\text{one-loop}}$$



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$$\frac{v}{f^2} \partial_\mu h \partial^\mu |\chi|^2 + \dots$$



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$$m_\chi^2 \chi^* \chi + 2v\lambda h \chi^* \chi + \dots$$



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$$\frac{m_t}{f^2} \bar{t} t \chi^* \chi$$



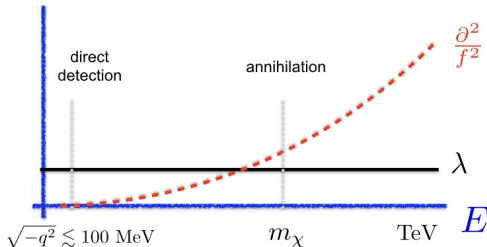
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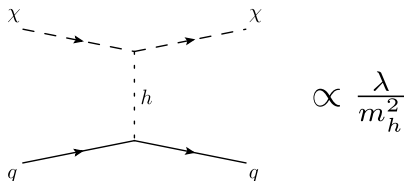
$$g_{\text{DM-SM}}^2(E)$$





Direct Detection

- Higgs exchange in t -channel dominates χN scattering



- Derivative interactions negligible $\frac{-q^2}{f^2} \ll 1$

- Spin-Independent χN cross-section determined by λ

$$\sigma_{\text{SI}}^{\chi N} \simeq \frac{f_N^2}{\pi} \frac{m_N^4 \lambda^2}{m_\chi^2 m_h^4} = 4 \times 10^{-46} \text{cm}^2 \left(\frac{\lambda}{0.03} \right)^2 \left(\frac{300 \text{ GeV}}{m_\chi} \right)^2$$

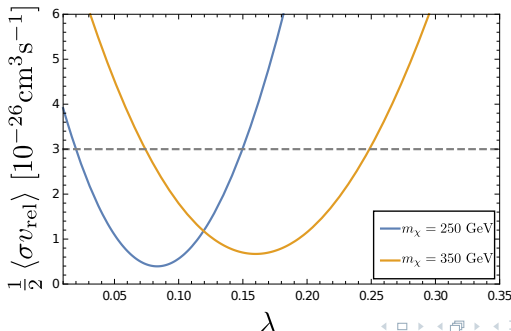


Relic Abundance

- Annihilation mainly through Higgs portal into WW, ZZ, hh, tt

- Interplay of derivative and portal coupling $\mathcal{L} \sim \frac{v}{f^2} \partial_\mu h \partial^\mu |\chi|^2 - 2v\lambda h |\chi|^2$

$$\sigma \propto \left(\frac{2m_\chi^2}{f^2} - \lambda \right)^2 \quad \text{fixed } m_\chi: \text{two values of } \lambda \text{ reproduce right RA}$$



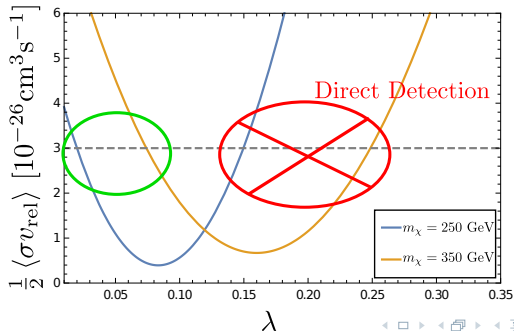


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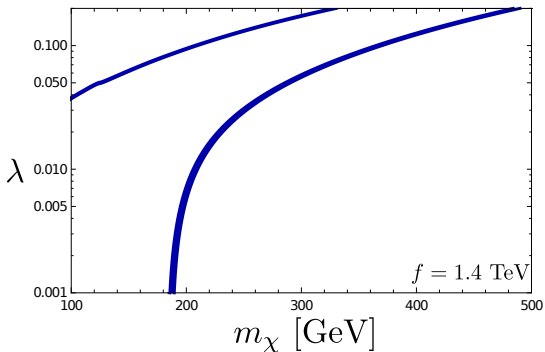
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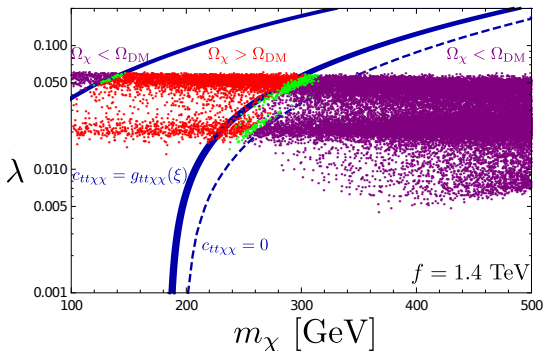


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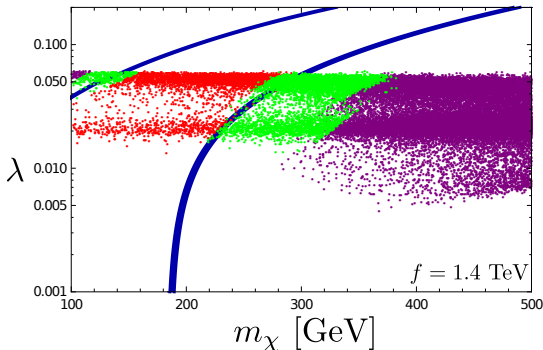


Green: $\Omega_\chi = \Omega_{\text{DM}} \pm 5\%$

All points satisfy collider constraints (next slide)



Relic Abundance



Caveat:

Large loop corrections to

$$\frac{1}{f^2} \partial_\mu |H|^2 \partial^\mu |\chi|^2$$

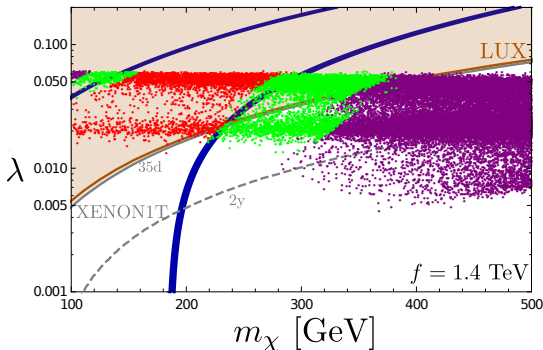
$\Rightarrow$ irreducible uncertainty of 50%

Green: $\Omega_\chi = \Omega_{\text{DM}} \pm 50\%$

All points satisfy collider constraints (next slide)



Relic Abundance



Caveat:

Large loop corrections to

$$\frac{1}{f^2} \partial_\mu |H|^2 \partial^\mu |\chi|^2$$

⇒ irreducible uncertainty of 50%

Green: $\Omega_\chi = \Omega_{\text{DM}} \pm 50\%$

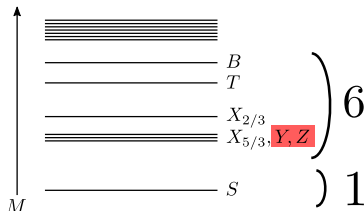
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Collider Phenomenology

- Top Partners are **colored** and can be produced at hadron colliders
- They come in complete $SO(6)$ representations

$$SO(7)/SO(6) : \quad \mathbf{7} = \mathbf{6} + \mathbf{1}$$



$U(1)_{\text{DM}}$ charged Top Partners $\mathbf{Y}, \mathbf{Z}$

stop-like Signatures $\mathbf{Y} \rightarrow t + \chi$

Current Bounds

$$M_S > 1 \text{ TeV}, \quad M_T > 1.2 \text{ TeV}, \quad M_{X_{5/3}} > 1.2 \text{ TeV}, \quad M_{Y,Z} > 1.4 \text{ TeV}$$



- DM as pNGB in $SO(7)/SO(6)$ Composite Higgs model
- Natural and UV-Safe symmetry protecting the DM
- Viable mass range 200 – 400 GeV will be fully tested by XENON1T
- New LHC signals from $U(1)_{\text{DM}}$ charged top partners

Outlook:

$U(1)_{\text{DM}}$ can be gauged $\rightarrow$ very different phenomenology (work in progress)



Backup



One-Loop Scalar Potential

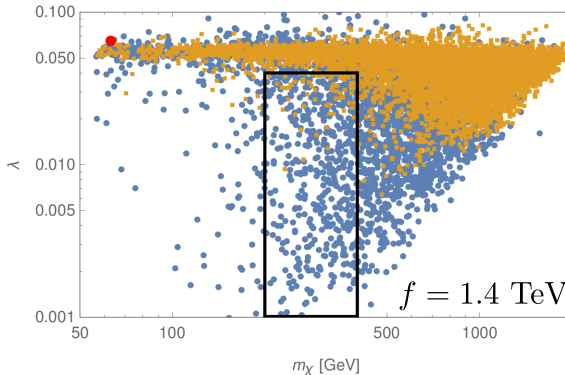
$$V(h, \chi) = \frac{1}{2} \mu_h^2 h^2 + \frac{\lambda_h}{4} h^4 + \mu_{\text{DM}}^2 \chi^* \chi + \lambda_{\text{DM}} (\chi^* \chi)^2 + \lambda h^2 \chi^* \chi$$

- Portal coupling λ important for DM-Pheno
- Integrate out vector and fermion resonances and compute Coleman-Weinberg Potential [S. R. Coleman, E. J. Weinberg '73](#)
→ Require sum rules on model parameters for finite potential

[Marzocca et al. '12, Pomarol, Riva '12](#)

- One layer of fermionic resonances: $m_\chi \sim \frac{m_h}{2}, \lambda \sim \frac{\lambda_h}{2} \sim 0.06$
Ruled out by direct detection bounds!

- Two Layers of fermionic resonances



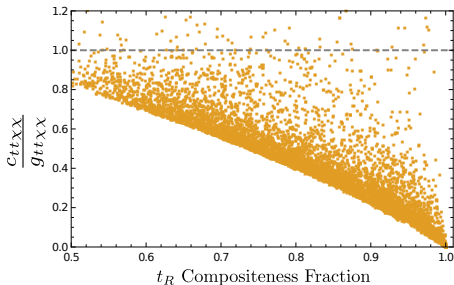
Orange: Top Partner masses above 1 TeV



Top interactions: $\mathcal{L}_t \sim -2c_{tt\chi\chi} \frac{m_t}{v^2} \bar{t} t \chi^* \chi \Rightarrow \sigma \propto \left(\frac{2m_\chi^2}{f^2} - \lambda \right)^2 + f(c_{tt\chi\chi})$

Mixing Contributions to $c_{tt\chi\chi}$ Azatov, Galloway '12

$$c_{tt\chi\chi} = g_{tt\chi\chi}(\xi) + \mathcal{O}(\xi (\epsilon_t/m_\psi)^2), \quad \xi = v^2/f^2$$



Fully composite t_R respects full $\mathcal{G}$ symmetry

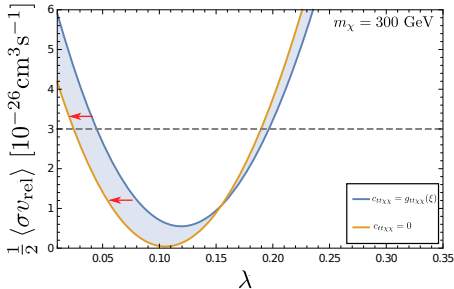
$\Rightarrow$ Non-derivative couplings vanish



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Shifts relic abundance to smaller λ

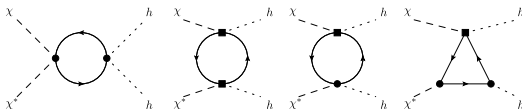


We use tree-level and one-loop couplings in V_{eff} for DM annihilation

BUT derivative interactions in $\mathcal{L}_{nl\sigma m}$

$$\mathcal{L}_{nl\sigma m} \supset \frac{c}{f^2} \partial_\mu |H|^2 \partial^\mu |\chi|^2$$

receive large corrections



$$c^{\text{tree}} = 1, \quad c^{1\text{-loop}} = \frac{N_c}{2\pi^2 f^2} \left(\epsilon_T^2 - \frac{\epsilon_Q^2}{8} \right) \log \frac{\Lambda^2}{m_\psi^2}$$

⇒ Assign 50% theoretical uncertainty to annihilation cross-section



Estimate of inverse Tuning (1 Layer Model)

$$\frac{1}{\Delta} \sim 2\xi \approx 6\% \text{ for } f = 1.4 \text{ TeV}$$

Numerical estimate: $\Delta = \max_i \left| \frac{\partial \log \xi}{\partial \log c_i} \right|$

