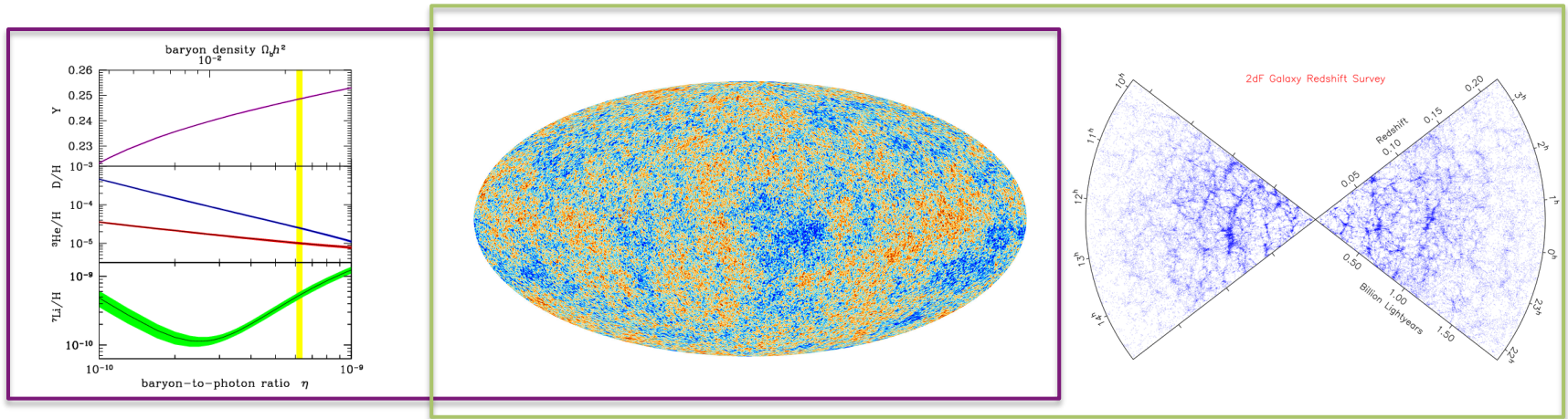


Sterile neutrinos with secret interactions

Maria Archidiacono
RWTH Aachen University

DESY Theory Workshop
26 – 29 September 2017
Fundamental physics in the cosmos:
The early, the large and the dark Universe

The case for sterile neutrinos



N_{eff} & flavour

N_{eff} & (Σm_ν)

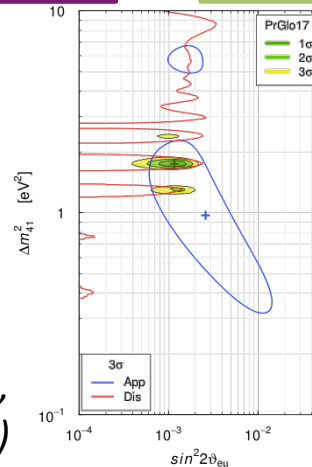
(N_{eff}) & Σm_ν

$$N_{\text{eff}} = 2.99 \pm 0.20 \text{ (68\%cl)}$$

$$\Sigma m_\nu < 0.13 \text{ eV (95\%cl)}$$

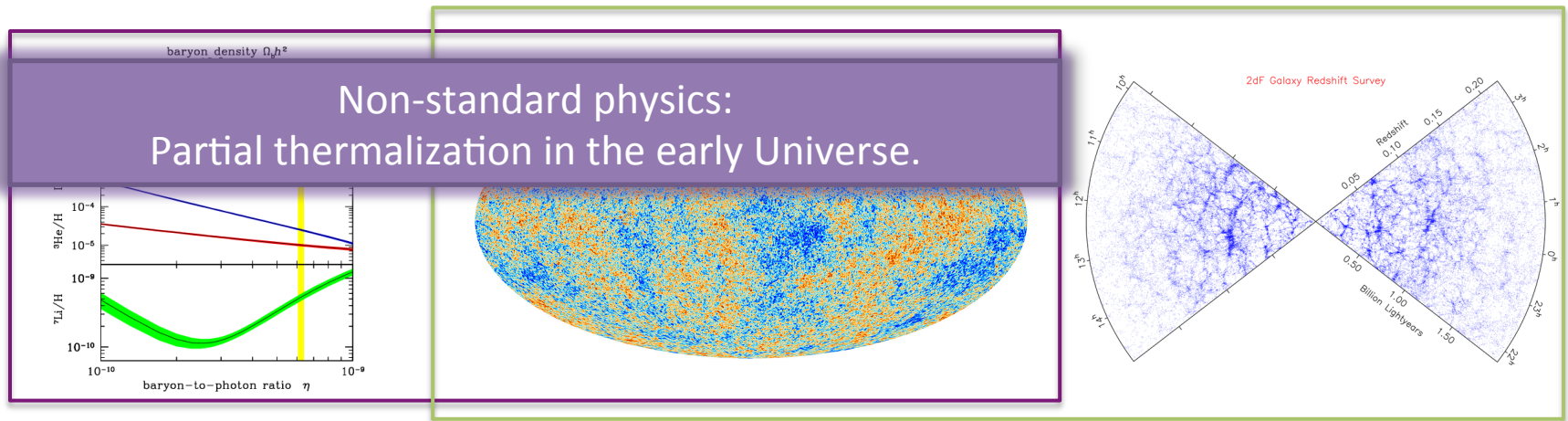
*Cuesta et al.,
Phys. Dark Univ (2016)*

*Gariazzo et al.,
Global fit (2017)*



$$N_{\text{eff}} = 4.0... \\ \Sigma m_\nu \approx 1 \text{ eV}$$

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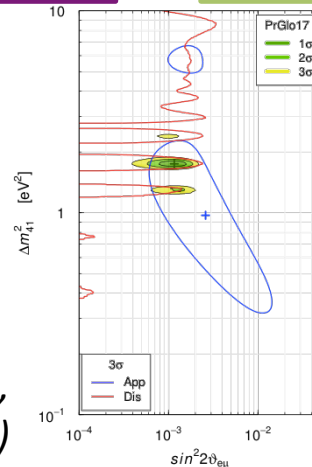
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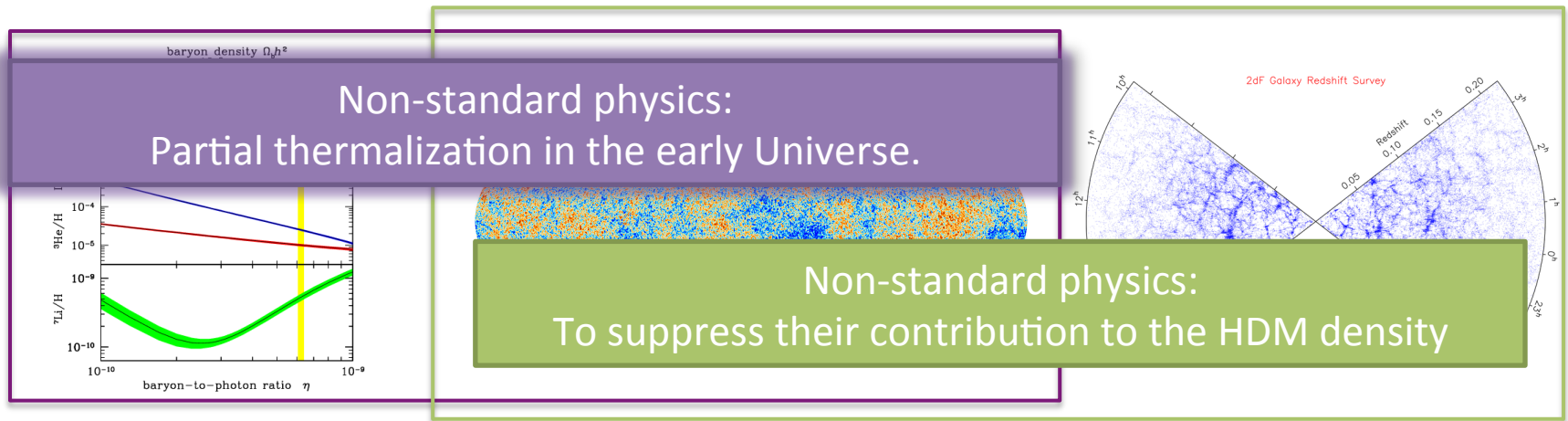
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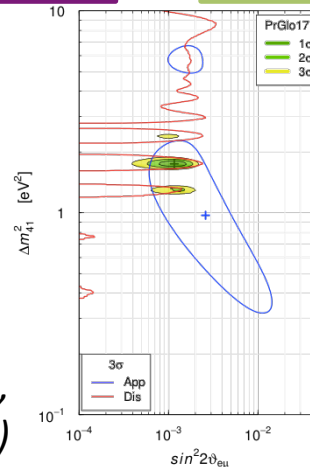
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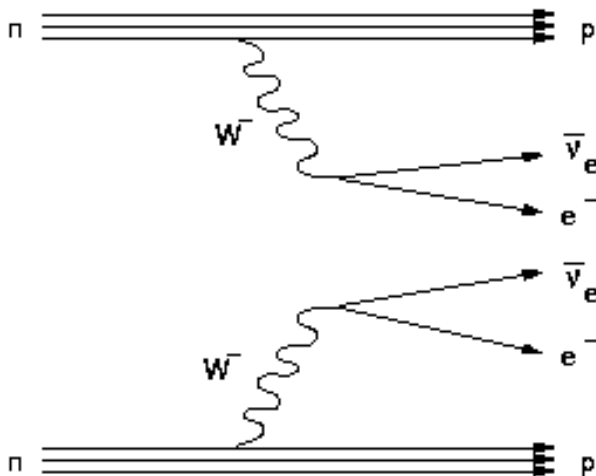
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ν_s secret interactions

The sterile neutrino is coupled to a new light pseudoscalar (Majoron models)

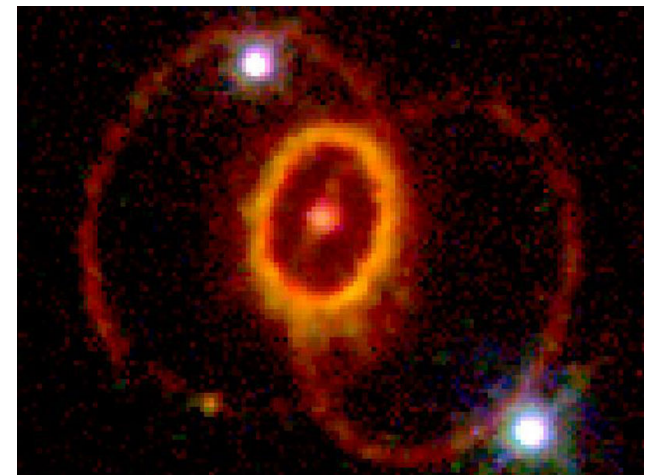
$$L_{\text{int}} \sim g_s \phi \bar{\nu}_s \gamma_5 \nu_s$$

Non-cosmological constraints:



$$\beta\beta g_e < 3 \times 10^{-5}$$

Bernatowicz et al, PRL (1992)



$$\text{SN energy loss } g_e < 4 \times 10^{-7}$$

$$g_s < 10^{-5}$$

Farzan, PRD (2003)

Flavour evolution

$$\rho(p,t) = \begin{pmatrix} \rho_{aa} & \rho_{as} \\ \rho_{sa} & \rho_{ss} \end{pmatrix} = \frac{f_0(p)}{2} [P_0(p,t) + \bar{\sigma} \times \bar{P}(p,t)];$$

$$\frac{d\bar{P}}{dt} = \bar{V} \times \bar{P} - D\bar{P}_T + \frac{R}{f_0} \hat{z}$$

$$\bar{V} = \bar{V}_{vacuum} + \bar{V}_{medium} + \bar{V}_s$$

$$V_{vacuum} = \frac{\Delta m^2}{2p}$$

$$V_{medium} \propto \frac{G_F}{M_z^2} n_a p T^4$$

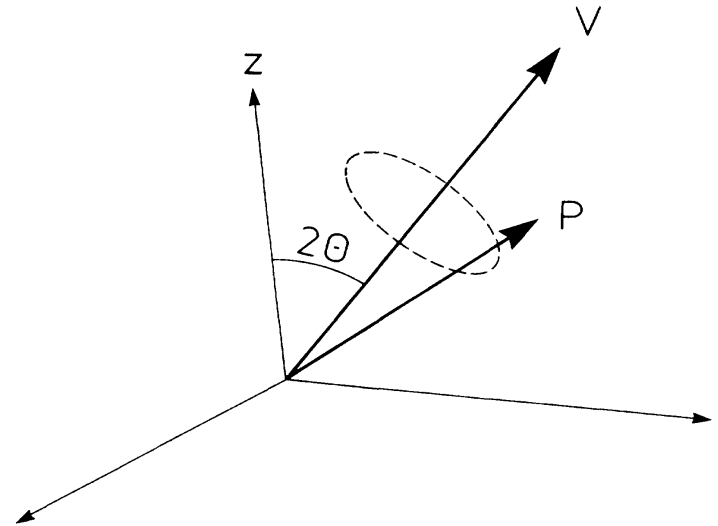
$$D = \frac{1}{2} \Gamma \quad \text{damping}$$

$$R = \Gamma \left(f_0 - \frac{f_0}{2} (P_0 + P_z) \right) \quad \text{repopulation}$$

$$\Gamma = \Gamma_a + \Gamma_s$$

$$\Gamma_a \propto G_F^2 p T^4$$

Stodolsky, PRD (1987)



Early time phenomenology:

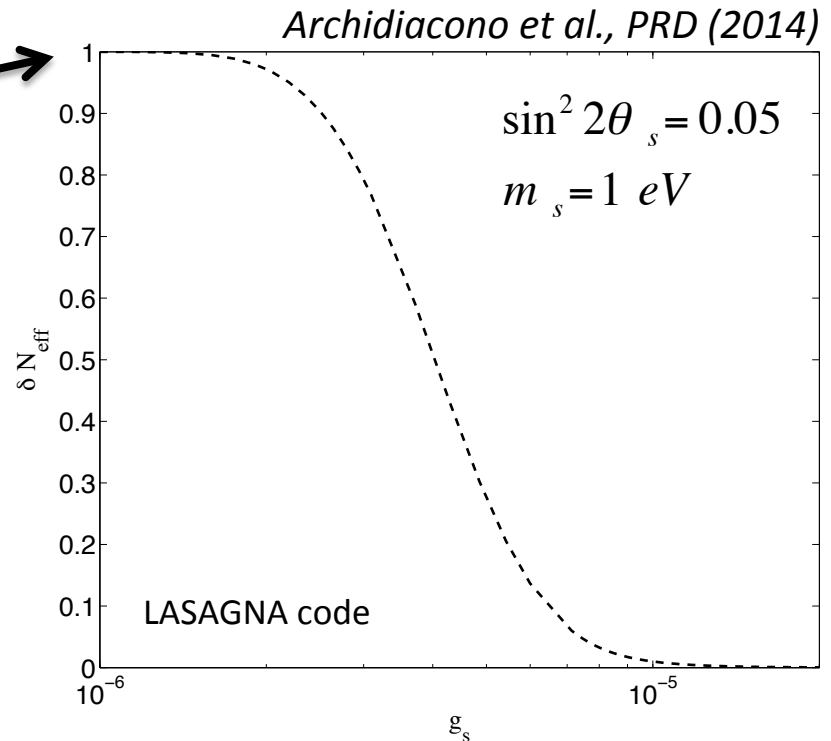
$$V_s(p_s) \sim 10^{-1} g_s^2 T_s$$

$$V_s > V_{vacuum} \quad \text{until} \quad \frac{\Gamma_a}{H} > 1 \quad (\approx 1 \text{ MeV})$$

N_{eff} at BBN

BBN bounds:
 $\Delta N_{\text{eff}} \leq 1$ (95% c.l.)

When sterile neutrinos are produced, they will create non-thermal distortions in the sterile neutrino distribution, and the sterile neutrino spectrum end up being somewhat non-thermal.

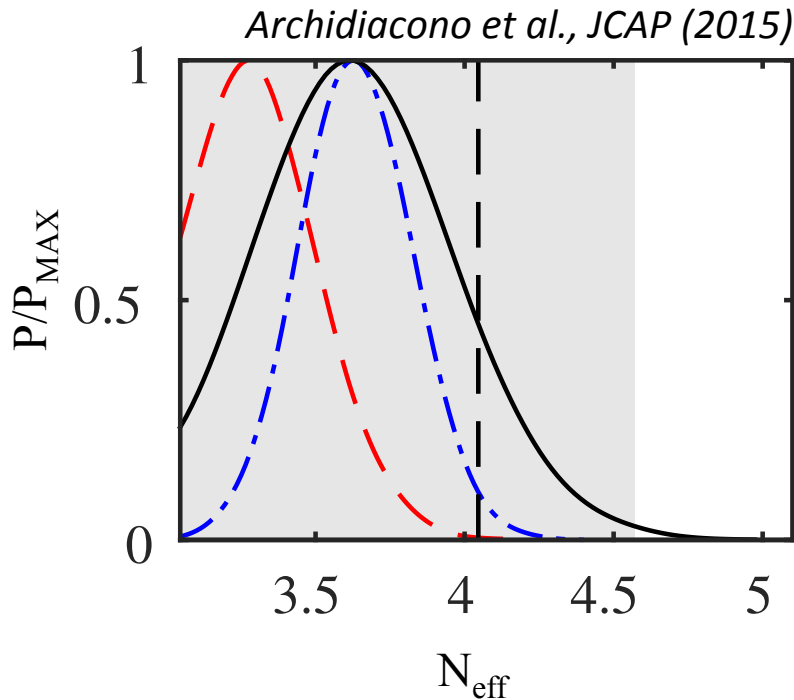


The transition between full thermalization and no thermalization occurs for coupling $10^{-6} < g_s < 10^{-5}$

N_{eff} at CMB

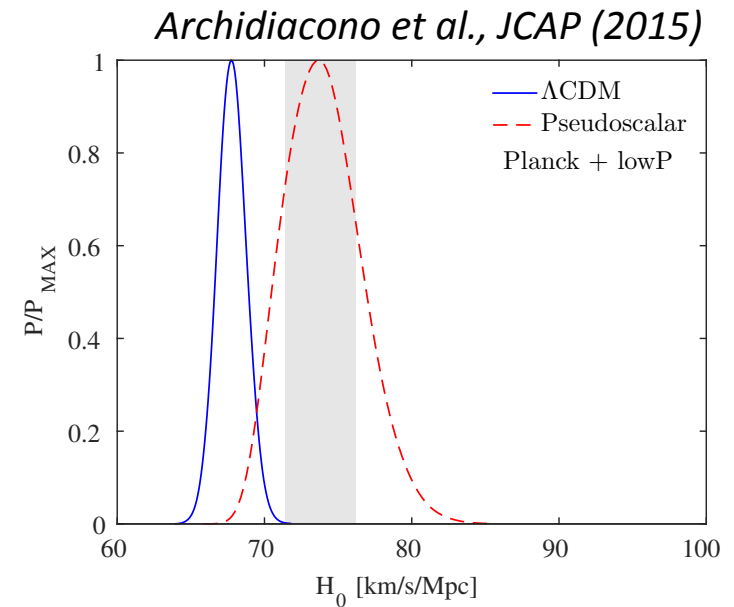
The $\nu_s - \phi$ fluid becomes strongly interacting before neutrinos go non-relativistic around recombination.

$$\Gamma_s = \frac{g_s^4}{4\pi T_s^2} n_s$$



Correlation between values of N_{eff} and values of g_s

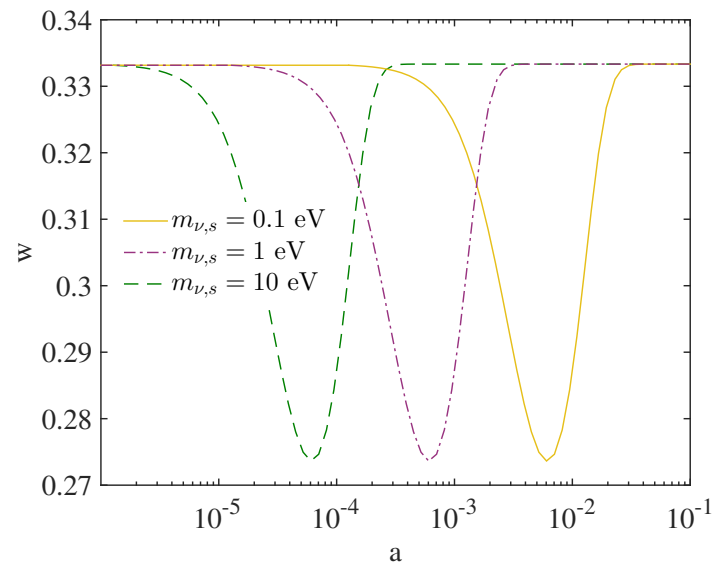
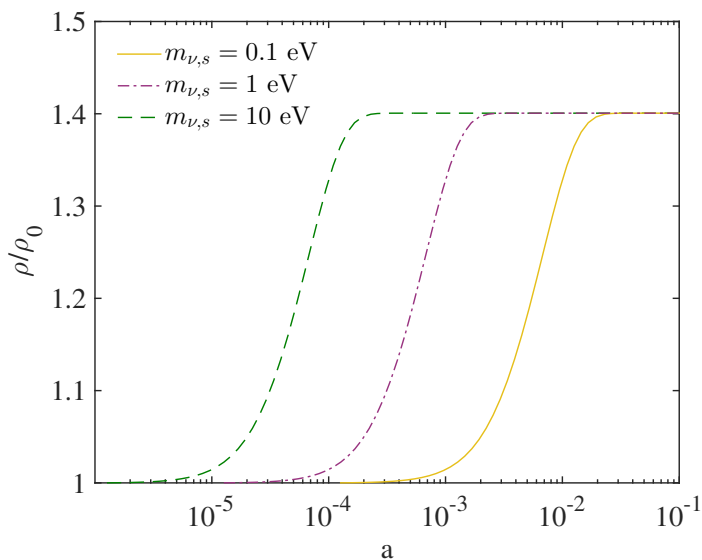
Consistent with HST



Σm_ν and LSS

As soon as sterile neutrinos go non-relativistic, they start annihilating into pseudoscalars.

Archidiacono et al., JCAP (2015)

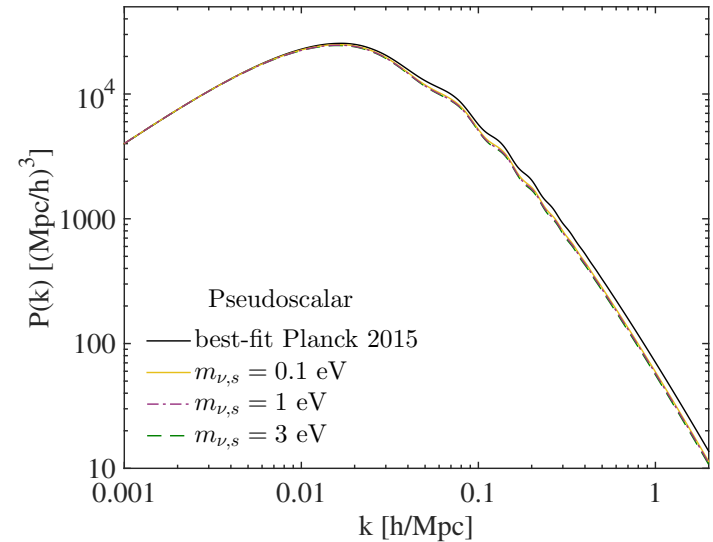
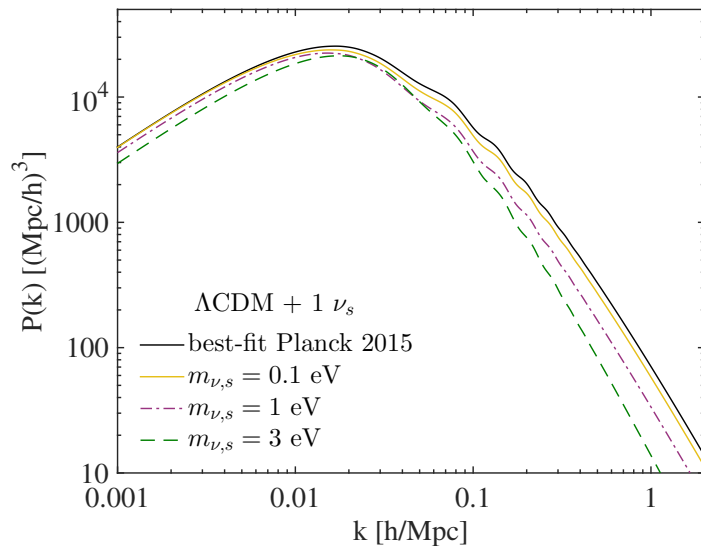


Σm_ν and LSS

Sterile neutrinos disappear from the cosmic neutrino background.

Neutrinoless Universe, Beacom et al., PRL (2004)

Archidiacono et al., JCAP (2015)

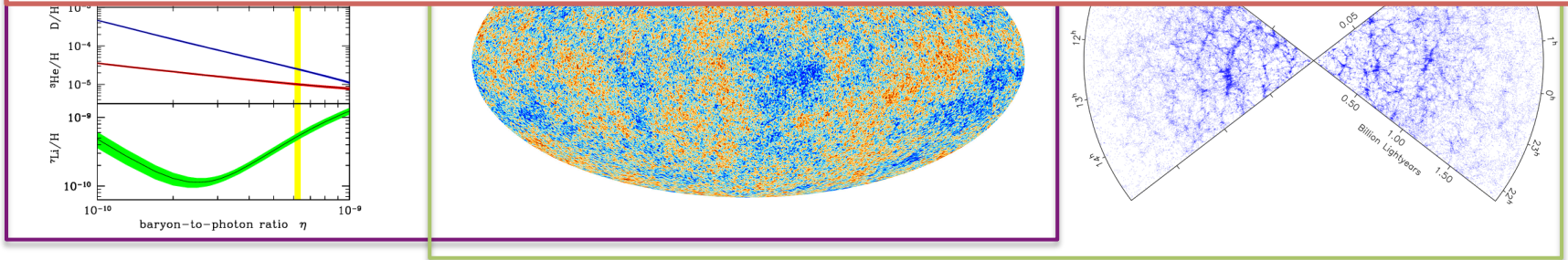


If the mediator is a massive MeV vector boson, then the late time phenomenology is different.

Hannestad et al., PRL(2013); Bringmann et al., JCAP (2014); Mirizzi et al., PRD (2014); Chu, Dasgupta, Kopp, JCAP (2015)

Conclusions

“Secret” sterile neutrino self-interactions mediated by a light pseudoscalar can accommodate one additional massive sterile state in cosmology without spoiling BBN, CMB and LSS constraints.



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$$N_{\text{eff}} = 2.99 \pm 0.20 \text{ (68\%cl)}$$

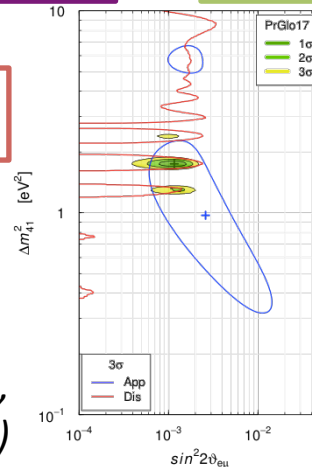
$$\Sigma m_\nu < 0.13 \text{ eV (95\%cl)}$$

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$$10^{-6} < g_s < 10^{-5}$$

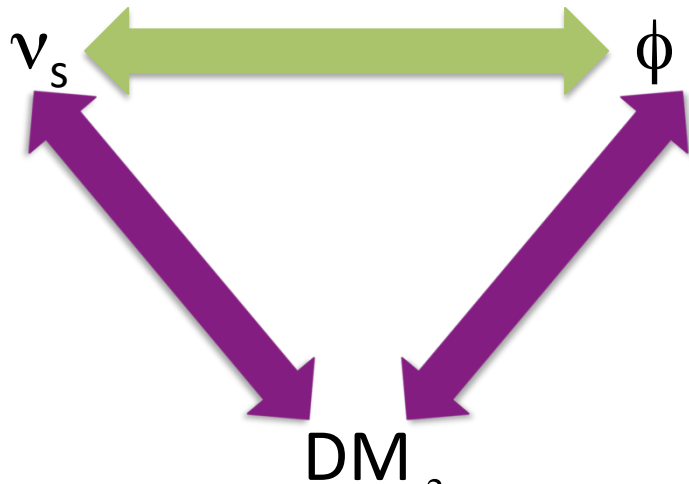
$$m_\phi \ll m_s$$

*Gariazzo et al.,
Global fit (2017)*



$$N_{\text{eff}} = 4.0... \\ \Sigma m_\nu \approx 1 \text{ eV}$$

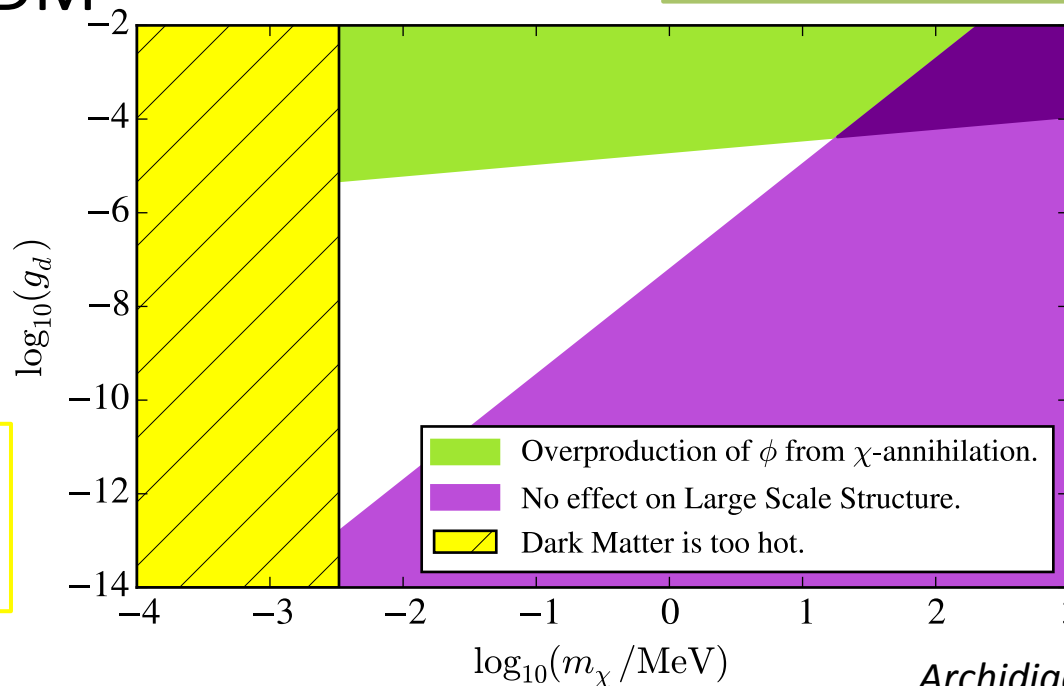
Outlook



Secret interactions might also solve the small scale problems of cold dark matter.

$$g_d \leq 2 \times 10^{-5} \left(\frac{m_\chi}{\text{MeV}} \right)^{1/4}$$

$$g_d \geq 6 \times 10^{-8} \left(\frac{m_\chi}{\text{MeV}} \right)^{9/4}$$



WDM limit
 $m_\chi > 3.3 \text{ keV}$
 Viel et al. (2013)

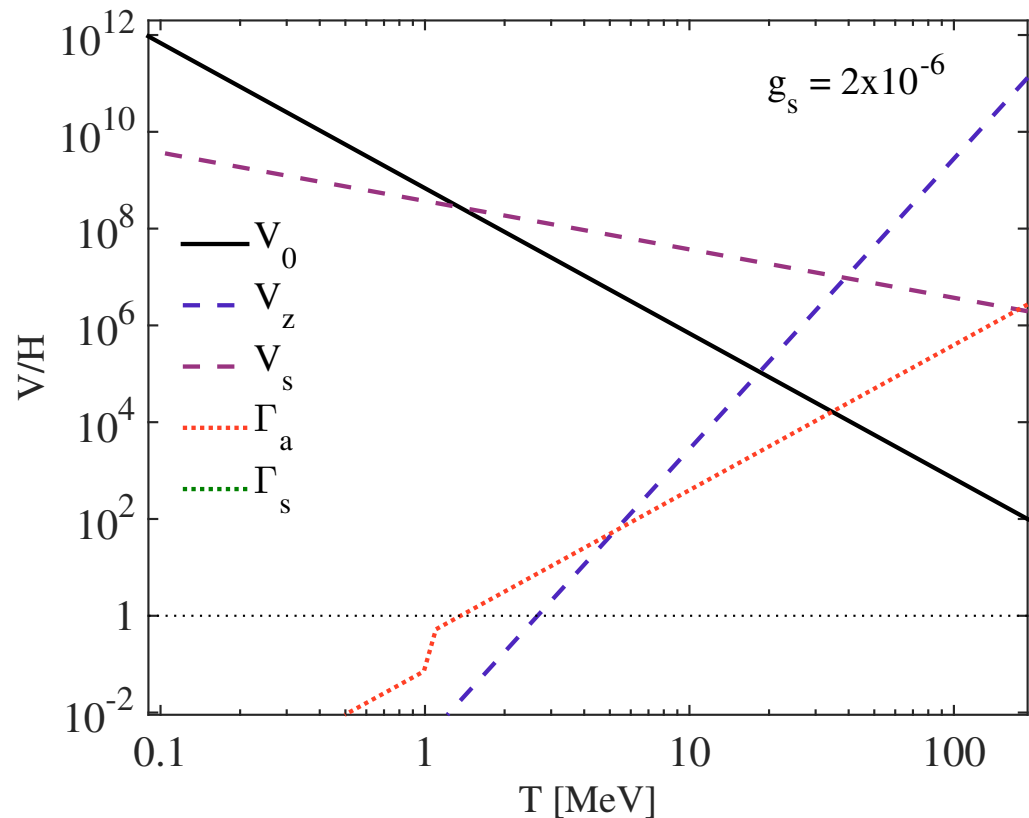
The scattering time scale of DM self interactions has to be less than the age of the Universe.

Archidiacono et al., PRD (2014)

Flavour oscillations

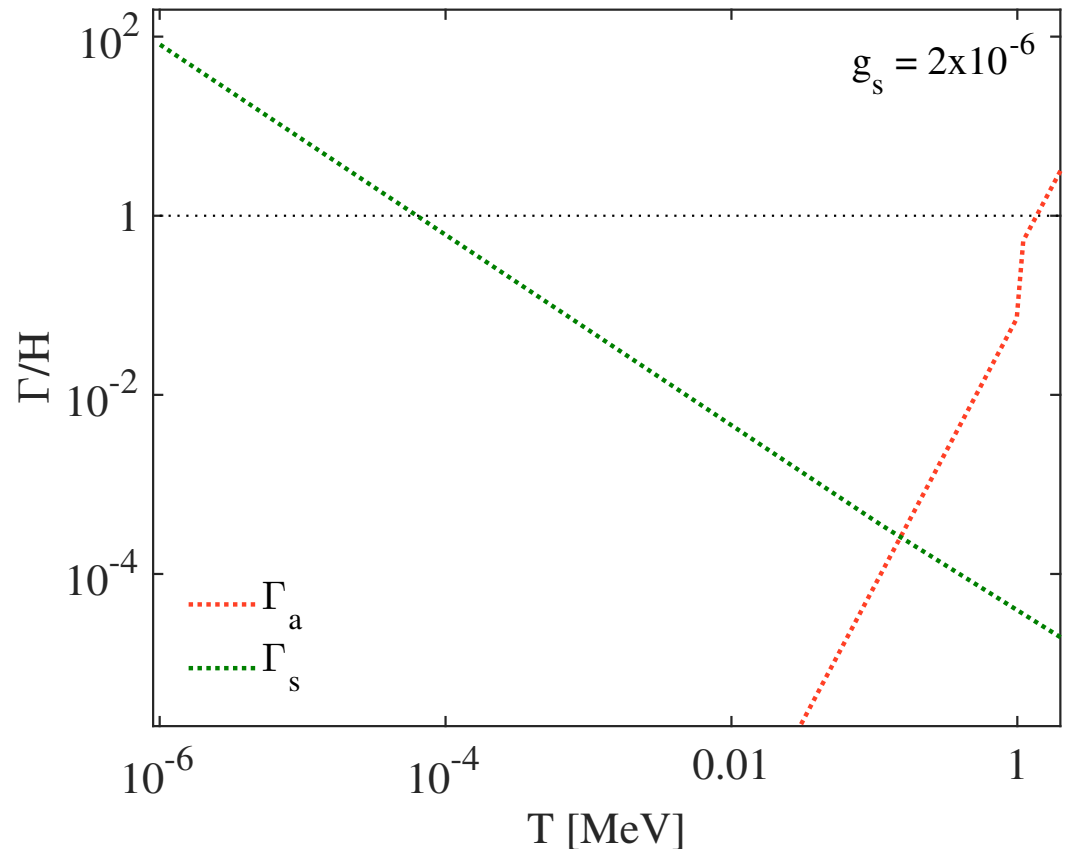
$$\sin^2 2\theta_m = \frac{\sin^2 2\theta_0}{\left(\cos 2\theta_0 + \frac{2E}{\Delta m^2} V_{\text{eff}}\right)^2 + \sin^2 2\theta_0}$$

$$V_s(p_s) = \frac{g_s^2}{8\pi^2 p_s} \int p dp (f_\phi + f_s) \sim 10^{-1} g_s^2 T_s$$

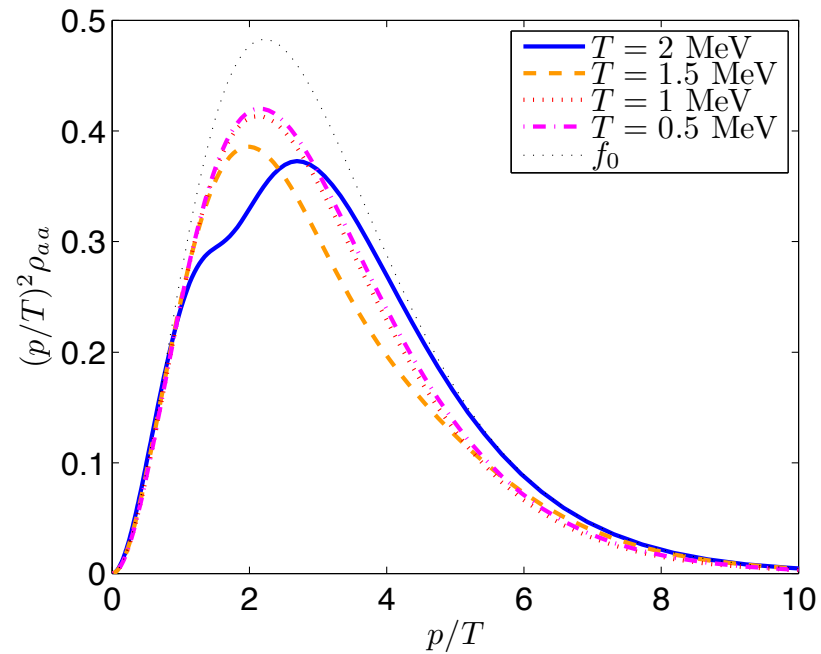


Collisional re-coupling

$$\Gamma_{\text{coll}} = n_{\nu_s} \sigma \sim \begin{cases} n_{\nu_s} e_s^4 \frac{E^2}{M^4} & \text{for } T_s \ll M \\ n_{\nu_s} e_s^4 \frac{1}{E^2} & \text{for } T_s \gg M \end{cases} \quad \Gamma_s \simeq \frac{1}{2} \sin^2 2\theta_m \times \frac{3}{4} n_{\nu_a}^{\text{SM}} \cdot \begin{cases} e_s^4 \frac{E^2}{M^4} & \text{for } T_s \ll M \\ e_s^4 \frac{1}{E^2} & \text{for } T_s \gg M \end{cases}$$



Neutrino PDF



Hannestad et al., PRL (2013)

DM-pseudoscalar

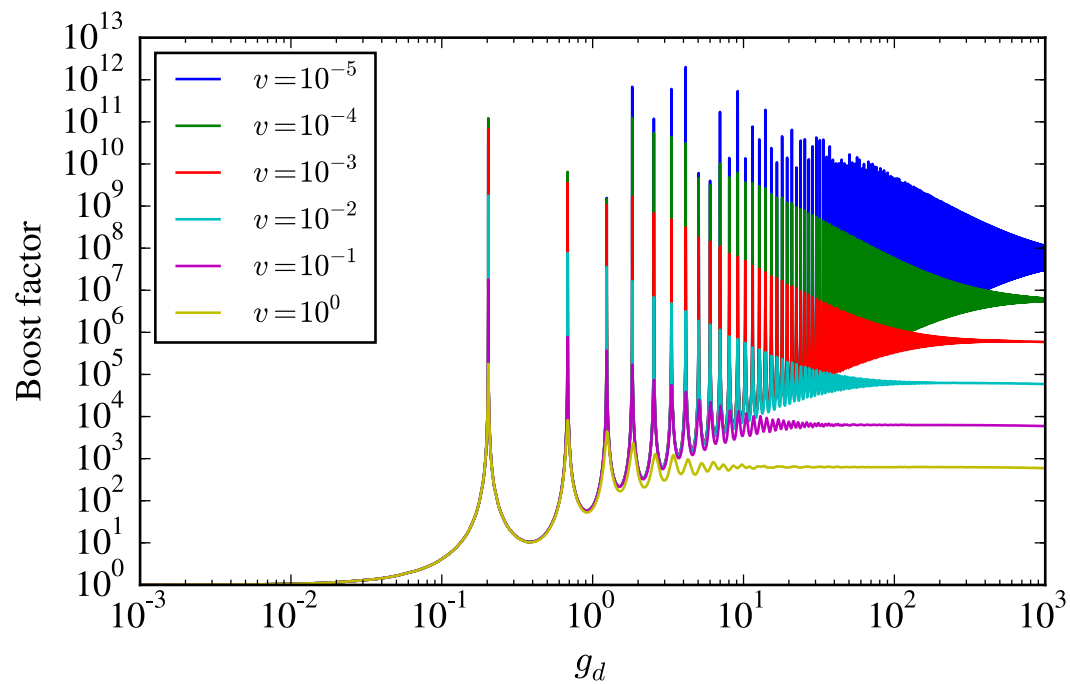
$$V(r) = -\frac{g_d^2}{m_\chi^2} \frac{e^{-m_\phi r}}{4\pi r^3} h(m_\phi r) \mathcal{S},$$

$$h(m_\phi, r) = 1 + m_\phi r + \frac{1}{3}(m_\phi r)^2$$

$$\Gamma_d(T) = \langle \sigma | v | \rangle n \propto g_d^4 T$$

$$\Gamma_d(T_{MAX} = m_\chi / 3) < H(T_{MAX} = m_\chi / 3)$$

Sommerfeld & pseudoscalar



Dark matter & pseudoscalar

$$\frac{\tau_{scat}}{\tau_{dyn}} = \frac{2R^2}{3N_\chi\sigma} \left\{ \begin{array}{l} \tau_{dyn} = \frac{2\pi R}{v} \\ \tau_{scat} = \frac{1}{n\langle\sigma|v|\rangle} \end{array} \right. \quad N_\chi = \frac{M_{gal}}{m_\chi}$$

Hard scattering $\sigma \sim 4\pi b^2$ $\frac{1}{2}m_\chi v^2 = \frac{\alpha_d}{m_\chi b^3}$ $\alpha_d = \frac{g_d^2}{4\pi}$

The condition for having observable consequences on galactic dynamics is that the scattering time scale of DM self interactions is less than the age of the Universe.

Milky Way:

$$g_d \geq 6 \times 10^{-8} \left(\frac{m_\chi}{MeV} \right)^{9/4}$$

Bellazzini et al., PRD (2013)
Ackerman et al., PRD (2009)

It is just a lower bound
 It requires further
 investigation