

# Anti-brane induced inflation

Marco Scalisi



based on

E. McDonough and MS - JCAP 1611, n.11, 028 (2016) [arXiv:1609.00364]

*supported by*

*Angelo Della Riccia Foundation  
FWO*

*European Union's Horizon 2020 Marie Skłodowska-Curie grant agreement N. 665501*



September 27<sup>th</sup>, 2017 - DESY Theory Workshop “Fundamental Physics in the Cosmos”- Hamburg

*M17 Nebula in Sagittarius - Gianni Benintende (Sicily - Italy)*

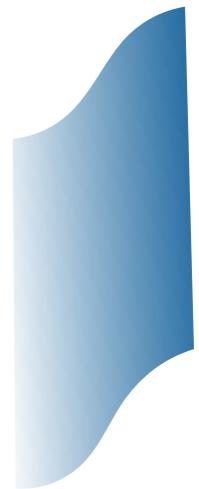
# Two Observations

de Sitter

$$V > 0$$



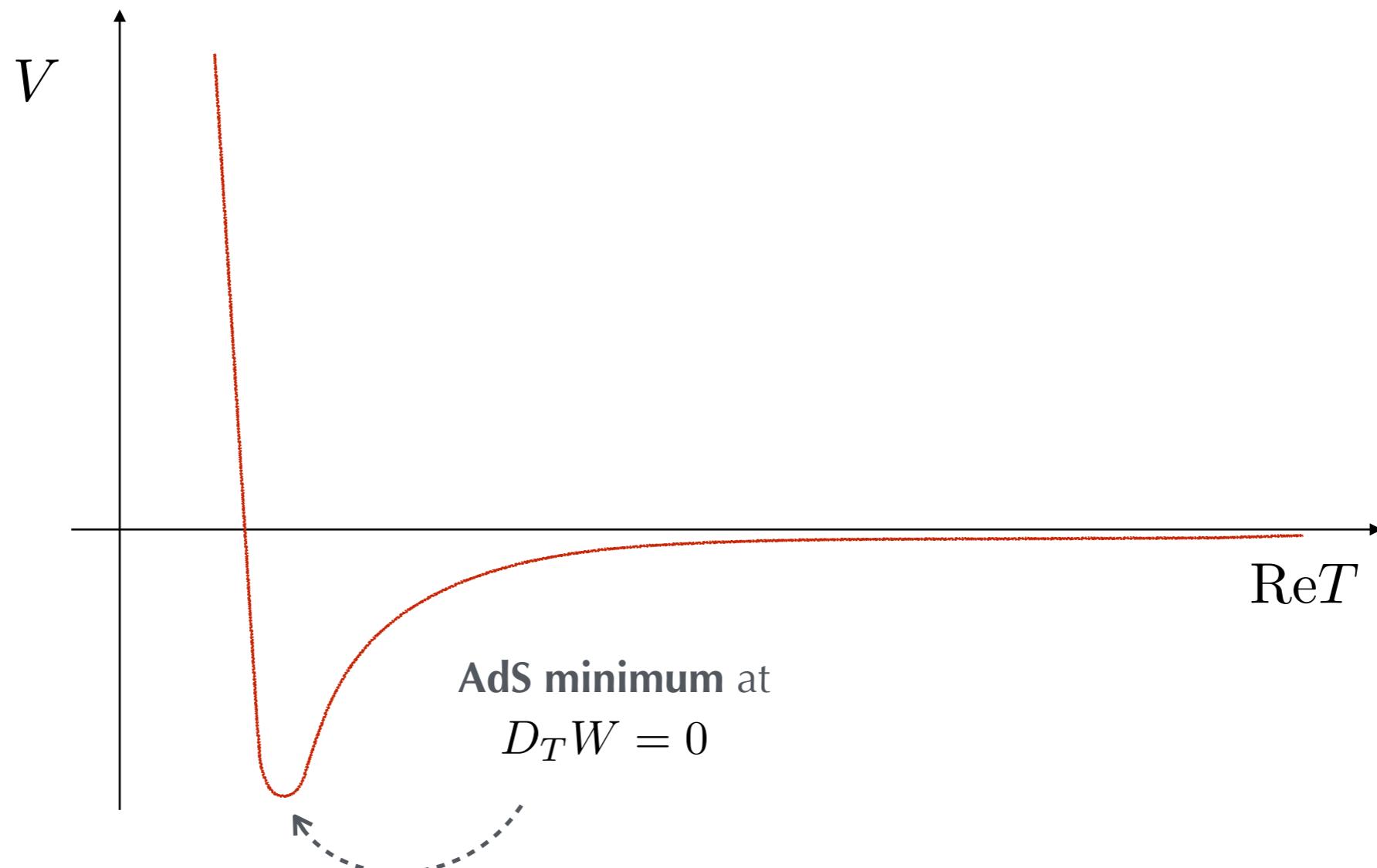
String Theory  
Anti-D3 Brane



$$K = -3 \log(T + \bar{T})$$

$$V = V_T$$

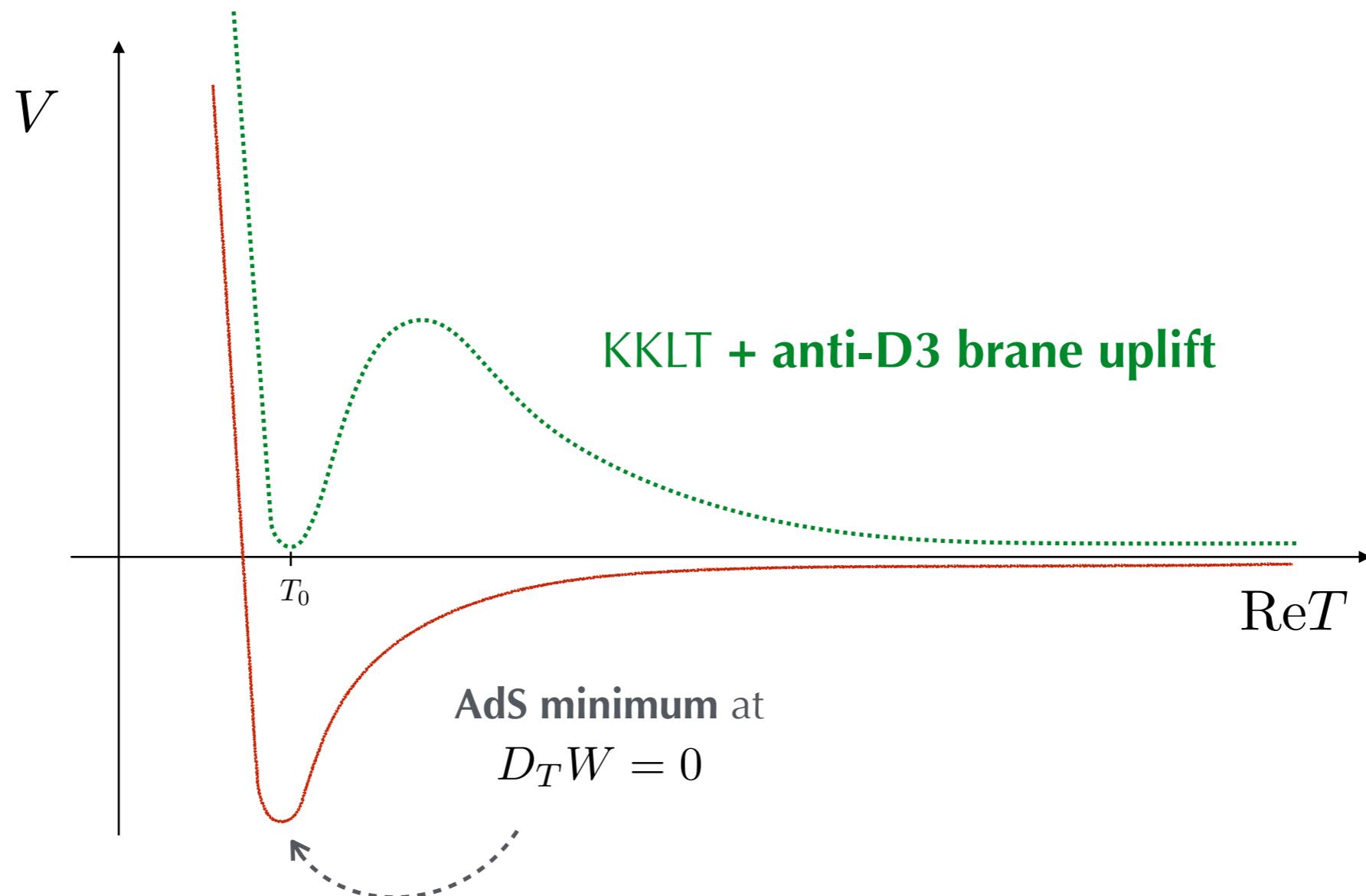
$$W = W_0 + A \exp(-aT)$$



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$$V = V_T + \frac{\mu^4}{(T + \bar{T})^2}$$

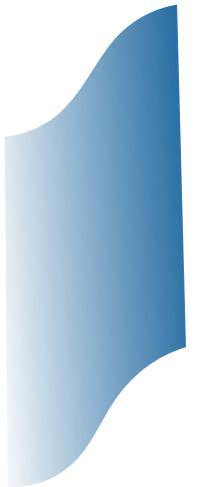


de Sitter

$$V > 0$$

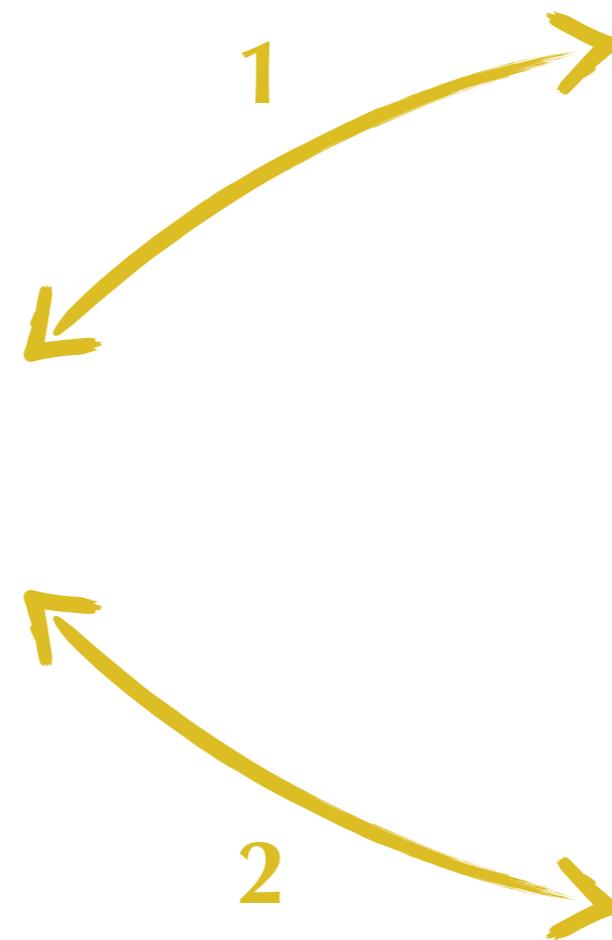


String Theory  
Anti-D3 Brane



de Sitter

$$V > 0$$



String Theory  
Anti-D3 Brane



Supergravity  
Nilpotent  
Superfield



*see Wrase's talk*

# The nilpotent superfield

$$S(x, \theta) = s(x) + \sqrt{2}\lambda(x)\theta + F(x)\theta^2$$

Volkov, Akulov 1972, 1973

Rocek; Ivanov, Kapustnikov 1978

$$S^2(x, \theta) = 0$$

Lindstrom, Rocek 1979

Casalbuoni, De Curtis, Dominici, Feruglio, Gatto 1989

Komargodski, Seiberg 2009



$$S(x, \theta) = \frac{\lambda\lambda}{2F} + \sqrt{2}\lambda\theta + F\theta^2$$

**no scalar!**

**just fermions!**

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- SUSY **non-linearly** realized
- SUSY **spontaneously broken**  $F \neq 0$

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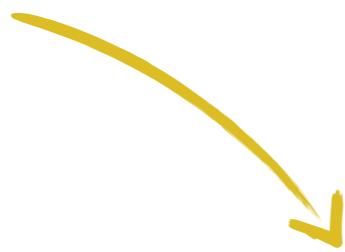
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when coupled to SUGRA  
**de Sitter Supergravity**

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Bergshoeff, Freedman, Kallosh, Van Proyen 2015

Hasegawa, Yamada 2015

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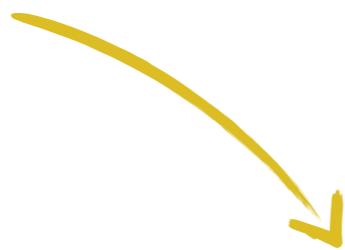
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Hasegawa, Yamada 2015

Farakos, Kehagias 2013

“Decoupling Limits of sGoldstino Modes in  
Global and Local Supersymmetry”  
**Appendix of arXiv:1302.0866v1**



# de Sitter supergravity

$$K = S \bar{S}$$

$$W = W_0 + M S \quad \text{most general } W \text{ as } S^2 = 0$$



at  $S = 0$

$$V = M^2 - 3W_0^2$$

no scalar degrees of freedom!

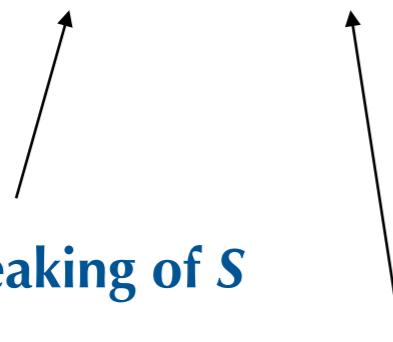
**'de Sitter supergravity'**

*Antoniadis, Dudas, Ferrara, Sagnotti 2014*

*Bergshoeff, Freedman, Kallosh, Van Proeyen 2015*

*Hasegawa, Yamada 2015*

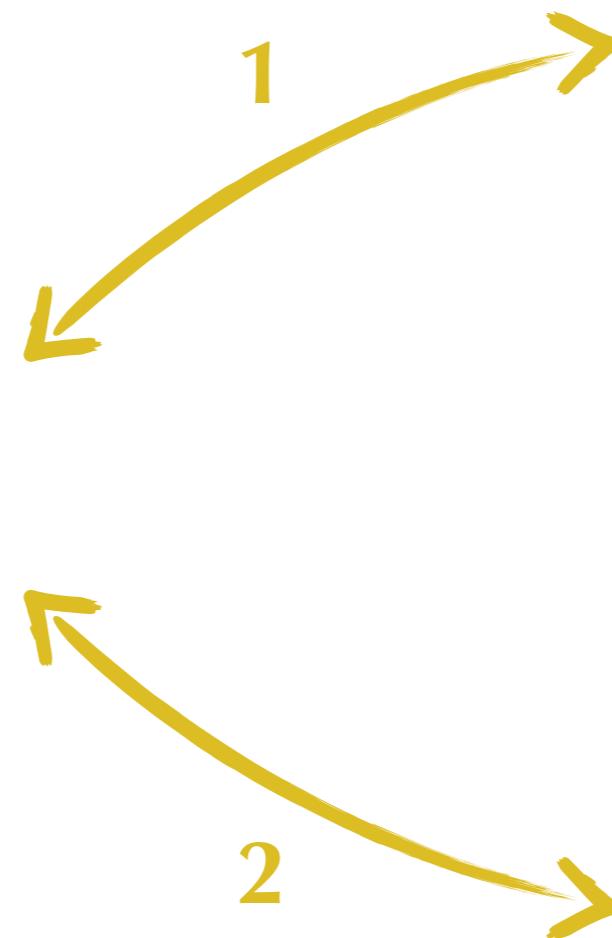
**SUSY breaking of  $S$**



**Gravitino mass**

de Sitter

$$V > 0$$



String Theory  
Anti-D3 Brane

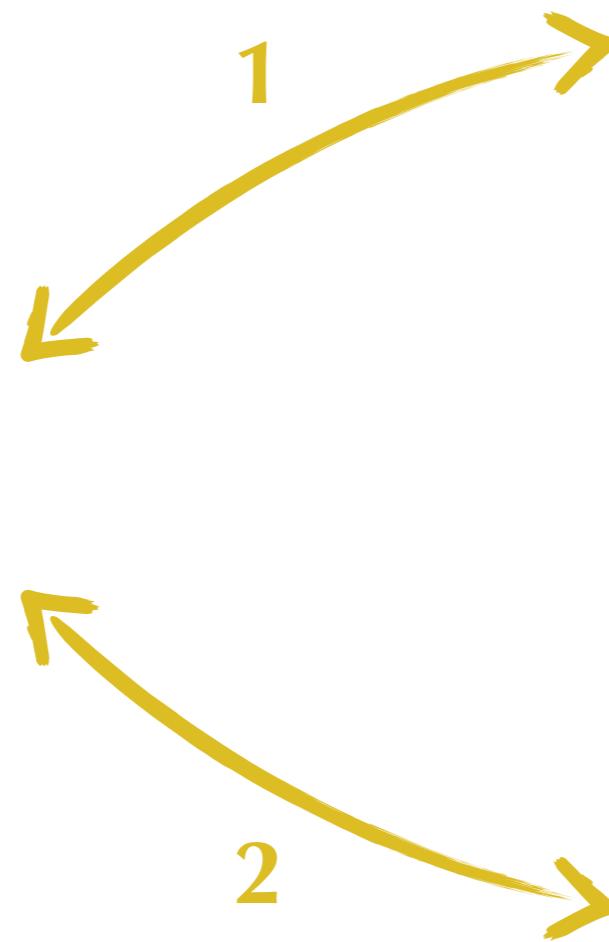


Supergravity  
Nilpotent  
Superfield



de Sitter

$$V > 0$$



String Theory  
Anti-D3 Brane



Supergravity  
Nilpotent  
Superfield





String Theory

Anti-D3 Brane



QUESTION #1

**Does the Anti-D3 brane  
break SUSY spontaneously?**

*McGuirk, Shiu, Ye 2012*

QUESTION #2

**Can we package the  
uplifting term into  $K$  and  $W$ ?**

*Ferrara, Kallosh, Linde 2014*

Supergravity

Nilpotent  
Superfield



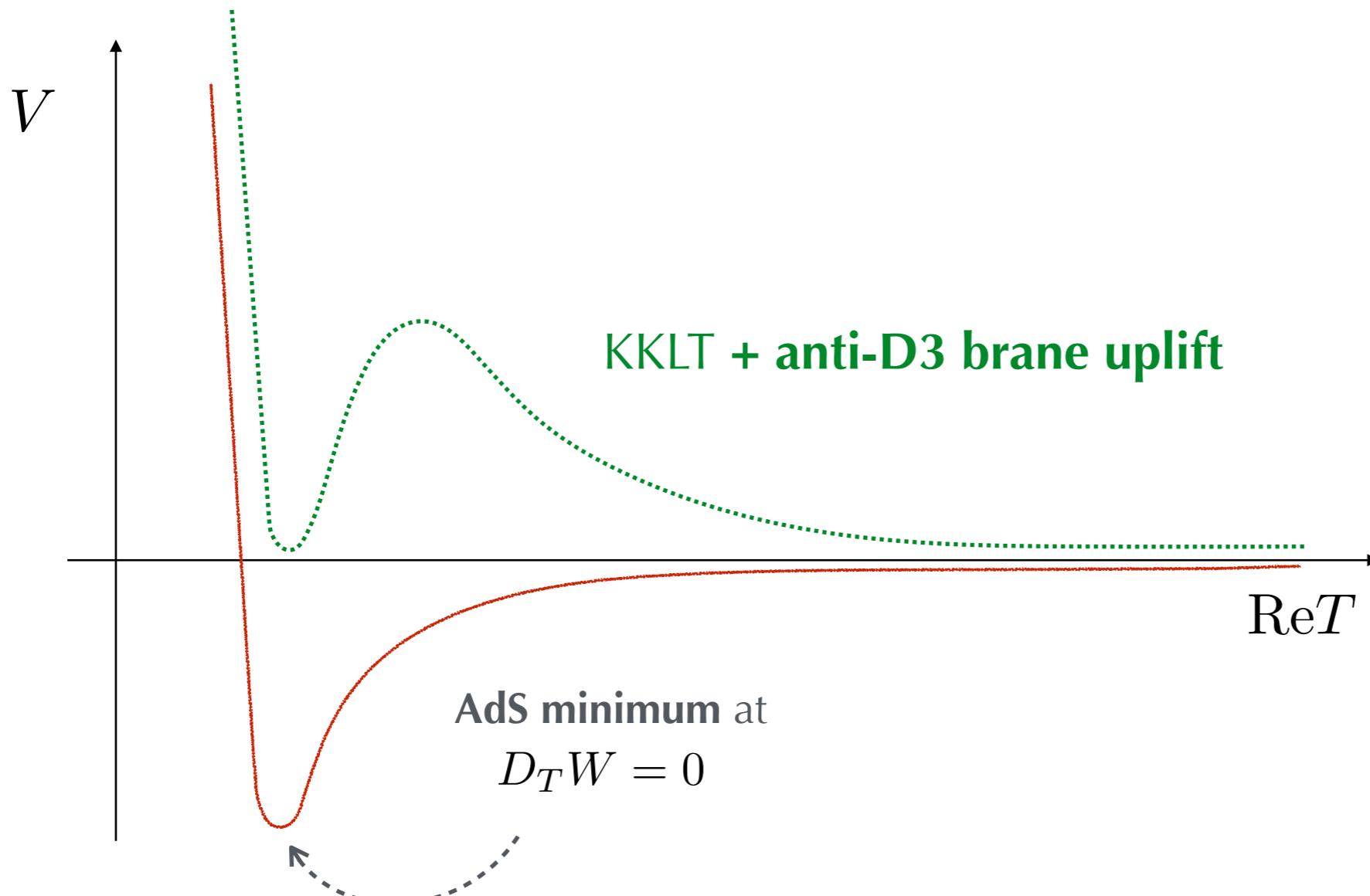
# KKLT + Nilpotent Superfield

Ferrara, Kallosh, Linde 2014

$$K = -3 \log(T + \bar{T} - S\bar{S})$$

$$W = W_0 + A \exp(-aT) + \mu^2 S$$

$$V = V_T + \frac{\mu^4}{3(T + \bar{T})^2}$$



## NILPOTENT SUPERFIELD $\longleftrightarrow$ D-BRANES

The **four dimensional** description of an **anti-D3 brane**  
in an  $N = 1$  flux background is  
a supergravity theory of a **nilpotent superfield**.

*Kallosh & Wrane 2014*

## NILPOTENT SUPERFIELD    $\longleftrightarrow$    D-BRANES

The **fermions** arising when one or more **anti-branes**, placed in certain geometries, break supersymmetry spontaneously can often be packaged into **constrained superfields**.

*McGuirk, Shiu, Ye 2012*

*Kallosh & Wrane 2014*

*Bergshoeff, Dasgupta, Kallosh, Van Proeyen & Wrane 2014*

*Kallosh, Quevedo & Uranga 2015*

*Bertolini, Musso, Papadimitriou & Raj 2015*

*Aparicio, Quevedo & Valandro 2015*

*Garca-Etxebarria, Quevedo & Valandro 2015*

*Dasgupta, Emelin and McDonough 2016*

*Vercnocke & Wrane 2016*

*Kallosh, Vercnocke & Wrane 2016*

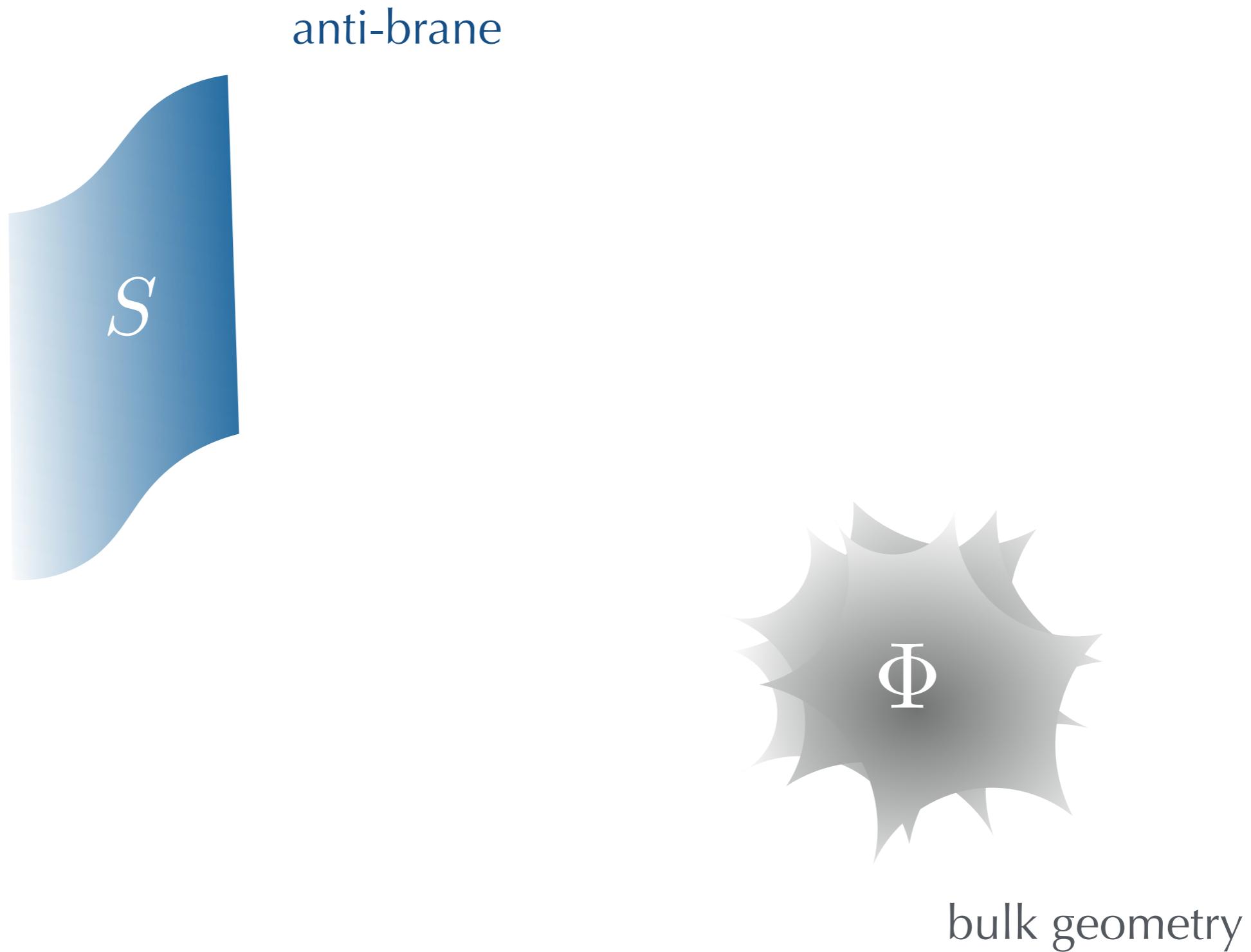
*Bandos, Heller, Kuzenko, Martucci and Sorokin 2016*

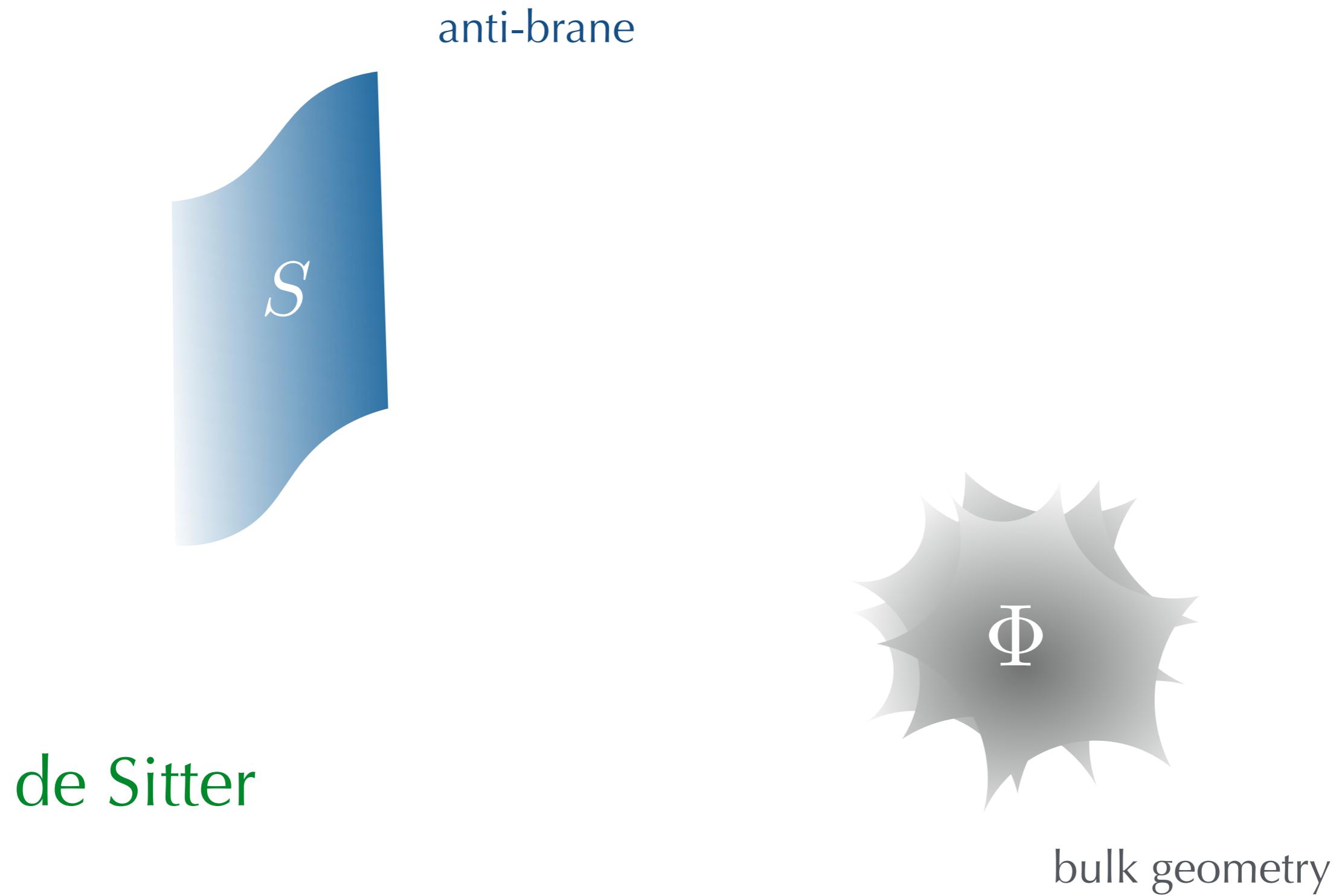
*Aalsma, van der Schaar, Vercnocke 2017*

*Garcia del Moral, Parameswaran, Quiroz, Zavala 2017*

# Our Work

*McDonough & MS 2016*





de Sitter



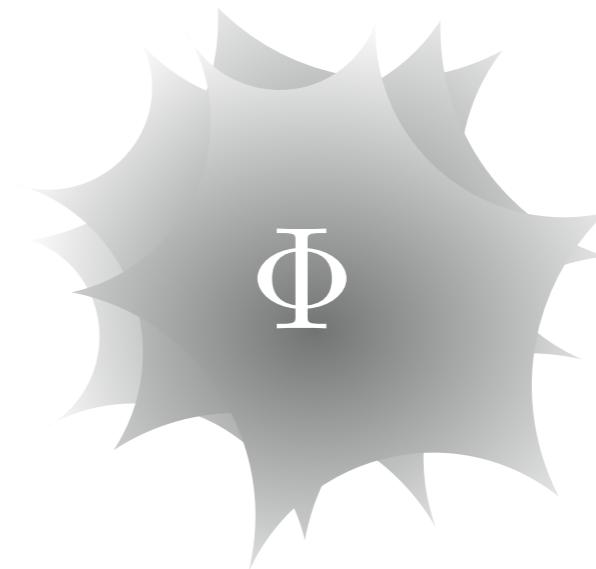
anti-brane

!

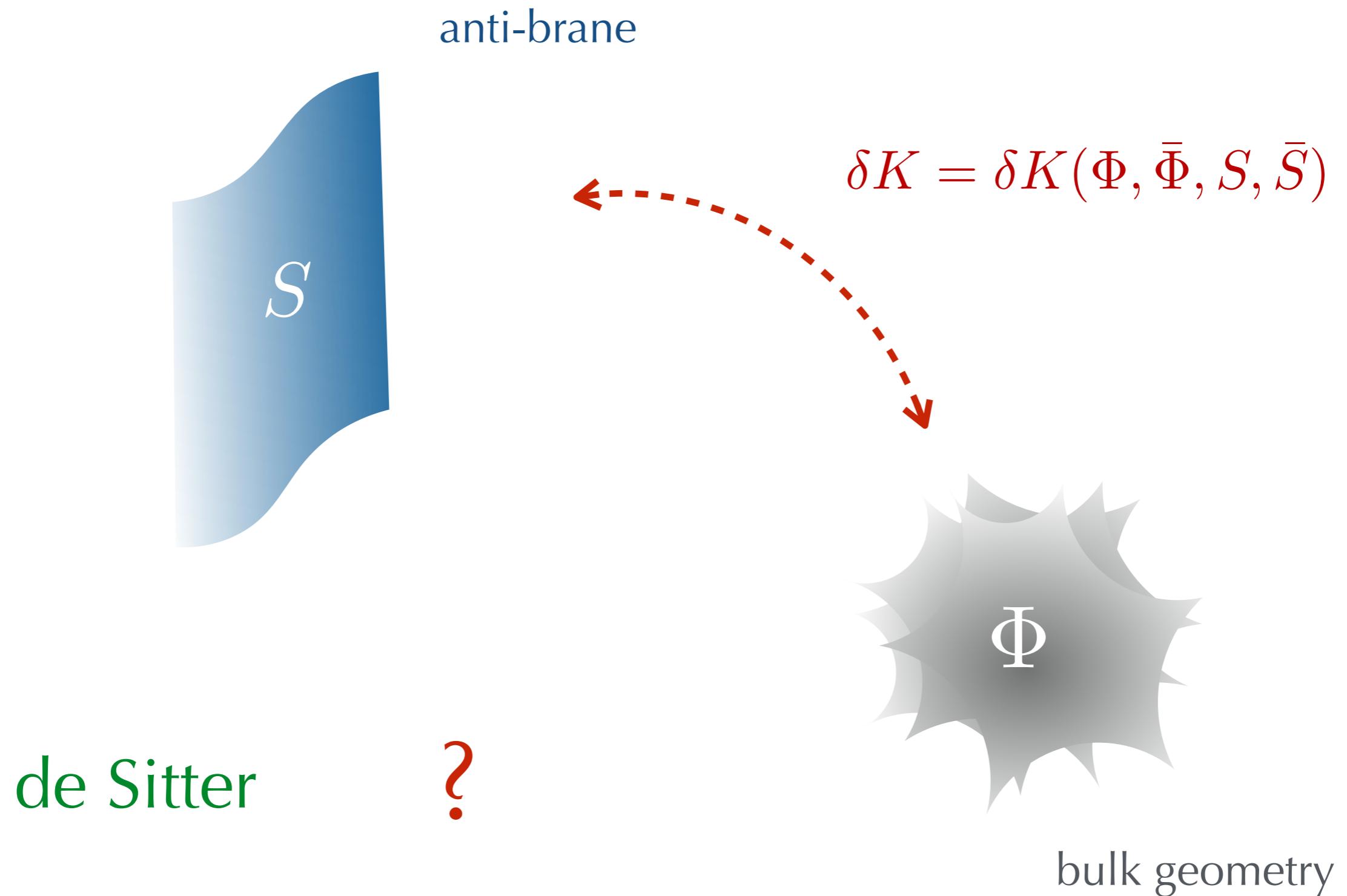
10D analysis of the backreaction

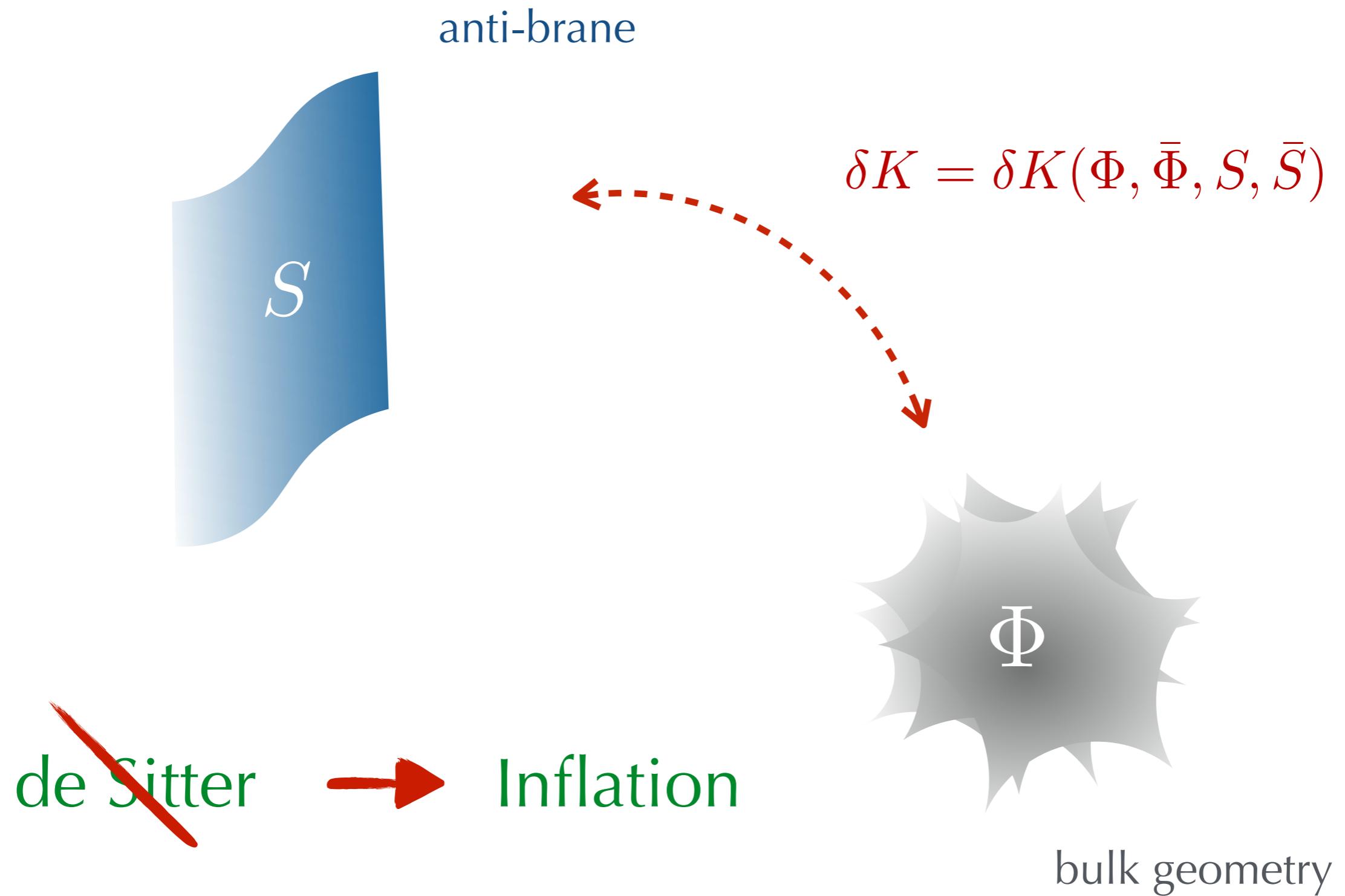
*Moritz, Retolaza, Westphal 2017*

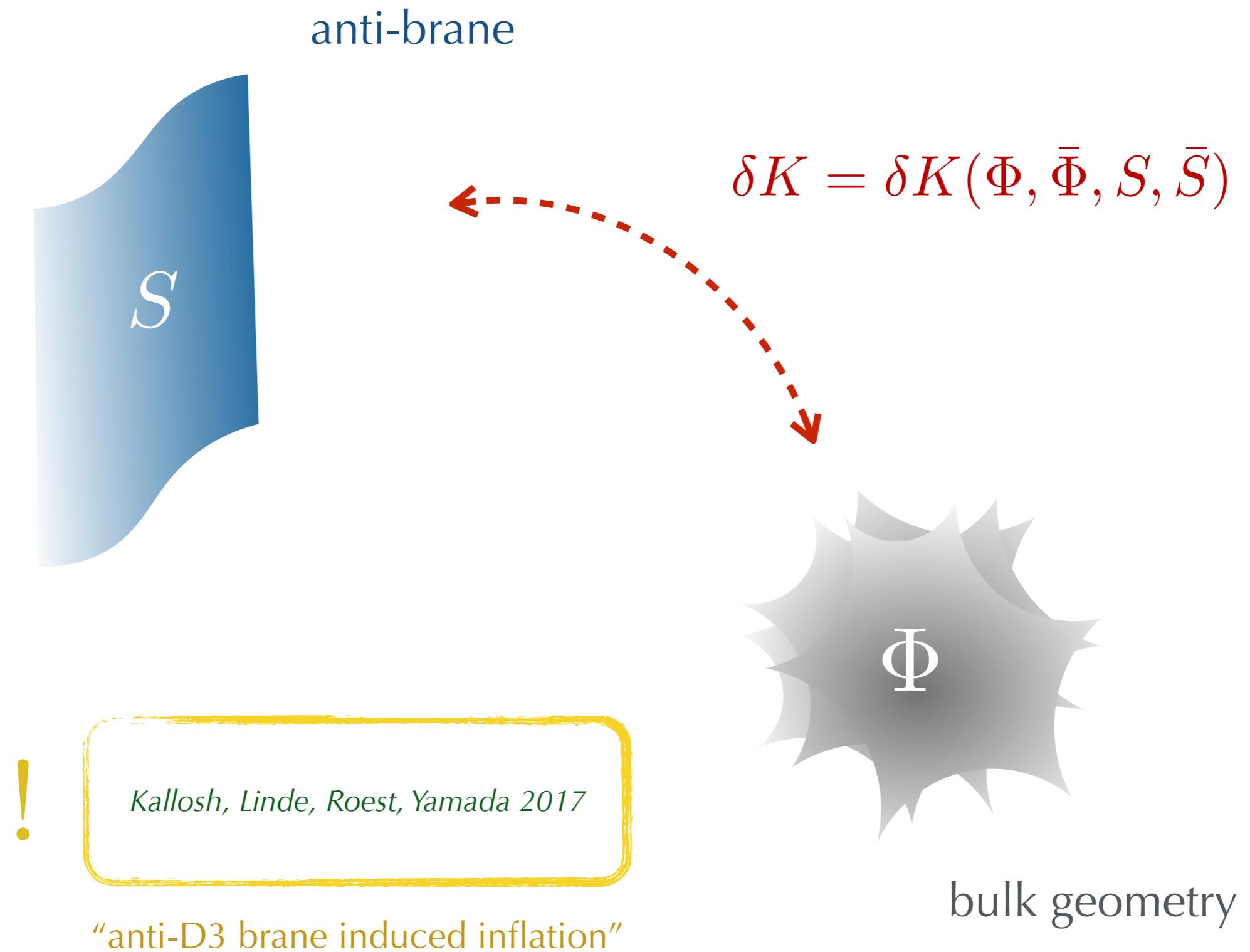
*see Moritz's talk*



bulk geometry







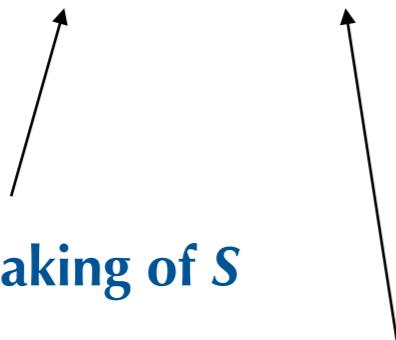
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$$W = W_0 + M S$$



at  $S = 0$

$$V = M^2 - 3W_0^2$$



**SUSY breaking of  $S$**

**Gravitino mass**

# Flat Kähler geometry

McDonough & MS 2016

$$K = S \bar{S} - \frac{1}{2} (\Phi - \bar{\Phi})^2$$

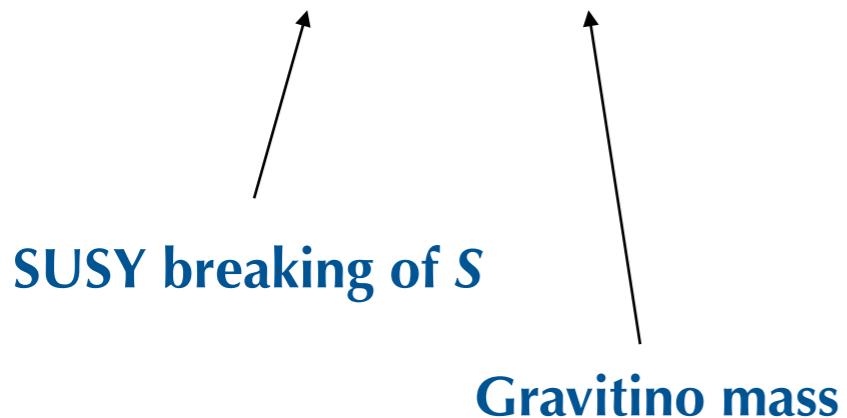
$$W = W_0 + M S$$



at  $S = 0$  and at  $Im\Phi = 0$

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**no inflation!** (shift symmetry not broken!)



# Flat Kähler geometry

McDonough & MS 2016

$$K = S\bar{S} - \frac{1}{2} (\Phi - \bar{\Phi})^2$$

**inflation!**

$$W = W_0 + M S$$



$$W = g(\Phi) + f(\Phi)S$$

Kallosch, Linde, MS 2014

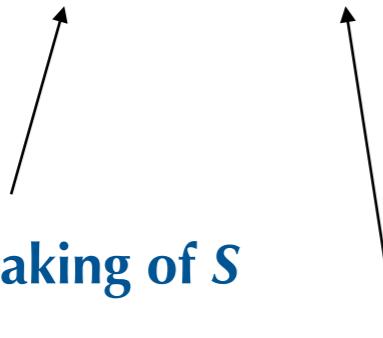


at  $S = 0$  and at  $Im\Phi = 0$

Kallosch, Linde 2014  
Dall'Agata, Zwirner 2014

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**SUSY breaking of  $S$**

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# Flat Kähler geometry

McDonough & MS 2016

$$K = S\bar{S} - \frac{1}{2} (\Phi - \bar{\Phi})^2 + f(\Phi, \bar{\Phi})S\bar{S} + g(\Phi, \bar{\Phi})S + \bar{g}(\Phi, \bar{\Phi})\bar{S}$$

$$W = W_0 + MS$$

$$\delta K$$

# Flat Kähler geometry

McDonough & MS 2016

$$K = -\frac{1}{2} (\Phi - \bar{\Phi})^2 + [1 + f(\Phi, \bar{\Phi})] S \bar{S}$$

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## Phenomenological Flexibility

- Arbitrary **Inflationary Potential**
- Controllable level of **SUSY breaking**
- Tunable level of the **CC**



## SUSY broken just in the $S$ direction

$$D_\Phi W = 0 \quad D_S W = M$$

## SUSY breaking of S

$$K^{S\bar{S}} |D_S W|^2$$

**Gravitino mass**

$$V = M^2 F(\Phi) - 3W_0^2$$

**Effects of Kähler corrections**

**“Original” SUSY breaking of S**

- ! **SUSY broken just in the S direction**
- !  $D_\Phi W = 0 \quad D_S W = M$

# Ex. Quadratic Inflation

McDonough & MS 2016

$$K = -\frac{1}{2} (\Phi - \bar{\Phi})^2 + S\bar{S} - \frac{m^2}{2M^2} \Phi\bar{\Phi} \cdot S\bar{S}$$

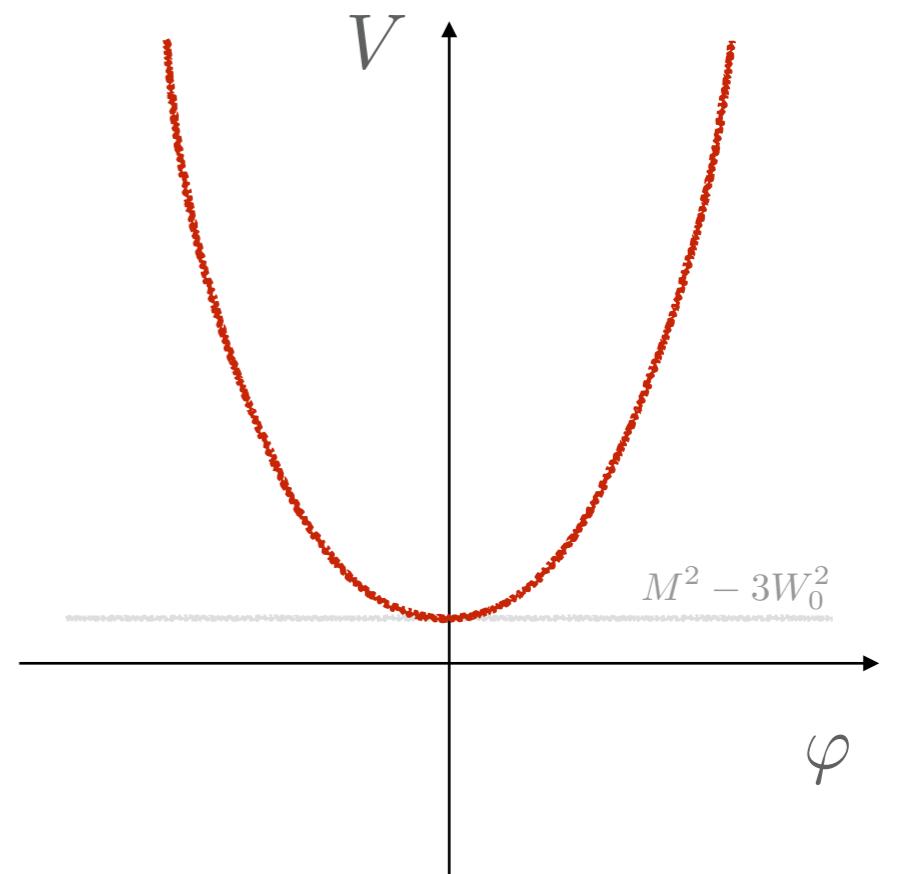
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at  $S = 0$  and at  $Im\Phi = 0$

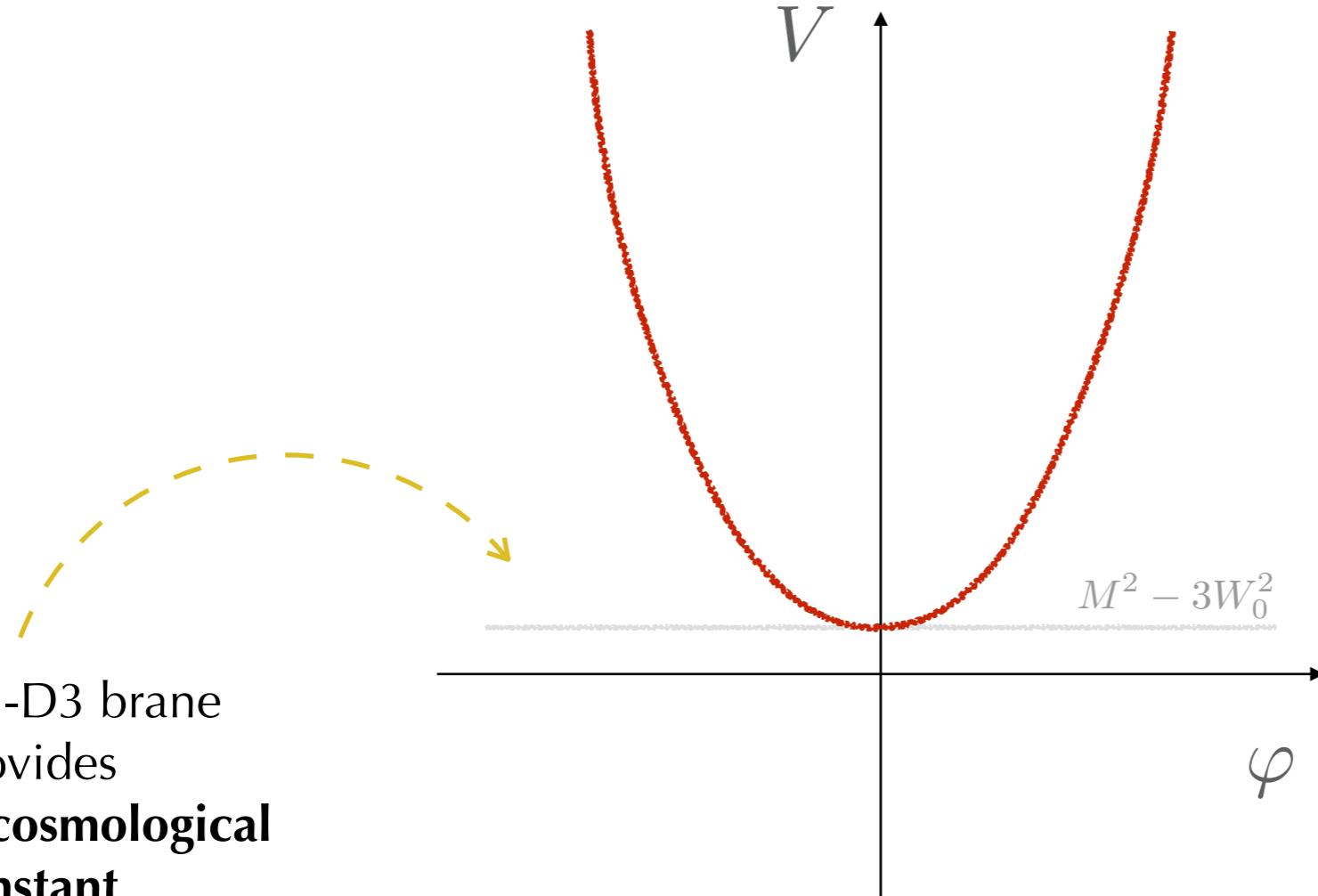
$$V = (M^2 - 3W_0^2) + \frac{1}{2}m^2\varphi^2$$

with  $\varphi = \sqrt{2}\text{Re}\Phi$



# Ex. Quadratic Inflation

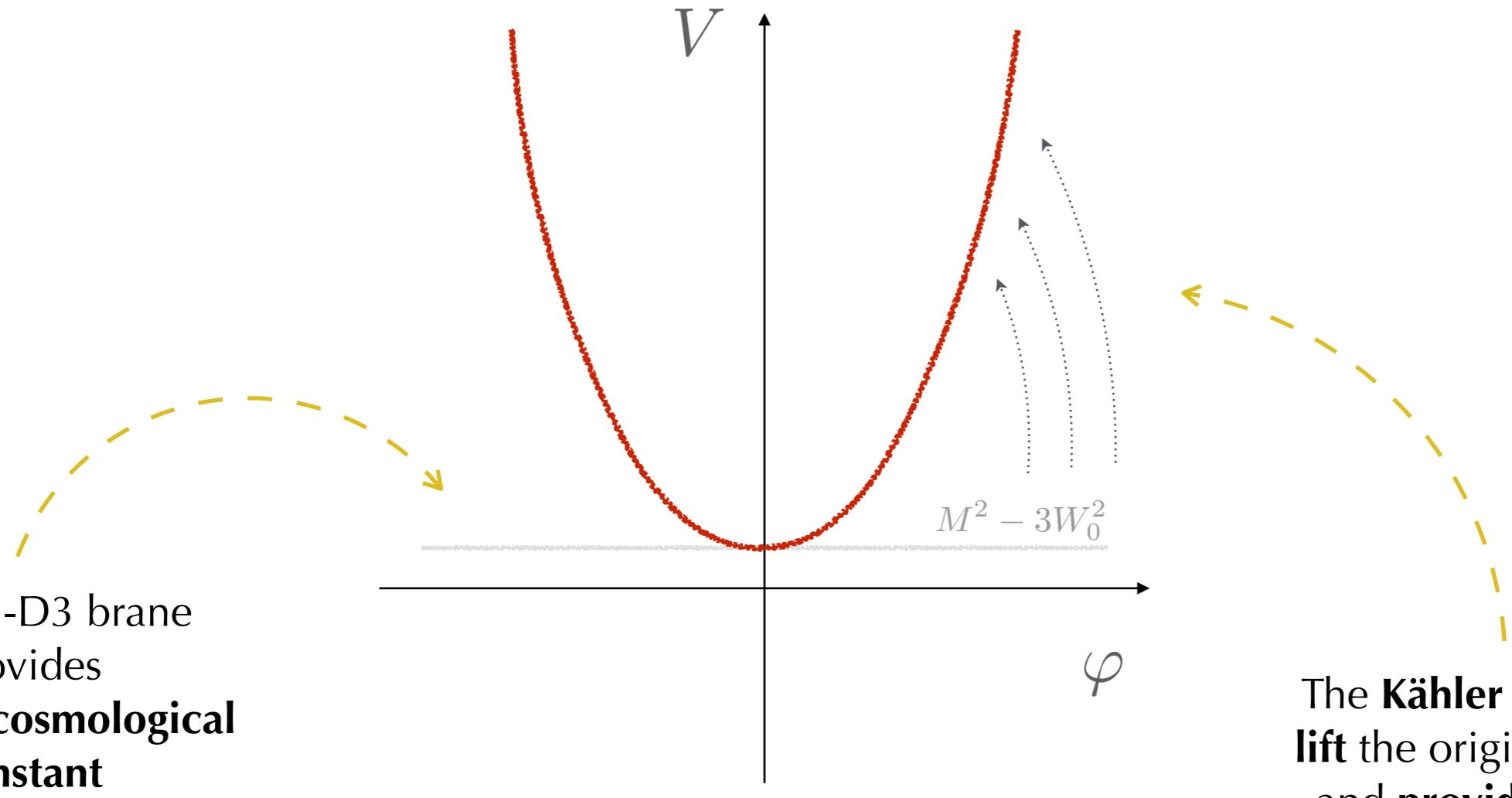
McDonough & MS 2016



The anti-D3 brane  
provides  
the **small cosmological  
constant**

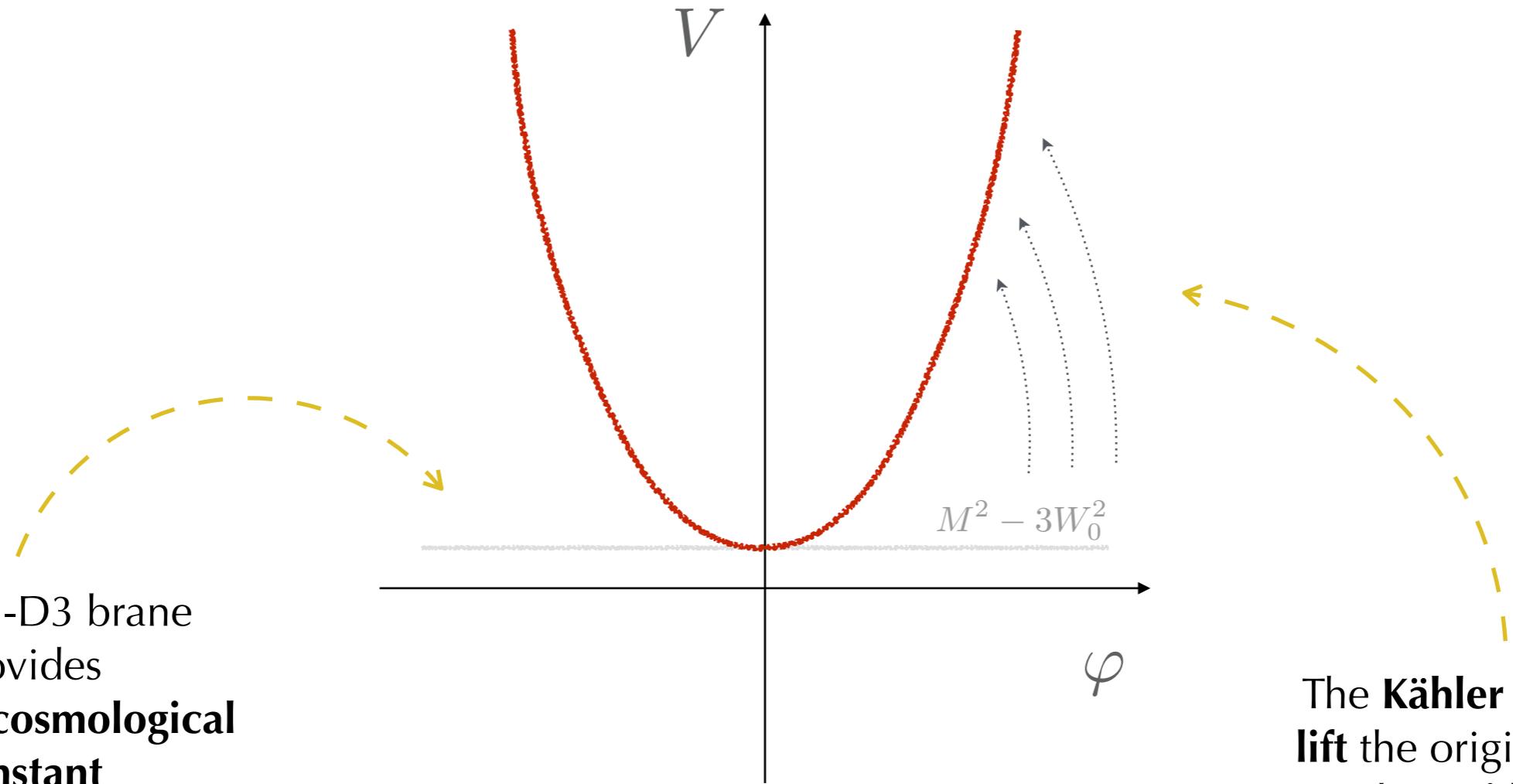
# Ex. Quadratic Inflation

McDonough & MS 2016



# Ex. Quadratic Inflation

McDonough & MS 2016



Inflation happens at large Kähler corrections

# Hyperbolic Kähler geometry

McDonough & MS 2016

# Hyperbolic Kähler geometry

McDonough & MS 2016

$$K = S\bar{S} - 3\alpha \log \left( \frac{\Phi + \bar{\Phi}}{2|\Phi|} \right)$$

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related to

$$-3\alpha \log (\Phi + \bar{\Phi})$$

by means of a  
Kähler transformation

# Hyperbolic Kähler geometry

McDonough & MS 2016

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related to

$$-3\alpha \log (\Phi + \bar{\Phi})$$

by means of a  
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at  $S = 0$  and at  $Im\Phi = 0$

$$V = M^2 - 3W_0^2$$

**no inflation!** (shift symmetry not broken!)

# Hyperbolic Kähler geometry

McDonough & MS 2016

$$K = S\bar{S} - 3\alpha \log \left( \frac{\Phi + \bar{\Phi}}{2|\Phi|} \right) + f(\Phi, \bar{\Phi})S\bar{S} + g(\Phi, \bar{\Phi})(S + \bar{S})$$
$$W = W_0 + MS$$

$\longleftrightarrow$   
 $\delta K$

at  $S = 0$  and at  $Im\Phi = 0$



# Hyperbolic Kähler geometry

McDonough & MS 2016

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$$W = W_0 + MS$$

$\longleftrightarrow$   
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at  $S = 0$  and at  $Im\Phi = 0$



**the same story as before!**

$$V = M^2 F(\Phi) - 3W_0^2$$

# Hyperbolic Kähler geometry

McDonough & MS 2016

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$$W = W_0 + MS$$

$\longleftrightarrow$   
 $\delta K$

at  $S = 0$  and at  $Im\Phi = 0$

$$f = \sum_{n=1}^{\infty} f_n |\Phi|^n \quad g = \sum_{n=1}^{\infty} g_n |\Phi|^n \quad \text{small perturbative corrections!}$$



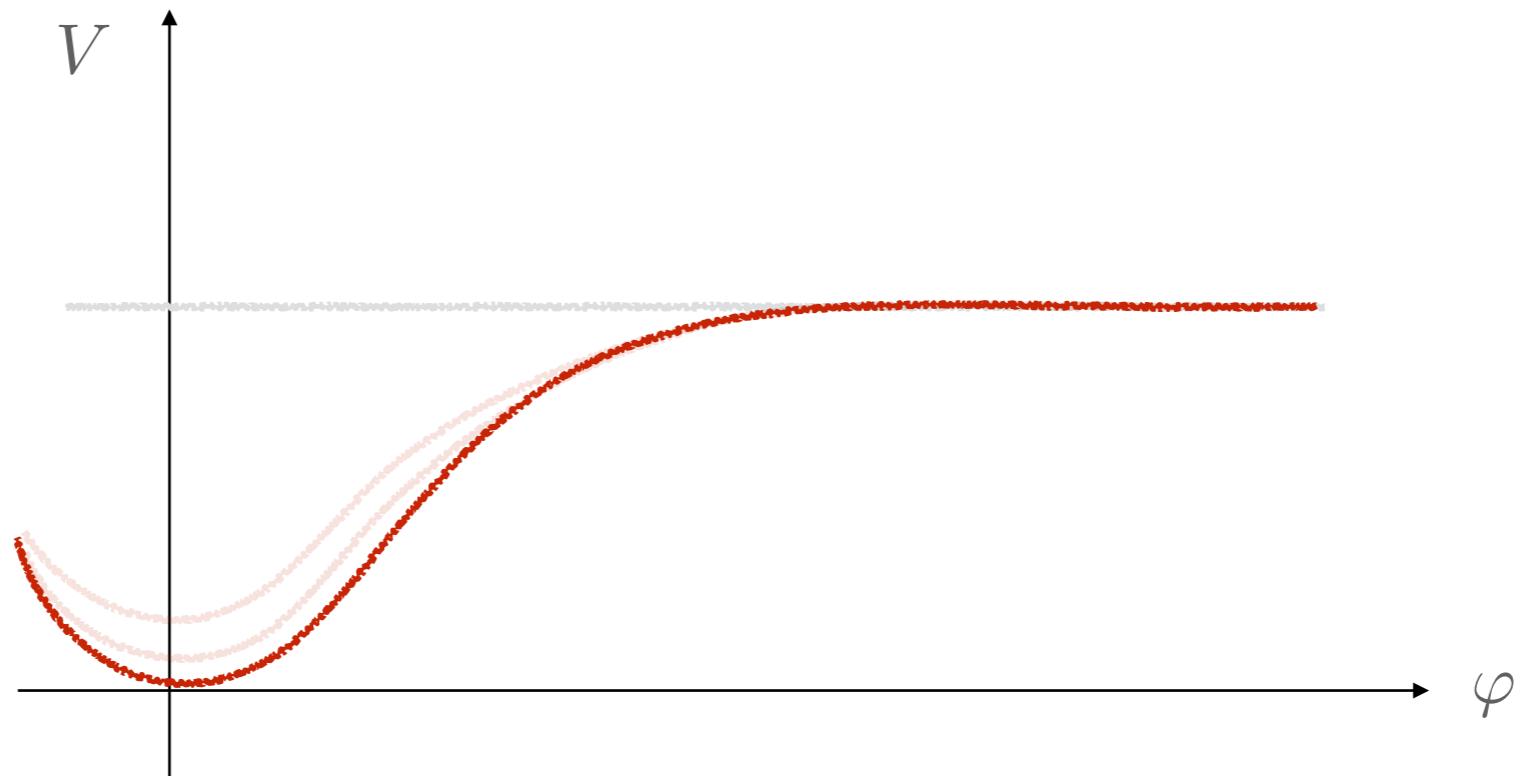
$$V = V_0 + V_1 \exp \left( -\sqrt{2/3\alpha} \varphi \right) + \dots$$



$$V_0 = M^2 - 3W_0^2$$

# Hyperbolic Kähler geometry

McDonough & MS 2016



$$V = V_0 + V_1 \exp\left(-\sqrt{2/3\alpha} \varphi\right) + \dots$$

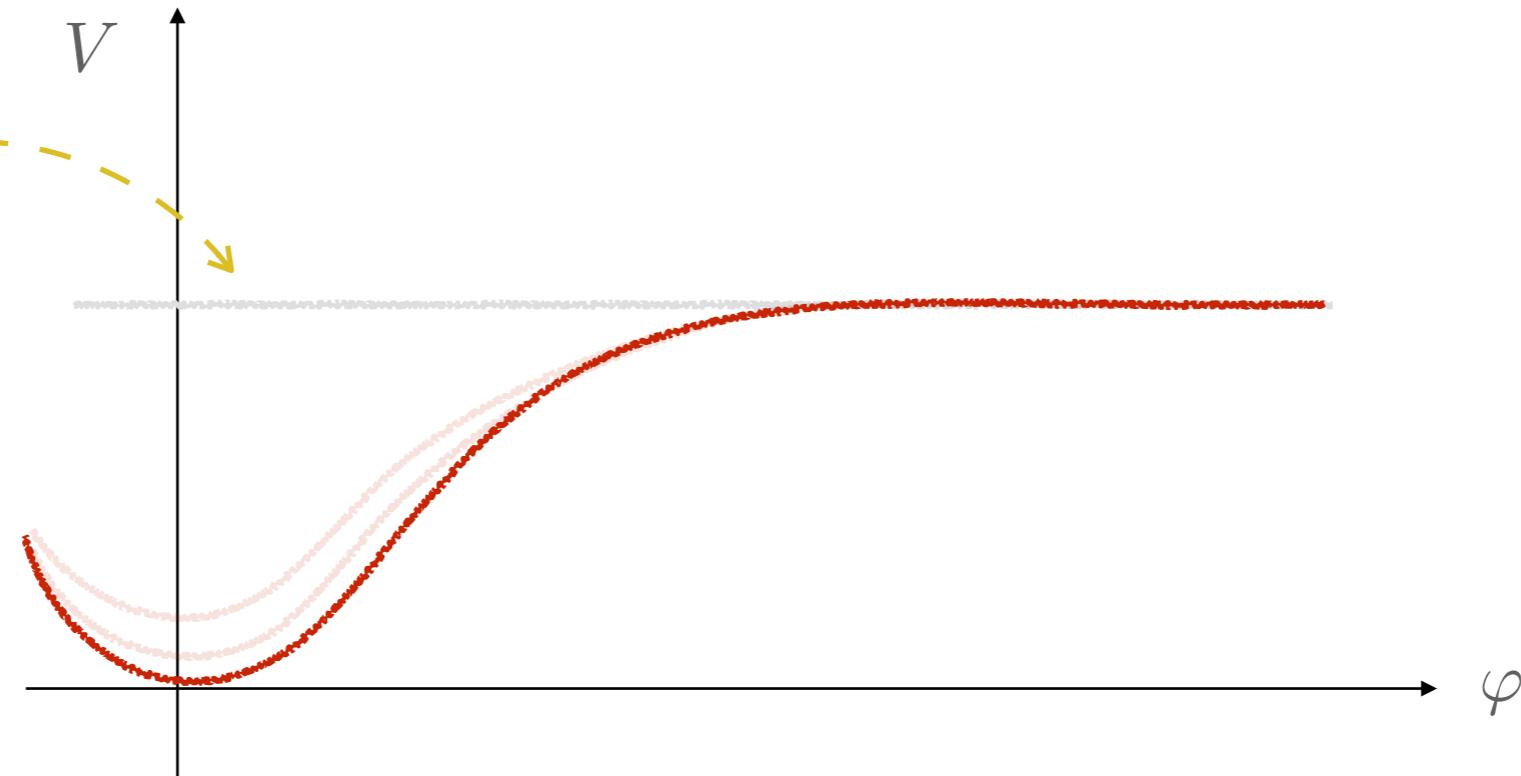


$$V_0 = M^2 - 3W_0^2$$

# Hyperbolic Kähler geometry

McDonough & MS 2016

The anti-D3 brane  
provides  
the **Hubble**  
**inflationary energy**



$$V = V_0 + V_1 \exp\left(-\sqrt{2/3\alpha} \varphi\right) + \dots$$

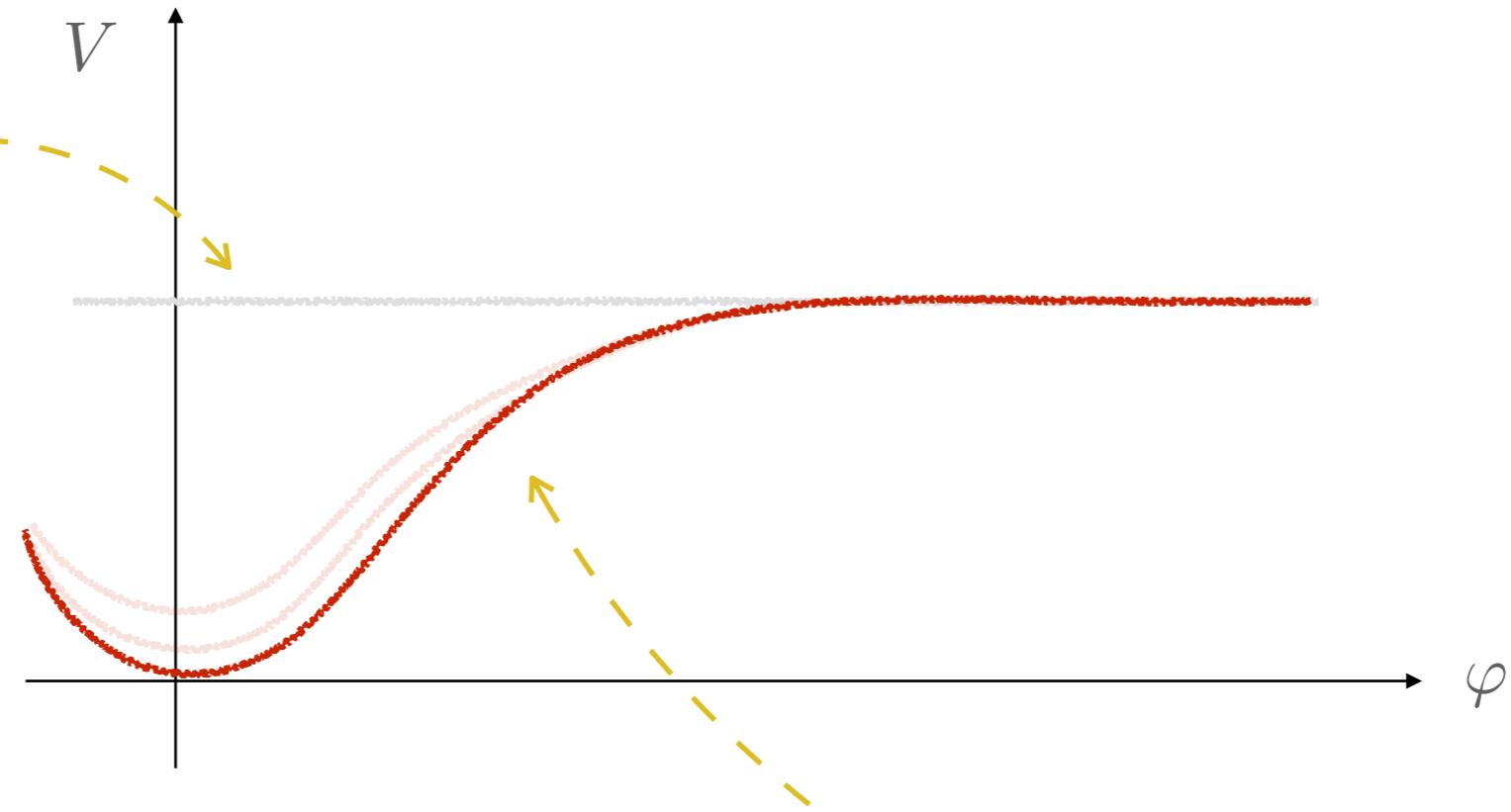


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McDonough & MS 2016

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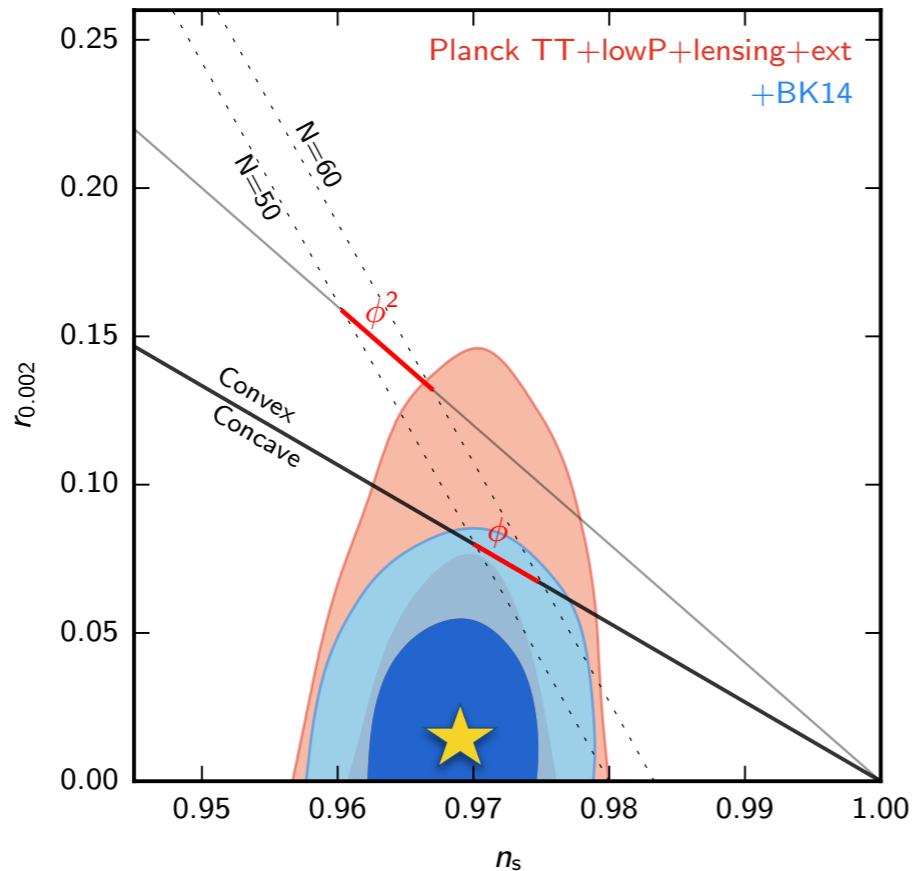
$$\uparrow$$
  
$$V_0 = M^2 - 3W_0^2$$

The **Kähler corrections**  
provide the **exponential**  
**fall-off** from dS

Inflation happens at  
**small Kähler corrections**

# Hyperbolic Kähler geometry

McDonough & MS 2016



## α-attractors

$$n_s = 1 - \frac{2}{N} \quad r = \frac{12\alpha}{N^2}$$

$$V = V_0 + V_1 \exp \left( -\sqrt{2/3\alpha} \varphi \right) + \dots$$



$$V_0 = M^2 - 3W_0^2$$

# Supersymmetry breaking scale

McDonough & MS 2016

$$V = M^2 F(\Phi) - 3W_0^2$$

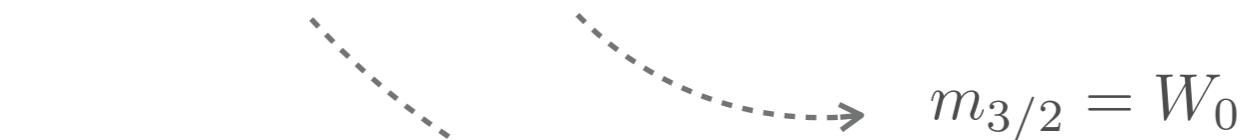

$$m_{3/2} = W_0$$


$$F(\Phi) \equiv \frac{1}{1 + f(\Phi, \Phi)}$$

# Supersymmetry breaking scale

McDonough & MS 2016

$$V = M^2 F(\Phi) - 3W_0^2$$


$$m_{3/2} = W_0$$

**at the minimum**

$$\Lambda = M^2 F(0) - 3W_0^2$$

$$F(\Phi) \equiv \frac{1}{1 + f(\Phi, \Phi)}$$

# Supersymmetry breaking scale

McDonough & MS 2016

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**large Kähler corrections**

$$|f| \gg 1 \quad |F| \ll 1$$

**small Gravitino mass**

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**large Kähler corrections**

$$|f| \gg 1 \quad |F| \ll 1$$

**small Gravitino mass**

**small Kähler corrections**

$$|f| \ll 1 \quad |F| \sim 1$$

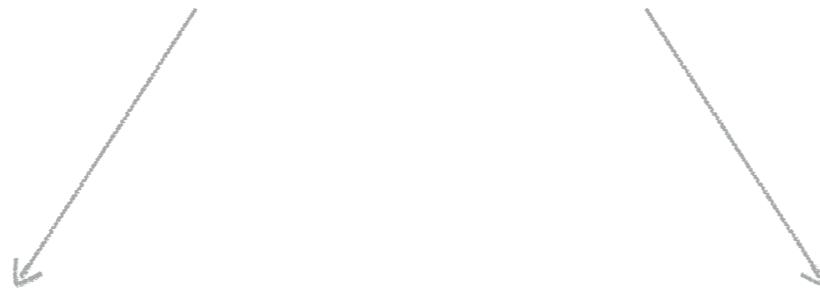
$$M \geq H$$

**large Gravitino mass**

# Stability

McDonough & MS 2016

$$m_{\text{Im}\Phi}^2 = -4W_0^2 + 4M^2 F(\Phi)$$



during inflation

$$\frac{m_{\text{Im}\Phi}^2}{H^2} = 12 + \frac{24W_0^2}{M^2 F(\Phi) - 3W_0^2}$$

at the minimum

$$m_{\text{Im}\Phi}^2 = 8W_0^2 + 4\Lambda$$

# Recap-Cartoon

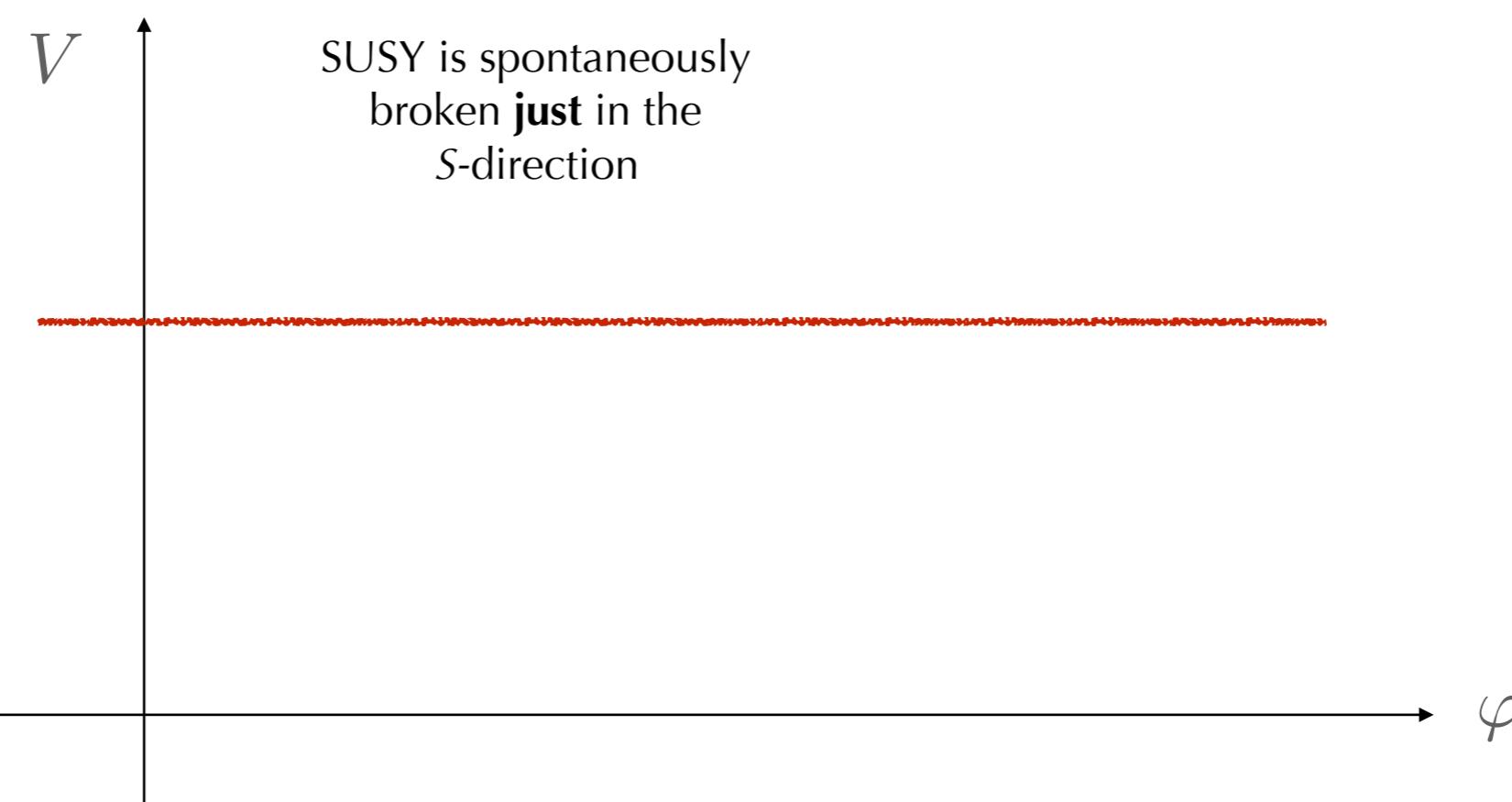


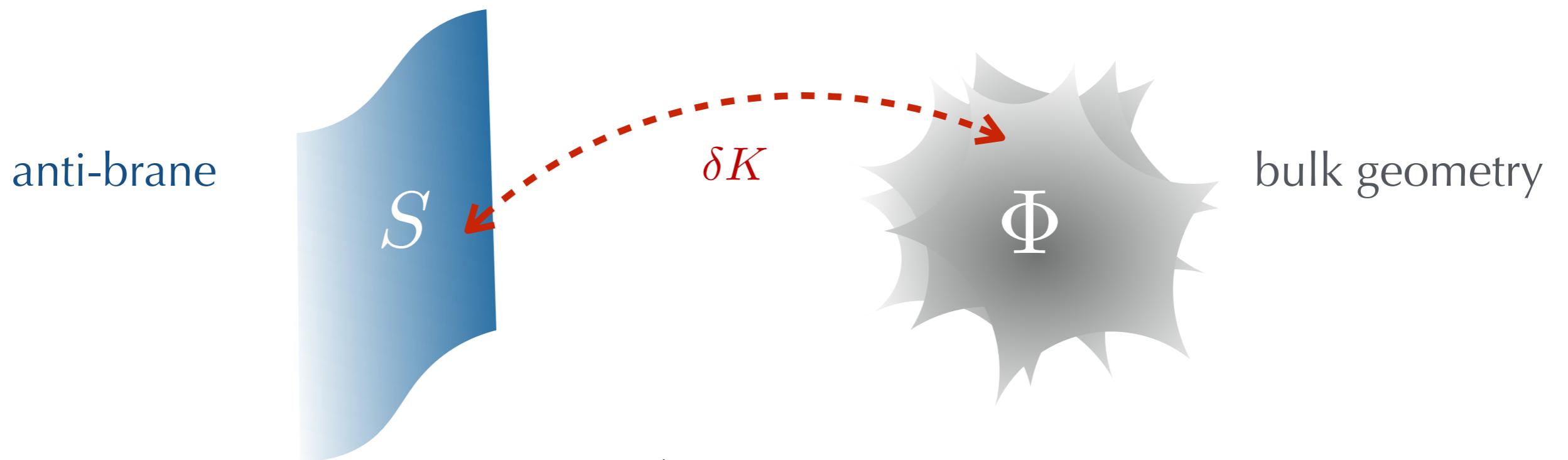
$$D_S W \neq 0$$

$$D_\Phi W = 0$$



pure de Sitter

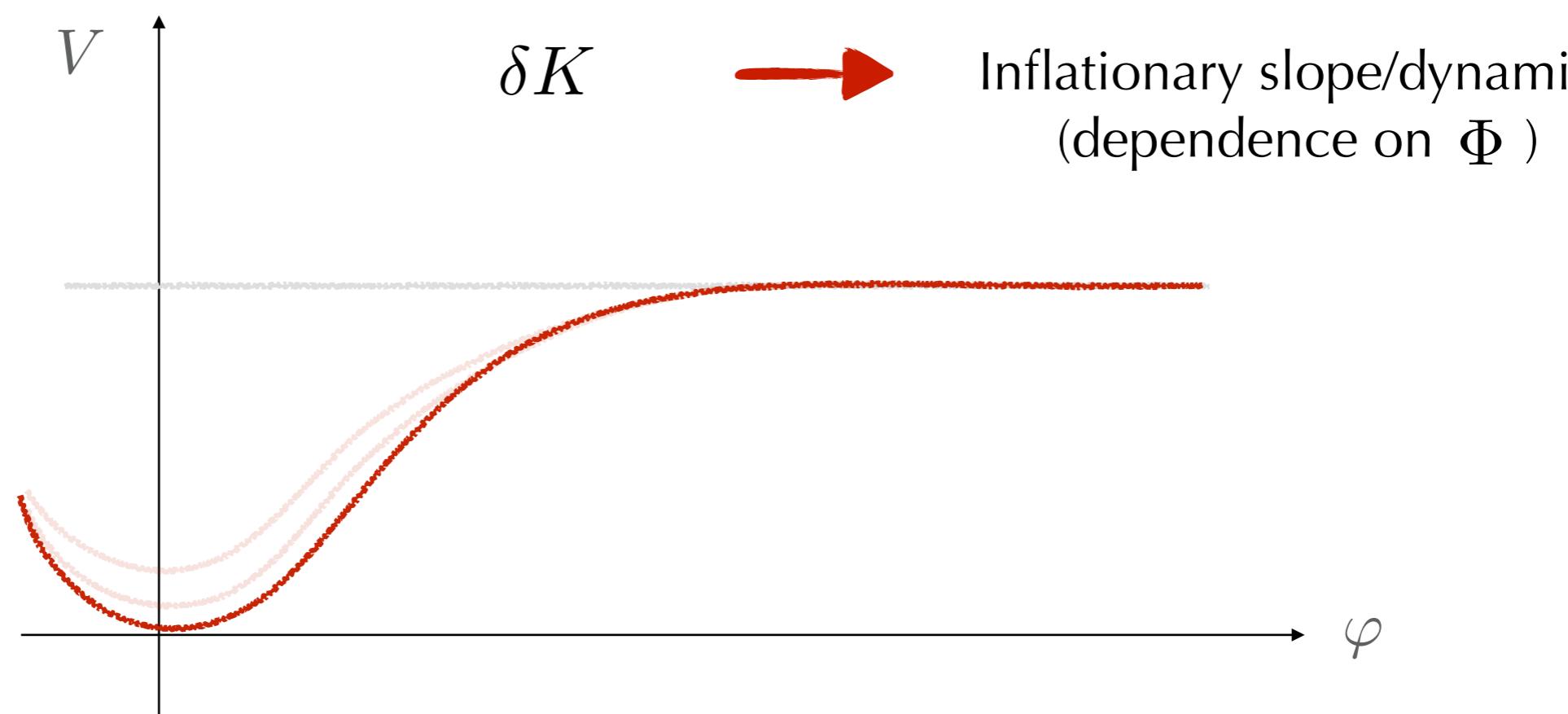


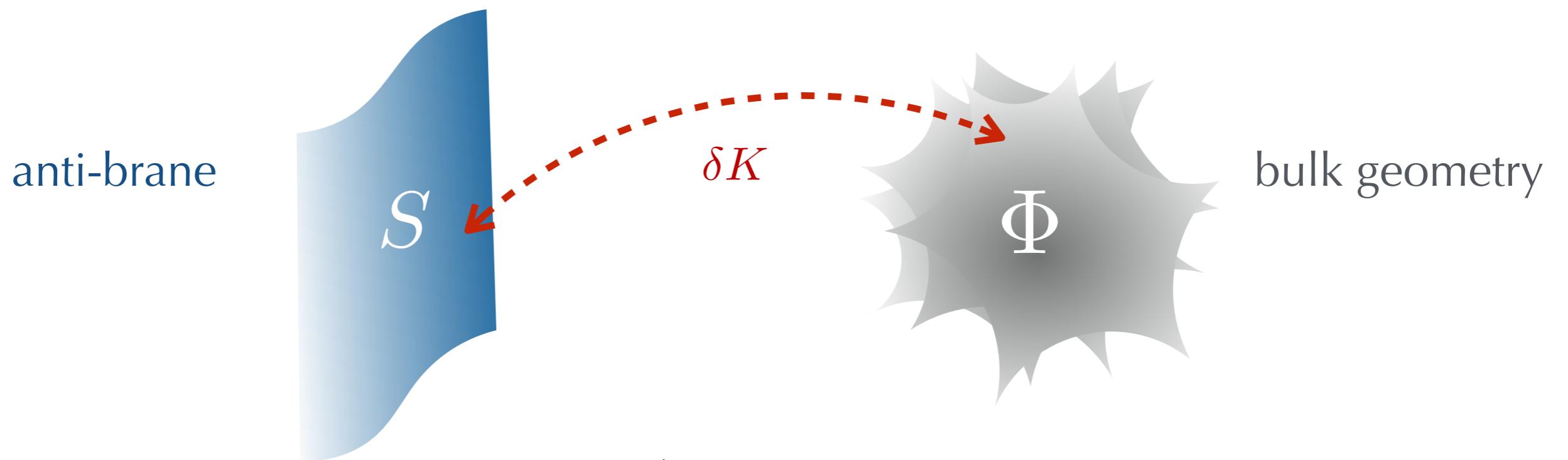


$$\begin{aligned} D_S W &\neq 0 \\ D_\Phi W &= 0 \end{aligned}$$



Hubble inflationary energy





$$\begin{aligned} D_S W &\neq 0 \\ D_\Phi W &= 0 \end{aligned}$$



Hubble inflationary energy

$$\delta K$$



Inflationary slope/dynamics  
(dependence on  $\Phi$ )

