

# Dealing with axion monodromy and other oscillating inflationary models

Deanna C. Hooper

*Based on DH and Lesgourges (in prep)*

**DESY THEORY WORKSHOP**  
**September 2017**

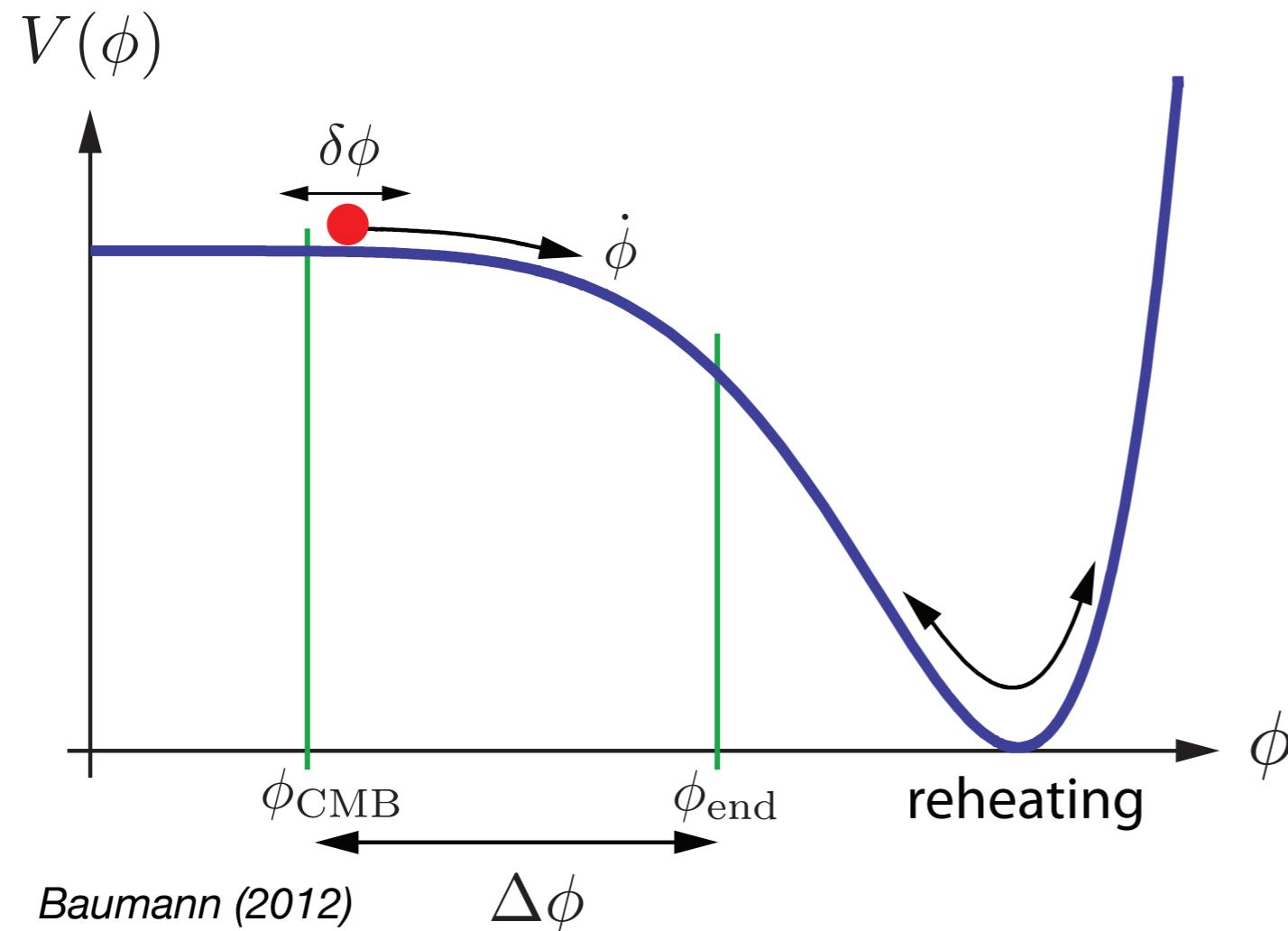


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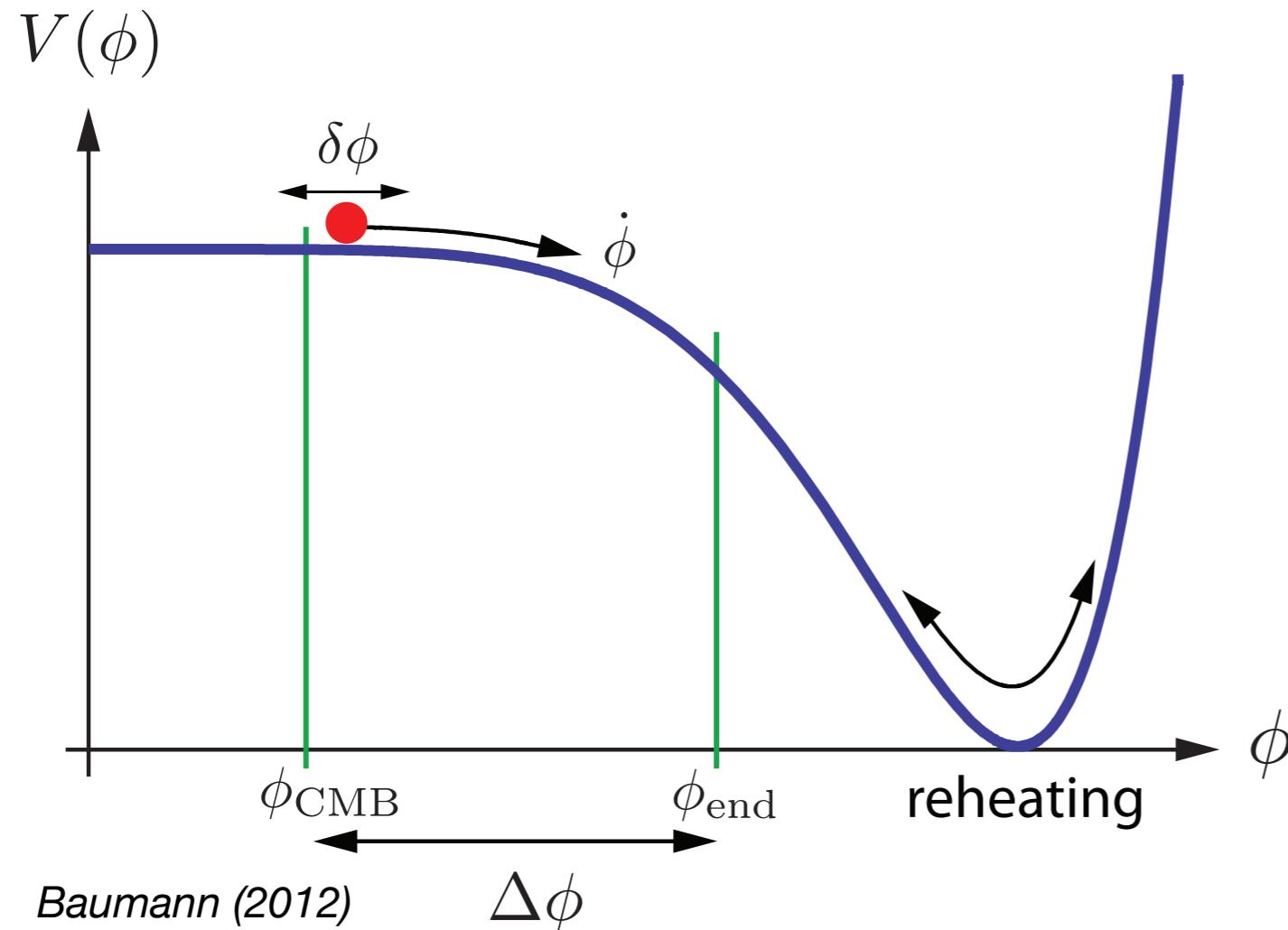
# Overview

- Backdrop
  - *Why is this relevant?*
- Axion Monodromy Inflation
  - *Why this model?*
  - *What has been done?*
  - *What are we doing?*
- Comparisons
  - *What did we find?*
  - *What next?*
- Summary
  - *Questions?*

# Slow Roll Inflation



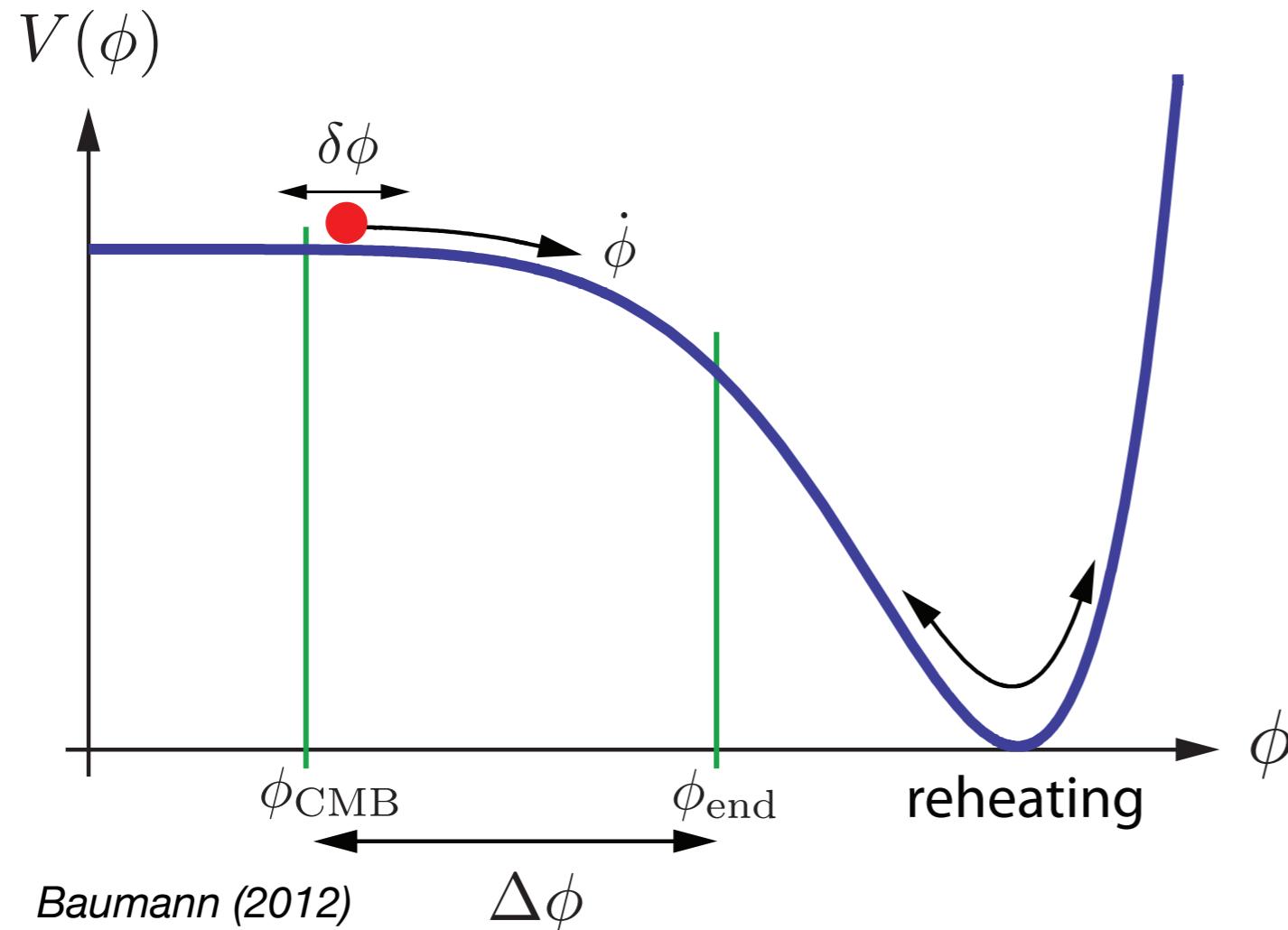
# Slow Roll Inflation



$$\epsilon_1 = \frac{M_{\text{Pl}}^2}{2} \left( \frac{V'}{V} \right)^2$$

$$\epsilon_2 = 2M_{\text{Pl}}^2 \left[ \left( \frac{V'}{V} \right)^2 - \frac{V''}{V} \right]$$

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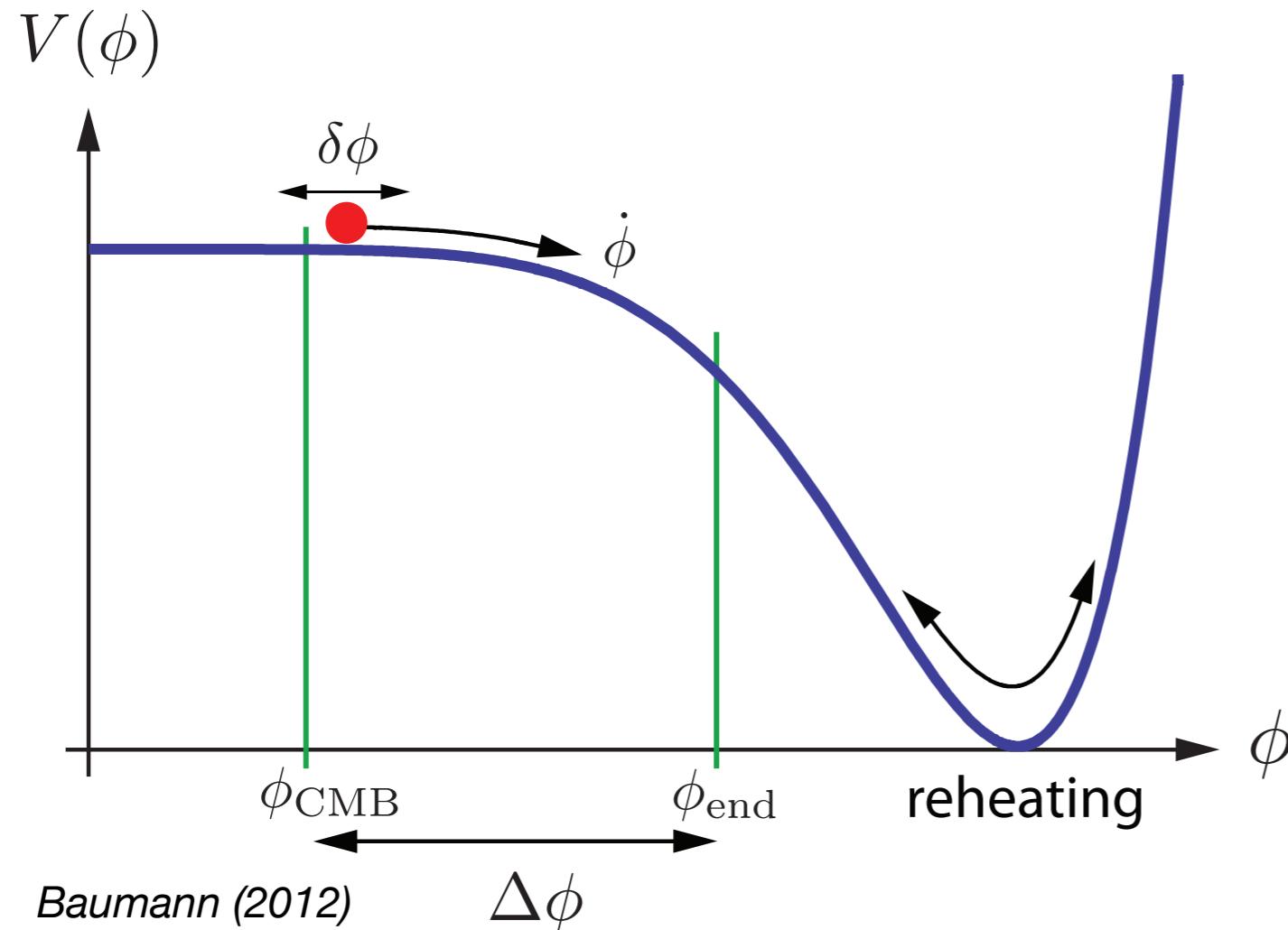
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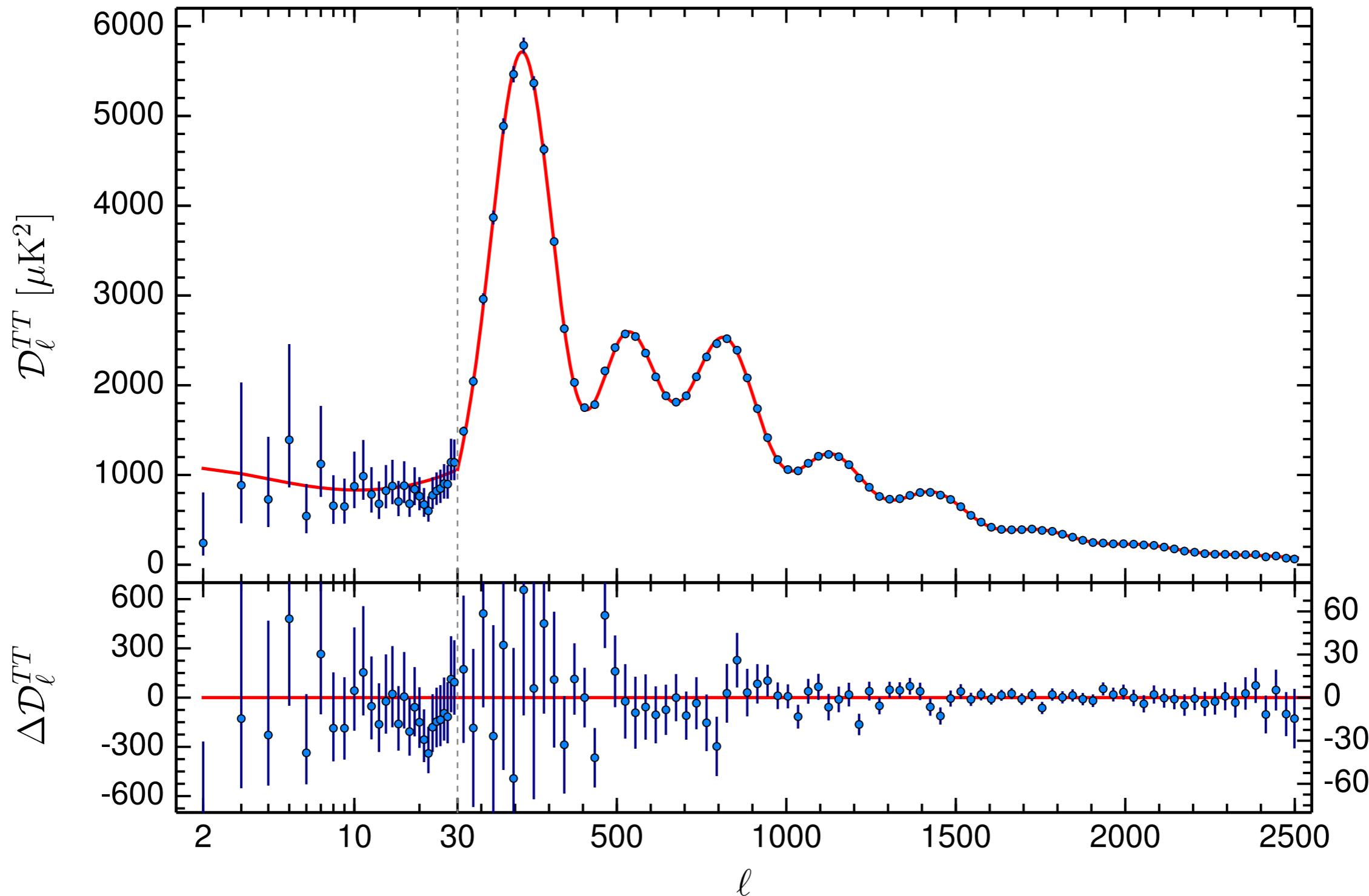
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$$\mathcal{P}_{\mathcal{R}}(k) = \left( \frac{V}{24\pi^2 \epsilon_1 M_{\text{Pl}}^4} \right) \cdot \left( 1 - 2(C+1)\epsilon_1 - C\epsilon_2 + \mathcal{O}(\epsilon_i^2) \right)$$

# CMB Power Spectra



Planck Collaboration (2015)

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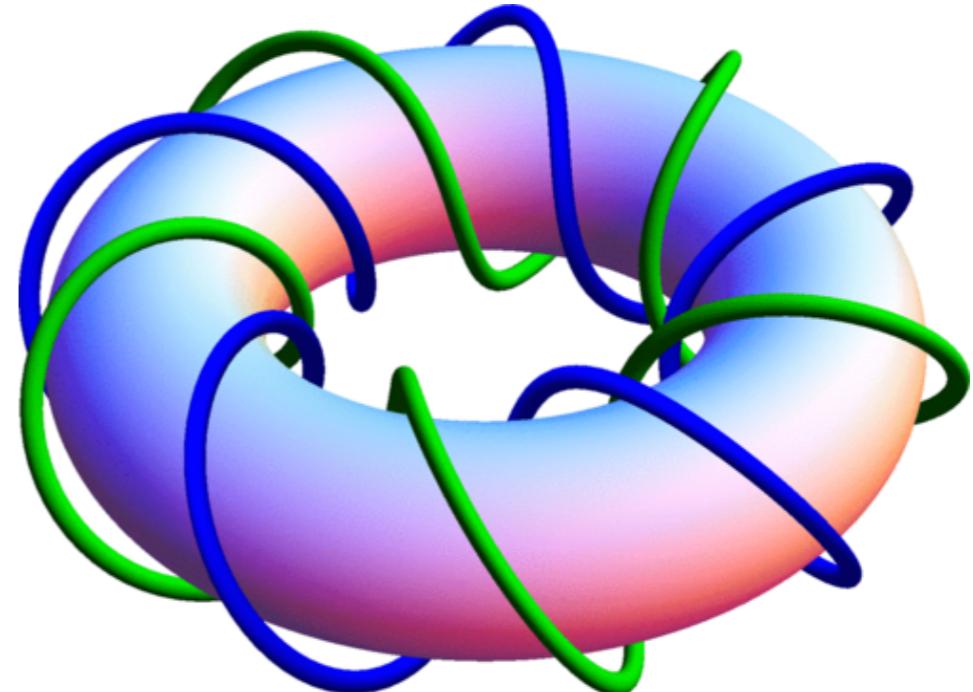
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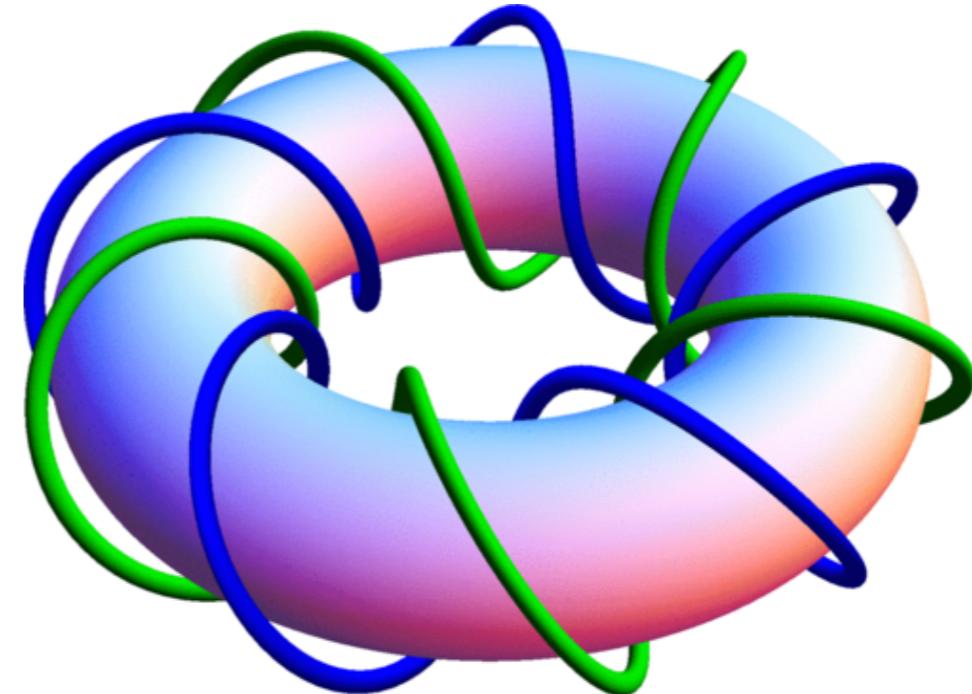
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- Energy scale of inflation unconstrained, could be GUT scale
- Discrete shift symmetry leads to sub-Planckian periodicity
- Monodromy explains slowly varying  $\phi$
- Axion field is wrapped around a higher-dimension brane
- AMI would produce clear primordial gravitational waves, can be constrained with tensor modes



# AMI Potential

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$$V(\phi) = \mu^{4-p} \phi^p + \Lambda_0^4 e^{-C_0 \left( \frac{\phi}{\phi_0} \right)^{p_\Lambda}} \cos \left( \gamma_0 + \frac{\phi_0}{f} \left( \frac{\phi}{\phi_0} \right)^{p_f+1} \right)$$

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*Flauger et al. (2009)*

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$$\mathcal{P}_{\mathcal{R}}(k) = \mathcal{P}_{\mathcal{R}}(k_\star) \left( \frac{k}{k_\star} \right)^{n_s-1} \left\{ 1 + \delta n_s \cos \left[ \frac{\phi_0}{f'} \left( \frac{\phi_k}{\phi_0} \right)^{p_f+1} + \Delta\phi \right] \right\}$$

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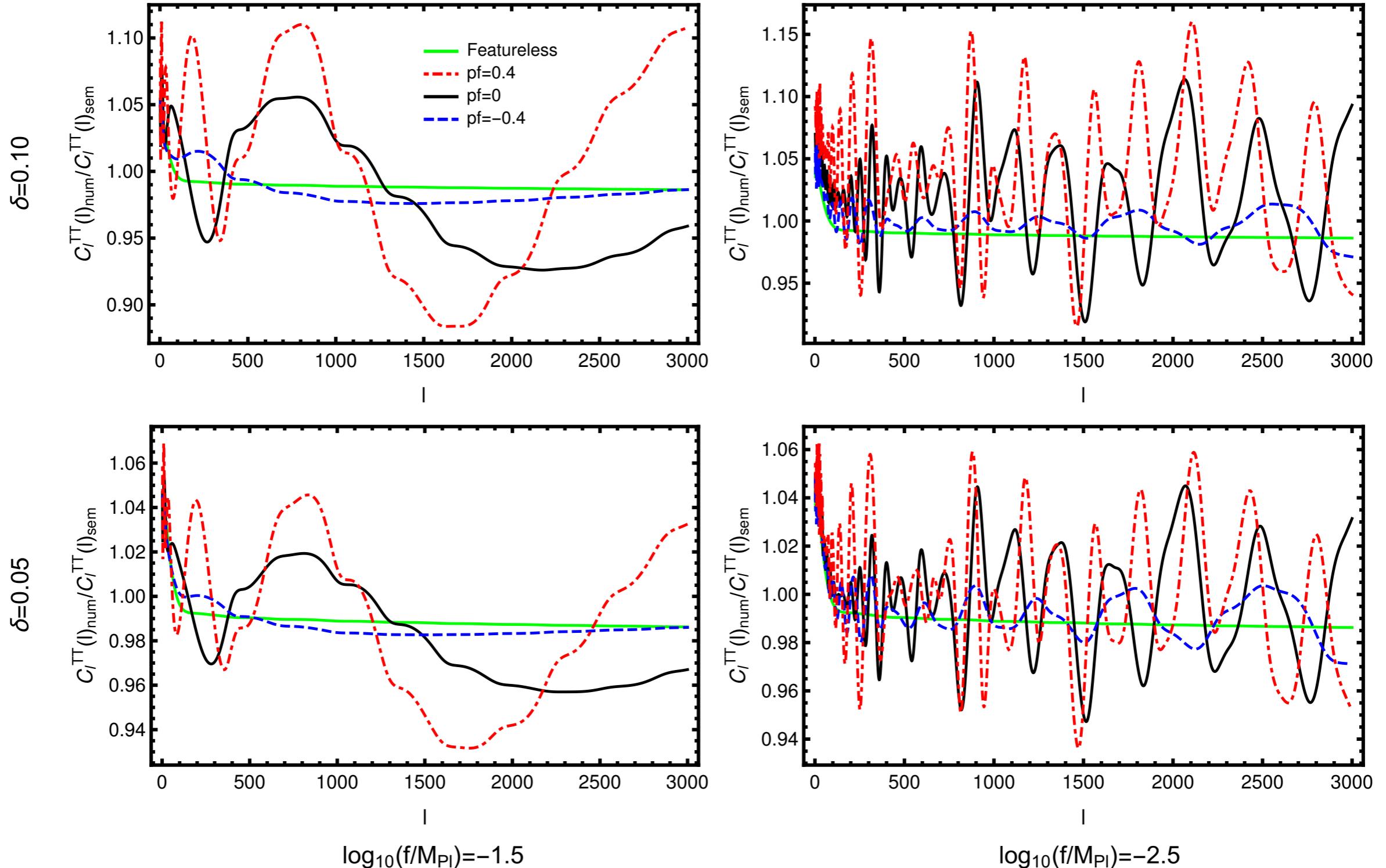
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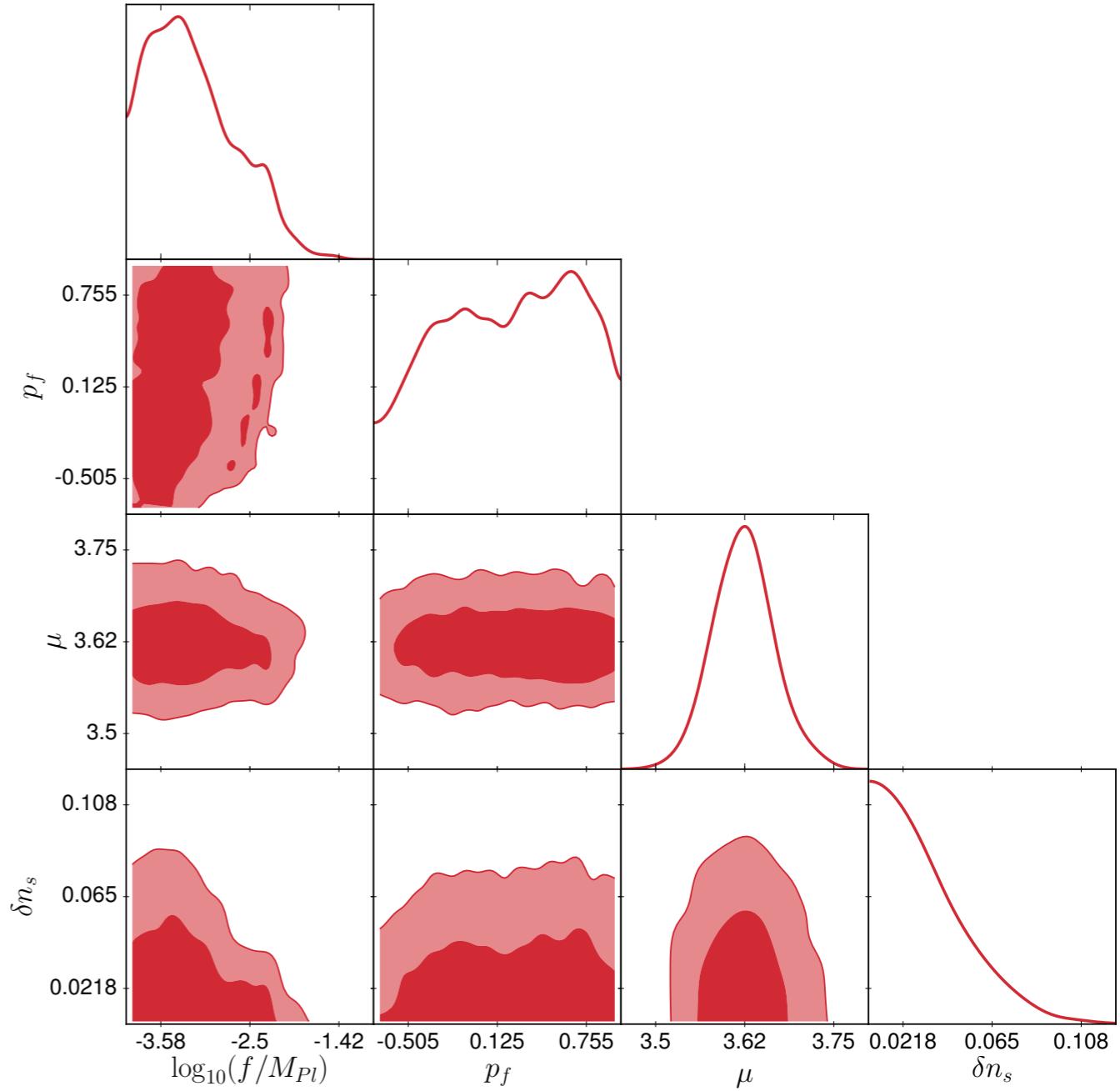
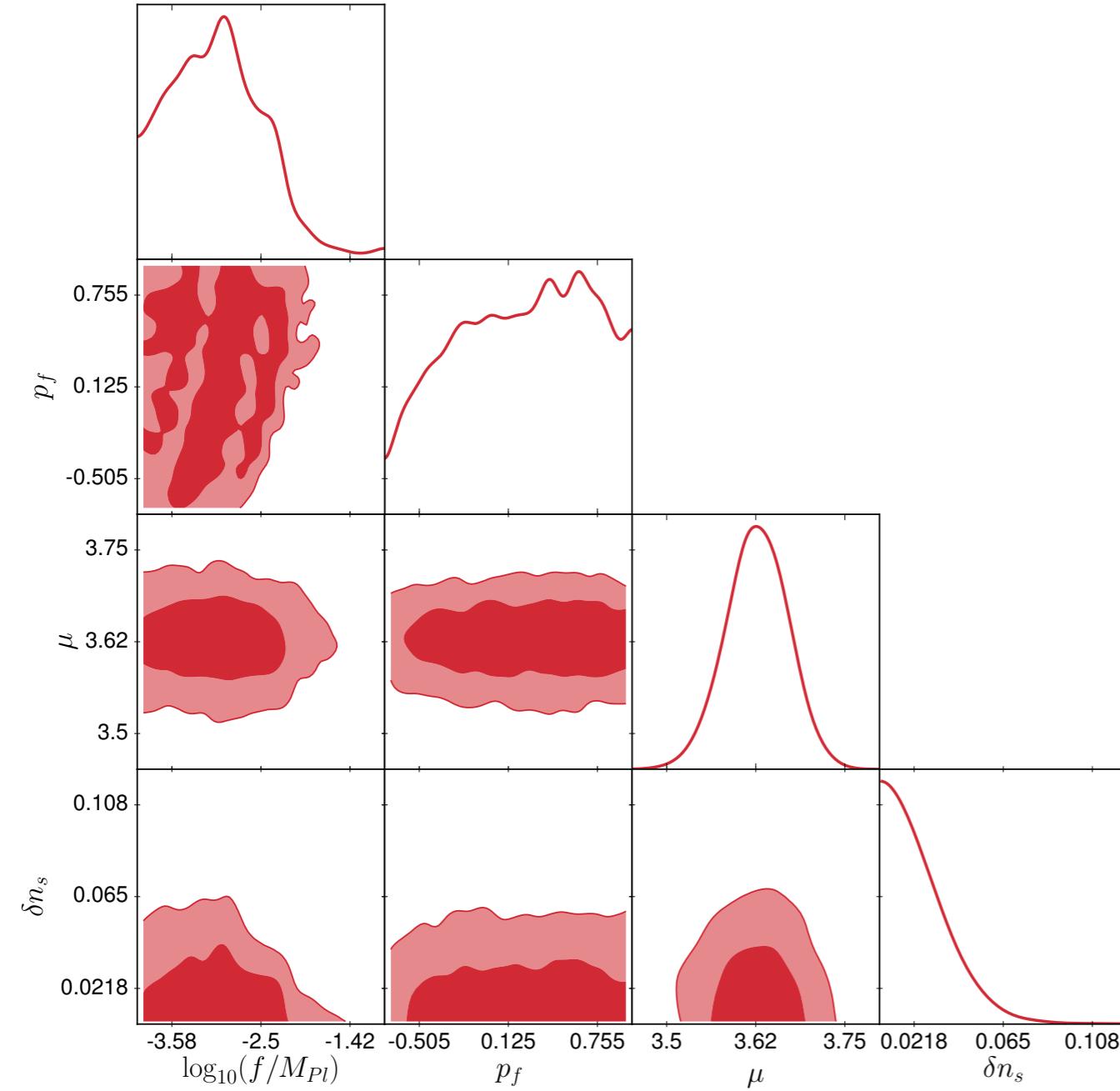
$$\mathcal{P}_{\mathcal{R}}(k) = \mathcal{P}_{\mathcal{R}}(k_\star) \left( \frac{k}{k_\star} \right)^{n_s-1} \left\{ 1 + \delta n_s \cos \left[ \frac{\phi_0}{f'} \left( \frac{\phi_k}{\phi_0} \right)^{p_f+1} + \Delta\phi \right] \right\}$$

$$\delta n_s = 3b \left( \frac{2\pi}{\alpha} \right)^{1/2} \rightarrow \begin{cases} \alpha = (1 + p_f) \frac{\phi_0}{2fN_0} \left( \frac{\sqrt{2pN_0}}{\phi_0} \right)^{1+p_f} \\ b = \frac{\Lambda_0^4}{p\mu^{4-p}\phi^{p-1}f} (1 + p_f) \left( \frac{\phi}{\phi_0} \right)^{p_f} \end{cases}$$

# First comparisons



# SAT fits

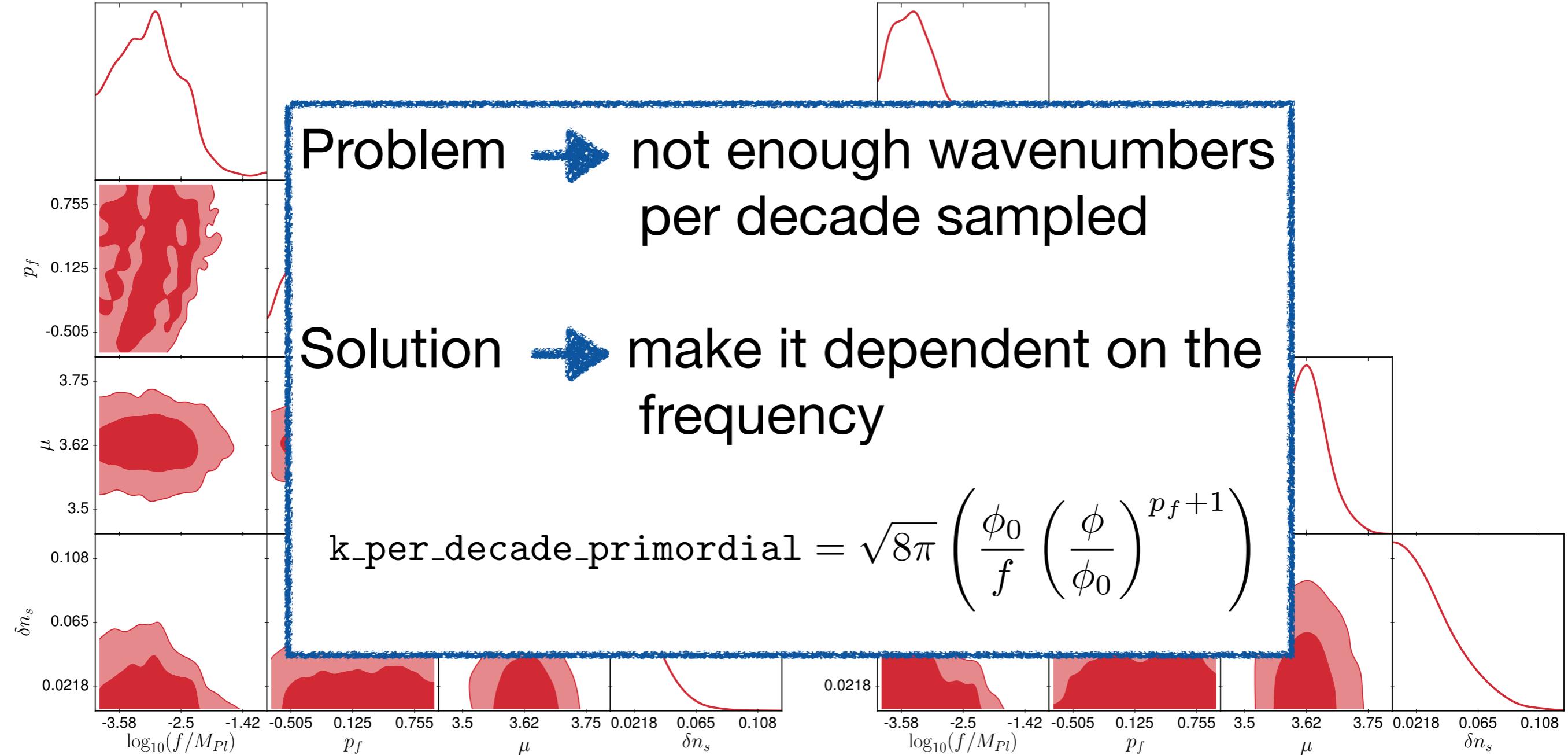


# SAT fits

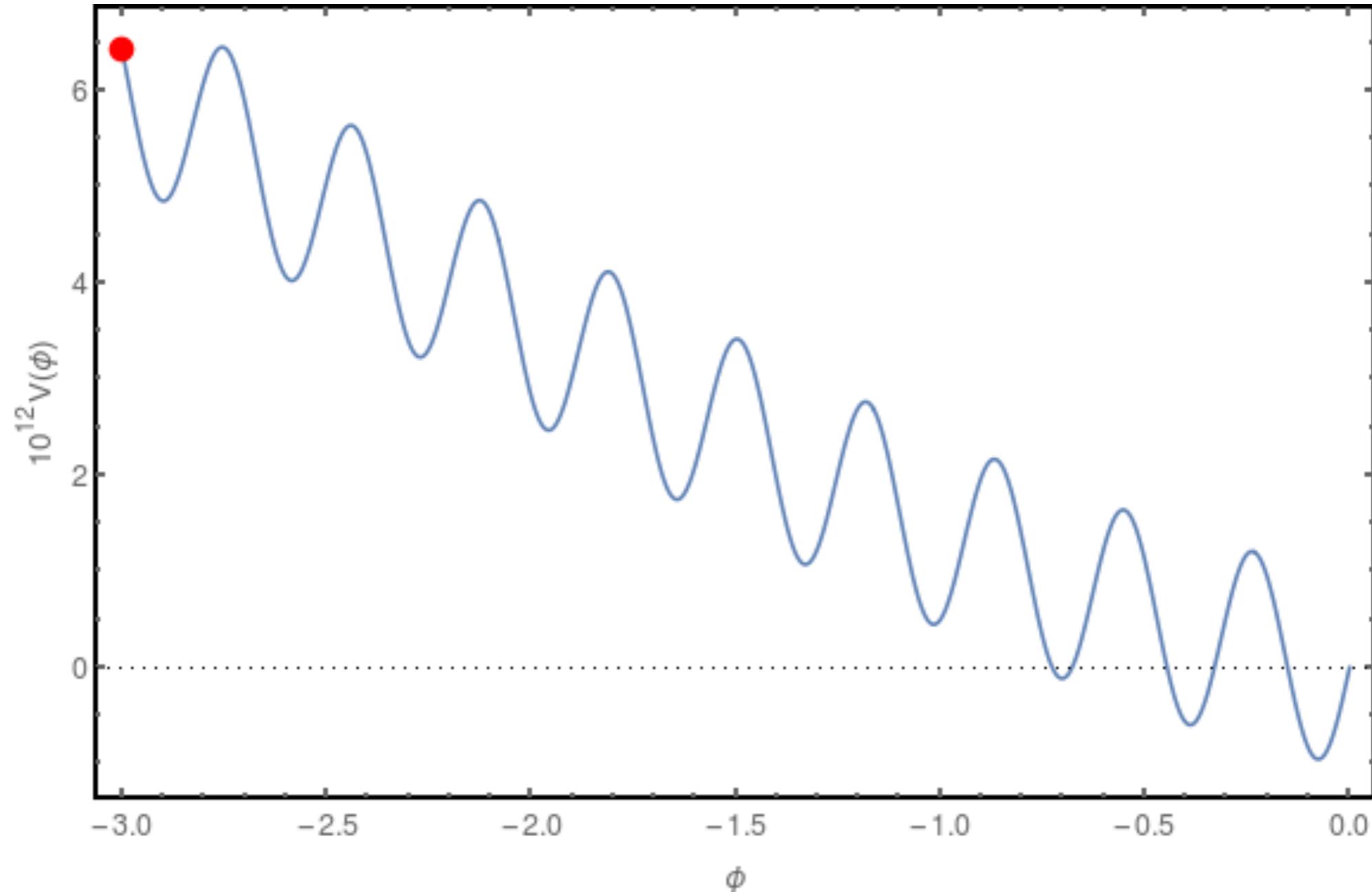
Problem → not enough wavenumbers per decade sampled

Solution → make it dependent on the frequency

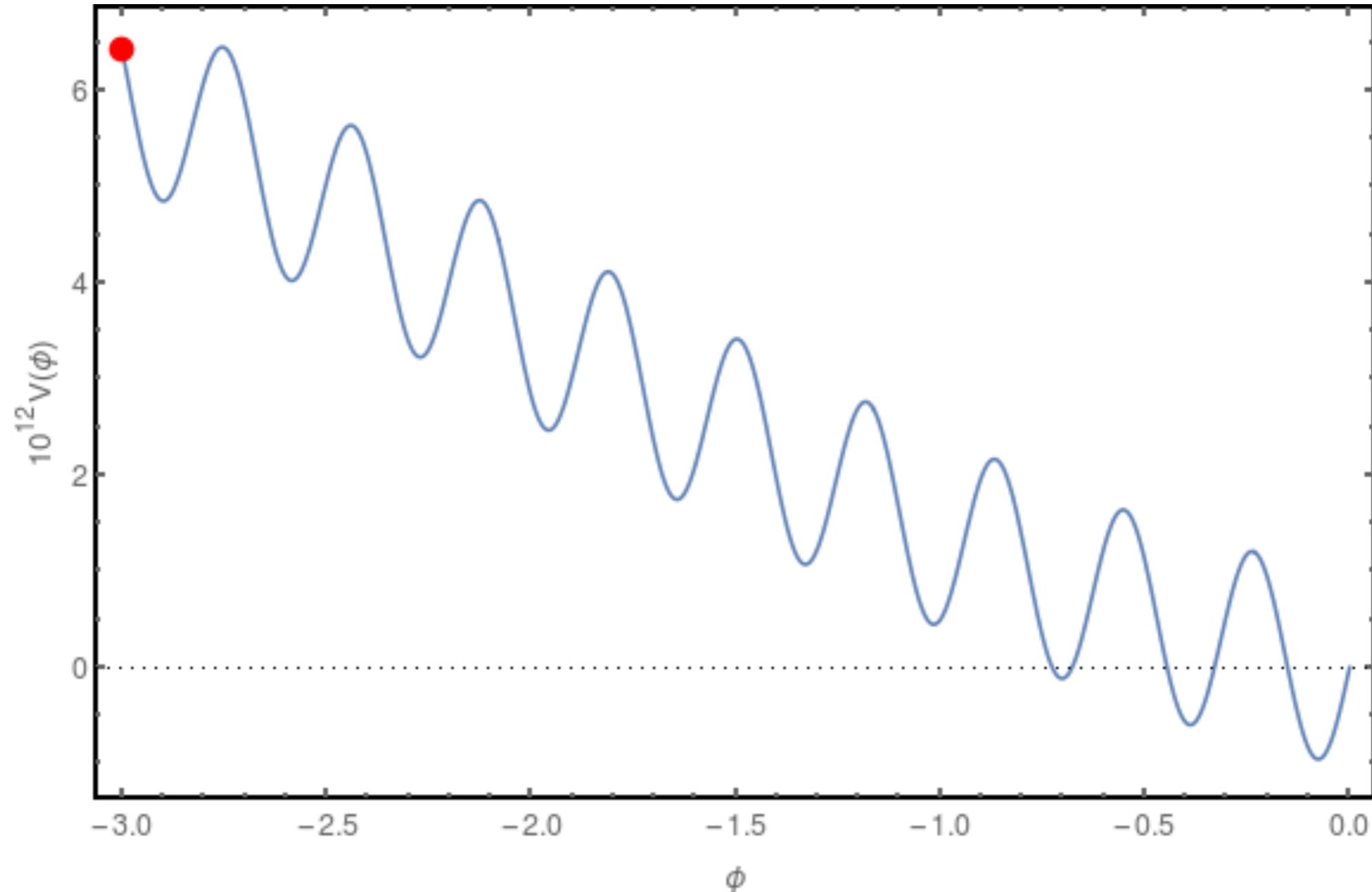
$$k_{\text{per\_decade\_primordial}} = \sqrt{8\pi} \left( \frac{\phi_0}{f} \left( \frac{\phi}{\phi_0} \right)^{p_f + 1} \right)$$



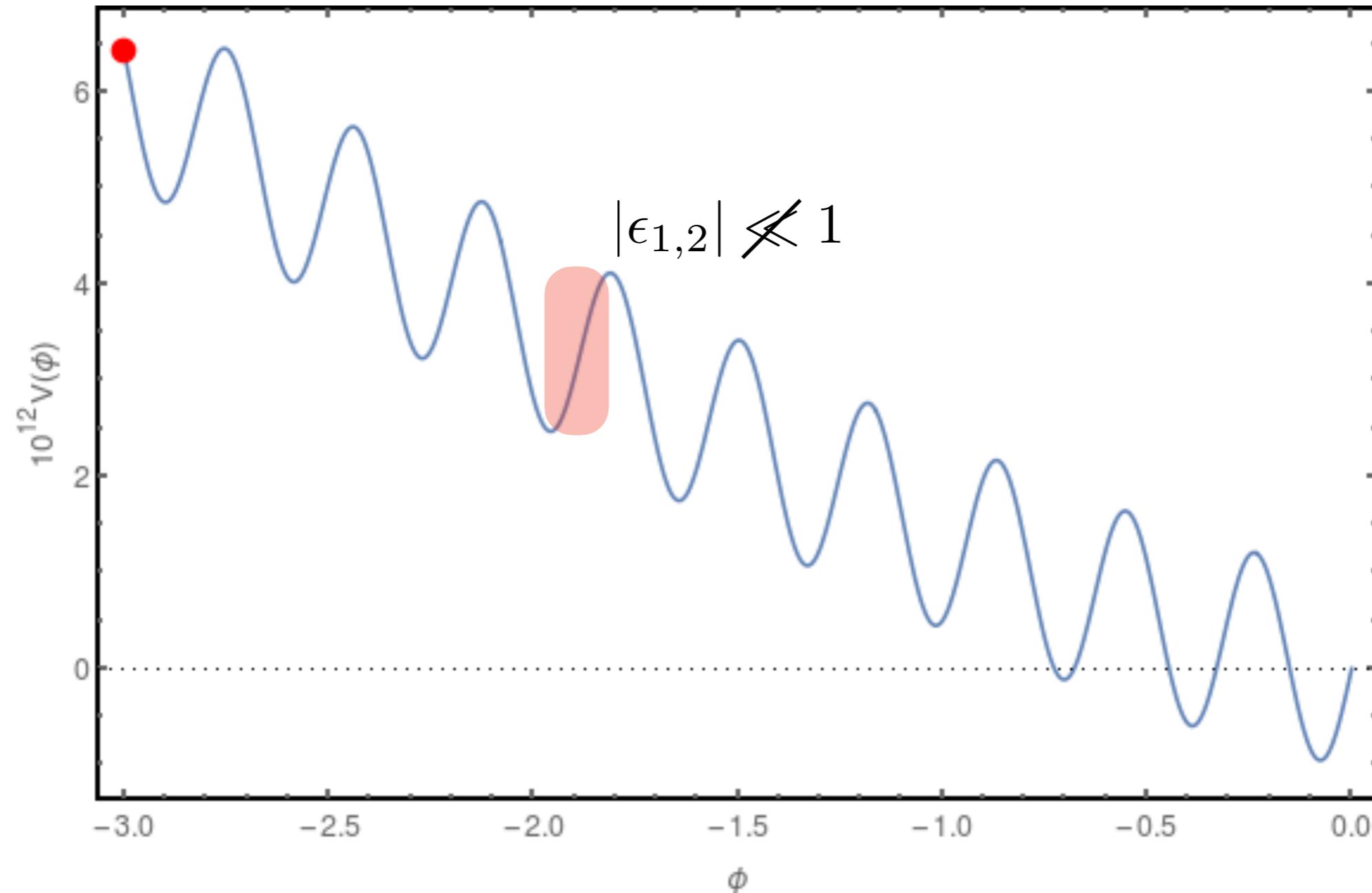
# Dealing with non-monotonic potentials



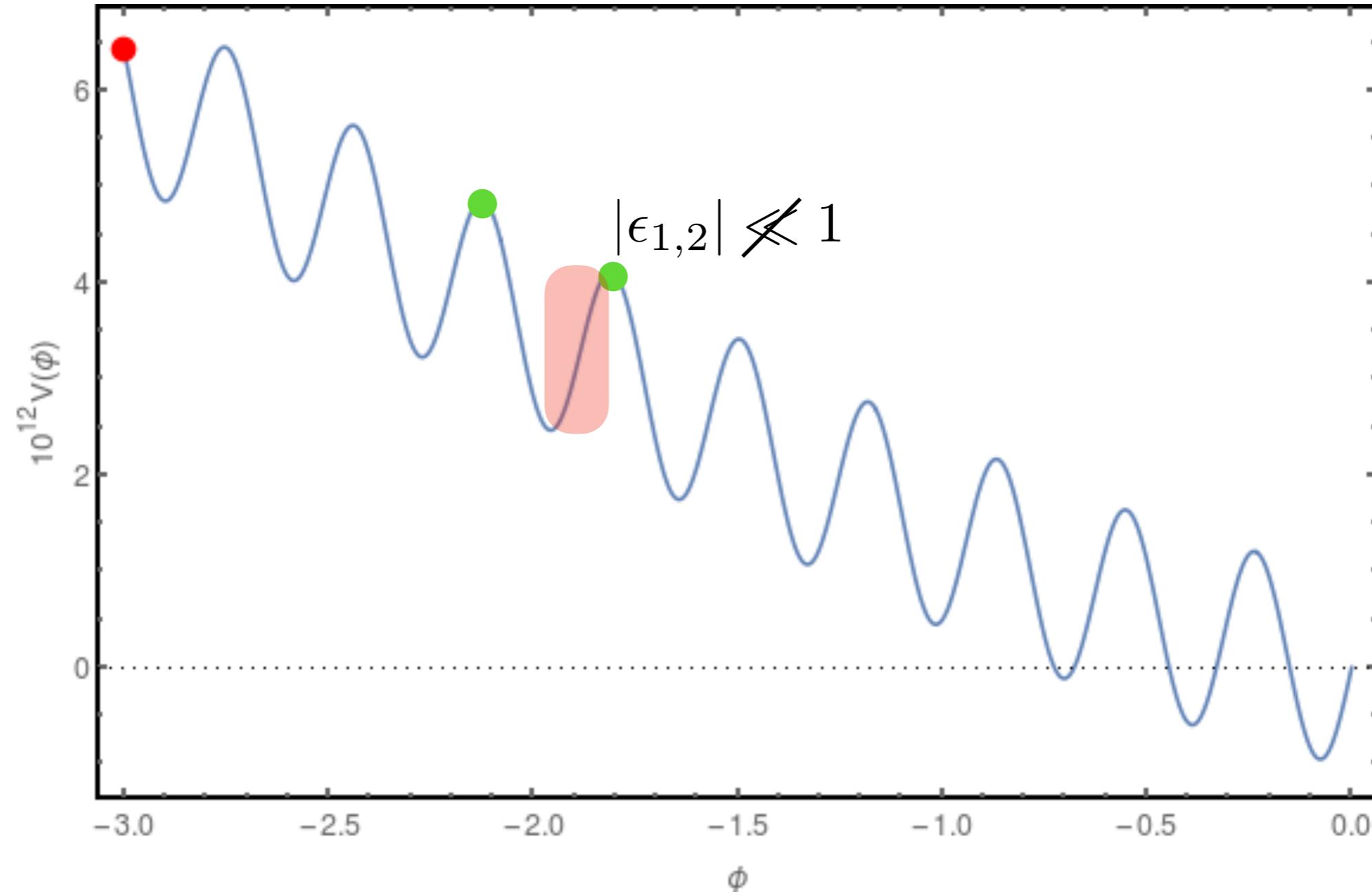
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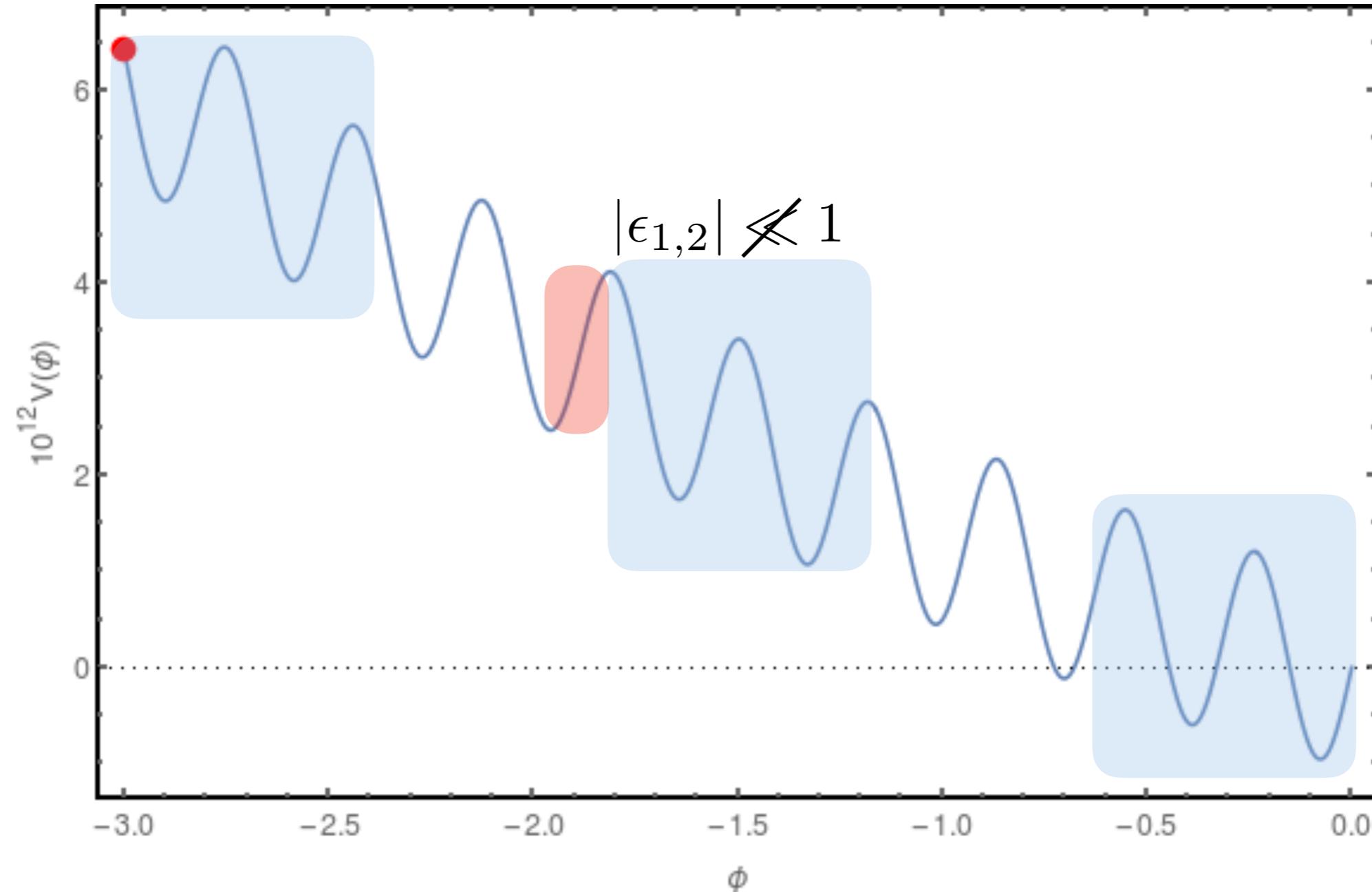
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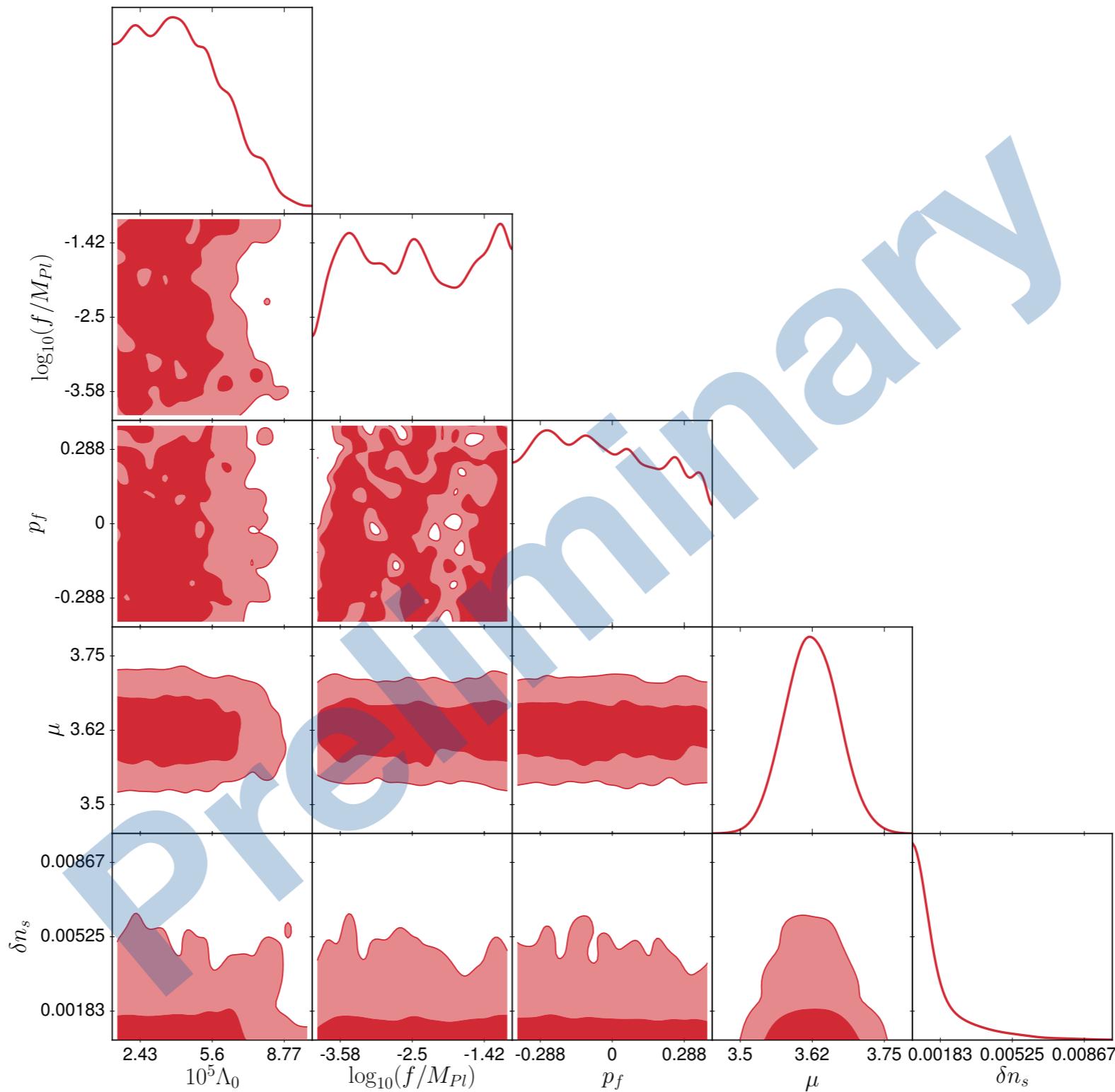
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# MultiNest Fits



# Summary

- The search for oscillating inflaton fields is motivated by high energy models
- These models can violate the SR conditions, and require special treatment
- We propose a method to treat oscillating potentials numerically
- Tested on the string-theory motivated AMI
- Strong framework in place to constrain these models with Planck data (and BKP)

**Thank you for your attention**