

Primordial black hole constraints for extended mass functions

B. Carr, M. Raidal, T. Tenkanen, VV and H. Veermäe, arXiv:1705.05567.

M. Raidal, VV and H. Veermäe, arXiv:1707.01480.

by

Ville Vaskonen

Outline

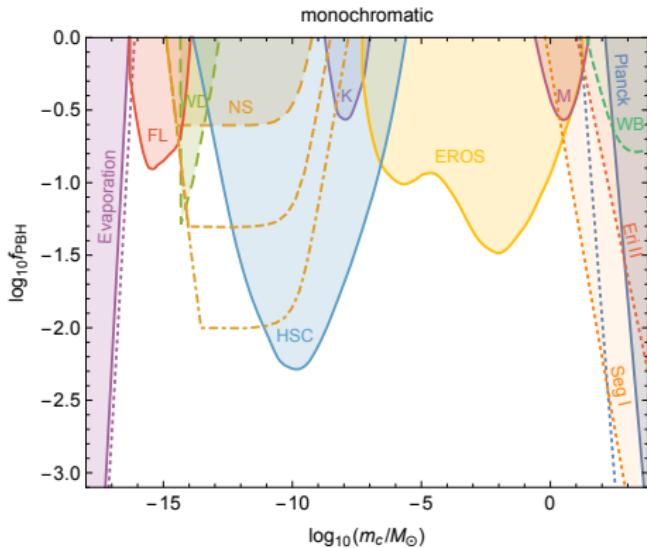
1. Extracting constraints for extended PBH mass functions
2. Constraint on PBH abundance from non-observation of the GW background by LIGO

1

Extracting constraints for
extended PBH mass functions

Introduction

- Monochromatic PBH mass function: $\psi(m) = f_{\text{PBH}} \delta(m - m_c)$,
 $f_{\text{PBH}} \equiv \Omega_{\text{PBH}} / \Omega_{\text{DM}}$.



Extended PBH mass function

- ▶ Typically PBHs are produced at various masses
 \Rightarrow extended mass function, $f_{\text{PBH}} = \int dm \psi(m)$.
- ▶ Lognormal mass function:

$$\psi(m) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma m} \exp\left(-\frac{\log^2(m/m_c)}{2\sigma^2}\right),$$

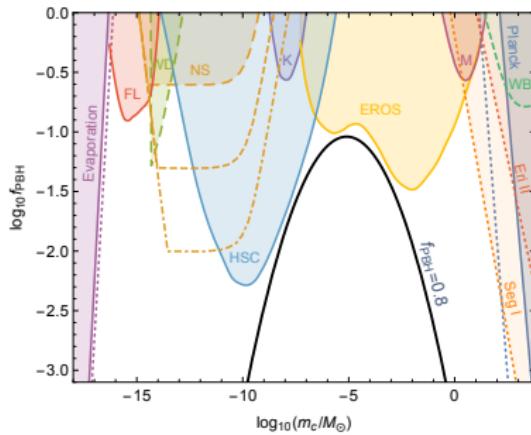
$\log m_c$ and σ^2 are the mean and variance of the $\log m$ distribution.

Extended PBH mass function

- ▶ Typically PBHs are produced at various masses
 \Rightarrow extended mass function, $f_{\text{PBH}} = \int dm \psi(m)$.
- ▶ Lognormal mass function:

$$\psi(m) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma m} \exp\left(-\frac{\log^2(m/m_c)}{2\sigma^2}\right),$$

$\log m_c$ and σ^2 are the mean and variance of the $\log m$ distribution.

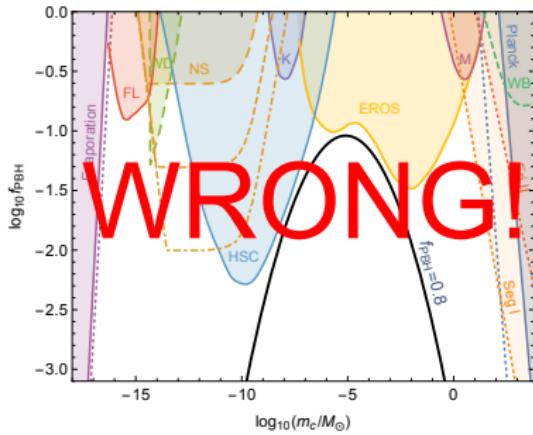


Extended PBH mass function

- ▶ Typically PBHs are produced at various masses
 \Rightarrow extended mass function, $f_{\text{PBH}} = \int dm \psi(m)$.
- ▶ Lognormal mass function:

$$\psi(m) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma m} \exp\left(-\frac{\log^2(m/m_c)}{2\sigma^2}\right),$$

$\log m_c$ and σ^2 are the mean and variance of the $\log m$ distribution.



Constraints for arbitrary ψ

- ▶ PBHs of different mass contribute independently to all observables A for which the constraints were shown in the first plot:

$$A = A_0 + \int dm \psi(m) K(m).$$

Constraints for arbitrary ψ

- ▶ PBHs of different mass contribute independently to all observables A for which the constraints were shown in the first plot:

$$A = A_0 + \int dm \psi(m) K(m).$$

- ▶ For a monochromatic mass function the constraint $A \leq A_{\text{exp}}$ can be written as

$$f_{\text{PBH}}(m) \leq \frac{A_{\text{exp}} - A_0}{K(m)} \equiv f_{\text{max}}(m),$$

Constraints for arbitrary ψ

- ▶ PBHs of different mass contribute independently to all observables A for which the constraints were shown in the first plot:

$$A = A_0 + \int dm \psi(m) K(m).$$

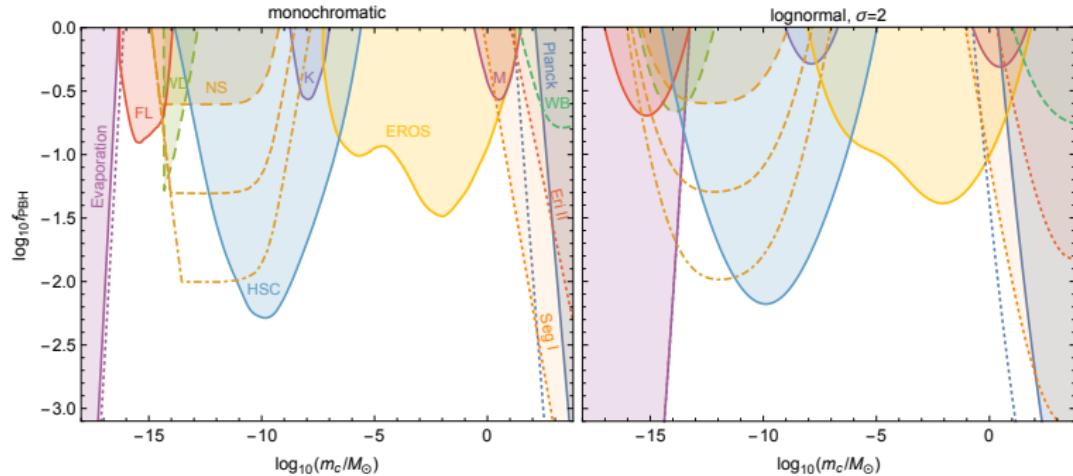
- ▶ For a monochromatic mass function the constraint $A \leq A_{\text{exp}}$ can be written as

$$f_{\text{PBH}}(m) \leq \frac{A_{\text{exp}} - A_0}{K(m)} \equiv f_{\text{max}}(m),$$

so for any mass function $A \leq A_{\text{exp}}$ implies

$$\int dm \frac{\psi(m)}{f_{\text{max}}(m)} \leq 1.$$

Constraints for lognormal ψ



- ▶ The wider the mass function is, the smaller is the maximal allowed PBH abundance.

2

Constraint on PBH abundance
from non-observation of the
GW background by LIGO

GW background

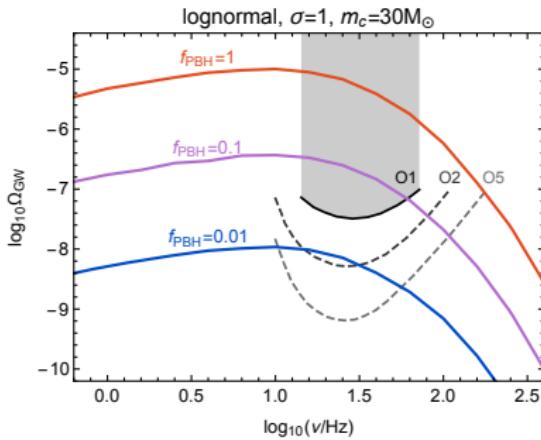
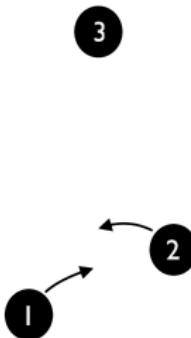
- ▶ The GWs from the unobservable events combine to create a stochastic GW background.
- ▶ Non-observation of the GW background constrains the PBH abundance.

GW background

- ▶ The GWs from the unobservable events combine to create a stochastic GW background.
- ▶ Non-observation of the GW background constrains the PBH abundance.
- ▶ PBH binaries are formed in the early Universe dominantly by three-body interactions.

Nakamura, Sasaki, Tanaka, Thorne, astro-ph/9708060.

Ali-Haïmoud, Kovetz, Kamionkowski, 1709.06576.



Constraints from LIGO

- ▶ The observable is of the form

$$A = \int dm_1 dm_2 dm_3 \psi(m_1) \psi(m_2) \psi(m_3) K(m_1, m_2, m_3) + \dots$$

⇒ the constraint must be calculated separately for each mass function.

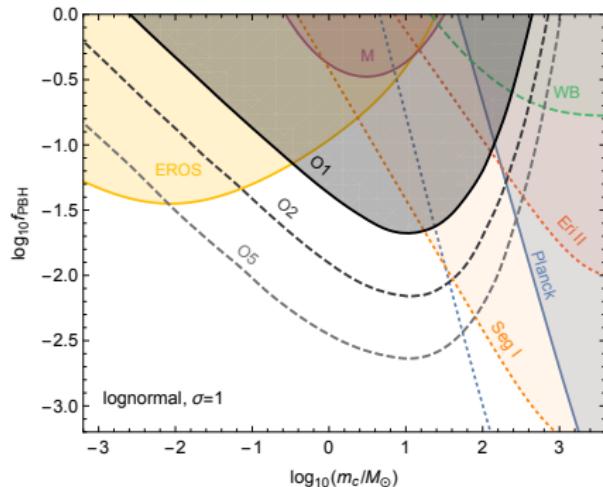
Constraints from LIGO

- The observable is of the form

$$A = \int dm_1 dm_2 dm_3 \psi(m_1) \psi(m_2) \psi(m_3) K(m_1, m_2, m_3) + \dots$$

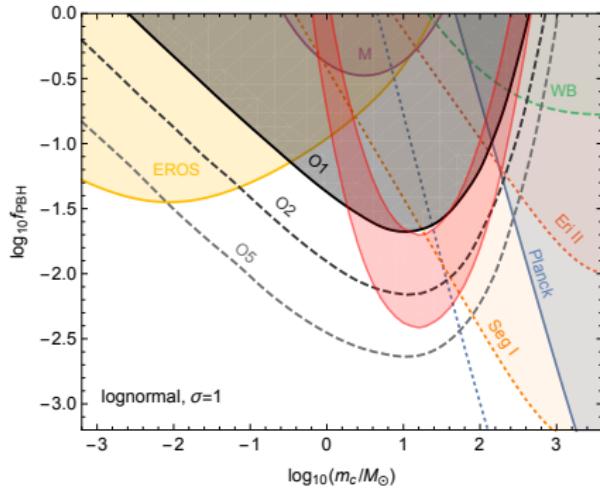
⇒ the constraint must be calculated separately for each mass function.

- LIGO gives the strongest constraint for $m_c \sim 1 - 10 M_\odot$:



LIGO observations

- PBH mergers are a viable explanation for the BH merger events observed by LIGO if the mass function is sufficiently narrow and peaked roughly around $30M_{\odot}$.



Summary

1. Each constraint on PBH abundance can be expressed in the form

$$\int dm \frac{\psi(m)}{f_{\max}(m)} \leq 1,$$

where $f_{\max}(m)$ is the constraint for monochromatic mass function.

2. The wider the mass function is, the smaller is the maximal allowed PBH abundance.
3. The non-observation on the GW background from PBH mergers by LIGO gives the dominant constraint for $m_c \sim 1 - 10M_\odot$.
4. The fraction of DM in PBHs needed to explain the observed BH merger events is $f_{\text{PBH}} \sim 1\%$.

Lognormal mass function

For the lognormal mass function

$$\psi(m) = \frac{f_{\text{PBH}}}{\sqrt{2\pi}\sigma m} \exp\left(-\frac{\log^2(m/m_c)}{2\sigma^2}\right)$$

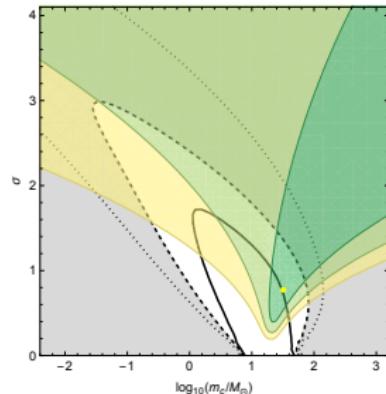
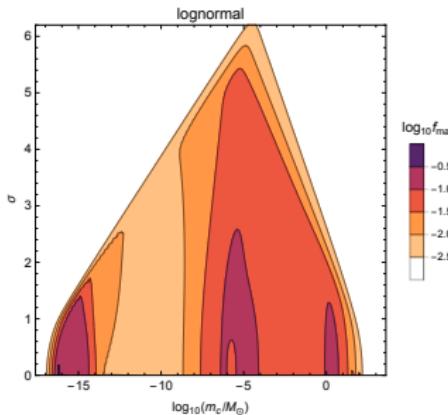
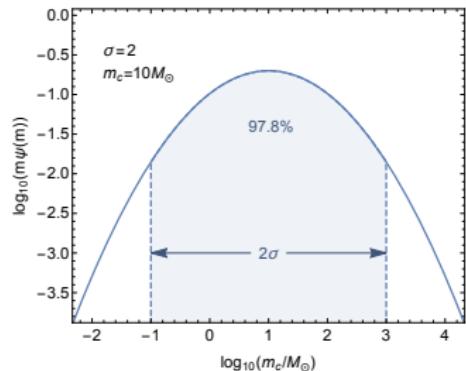
$\log m_c$ and σ^2 are the mean and variance of the $\log m$ distribution:

$$\log m_c \equiv \langle \log m \rangle_\psi, \quad \sigma^2 \equiv \langle \log^2 m \rangle_\psi - \langle \log m \rangle_\psi^2,$$

where

$$\langle X \rangle_\psi \equiv \frac{1}{f_{\text{PBH}}} \int dm \psi(m) X(m).$$

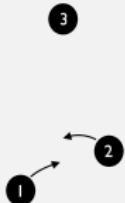
More figures



PBH binary formation

Early Universe

$$\frac{R_3(t_0)}{\text{Gpc}^{-3}\text{yr}^{-1}} \approx 5.1 \times 10^4 \delta_{\text{dc}}^{\frac{16}{37}} f_{\text{PBH}}^{\frac{53}{37}} \left(\frac{m_c}{30M_\odot} \right)^{-\frac{32}{37}}$$



Late Universe

$$\frac{R_2(t_0)}{\text{Gpc}^{-3}\text{yr}^{-1}} \approx 3.7 \times 10^{-7} \delta_0 f_{\text{PBH}}^2 \left(\frac{v_{\text{PBH}}}{10 \text{ km/s}} \right)^{-\frac{11}{7}}$$



$\implies R_3$ dominates over R_2 unless $\delta_0 \gtrsim 10^{11}$.