

AKK update: AKK08 fragmentation functions

Simon Albino

¹II. Institute of Theoretical Physics, University of Hamburg

May 22, 2009

Outline

- 1 Single hadron inclusive production and pQCD
 - Introduction
 - Measurements of single hadron inclusive production
 - Outline of factorization
- 2 FFs at large x
 - Cross section in the fixed order approach
 - Global fits
 - Predictions using FFs from e^+e^-
 - Theoretical improvements
 - Latest global fit: AKK08 (Nucl. Phys. B)
- 3 Improving the small x region
 - Soft gluon logarithms in splitting functions
 - Unified approach
 - Fits to data

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
- x : fraction of the process's available energy / momentum taken away by h

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
- } final state

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
 - E_S : energy scale of process
- } final state

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
 - E_S : energy scale of process
- } final state

Parton (quark/gluon) fragmentation “model” from probability:

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
 - E_S : energy scale of process
- } final state

Parton (quark/gluon) fragmentation “model” from probability:

$$XS(? \rightarrow h + X) = XS(? \rightarrow i + X) \times \text{FF}(i \rightarrow h + X)$$

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
 - E_S : energy scale of process
- } final state

Parton (quark/gluon) fragmentation “model” from probability:

$$XS(? \rightarrow h + X) = XS(? \rightarrow i + X) \times \text{FF}(i \rightarrow h + X)$$

FFs: probability for parton i to *fragment* to h

General cross section and parton fragmentation

$? \rightarrow h + X$ ($X = \text{anything}$) takes the form $\frac{d\sigma^h}{dx}(x, E_S^2)$

(Initial states “?”: e^+e^- , γ^*p , pp)

- h : detected hadron
 - x : fraction of the process's available energy / momentum taken away by h
 - E_S : energy scale of process
- } final state

Parton (quark/gluon) fragmentation “model” from probability:

$$XS(? \rightarrow h + X) = XS(? \rightarrow i + X) \times FF(i \rightarrow h + X)$$

FF s: probability for parton i to *fragment* to h

$XS(? \rightarrow i + X)$: partonic final state, pQCD at high E_S
 \rightarrow series in QCD coupling $a_s(E_S) \sim \frac{1}{\ln E_S/\Lambda_{\text{QCD}}}$

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^* p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp
 - $\mathcal{L} = 0.5 \text{ fb}^{-1}$ (1996-2007), theory \leftrightarrow exp

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp
 - $\mathcal{L} = 0.5 \text{ fb}^{-1}$ (1996-2007), theory \leftrightarrow exp

$$pp(\bar{p}) \rightarrow h + X$$

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp
 - $\mathcal{L} = 0.5 \text{ fb}^{-1}$ (1996-2007), theory \leftrightarrow exp

$$pp(\bar{p}) \rightarrow h + X$$

- CDF (2005) @ Tevatron ($p\bar{p}$, $\sqrt{s} = 630, 1800 \text{ GeV}$)
no norm \rightarrow FF shape

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^*p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp
 - $\mathcal{L} = 0.5 \text{ fb}^{-1}$ (1996-2007), theory \leftrightarrow exp

$$pp(\bar{p}) \rightarrow h + X$$

- CDF (2005) @ Tevatron ($p\bar{p}$, $\sqrt{s} = 630, 1800 \text{ GeV}$)
no norm \rightarrow FF shape
- BRAHMS, PHENIX, STAR (2006,2007) @ RHIC (pp , $\sqrt{s} = 200 \text{ GeV}$)
 \rightarrow remaining FFs (gluon, valence quark, ...)

Measurements from several processes

$$e^+e^- \rightarrow h + X$$

- ALEPH, DELPHI, TASSO, OPAL, ... — many data points
- Most precise data for (charge-summed) FF extraction

$$\gamma^* p \rightarrow h + X$$

- $\sim PDFs \times e^+e^- \rightarrow h + X \rightarrow$ different FF flavour weighting
- H1+ZEUS@HERA
 - $\mathcal{L} = 0.5 - 2.5 \text{ pb}^{-1}$ (1993,1994), theory \leftrightarrow exp
 - $\mathcal{L} = 0.5 \text{ fb}^{-1}$ (1996-2007), theory \leftrightarrow exp

$$pp(\bar{p}) \rightarrow h + X$$

- CDF (2005) @ Tevatron ($p\bar{p}$, $\sqrt{s} = 630, 1800 \text{ GeV}$)
no norm \rightarrow FF shape
- BRAHMS, PHENIX, STAR (2006,2007) @ RHIC (pp , $\sqrt{s} = 200 \text{ GeV}$)
 \rightarrow remaining FFs (gluon, valence quark, ...)
- LHC (pp , $\sqrt{s} = 14 \text{ TeV}$)

Factorize into hard + soft

- Emerging picture:

- Purpose:

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = d\sigma_{q,g}(E_S, M_f) \otimes D_{q,g}^h(M_f)$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard"}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft"}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft"}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$
- processes $< O(M_f)$ into FFs (so choose $M_f = O(E_S)$)

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$
- processes $< O(M_f)$ into FFs (so choose $M_f = O(E_S)$)

Allows for perturbative calculations dependent on initial state:

$$d\sigma_{q,g}(E_S, M_f) = \sum_n a_s^n(M_f) d\sigma_{q,g}^{(n)}(E_S, M_f)$$

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$
- processes $< O(M_f)$ into FFs (so choose $M_f = O(E_S)$)

Allows for perturbative calculations dependent on initial state:

$$d\sigma_{q,g}(E_S, M_f) = \sum_n a_s^n(M_f) d\sigma_{q,g}^{(n)}(E_S, M_f)$$

- M_f dependence of FFs also perturbative \rightarrow DGLAP:

$$\frac{d}{d \ln M_f} D_{q,g}^h(M_f) = P_{q,g} \otimes_{q,g}(a_s(M_f)) \otimes D_{q,g}^h(M_f)$$

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$
- processes $< O(M_f)$ into FFs (so choose $M_f = O(E_S)$)

Allows for perturbative calculations dependent on initial state:

$$d\sigma_{q,g}(E_S, M_f) = \sum_n a_s^n(M_f) d\sigma_{q,g}^{(n)}(E_S, M_f)$$

- M_f dependence of FFs also perturbative \rightarrow DGLAP:

$$\frac{d}{d \ln M_f} D_{q,g}^h(M_f) = P_{q,g} \otimes_{q,g}(a_s(M_f)) \otimes D_{q,g}^h(M_f)$$

$$\text{where } P(a_s) = \sum_n a_s^n P^{(n-1)} \text{ calculable}$$

Factorize into hard + soft

- Emerging picture:

$$d\sigma^h(E_S) = \underbrace{d\sigma_{q,g}(E_S, M_f)}_{\text{"hard": pQCD, PDFs}} \otimes \underbrace{D_{q,g}^h(M_f)}_{\text{"soft": FFs}}$$

($\otimes \equiv$ sum over partons, spin, integral over momentum etc.)

- Purpose:

- processes $< O(E_S)$ spoil a_s -series for $d\sigma_{q,g}$
- processes $< O(M_f)$ into FFs (so choose $M_f = O(E_S)$)

Allows for perturbative calculations dependent on initial state:

$$d\sigma_{q,g}(E_S, M_f) = \sum_n a_s^n(M_f) d\sigma_{q,g}^{(n)}(E_S, M_f)$$

- M_f dependence of FFs also perturbative \rightarrow DGLAP:

$$\frac{d}{d \ln M_f} D_{q,g}^h(M_f) = P_{q,g} \otimes D_{q,g}^h(M_f)$$

$$\text{where } P(a_s) = \sum_n a_s^n P^{(n-1)} \text{ calculable}$$

- FFs at $M_f = M_0$ from exp, then this approach predicts any XS

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz \, d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz \, d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

\longrightarrow PDF evolution techniques can be used:

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

→ PDF evolution techniques can be used:

- Parameterize FFs at $M_f = M_0$, $D_i^h(z, M_0^2) = N_i z^{\alpha_i} (1 - z)^{\beta_i}$

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

→ PDF evolution techniques can be used:

- Parameterize FFs at $M_f = M_0$, $D_i^h(z, M_0^2) = N_i z^{\alpha_i} (1-z)^{\beta_i}$
- Solve DGLAP in Mellin space $D^h(\omega, M_f^2) = \int_0^1 dz z^\omega D^h(z, M_f^2)$
convolution $\otimes \rightarrow$ product \times :

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d\ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d\ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

→ PDF evolution techniques can be used:

- Parameterize FFs at $M_f = M_0$, $D_i^h(z, M_0^2) = N_i z^{\alpha_i} (1-z)^{\beta_i}$
- Solve DGLAP in Mellin space $D^h(\omega, M_f^2) = \int_0^1 dz z^\omega D^h(z, M_f^2)$
convolution $\otimes \rightarrow$ product \times :

$$\frac{d}{d\ln M_f^2} D_i^h(\omega, M_f^2) = P_{ij}(\omega, a_s(M_f^2)) D_j^h(\omega, M_f^2)$$

$$\text{gives } D_i^h(\omega, M_f^2) = E_{ij}(\omega, a_s(M_f^2), a_s(M_0^2)) D_j^h(\omega, M_0^2)$$

Outline of factorization

Convolution \otimes in $d\sigma^h = d\sigma_{q,g} \otimes D_{q,g}^h$, $\frac{d}{d \ln M_f} D_{q,g}^h = P_{q,gq,g} \otimes D_{q,g}^h$

$$d\sigma^h(x, E_S) = \sum_i \int_x^1 dz d\sigma_i\left(\frac{x}{z}, E_S, M_f\right) D_i^h(z, M_f)$$

$$\frac{d}{d \ln M_f^2} D_i^h(z, M_f^2) = \sum_j \int_z^1 \frac{dz'}{z'} P_{ij}\left(\frac{z}{z'}, a_s(M_f^2)\right) D_j^h(z', M_f^2)$$

x = produced hadron momentum / available momentum

z = produced hadron momentum / fragmenting parton momentum

\sim inclusive hadron-initiated process (PDF convolution)

→ PDF evolution techniques can be used:

- Parameterize FFs at $M_f = M_0$, $D_i^h(z, M_0^2) = N_i z^{\alpha_i} (1 - z)^{\beta_i}$
- Solve DGLAP in Mellin space $D^h(\omega, M_f^2) = \int_0^1 dz z^\omega D^h(z, M_f^2)$
convolution $\otimes \rightarrow$ product \times :

$$\frac{d}{d \ln M_f^2} D_i^h(\omega, M_f^2) = P_{ij}(\omega, a_s(M_f^2)) D_j^h(\omega, M_f^2)$$

$$\text{gives } D_i^h(\omega, M_f^2) = E_{ij}(\omega, a_s(M_f^2), a_s(M_0^2)) D_j^h(\omega, M_0^2)$$

- $d\sigma^h(\omega, E_S) = d\sigma_i(\omega, E_S, M_f) D_i^h(\omega, M_f^2)$

Outline

- 1 Single hadron inclusive production and pQCD
 - Introduction
 - Measurements of single hadron inclusive production
 - Outline of factorization
- 2 FFs at large x
 - Cross section in the fixed order approach
 - Global fits
 - Predictions using FFs from e^+e^-
 - Theoretical improvements
 - Latest global fit: AKK08 (Nucl. Phys. B)
- 3 Improving the small x region
 - Soft gluon logarithms in splitting functions
 - Unified approach
 - Fits to data

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv \text{large } \omega$

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv \text{large } \omega$
- Further improvement possible for very large x (see later)

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv$ large ω
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv$ small ω (see later)

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv$ large ω
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv$ small ω (see later)

Other corrections include ($E_S = \sqrt{s}$, Q , ...)

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv \text{large } \omega$
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv \text{small } \omega$ (see later)

Other corrections include ($E_S = \sqrt{s}$, Q , ...)

- hadron mass, $O\left(\frac{m_h^2}{E_S^2}\right)$ (later)

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv \text{large } \omega$
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv \text{small } \omega$ (see later)

Other corrections include ($E_S = \sqrt{s}$, Q , ...)

- hadron mass, $O\left(\frac{m_h^2}{E_S^2}\right)$ (later)
- quark mass, $O\left(\frac{m_q^2}{E_S^2}\right)$

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv$ large ω
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv$ small ω (see later)

Other corrections include ($E_S = \sqrt{s}$, Q , ...)

- hadron mass, $O\left(\frac{m_h^2}{E_S^2}\right)$ (later)
- quark mass, $O\left(\frac{m_q^2}{E_S^2}\right)$
- higher twist, $O\left(\frac{\Lambda_{\text{QCD}}^p}{E_S^p}\right)$ (beyond factorization)

Basic approach

- Fixed order pQCD

$$P(\omega, a_s) = \sum_{n=1}^{\infty} a_s^n P^{(n-1)}(\omega) = \underbrace{a_s P^{(0)}}_{\text{LO}} + \underbrace{a_s^2 P^{(1)}}_{\text{NLO}} + \dots$$

$$d\sigma_i(\omega, E_S) = \sum_{n=0}^{\infty} a_s^n(E_S) d\sigma_i^{(n)}(\omega) \quad (M_f = O(E_S) \text{ fixed})$$

- Works well for sufficiently large $x \equiv \text{large } \omega$
- Further improvement possible for very large x (see later)
- ... and significantly so for small $x \equiv \text{small } \omega$ (see later)

Other corrections include ($E_S = \sqrt{s}$, Q , ...)

- hadron mass, $O\left(\frac{m_h^2}{E_S^2}\right)$ (later)
 - quark mass, $O\left(\frac{m_q^2}{E_S^2}\right)$
 - higher twist, $O\left(\frac{\Lambda_{\text{QCD}}^p}{E_S^p}\right)$ (beyond factorization)
- } small x

Data

$$e^+e^-$$

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)
- (heavy) quark flavour tagging \rightarrow flavour separation of FFs

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)
- (heavy) quark flavour tagging \rightarrow flavour separation of FFs
- OPAL tagging probabilities separate light flavours

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)
- (heavy) quark flavour tagging \rightarrow flavour separation of FFs
- OPAL tagging probabilities separate light flavours
- exclude $x \lesssim 0.05$ due to failure of FO series

$$pp(\bar{p})$$

Data

$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)
- (heavy) quark flavour tagging \rightarrow flavour separation of FFs
- OPAL tagging probabilities separate light flavours
- exclude $x \lesssim 0.05$ due to failure of FO series

$$pp(\bar{p})$$

- BRAHMS, CDF, PHENIX, STAR

Data

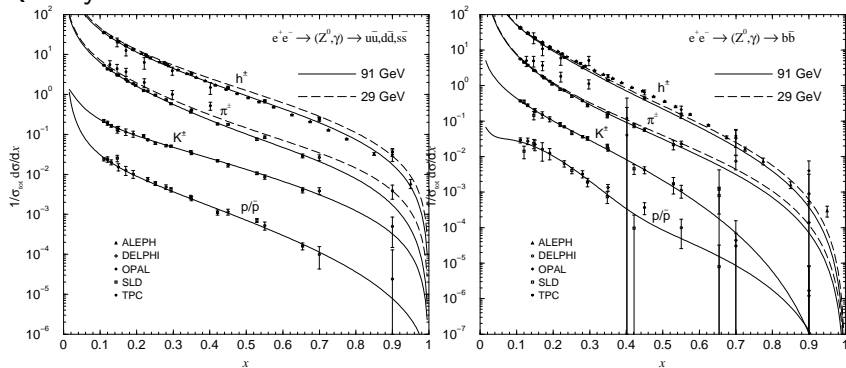
$$e^+e^-$$

- Hadrons: Light charged ($\pi^\pm(u, d)$, $K^\pm(u, s)$, $p/\bar{p}(u, u, d)$), neutral strange ($K_S^0(d, s)$, $\Lambda(u, d, s)$)
- Σ Hadron (quark) spins, charge (h^+ , h^- not distinguished)
- much accurate data @ $\sqrt{s} = 91 \text{ GeV} \gg \Lambda_{\text{QCD}} \simeq 0.2 \text{ GeV}$ (e.g. DELPHI, SLD)
- some exp. constraints @ $\sqrt{s} < 91 \text{ GeV}$ (e.g. TPC @ 29 GeV)
- (heavy) quark flavour tagging \rightarrow flavour separation of FFs
- OPAL tagging probabilities separate light flavours
- exclude $x \lesssim 0.05$ due to failure of FO series

$$pp(\bar{p})$$

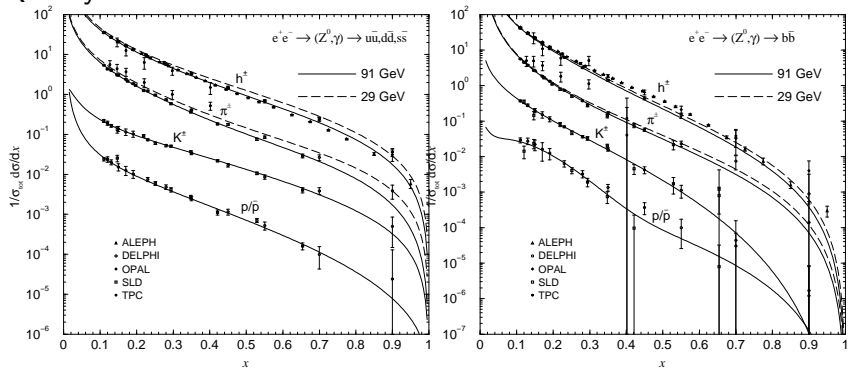
- BRAHMS, CDF, PHENIX, STAR
- h^+ , h^- separately

Quality of fits



From Albino-Kniehl-Kramer (AKK) fit of 2005
(no unidentified particle data, maybe contaminated with e, μ, \dots)

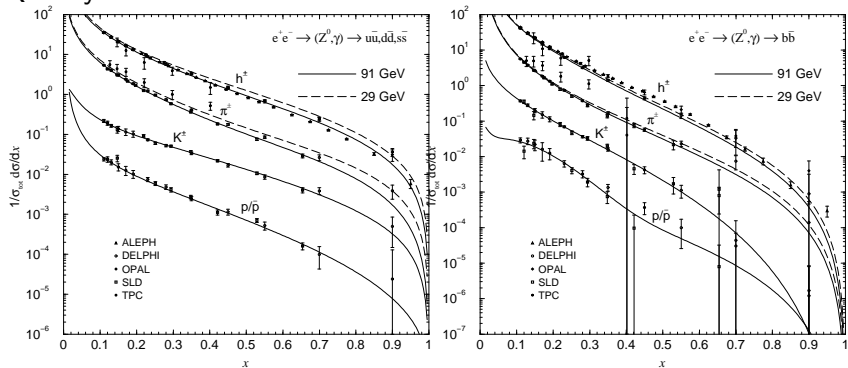
Quality of fits



From Albino-Kniehl-Kramer (AKK) fit of 2005
(no unidentified particle data, maybe contaminated with e, μ, \dots)

Such fits give e.g. $\alpha_s(M_Z) = 0.117^{+0.005}_{-0.007}$ (KKPötter)

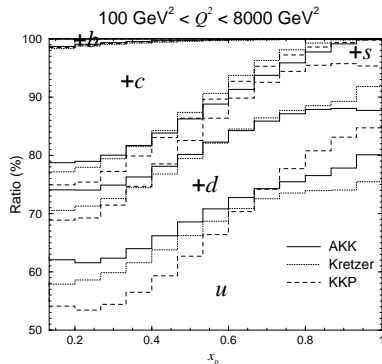
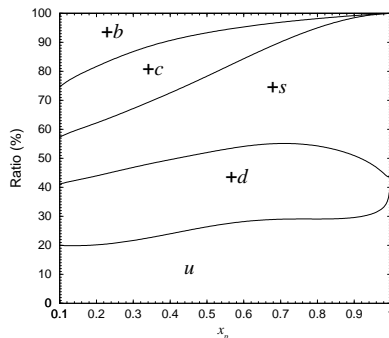
Quality of fits



From Albino-Kniehl-Kramer (AKK) fit of 2005
(no unidentified particle data, maybe contaminated with e, μ, \dots)

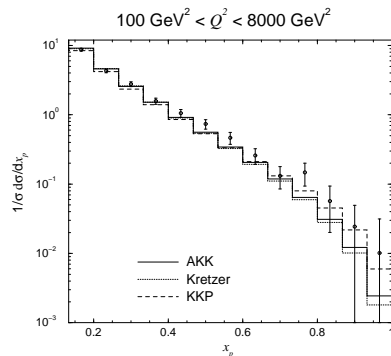
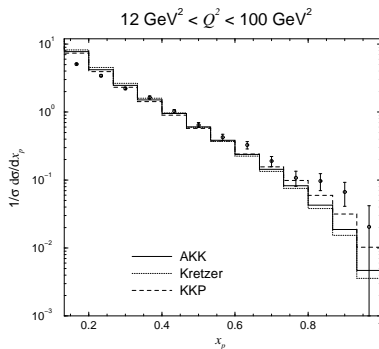
Such fits give e.g. $\alpha_s(M_Z) = 0.117^{+0.005}_{-0.007}$ (KKPötter)

More recently: HKNS, DSS, AKK08

$\gamma^* p \rightarrow h + X$ data from HERARatios of quark fragmentation contributions in $\gamma \rightarrow h + X$  ep : High Q @ H1 e^+e^- : $\sqrt{s} = 91 \text{ GeV}$ (This is for unidentified hadrons \simeq charged pions)

$\gamma^* p \rightarrow h + X$ data from HERA

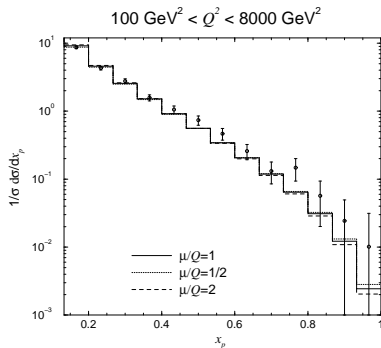
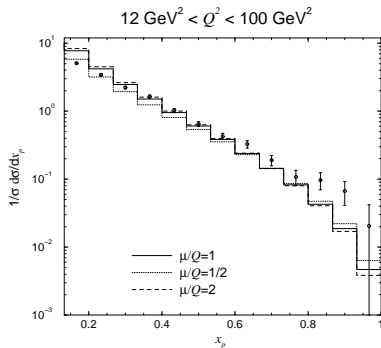
Albino et al. 2007



If good constraints on all pion FFs \rightarrow agreement at large Q ,
 then disagreement at small Q from theoretical errors
 (e.g. large/small x resummation, detected hadron mass, etc.)

$\gamma^* p \rightarrow h + X$ data from HERA

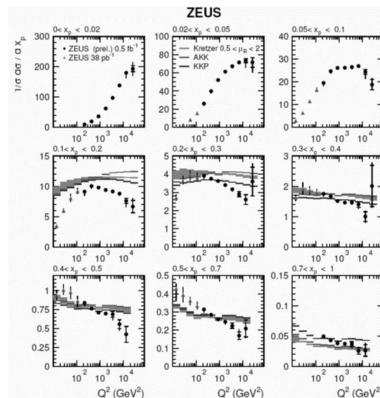
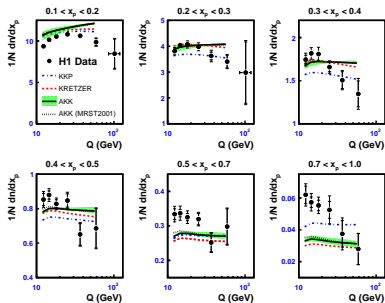
Perturbative errors from scale variation



Largest at small Q and small/large $x \rightarrow$ resummation may help

$\gamma^* p \rightarrow h + X$ data from HERA

New improved data from HERA



New improved FFs needed

Detected hadron mass effects (Albino et al., 2006-2008)

Detected hadron mass effects (Albino et al., 2006-2008)

- $$\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$$

Detected hadron mass effects (Albino et al., 2006-2008)

- $\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$
- identify x, z as ratios of light-like momenta

Detected hadron mass effects (Albino et al., 2006-2008)

- $\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$
- identify x, z as ratios of light-like momenta
- find $x = x(x_{\text{measured}})$

Detected hadron mass effects (Albino et al., 2006-2008)

- $\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$
- identify x, z as ratios of light-like momenta
- find $x = x(x_{\text{measured}})$
- take $\frac{d\sigma^h}{dx_{\text{measured}}} = \frac{dx}{dx_{\text{measured}}} \frac{d\sigma^h}{dx}$

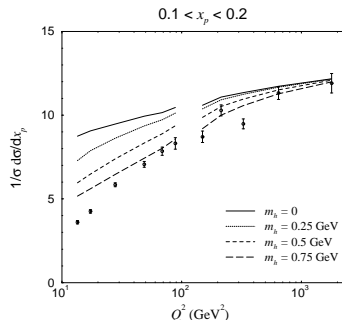
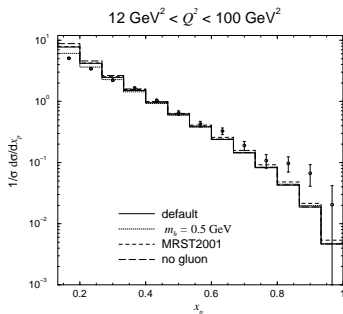
Detected hadron mass effects (Albino et al., 2006-2008)

- $\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$
- identify x, z as ratios of light-like momenta
- find $x = x(x_{\text{measured}})$
- take $\frac{d\sigma^h}{dx_{\text{measured}}} = \frac{dx}{dx_{\text{measured}}} \frac{d\sigma^h}{dx}$
- Modification at low x, E_S — e.g. HERA data:

Theoretical improvements

Detected hadron mass effects (Albino et al., 2006-2008)

- $\frac{d\sigma^h}{dx}(x, E_S^2) = \sum_{i=q,g} \int_x^1 \frac{dz}{z} \frac{d\sigma^i}{d(x/z)} \left(\frac{x}{z}, \frac{E_S^2}{M_f^2}, a_s(M_f^2) \right) D_i^h(z, M_f^2)$
- identify x, z as ratios of light-like momenta
- find $x = x(x_{\text{measured}})$
- take $\frac{d\sigma^h}{dx_{\text{measured}}} = \frac{dx}{dx_{\text{measured}}} \frac{d\sigma^h}{dx}$
- Modification at low x, E_S — e.g. HERA data:



Large x resummation

Large x resummation

- $\frac{d\sigma^i}{dx}$, large x (\equiv large ω), $\simeq 1 + Aa_s \ln \omega + Ba_s^2 \ln^2 \omega + \dots$
→ perturbation theory fails

Large x resummation

- $\frac{d\sigma^i}{dx}$, large x (\equiv large ω), $\simeq 1 + Aa_s \ln \omega + Ba_s^2 \ln^2 \omega + \dots$
→ perturbation theory fails
- **resum** to get $f(a_s \ln \omega)$

Large x resummation

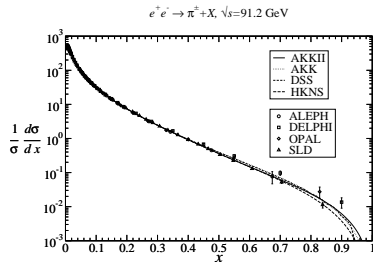
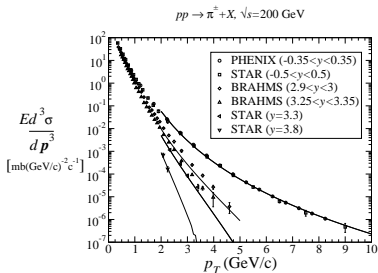
- $\frac{d\sigma^i}{dx}$, large x (\equiv large ω), $\simeq 1 + Aa_s \ln \omega + Ba_s^2 \ln^2 \omega + \dots$
→ perturbation theory fails
- **resum** to get $f(a_s \ln \omega)$
- Likewise, subleading $\bar{A}a_s + \bar{B}a_s(a_s \ln \omega) + \bar{C}a_s(a_s \ln \omega)^2 + \dots$
→ $a_s \bar{f}(a_s \ln \omega)$

Large x resummation

- $\frac{d\sigma^i}{dx}$, large x (\equiv large ω), $\simeq 1 + Aa_s \ln \omega + Ba_s^2 \ln^2 \omega + \dots$
→ perturbation theory fails
- **resum** to get $f(a_s \ln \omega)$
- Likewise, subleading $\bar{A}a_s + \bar{B}a_s(a_s \ln \omega) + \bar{C}a_s(a_s \ln \omega)^2 + \dots$
→ $a_s \bar{f}(a_s \ln \omega)$
- Albino et al. (2008): resummation in DGLAP evolution
→ modifies theoretical errors

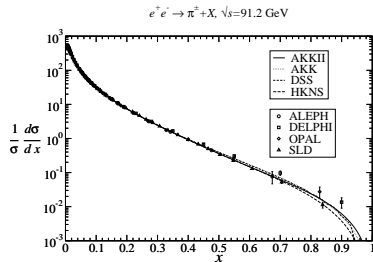
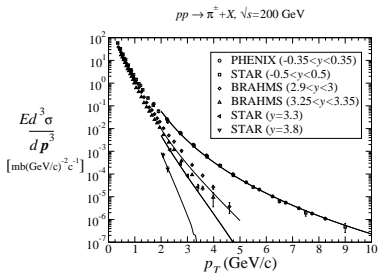
Detected hadron mass effects

“Naive” FO calculation is sufficient



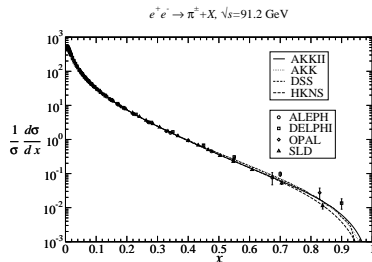
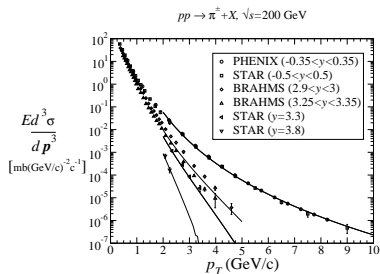
Detected hadron mass effects

“Naive” FO calculation is sufficient

Main caveat: **mass effects** for heavier particles

Detected hadron mass effects

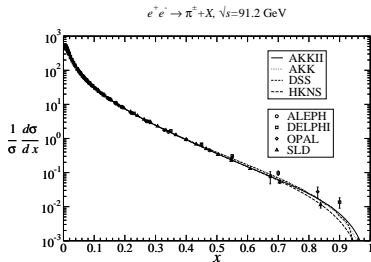
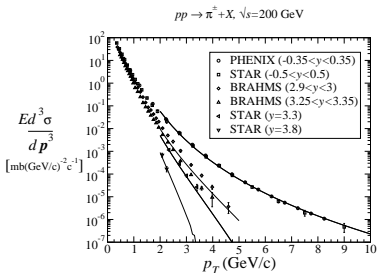
“Naive” FO calculation is sufficient



Main caveat: **mass effects** for heavier particles
 (Second caveat: **Complex decay patterns**)

Detected hadron mass effects

“Naive” FO calculation is sufficient



Main caveat: **mass effects** for heavier particles
 (Second caveat: **Complex decay patterns**)

Fitting m_h subtracts unaccounted-for small x, E_S effects

Detected hadron mass effects

AKK08: Fit m_h to e^+e^- data

Particle	Fitted mass (MeV)	True mass (MeV)
π^\pm	154.6	139.6
K^\pm	337.0	493.7
p/\bar{p}	948.8	938.3
K_S^0	343.0	497.6
$\Lambda/\bar{\Lambda}$	1127.0	1115.7

Detected hadron mass effects

AKK08: Fit m_h to e^+e^- data

Particle	Fitted mass (MeV)	True mass (MeV)
π^\pm	154.6	139.6
K^\pm	337.0	493.7
p/\bar{p}	948.8	938.3
K_S^0	343.0	497.6
$\Lambda/\bar{\Lambda}$	1127.0	1115.7

- π^\pm : overshoot \longrightarrow e.g. $\rho(770) \rightarrow \pi^+ + \pi^-$

Detected hadron mass effects

AKK08: Fit m_h to e^+e^- data

Particle	Fitted mass (MeV)	True mass (MeV)
π^\pm	154.6	139.6
K^\pm	337.0	493.7
p/\bar{p}	948.8	938.3
K_S^0	343.0	497.6
$\Lambda/\bar{\Lambda}$	1127.0	1115.7

- π^\pm : overshoot \longrightarrow e.g. $\rho(770) \rightarrow \pi^+ + \pi^-$
- K s: undershoot \longrightarrow complicated decays ($K^* \rightarrow \pi + K, \dots$)
undershoot for $K_S^0 \simeq$ for $K^\pm \rightarrow \text{SU}(2)$ isospin

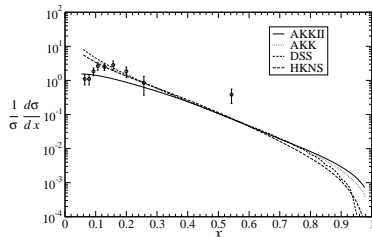
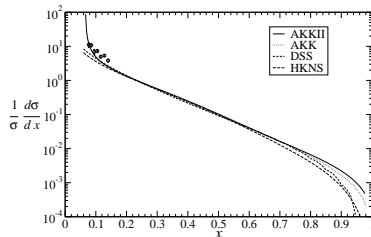
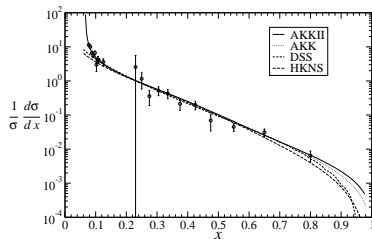
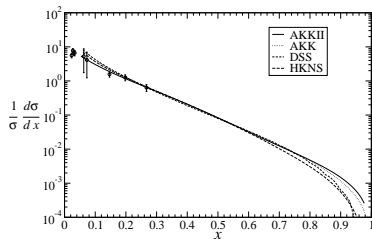
Detected hadron mass effects

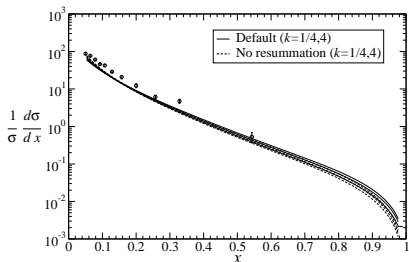
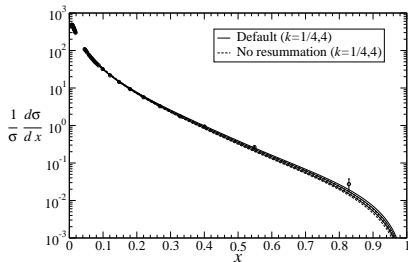
AKK08: Fit m_h to e^+e^- data

Particle	Fitted mass (MeV)	True mass (MeV)
π^\pm	154.6	139.6
K^\pm	337.0	493.7
p/\bar{p}	948.8	938.3
K_S^0	343.0	497.6
$\Lambda/\bar{\Lambda}$	1127.0	1115.7

- π^\pm : overshoot \longrightarrow e.g. $\rho(770) \rightarrow \pi^+ + \pi^-$
- K s: undershoot \longrightarrow complicated decays ($K^* \rightarrow \pi + K, \dots$)
undershoot for $K_S^0 \simeq$ for $K^\pm \rightarrow \text{SU}(2)$ isospin
- baryons ($p/\bar{p}, \Lambda/\bar{\Lambda}$): $\simeq +1\%$ \longrightarrow resonances slightly heavier
good environment to study partonic fragmentation

Detected hadron mass effects

 $e^+e^- \rightarrow p/\bar{p}+X$, TASSO, $\sqrt{s}=14$ GeV $e^+e^- \rightarrow p/\bar{p}+X$, HRS, $\sqrt{s}=29$ GeV $e^+e^- \rightarrow p/\bar{p}+X$, TPC, $\sqrt{s}=29$ GeV $e^+e^- \rightarrow p/\bar{p}+X$, TOPAZ, $\sqrt{s}=58$ GeV

Large x resummation $e^+e^- \rightarrow \pi^\pm + X$, TASSO, $\sqrt{s}=14$ GeV $e^+e^- \rightarrow \pi^\pm + X$, OPAL, $\sqrt{s}=91$ GeV

AKK08:

H	χ^2	
	Main fit	Unres. fit
π^\pm	518.7	519.0
K^\pm	416.6	439.4
p/\bar{p}	525.2	538.0
K_S^0	317.2	318.7
$\Lambda/\bar{\Lambda}$	273.1	325.7

Systematic errors

Systematic errors

$$\text{No systematic effects: } P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$$

Systematic errors

No systematic effects: $P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$

effect of K th source of systematic error: $f_i^e \rightarrow f_i^e + \lambda_K \sigma_i^K$,
 where $P(\lambda_K) \propto \exp \left[-\frac{1}{2} \lambda_K^2 \right]$ (i.e. $\langle \lambda_K^2 \rangle = 1$)

Systematic errors

No systematic effects: $P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$

effect of K th source of systematic error: $f_i^e \rightarrow f_i^e + \lambda_K \sigma_i^K$,

where $P(\lambda_K) \propto \exp \left[-\frac{1}{2} \lambda_K^2 \right]$ (i.e. $\langle \lambda_K^2 \rangle = 1$)

$$\chi^2 = \sum_i \left(\frac{f_i^t - (f_i^e + \sum_K \lambda_K \sigma_i^K)}{\sigma_i} \right)^2 + \sum_K \lambda_K^2$$

Systematic errors

No systematic effects: $P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$

effect of K th source of systematic error: $f_i^e \rightarrow f_i^e + \lambda_K \sigma_i^K$,

where $P(\lambda_K) \propto \exp \left[-\frac{1}{2} \lambda_K^2 \right]$ (i.e. $\langle \lambda_K^2 \rangle = 1$)

$$\chi^2 = \sum_i \left(\frac{f_i^t - (f_i^e + \sum_K \lambda_K \sigma_i^K)}{\sigma_i} \right)^2 + \sum_K \lambda_K^2$$

Choose λ_K such that $\frac{\partial \chi^2}{\partial \lambda_K} = 0$

Systematic errors

No systematic effects: $P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$

effect of K th source of systematic error: $f_i^e \rightarrow f_i^e + \lambda_K \sigma_i^K$,
 where $P(\lambda_K) \propto \exp \left[-\frac{1}{2} \lambda_K^2 \right]$ (i.e. $\langle \lambda_K^2 \rangle = 1$)

$$\chi^2 = \sum_i \left(\frac{f_i^t - (f_i^e + \sum_K \lambda_K \sigma_i^K)}{\sigma_i} \right)^2 + \sum_K \lambda_K^2$$

Choose λ_K such that $\frac{\partial \chi^2}{\partial \lambda_K} = 0$

$$\rightarrow \chi^2 = \sum_{ij} (f_i^t - f_i^e) (C^{-1})_{ij} (f_j^t - f_j^e) \text{ and } \lambda_K = \dots$$

Systematic errors

No systematic effects: $P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[-\frac{1}{2} \chi^2 = \sum_i \left(\frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$

effect of K th source of systematic error: $f_i^e \rightarrow f_i^e + \lambda_K \sigma_i^K$,

where $P(\lambda_K) \propto \exp \left[-\frac{1}{2} \lambda_K^2 \right]$ (i.e. $\langle \lambda_K^2 \rangle = 1$)

$$\chi^2 = \sum_i \left(\frac{f_i^t - (f_i^e + \sum_K \lambda_K \sigma_i^K)}{\sigma_i} \right)^2 + \sum_K \lambda_K^2$$

Choose λ_K such that $\frac{\partial \chi^2}{\partial \lambda_K} = 0$

$$\rightarrow \chi^2 = \sum_{ij} (f_i^t - f_i^e) (C^{-1})_{ij} (f_j^t - f_j^e) \text{ and } \lambda_K = \dots$$

Unknown systematic effects may give $|\lambda_K| \gg 1$

Systematic errors

 π^\pm production dataExpect $|\lambda_K| \simeq 1$

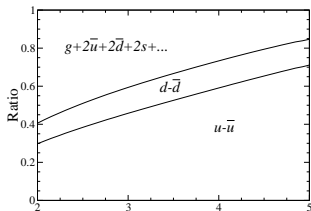
AKK08:

Collaboration	\sqrt{s} (GeV)	# data	Norm. (%)	χ^2_{DF}	λ_K
TASSO	12	5	20	0.50	0.21
TASSO	14	10	8.5	0.92	-1.26
TASSO	22	1	6.3	0.01	-0.08
TASSO	30	4	20	0.57	0.69
TASSO	34	10	6	1.07	0.62
TASSO	44	7	6	1.99	0.66
ALEPH	91.2	22	3	0.61	-0.55
BRAHMS, $y \in [2.9, 3]$ $y \in [3.25, 3.35]$	200	8	11,7,8(13), 2,1(3)	0.96	-1.76, -1.12, -1.22, -0.32, -0.13
		7		2.68	-2.01, -1.28, -1.80, -0.37, -0.32
PHENIX (π^0), $ \eta < 0.35$	200	13	9.7	0.54	-0.48
STAR (π^0), $\eta = 3.3$	200	4	16	0.70	-0.70
STAR (π^0), $\eta = 3.8$	200	2	16	0.57	-0.31
STAR, $ y < 0.5$	200	10	11.7	0.49	-0.34

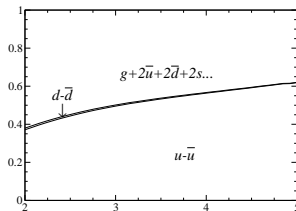
$pp \rightarrow h + X$ data from RHIC

Ratios of proton's valence/sea quark fragmentation contributions
(can be given scheme / scale independent definition)

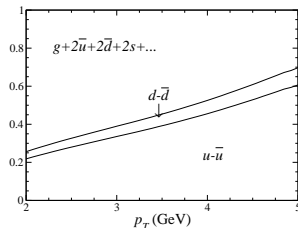
$pp \rightarrow \pi^\pm + X$ ($2.9 < y < 3$), $M_{\bar{F}} = p_T$, $\sqrt{s} = 200$ GeV



$pp \rightarrow K^\pm + X$ ($2.9 < y < 3$), $M_{\bar{F}} = p_T$, $\sqrt{s} = 200$ GeV



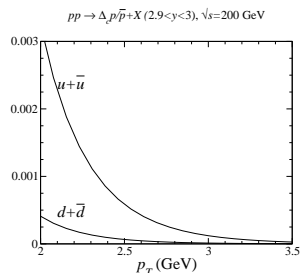
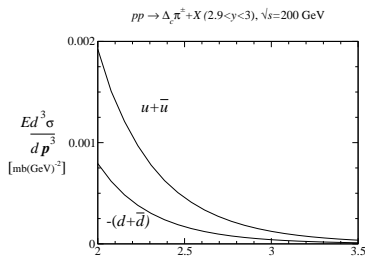
$pp \rightarrow p\bar{p} + X$ ($2.9 < y < 3$), $M_{\bar{F}} = p_T$, $\sqrt{s} = 200$ GeV



$pp \rightarrow h + X$ data from RHIC

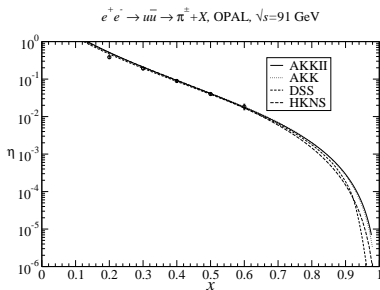
Proton's valence quark fragmentation contributions

$$\text{to } d\sigma^{h^+} - d\sigma^{h^-}$$



Comparison of FF sets

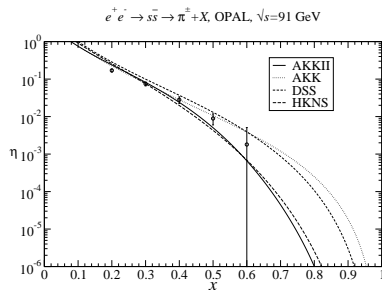
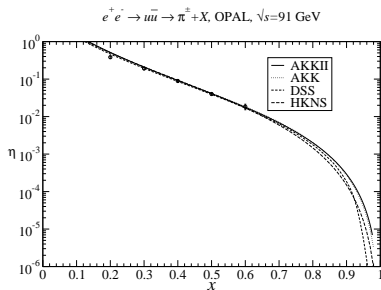
$u \rightarrow \pi$ favoured \therefore well constrained by e^+e^-



Comparison of FF sets

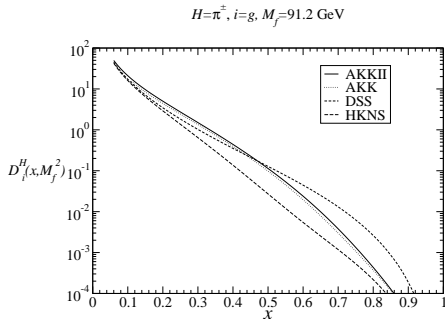
$u \rightarrow \pi$ favoured \therefore well constrained by e^+e^-

$s \rightarrow \pi$ unfavoured



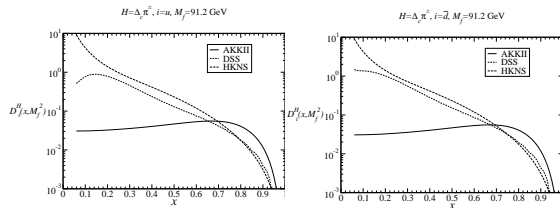
Comparison of FF sets

$g \rightarrow \pi$ constrained mainly by pp



Comparison of FF sets

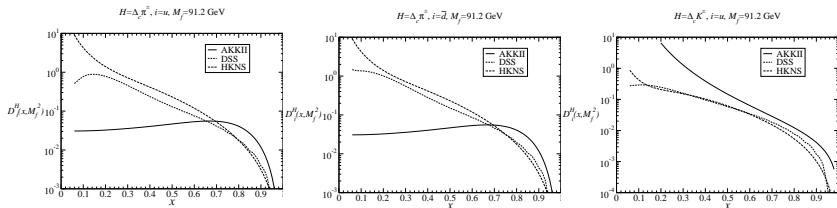
$$D_{u,d}^{h^+} - D_{u,d}^{h^-}$$



HKNS: no $d\sigma^{h^+} - d\sigma^{h^-}$ data, but \sim FF assumptions as DSS

Comparison of FF sets

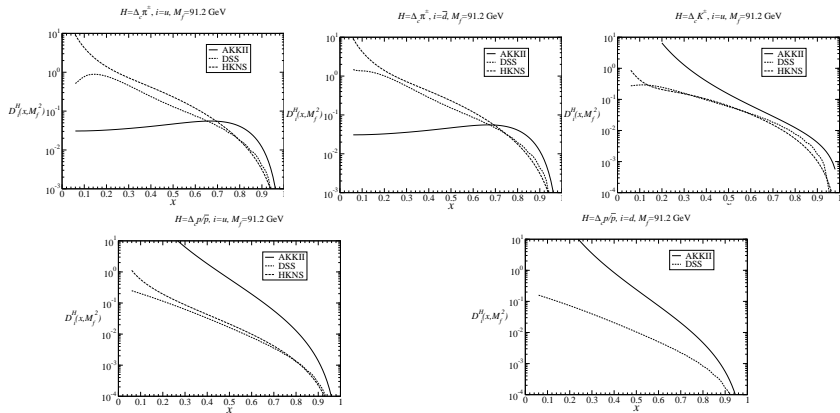
$$D_{u,d}^{h^+} - D_{u,d}^{h^-}$$



HKNS: no $d\sigma^{h^+} - d\sigma^{h^-}$ data, but \sim FF assumptions as DSS

Comparison of FF sets

$$D_{u,d}^{h^+} - D_{u,d}^{h^-}$$

BRAHMS \longrightarrow higher FF

(HKNS: negative)

HKNS: no $d\sigma^{h^+} - d\sigma^{h^-}$ data, but \sim FF assumptions as DSS

Outline

- 1 Single hadron inclusive production and pQCD
 - Introduction
 - Measurements of single hadron inclusive production
 - Outline of factorization
- 2 FFs at large x
 - Cross section in the fixed order approach
 - Global fits
 - Predictions using FFs from e^+e^-
 - Theoretical improvements
 - Latest global fit: AKK08 (Nucl. Phys. B)
- 3 Improving the small x region
 - Soft gluon logarithms in splitting functions
 - Unified approach
 - Fits to data

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

Type	FFs	PDFs
$\begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{NLO term}} = a_s^2 P^{(1)}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{\ln^2 z}{z} \\ \frac{1}{z} & \frac{\ln^2 z}{z} \end{pmatrix}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{1}{z} \\ \frac{1}{z} & \frac{1}{z} \end{pmatrix}$
pert. approach	$x > 0.1$	$x_B > 10^{-4}$

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

Type	FFs	PDFs
$\begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{NLO term}} = a_s^2 P^{(1)}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{\ln^2 z}{z} \\ \frac{1}{z} & \frac{\ln^2 z}{z} \end{pmatrix}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{1}{z} \\ \frac{1}{z} & \frac{1}{z} \end{pmatrix}$
pert. approach	$x > 0.1$	$x_B > 10^{-4}$

P for FFs contains **double logs**

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

Type	FFs	PDFs
$\begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{NLO term}} = a_s^2 P^{(1)}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{\ln^2 z}{z} \\ \frac{1}{z} & \frac{\ln^2 z}{z} \end{pmatrix}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{1}{z} \\ \frac{1}{z} & \frac{1}{z} \end{pmatrix}$
pert. approach	$x > 0.1$	$x_B > 10^{-4}$

P for FFs contains **double logs**

At $O(a_s^n)$, DL has form $(1/z)(a_s \ln z)^2 (a_s \ln^2 z)^{n-1}$

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

Type	FFs	PDFs
$\begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{NLO term}} = a_s^2 P^{(1)}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{\ln^2 z}{z} \\ \frac{1}{z} & \frac{\ln^2 z}{z} \end{pmatrix}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{1}{z} \\ \frac{1}{z} & \frac{1}{z} \end{pmatrix}$
pert. approach	$x > 0.1$	$x_B > 10^{-4}$

P for FFs contains **double logs**

At $O(a_s^n)$, DL has form $(1/z)(a_s \ln z)^2 (a_s \ln^2 z)^{n-1}$

→ small x inaccurate at $\ln(1/x) \gtrsim O(a_s^{-1/2})$

Small x : FFs vs. PDFs

$$\text{Small } z: \quad \begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{LO term}} = a_s P^{(0)} = a_s \begin{pmatrix} 0 & \frac{1}{z} \\ 0 & \frac{1}{z} \end{pmatrix}$$

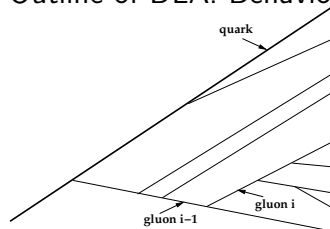
Type	FFs	PDFs
$\begin{pmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{pmatrix}_{\text{NLO term}} = a_s^2 P^{(1)}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{\ln^2 z}{z} \\ \frac{1}{z} & \frac{\ln^2 z}{z} \end{pmatrix}$	$a_s^2 \begin{pmatrix} \frac{1}{z} & \frac{1}{z} \\ \frac{1}{z} & \frac{1}{z} \end{pmatrix}$
pert. approach	$x > 0.1$	$x_B > 10^{-4}$

P for FFs contains **double logs**

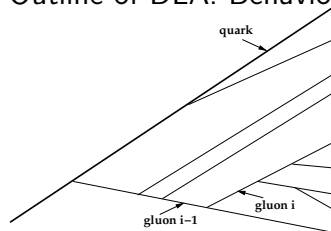
At $O(a_s^n)$, DL has form $(1/z)(a_s \ln z)^2 (a_s \ln^2 z)^{n-1} \left(\frac{a_s}{\omega} \left(\frac{a_s}{\omega^2} \right)^{n-1} \right)$

\rightarrow small x inaccurate at $\ln(1/x) \gtrsim O(a_s^{-1/2})$ ($\omega \lesssim O(\sqrt{a_s})$)

Outline of DLA: Behaviour at small momentum fractions

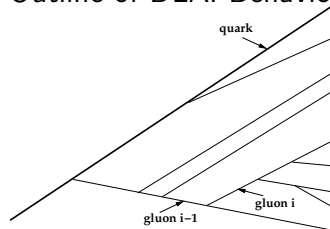


Outline of DLA: Behaviour at small momentum fractions



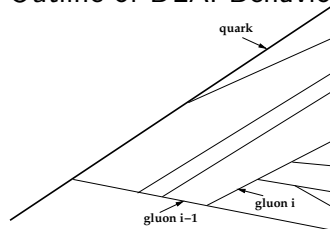
$$\bullet \mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$$

Outline of DLA: Behaviour at small momentum fractions



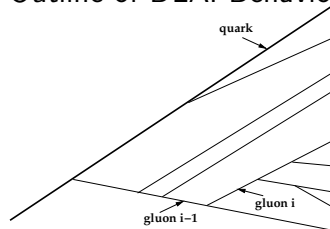
- $\mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$
- Strong ordering: $(E, \theta)_i \ll (E, \theta)_{i-1}$

Outline of DLA: Behaviour at small momentum fractions



- $\mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$
- Strong ordering: $(E, \theta)_i \ll (E, \theta)_{i-1}$
- Collinear part \in FFs:
 $E_T = E_{i-1} \theta_{i-1} > M_f$
 $\therefore E_i \theta_i > z' M_f$

Outline of DLA: Behaviour at small momentum fractions



- $\mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$

- Strong ordering: $(E, \theta)_i \ll (E, \theta)_{i-1}$

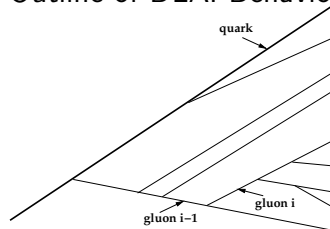
- Collinear part \in FFs:

$$E_T = E_{i-1} \theta_{i-1} > M_f$$

$$\therefore E_i \theta_i > z' M_f$$

$$\rightarrow \frac{d}{d \ln M_f^2} D(z, M_f^2) = \int_z^1 \frac{dz'}{z'} a_s(z'^2 M_f^2) P^{(0)}(z') D\left(\frac{z}{z'}, z'^2 M_f^2\right)$$

Outline of DLA: Behaviour at small momentum fractions

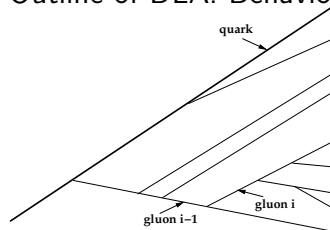


- $\mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$
- Strong ordering: $(E, \theta)_i \ll (E, \theta)_{i-1}$
- Collinear part \in FFs:
 $E_T = E_{i-1} \theta_{i-1} > M_f$
 $\therefore E_i \theta_i > \textcolor{red}{z}' M_f$

$$\rightarrow \frac{d}{d \ln M_f^2} D(z, M_f^2) = \int_z^1 \frac{dz'}{z'} a_s(\textcolor{red}{z}'^2 M_f^2) P^{(0)}(z') D\left(\frac{z}{z'}, \textcolor{red}{z}'^2 M_f^2\right)$$

In $P^{(0)}(z)$, only take small z sing $\simeq 2C_A A/z$ where $A^2 = A$

Outline of DLA: Behaviour at small momentum fractions



- $\mathcal{P}(G_{i-1} \rightarrow G_i + \dots) \propto a_s \frac{d\theta_i}{\theta_i} \frac{dE_i}{E_i} \simeq a_s \ln^2$
- Strong ordering: $(E, \theta)_i \ll (E, \theta)_{i-1}$
- Collinear part \in FFs:
 $E_T = E_{i-1} \theta_{i-1} > M_f$
 $\therefore E_i \theta_i > z' M_f$

$$\rightarrow \frac{d}{d \ln M_f^2} D(z, M_f^2) = \int_z^1 \frac{dz'}{z'} a_s(z'^2 M_f^2) P^{(0)}(z') D\left(\frac{z}{z'}, z'^2 M_f^2\right)$$

In $P^{(0)}(z)$, only take small z sing $\simeq 2C_A A/z$ where $A^2 = A$

\rightarrow DLA equation, contains all DLs

$$\frac{d}{d \ln M_f^2} D(z, M_f^2) = \int_z^1 \frac{dz'}{z'} \frac{2C_A}{z'} A a_s(z'^2 M_f^2) D\left(\frac{z}{z'}, z'^2 M_f^2\right)$$

DLA improved DGLAP (Albino et al. 2005/6)

DLA improved DGLAP (Albino et al. 2005/6)

- Use DGLAP equation *also* at small x

DLA improved DGLAP (Albino et al. 2005/6)

- Use DGLAP equation *also* at small x

- But $a_s P^{(0)}(z) \longrightarrow \underbrace{a_s \bar{P}^{(0)}(z)}_{a_s P^0 \text{ with DLs subtracted}} + \underbrace{P^{\text{DL}}(z, a_s)}_{\text{DLs to all orders}}$

→ Good approximation for P for large *and* small z

DLA improved DGLAP (Albino et al. 2005/6)

- Use DGLAP equation *also* at small x

- But $a_s P^{(0)}(z) \longrightarrow \underbrace{a_s \bar{P}^{(0)}(z)}_{a_s P^0 \text{ with DLs subtracted}} + \underbrace{P^{\text{DL}}(z, a_s)}_{\text{DLs to all orders}}$

→ Good approximation for P for large *and* small z

- P^{DL} (DL-resummed P) is small z behaviour of P from DLA

$$P^{\text{DL}}(\omega, a_s) = \frac{A}{4} \left(-\omega + \sqrt{\omega^2 + 16C_A a_s} \right)$$

DLA improved DGLAP (Albino et al. 2005/6)

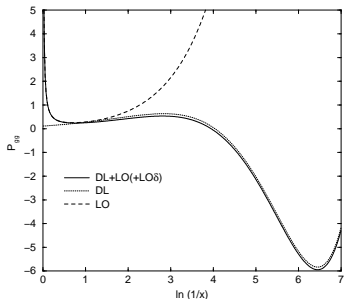
- Use DGLAP equation *also* at small x

- But $a_s P^{(0)}(z) \longrightarrow \underbrace{a_s \bar{P}^{(0)}(z)}_{a_s P^0 \text{ with DLs subtracted}} + \underbrace{P^{\text{DL}}(z, a_s)}_{\text{DLs to all orders}}$

→ Good approximation for P for large *and* small z

- P^{DL} (DL-resummed P) is small z behaviour of P from DLA

$$P^{\text{DL}}(\omega, a_s) = \frac{A}{4} \left(-\omega + \sqrt{\omega^2 + 16C_A a_s} \right)$$



DLA improved DGLAP (Albino et al. 2005/6)

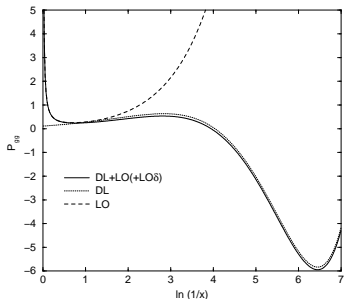
- Use DGLAP equation *also* at small x

- But $a_s P^{(0)}(z) \longrightarrow \underbrace{a_s \bar{P}^{(0)}(z)}_{a_s P^0 \text{ with DLs subtracted}} + \underbrace{P^{\text{DL}}(z, a_s)}_{\text{DLs to all orders}}$

→ Good approximation for P for large *and* small z

- P^{DL} (DL-resummed P) is small z behaviour of P from DLA

$$P^{\text{DL}}(\omega, a_s) = \frac{A}{4} \left(-\omega + \sqrt{\omega^2 + 16C_A a_s} \right)$$



- Approach definable \forall SGLs

DLA improved DGLAP (Albino et al. 2005/6)

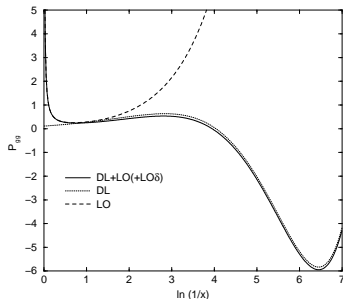
- Use DGLAP equation *also* at small x

- But $a_s P^{(0)}(z) \longrightarrow \underbrace{a_s \bar{P}^{(0)}(z)}_{a_s P^0 \text{ with DLs subtracted}} + \underbrace{P^{\text{DL}}(z, a_s)}_{\text{DLs to all orders}}$

→ Good approximation for P for large *and* small z

- P^{DL} (DL-resummed P) is small z behaviour of P from DLA

$$P^{\text{DL}}(\omega, a_s) = \frac{A}{4} \left(-\omega + \sqrt{\omega^2 + 16C_A a_s} \right)$$



- Approach definable \forall SGLs
- Consistent with MLLA

Procedure

Procedure

- Use hadron-unidentified data

Procedure

- Use hadron-unidentified data
- Large $M_0 \rightarrow$ assumptions for $D_i(z, M_0^2)$ from DLA:

Procedure

- Use hadron-unidentified data
- Large $M_0 \rightarrow$ assumptions for $D_i(z, M_0^2)$ from DLA:
 - Behaves like Gaussian in $\xi = \ln 1/z$ at small z

Procedure

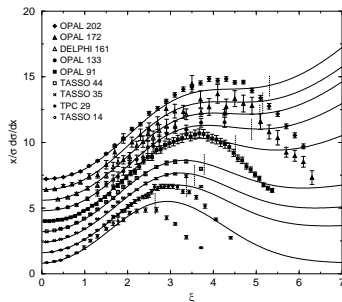
- Use hadron-unidentified data
- Large $M_0 \rightarrow$ assumptions for $D_i(z, M_0^2)$ from DLA:
 - Behaves like Gaussian in $\xi = \ln 1/z$ at small z
 - $\sum_q D_{q,\bar{q}} \simeq \frac{C_F}{C_A} D_g$

Procedure

- Use hadron-unidentified data
- Large $M_0 \rightarrow$ assumptions for $D_i(z, M_0^2)$ from DLA:
 - Behaves like Gaussian in $\xi = \ln 1/z$ at small z
 - $\sum_q D_{q,\bar{q}} \simeq \frac{C_F}{C_A} D_g$
- Fit Λ_{QCD} and m_h

Results

● $P = a_s P^{(0)}$ (old)



$$\chi^2_{\text{DF}} \simeq 3$$

$$\Lambda_{\text{QCD}} = 400 \text{ MeV}$$

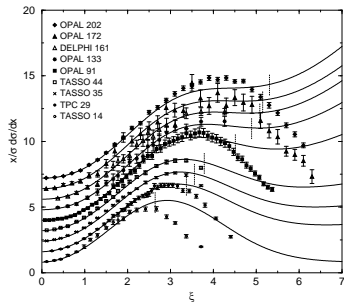
$$(\chi^2_{\text{DF}} \simeq 2)$$

$$(\Lambda_{\text{QCD}} = 1300 \text{ MeV})$$

$$(m_h \simeq 300 \text{ MeV})$$

Results

● $P = a_s P^{(0)}$ (old)



$$\chi^2_{DF} \simeq 3$$

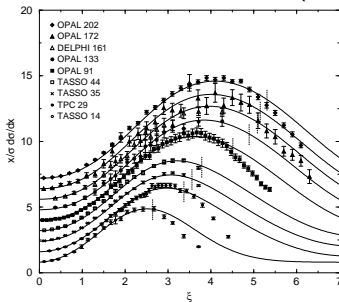
$$\Lambda_{QCD} = 400 \text{ MeV}$$

$$(\chi^2_{DF} \simeq 2)$$

$$(\Lambda_{QCD} = 1300 \text{ MeV})$$

$$(m_h \simeq 300 \text{ MeV})$$

● $P = a_s \bar{P}^{(0)} + P^{DL}$ (new)



$$\chi^2_{DF} \simeq 2$$

$$\Lambda_{QCD} = 800 \text{ MeV}$$

$$(\chi^2_{DF} \simeq 2)$$

$$(\Lambda_{QCD} = 400 \text{ MeV})$$

$$(m_h \simeq 300 \text{ MeV})$$

Summary

Summary

- FFs well constrained by exp data, particularly e^+e^-

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality
- Valence quark / gluon FFs weakly constrained by RHIC pp

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality
- Valence quark / gluon FFs weakly constrained by RHIC pp
- Large x resummation can make significant improvement

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality
- Valence quark / gluon FFs weakly constrained by RHIC pp
- Large x resummation can make significant improvement
- Fitted baryon masses good,
deviation of fitted masses may measure indirect fragmentation

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality
- Valence quark / gluon FFs weakly constrained by RHIC pp
- Large x resummation can make significant improvement
- Fitted baryon masses good,
deviation of fitted masses may measure indirect fragmentation
- Full treatment of exp data errors

Summary

- FFs well constrained by exp data, particularly e^+e^-
- Predictions for pp/ep data consistent with FF universality
- Valence quark / gluon FFs weakly constrained by RHIC pp
- Large x resummation can make significant improvement
- Fitted baryon masses good,
deviation of fitted masses may measure indirect fragmentation
- Full treatment of exp data errors
- Small x resummation within standard framework

Future

Future

- Theory improvements (heavy quarks, NNLO, small x logs)

Future

- Theory improvements (heavy quarks, NNLO, small x logs)
- Accurate data from HERA, RHIC, LEP, CLEO, Belle, BaBar
- . . .

Future

- Theory improvements (heavy quarks, NNLO, small x logs)
- Accurate data from HERA, RHIC, LEP, CLEO, Belle, BaBar
- . . .
- Development of FF error determination