

Axion/ALPs in Astrophysics and Cosmology.

Andreas Ringwald

SFB Lecture
DESY
Hamburg, D
7 July 2017

Reminder

- > At very low energies, axion/ALP interactions with photons, electrons, nucleons described by

$$\mathcal{L} = \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} m_a^2 a^2 - \frac{\alpha}{8\pi} \frac{C_{a\gamma}}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} \frac{C_{af}}{f_a} \partial_\mu a \bar{\psi}_f \gamma^\mu \gamma_5 \psi_f$$

- > In case of axion:

$$m_A = 57.0(7) \left(\frac{10^{11} \text{ GeV}}{f_A} \right) \mu\text{eV}$$

$$C_{A\gamma} = \frac{E}{N} - 1.92(4)$$

$$C_{Ap} = -0.47(3) + 0.88(3)C_{Au} - 0.39(2)C_{Ad} - 0.038(5)C_{As} \\ - 0.012(5)C_{Ac} - 0.009(2)C_{Ab} - 0.0035(4)C_{At},$$

$$C_{An} = -0.02(3) + 0.88(3)C_{Ad} - 0.39(2)C_{Au} - 0.038(5)C_{As} \\ - 0.012(5)C_{Ac} - 0.009(2)C_{Ab} - 0.0035(4)C_{At}$$

- Electron coupling very model-dependent

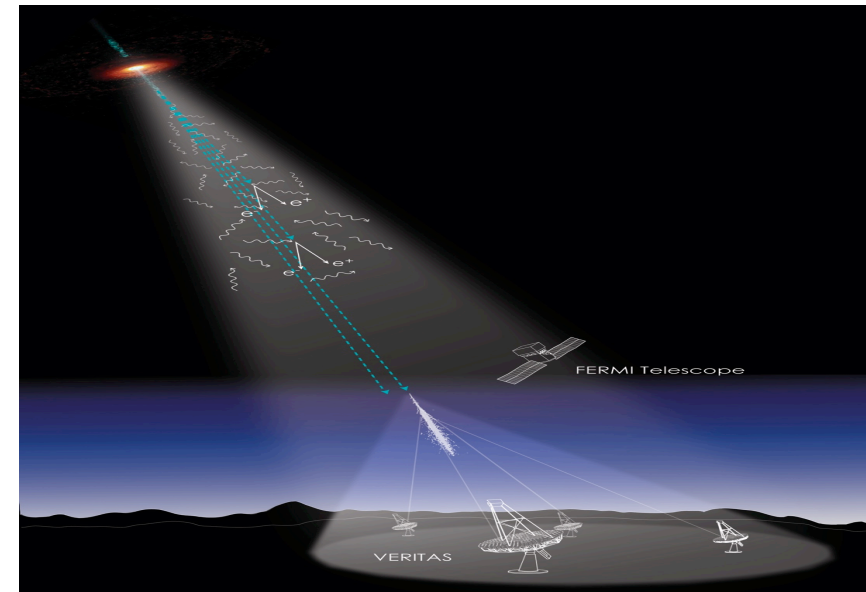
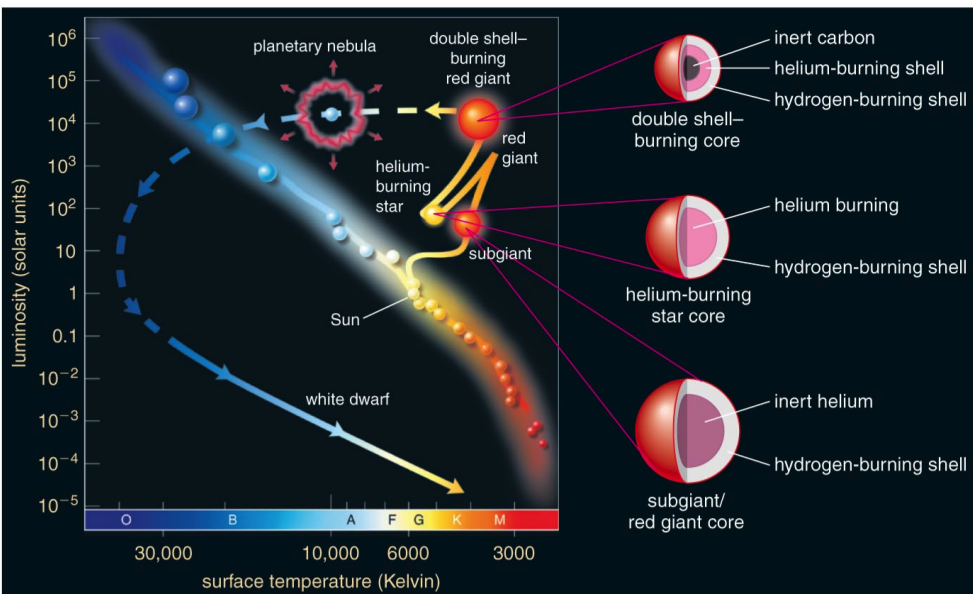
- > In case of ALP: mass and couplings very model-dependent



Astrophysical Signatures for Axion/ALPs

> Axion/ALPs may

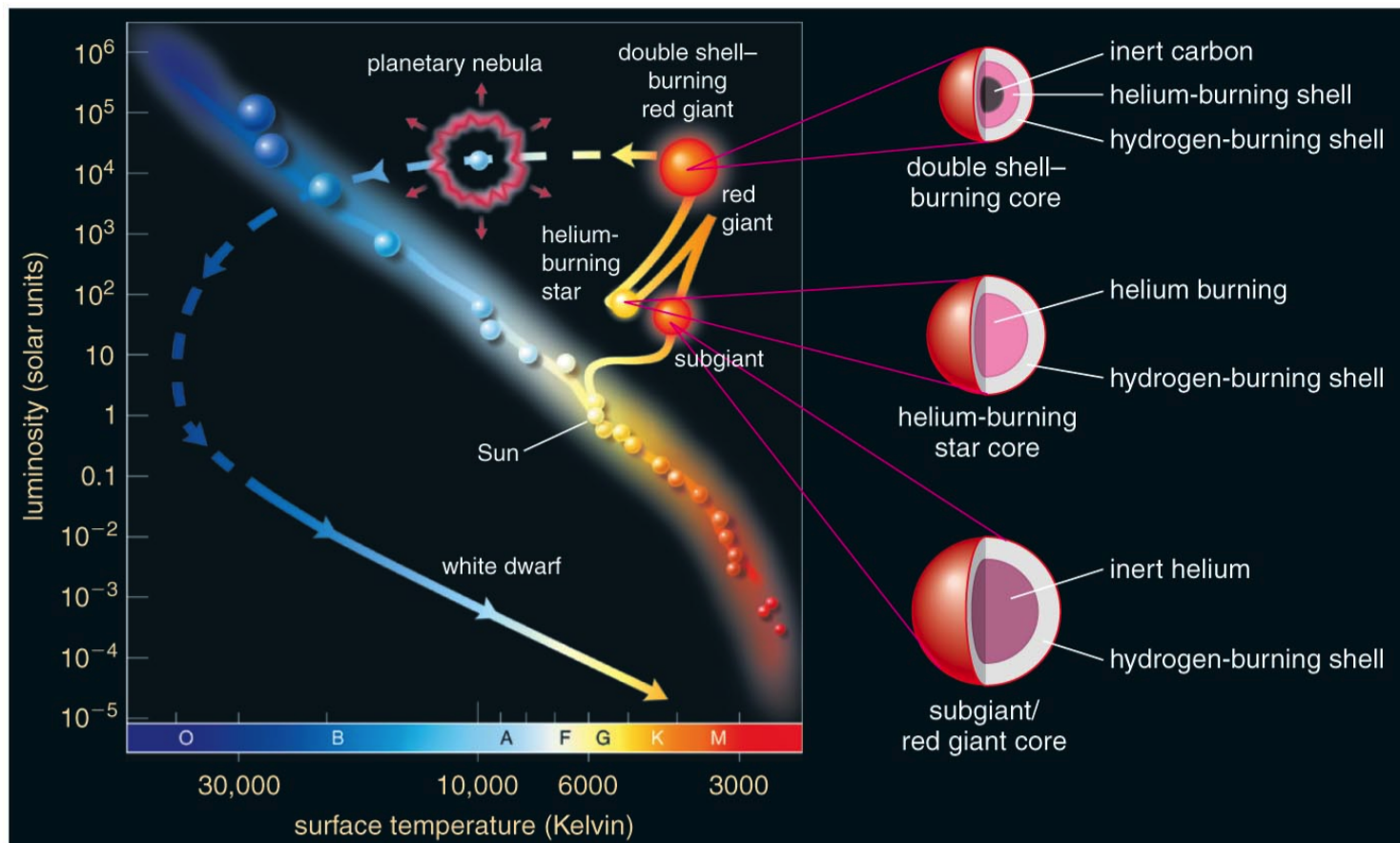
- lead to excessive energy losses of stars in various evolutionary stages
- may convert to photons (or vice-versa) in astrophysical magnetic fields



[Copyright Addison Wesley]

Searching for Axion/ALP Energy Losses of Stars

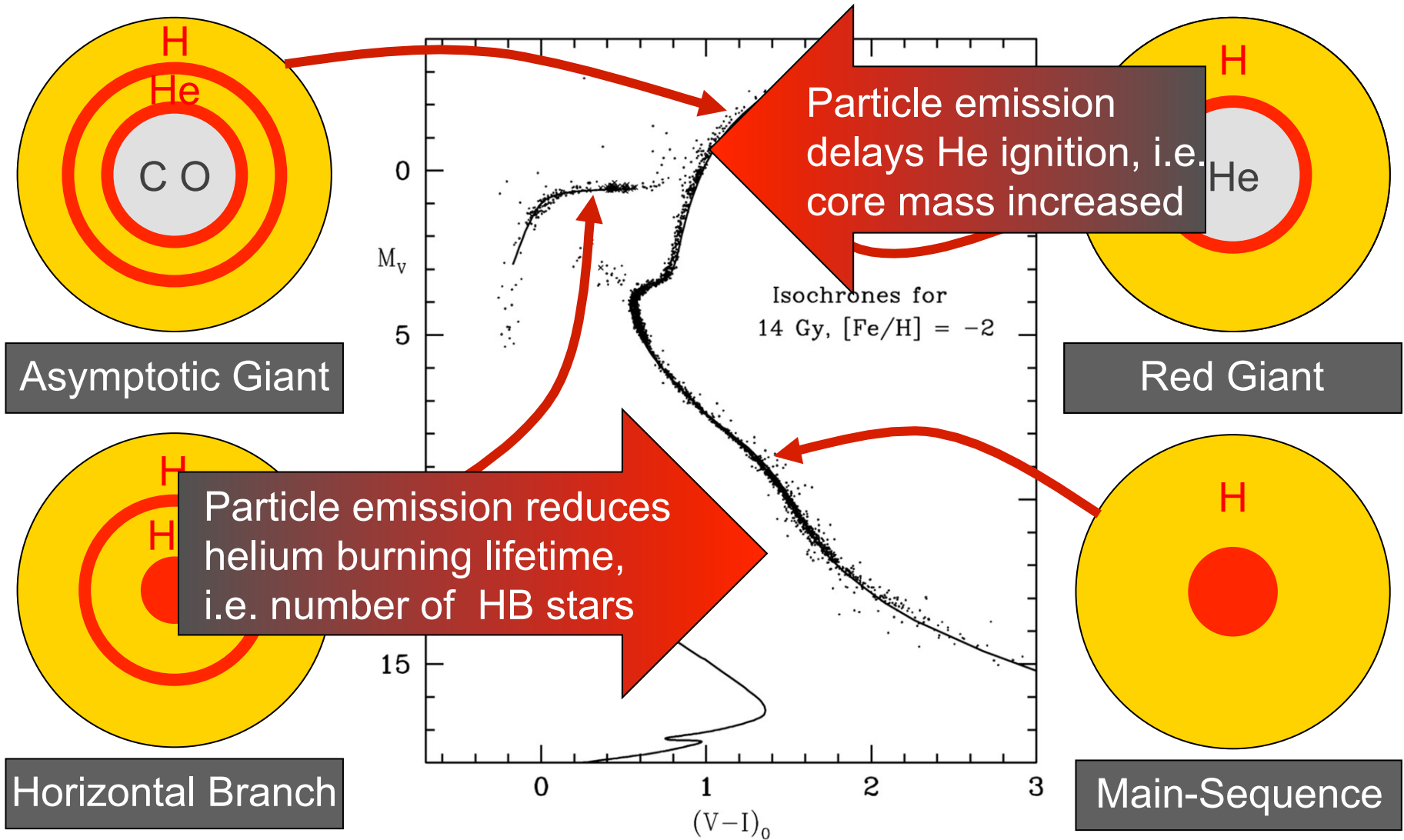
- Evolution of stars (Main Sequence – Red-Giant (RG) – Helium Burning (HB) – White Dwarf (WD)) sensitive to additional energy losses



[Copyright Addison Wesley]



Searching for Axion/ALP Energy Losses of Stars

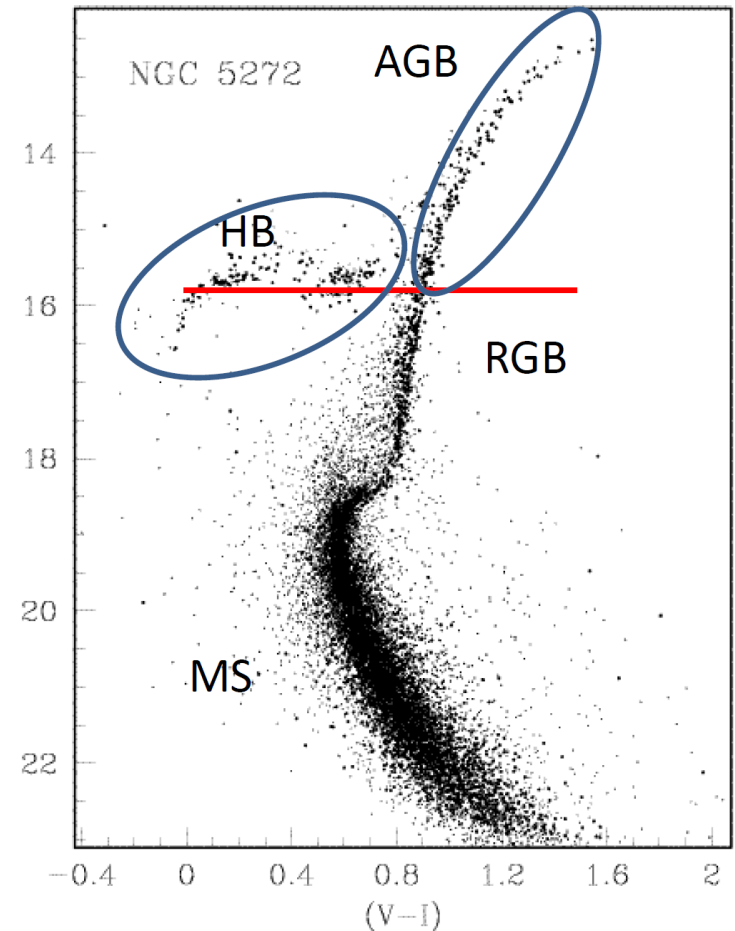


[Raffelt 14]



Searching for Axion/ALP Energy Losses of Stars

- > RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster
[Viaux et al. 13]
- > HB cooling rate: Number of HB stars vs. number of RGs in color-magnitude diagram of globular cluster
[Ayala et al. 14]



[Giannotti '16]

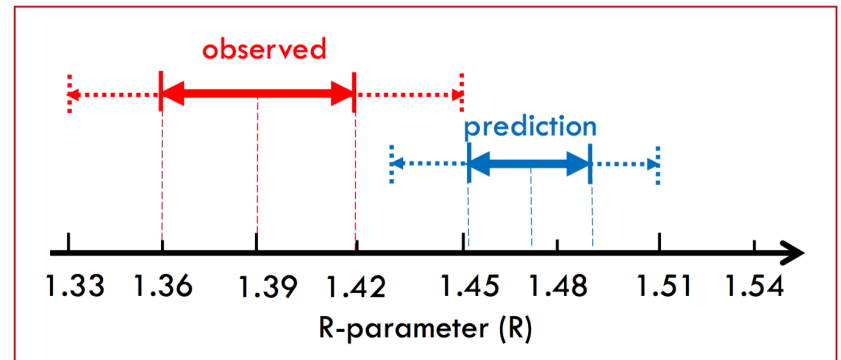
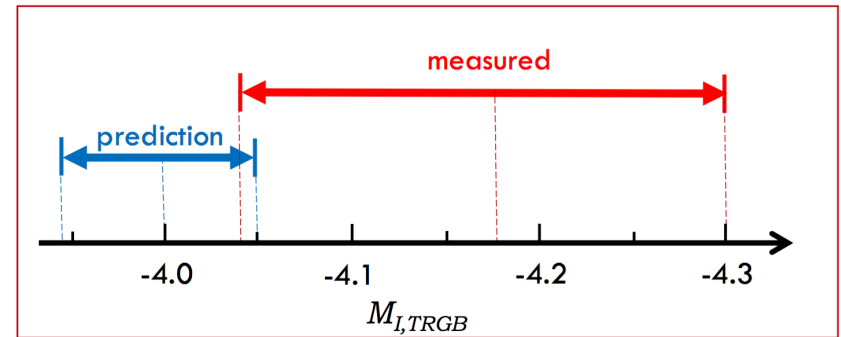
Searching for Axion/ALP Energy Losses of Stars

- > RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster

[Viaux et al. 13]

- > HB cooling rate: Number of HB stars vs. number of RGs in color-magnitude diagram of globular cluster

[Ayala et al. 14]



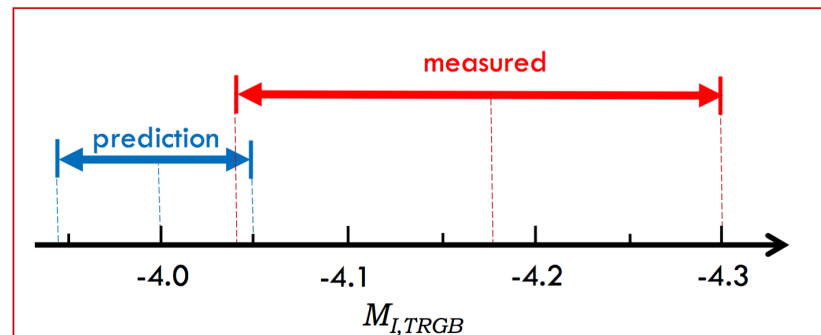
[Giannotti 15]



Searching for Axion/ALP Energy Losses of Stars

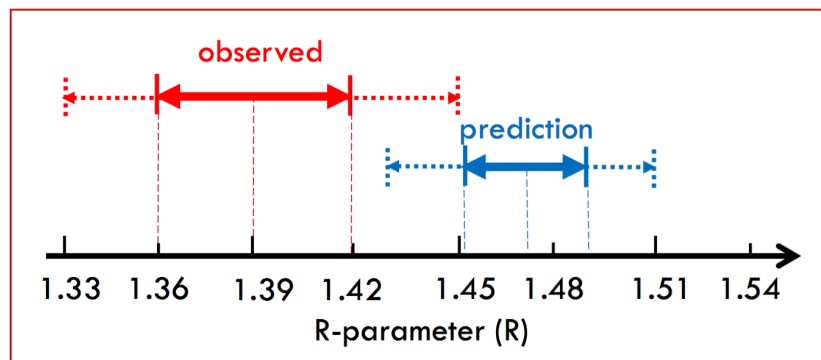
- > RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster

[Viaux et al. 13]



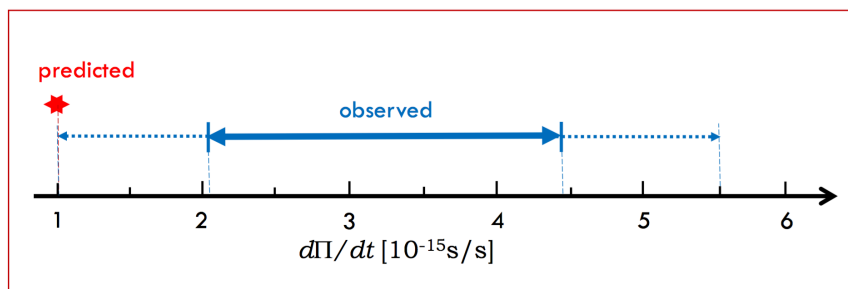
- > HB cooling rate: Number of HB stars vs. number of RGs in color-magnitude diagram of globular cluster

[Ayala et al. 14]



- > WD cooling rate:

- Period decrease of variable WDs [Kepler et al. 91,...]

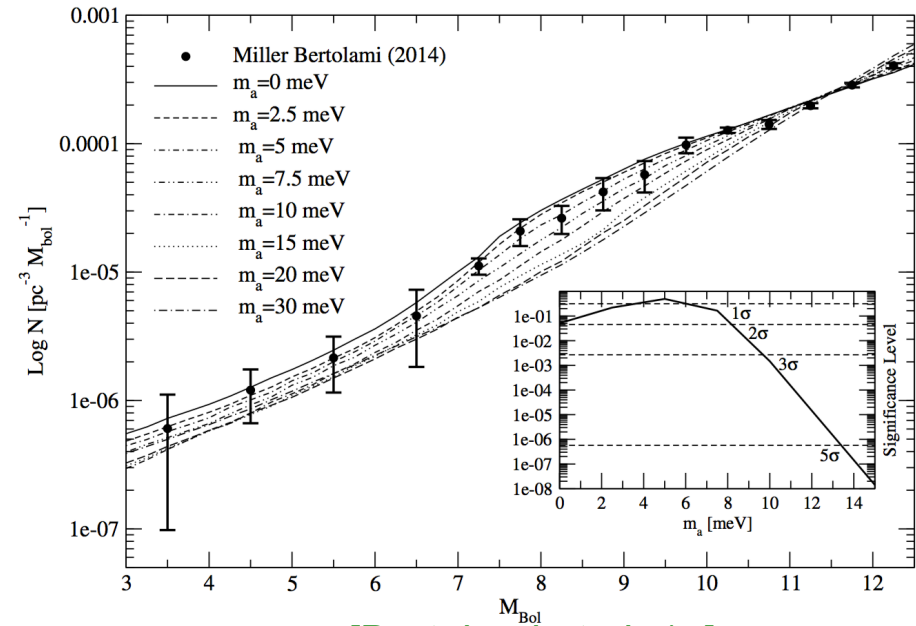


[Giannotti 15]



Searching for Axion/ALP Energy Losses of Stars

- > RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster
[Viaux et al. 13]
- > HB cooling rate: Number of HB stars vs. number of RGs in color-magnitude diagram of globular cluster
[Ayala et al. 14]
- > WD cooling rate:
 - Period decrease of variable WDs
[Kepler et al. 91,...]
 - White dwarf luminosity function (WDLF)
[Isern et al. 08-12]



[Bertolami et al. 15]



Searching for Axion/ALP Energy Losses of Stars

- > RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster

[Viaux et al. 13]

$$g_{ae} \equiv \frac{m_e}{f_a} |C_{ae}| < 4.3 \times 10^{-13}$$

- > HB cooling rate: Number of HB stars vs. number of RGs in color-magnitude diagram of globular cluster

[Ayala et al. 14]

$$g_{a\gamma} \equiv \frac{\alpha}{2\pi f_a} |C_{a\gamma}| < 6.6 \times 10^{-11} \text{ GeV}^{-1}$$

- > WD cooling rate:

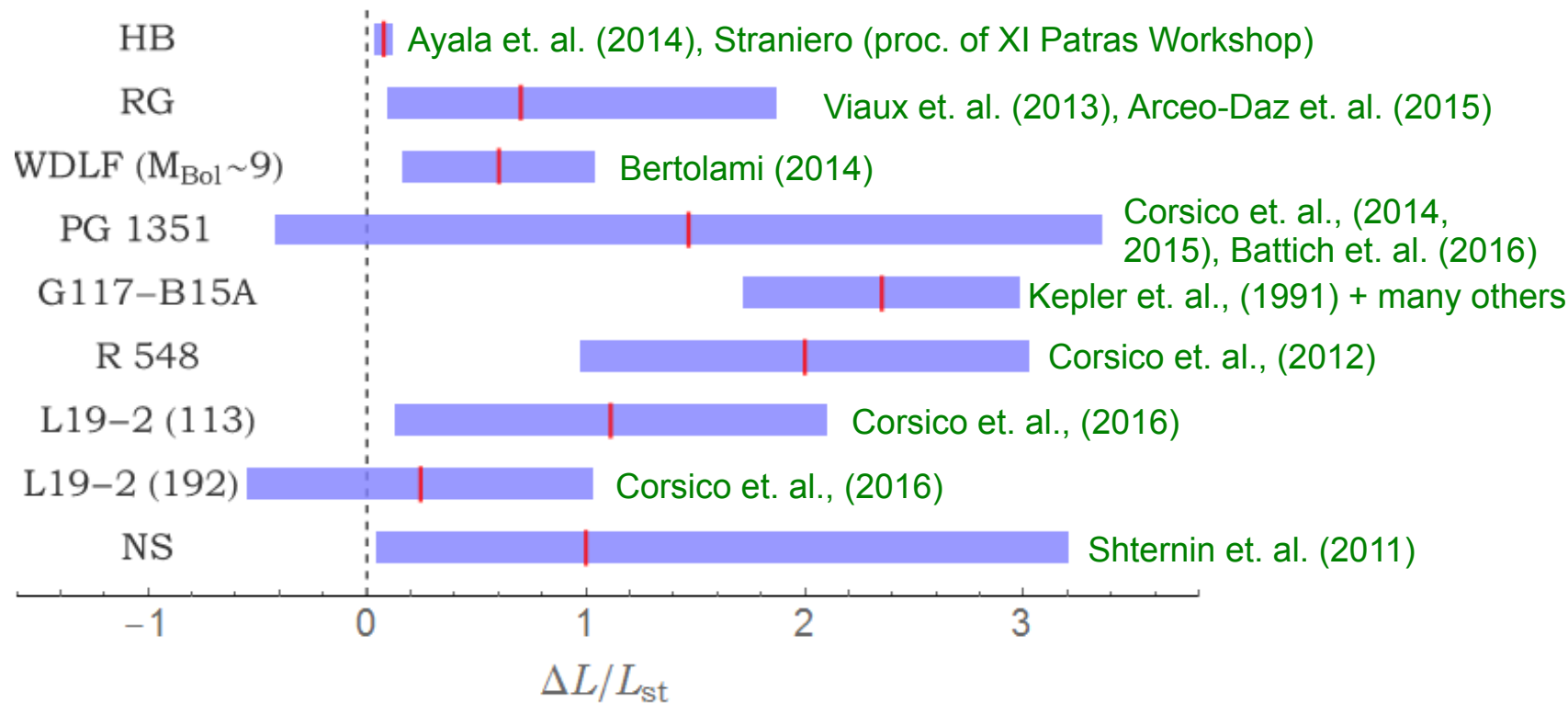
- Period decrease of variable WDs [Kepler et al. 91,...]
- White dwarf luminosity function (WDLF) [Isern et al. 08-12]

$$g_{ae} \equiv \frac{m_e}{f_a} |C_{ae}| < 2.8 \times 10^{-13}$$



Searching for Axion/ALP Energy Losses of Stars

- However, practically every stellar systems seems to be cooling a bit faster than predicted by models based on SM:

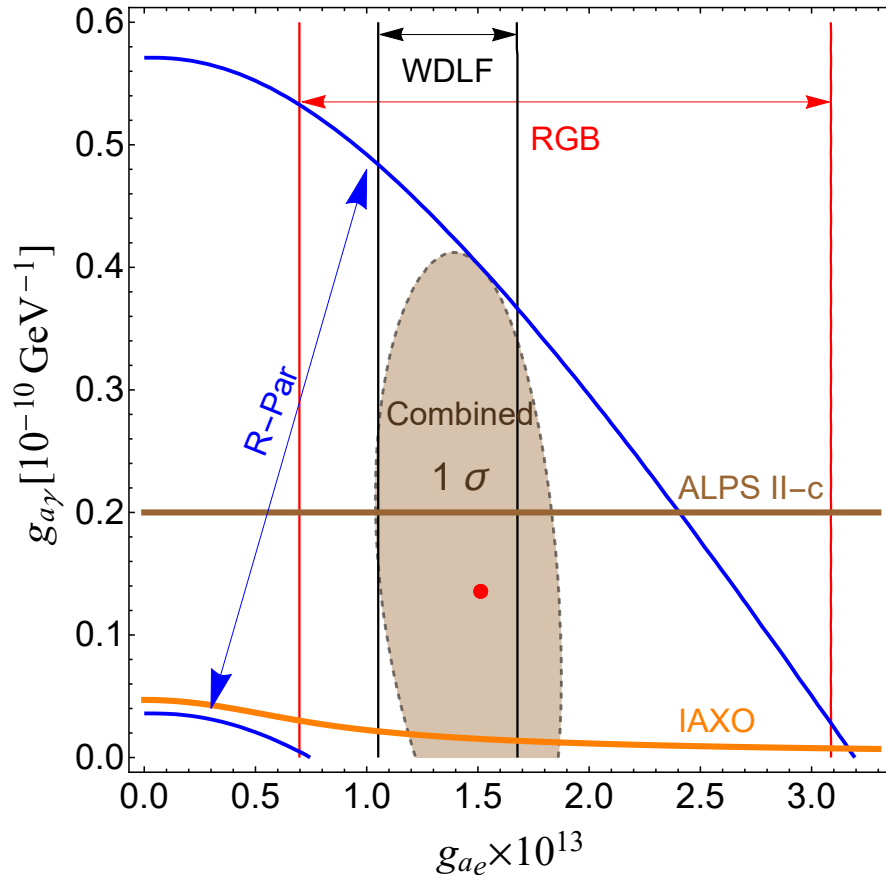


[Giannotti, Irastorza, Redondo, AR (2015); Giannotti, Irastorza, Redondo, AR (in preparation)]

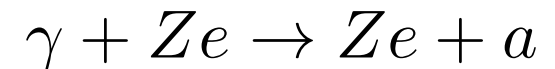


Searching for Axion/ALP Energy Losses of Stars

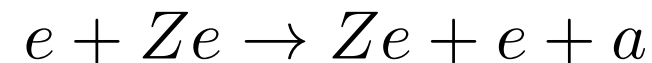
- Excessive energy losses of HBs, RG, WDs can be explained at one stroke by production of axion/ALP with coupling to photons and electrons:



$$g_{a\gamma} = C_{a\gamma} \alpha / (2\pi f_a)$$



$$g_{ai} = C_{ai} m_i / f_a$$

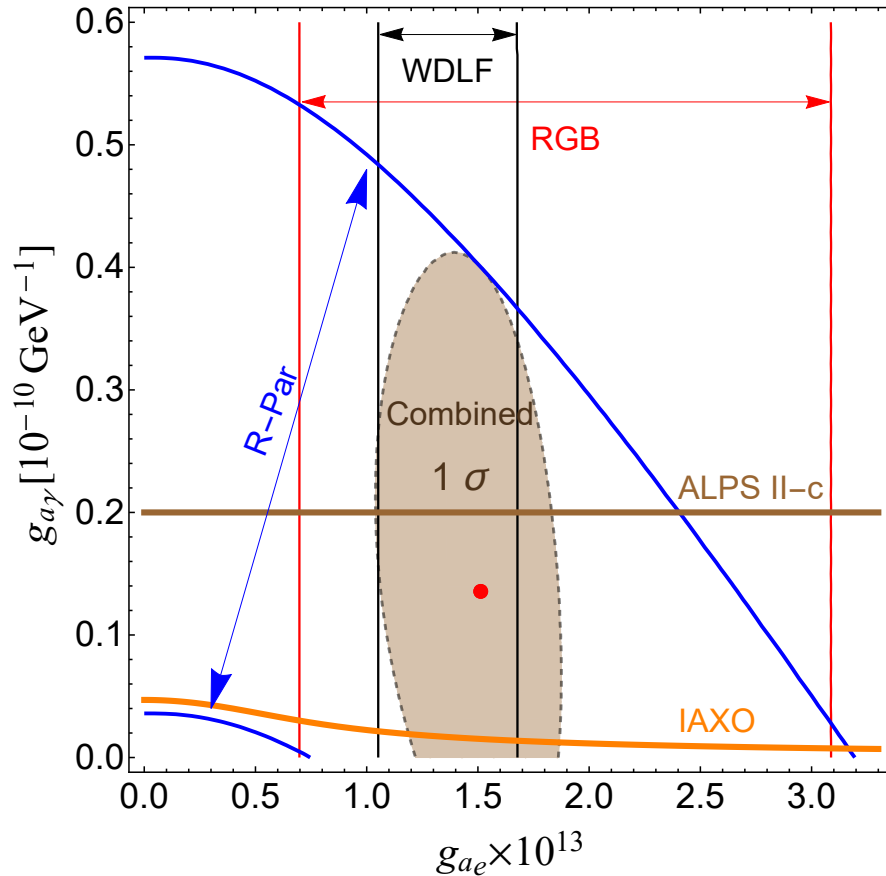


[Giannotti, Irastorza, Redondo, AR, Saikawa (in preparation)]

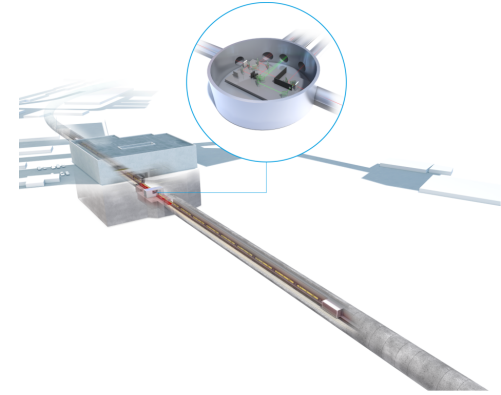


Searching for Axion/ALP Energy Losses of Stars

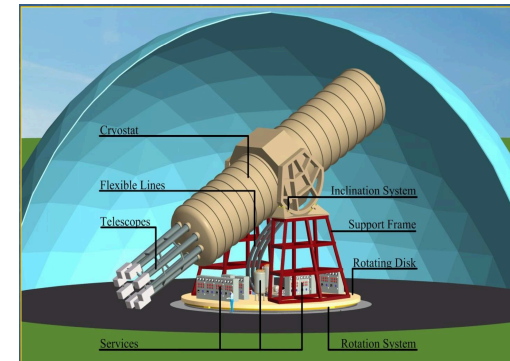
- Excessive energy losses of HBs, RG, WDs can be explained at one stroke by production of axion/ALP with coupling to photons and electrons:



ALPS II



IAXO

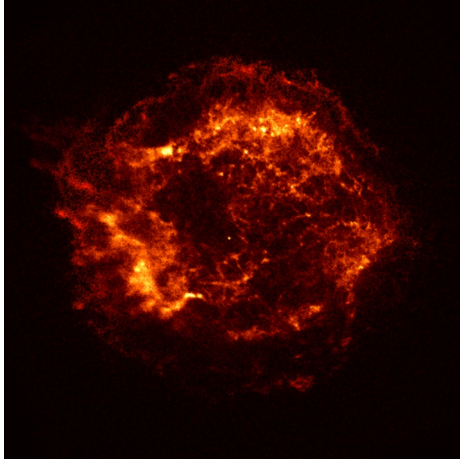


[Giannotti,Irastorza,Redondo,AR,Saikawa (in preparation)]

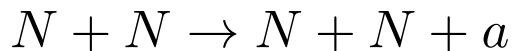


Searching for Axion/ALP Energy Losses of Stars

> Neutron star in Cas A:

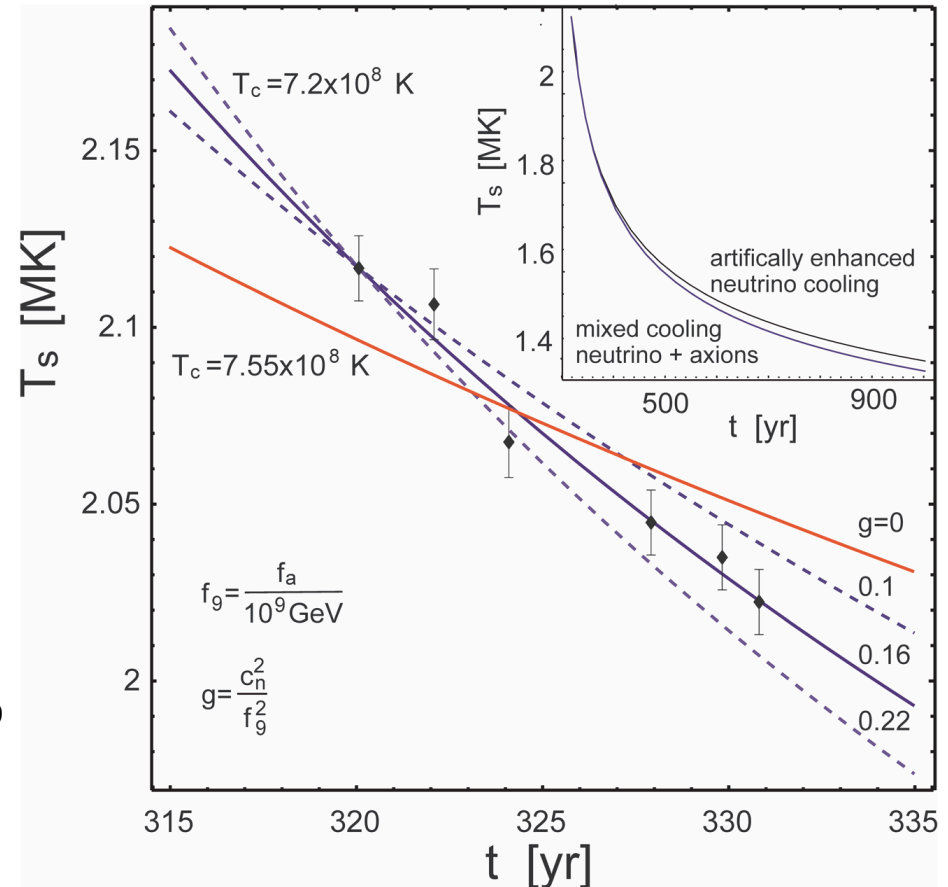


- Measured surface temperature reveals unusually fast cooling rate
- Hint on extra cooling by axion/ALP due to nucleon bremsstrahlung



- Required coupling to neutron:

$$g_{an} \equiv \frac{m_n}{f_a} |C_{an}| = (3.8 \pm 3) \times 10^{-10}$$



[Leinson 14]



Searching for Axion/ALP Energy Losses of Stars

> SN 1987 A:

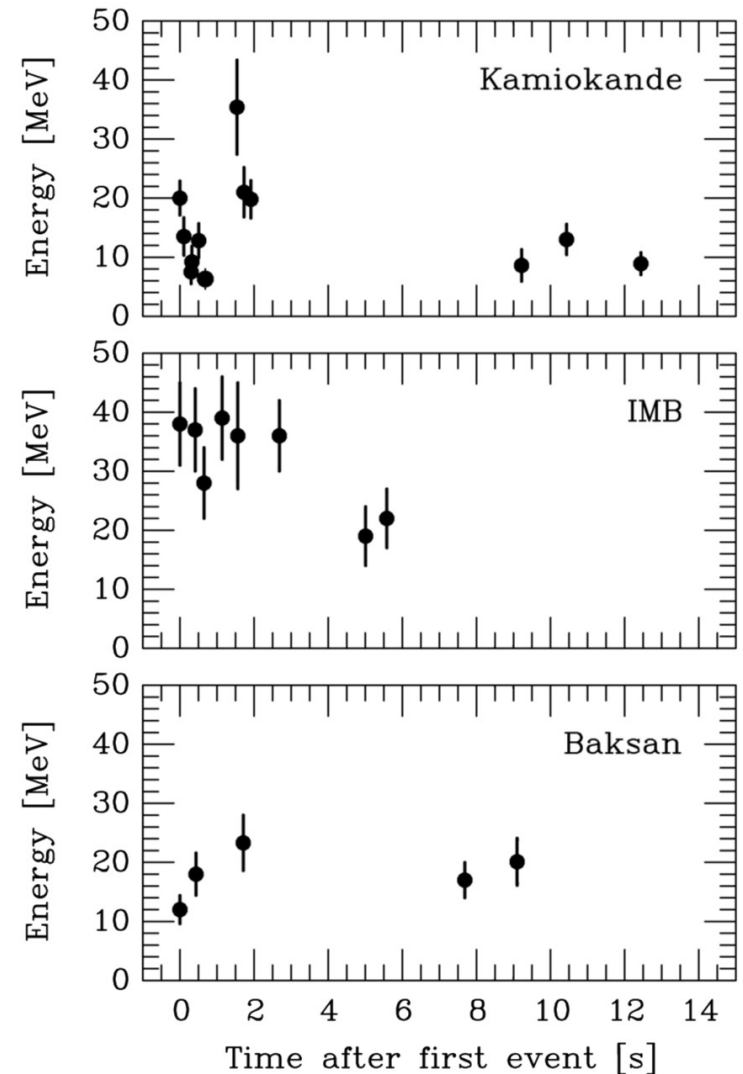


- Emission of axion/ALPs would take away energy from neutrino burst and shorten it

$$g_{ap}^2 + g_{an}^2 < 3.6 \times 10^{-19}$$

[Raffelt 08; Fischer et al. 16; Giannotti et al. (in prep.)]

- Not very solid bound: sparse data and interaction difficult to model



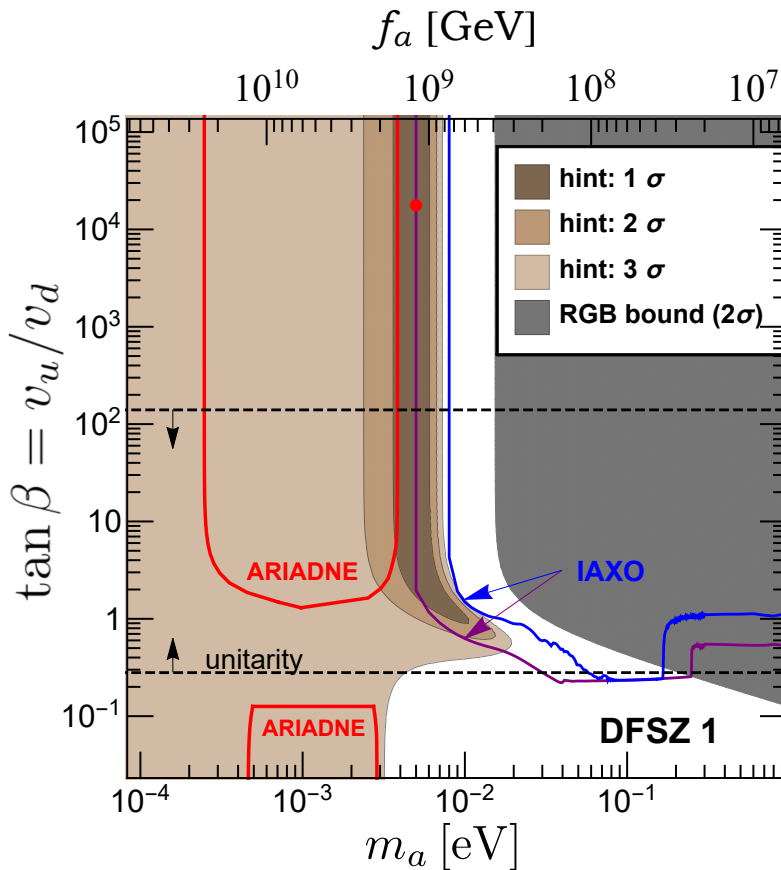
[Raffelt]



Searching for Axion/ALP Energy Losses of Stars

- Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for DFSZ axion model:

$$f_a = \frac{v_{PQ}}{6},$$



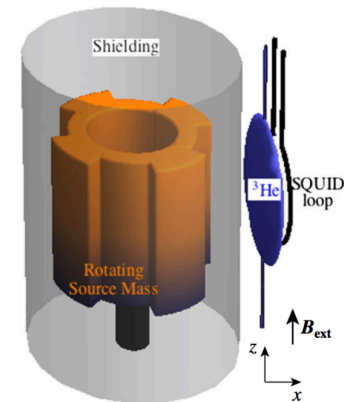
$$C_{ae}^{\text{DFSZ I}} = \frac{1}{3} \sin^2 \beta, \quad C_{ae}^{\text{DFSZ II}} = \frac{1}{3} (1 - \sin^2 \beta)$$

$$C_{a\gamma}^{\text{DFSZ I}} = \frac{8}{3} - 1.92(4), \quad C_{a\gamma}^{\text{DFSZ II}} = \frac{2}{3} - 1.92(4),$$

$$C_{Ap} = -0.435 \sin^2 \beta + (-0.182 \pm 0.025),$$

$$C_{An} = 0.414 \sin^2 \beta + (-0.16 \pm 0.025).$$

ARIADNE



[Giannotti,Irastorza,Redondo,AR,Saikawa (in preparation)]



Searching for Axion/ALP Energy Losses of Stars

- Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for DFSZ axion model:

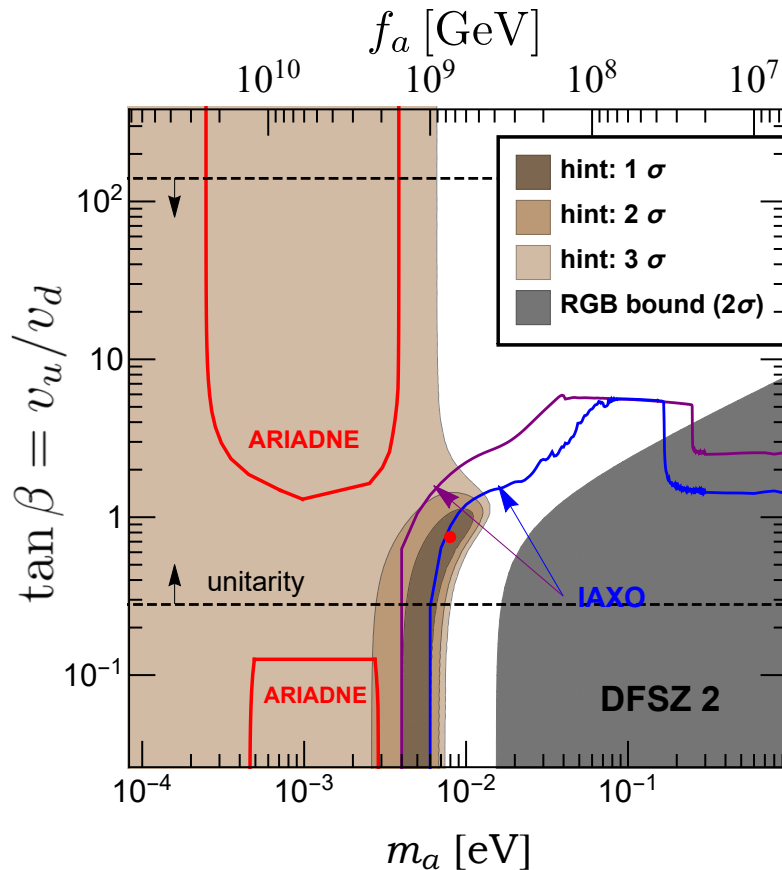
$$f_a = \frac{v_{PQ}}{6},$$

$$C_{ae}^{\text{DFSZ I}} = \frac{1}{3} \sin^2 \beta, \quad C_{ae}^{\text{DFSZ II}} = \frac{1}{3} (1 - \sin^2 \beta)$$

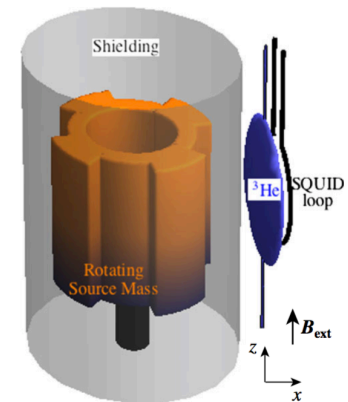
$$C_{a\gamma}^{\text{DFSZ I}} = \frac{8}{3} - 1.92(4), \quad C_{a\gamma}^{\text{DFSZ II}} = \frac{2}{3} - 1.92(4),$$

$$C_{Ap} = -0.435 \sin^2 \beta + (-0.182 \pm 0.025),$$

$$C_{An} = 0.414 \sin^2 \beta + (-0.16 \pm 0.025).$$



ARIADNE

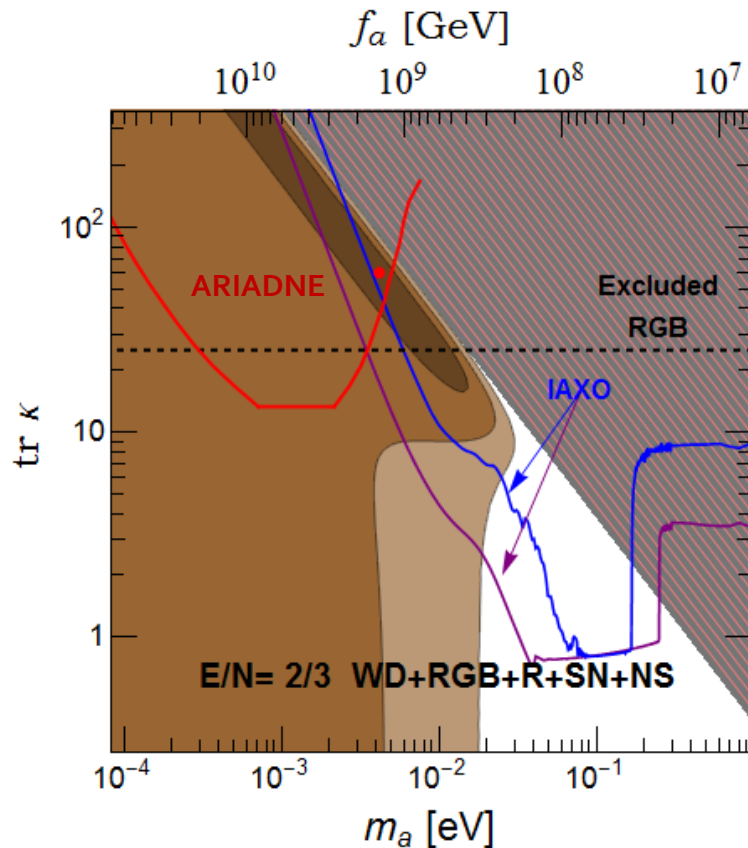


[Giannotti,Irastorza,Redondo,AR,Saikawa (in preparation)]



Searching for Axion/ALP Energy Losses of Stars

- Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for KSVZ axion/majoron model (as in SMASH):



$$\mathcal{L}_Y \supset -\bar{L}_i F_{ij} N_j H - \frac{1}{2} \bar{N}_i^c Y_{ij} N_j \sigma + \text{h.c.}$$

$$C_{ae} \simeq \frac{1}{8\pi^2 N} \left(\kappa_{ee} - \frac{1}{2} \text{tr} \kappa \right)$$

$$C_{aq} \simeq \frac{1}{8\pi^2 N} T_3^q \text{tr} \kappa$$

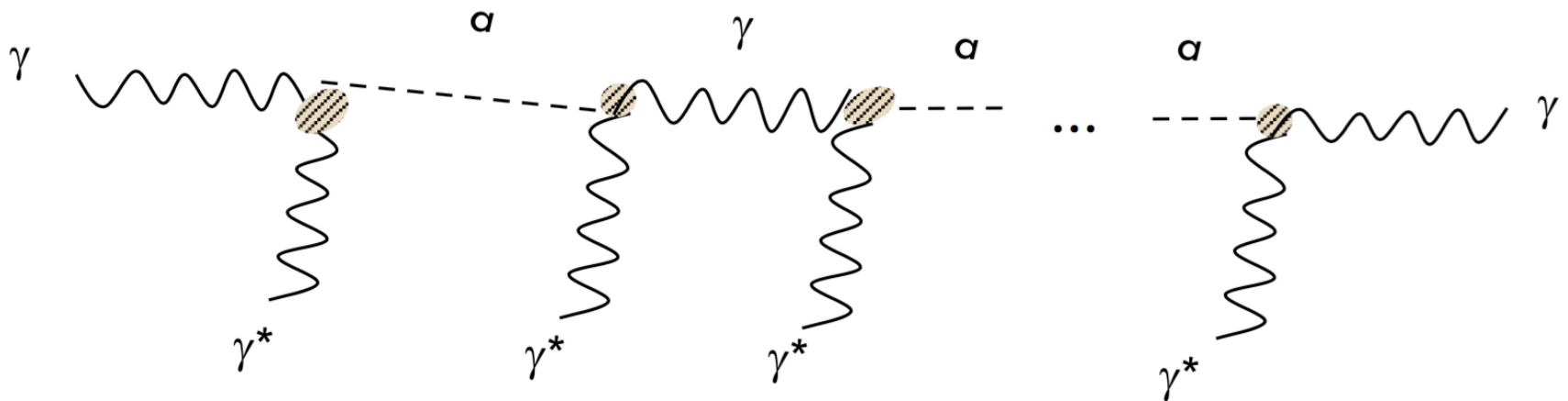
$$\kappa \equiv \frac{m_D m_D^\dagger}{v^2} = \frac{F F^\dagger}{2}$$

[Giannotti,Irastorza,Redondo,AR,Saikawa (in preparation)]



Photon – ALP Conversion in Cosmic Magnetic Fields?

- Photons propagating in magnetic field experience photon-ALP oscillations:



Photon – ALP Conversion in Cosmic Magnetic Fields?

- Photons propagating in magnetic field experience photon-ALP oscillations:

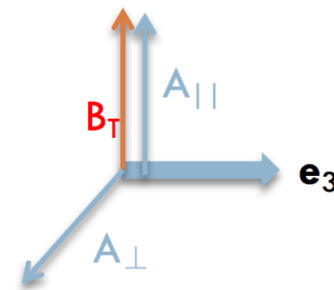
$$\left[E - i \frac{\partial}{\partial x_3} + \begin{pmatrix} \Delta_{\text{pl}} + \Delta_{\text{CMB}} - \frac{i\Gamma_{\text{abs}}}{2} & 0 & 0 \\ 0 & \Delta_{\text{pl}} + \Delta_{\text{CMB}} - i \frac{i\Gamma_{\text{abs}}}{2} & \Delta_{a\gamma} \\ 0 & \Delta_{a\gamma} & \Delta_a \end{pmatrix} \right] \begin{pmatrix} A_{\perp} \\ A_{\parallel} \\ a \end{pmatrix} = 0$$

▶ $\Delta_{a\gamma} \equiv \frac{g_{a\gamma} B_T}{2} \simeq 1.52 \times 10^{-8} \left(\frac{g_{a\gamma}}{10^{-17} \text{GeV}^{-1}} \right) \left(\frac{B_T}{10^{-9} \text{G}} \right) \text{Mpc}^{-1}$

▶ $\Delta_a \equiv -\frac{m_a^2}{2E} \simeq -7.8 \times 10^5 \left(\frac{m_a}{10^{-10} \text{eV}} \right)^2 \left(\frac{E}{\text{keV}} \right)^{-1} \text{Mpc}^{-1}$

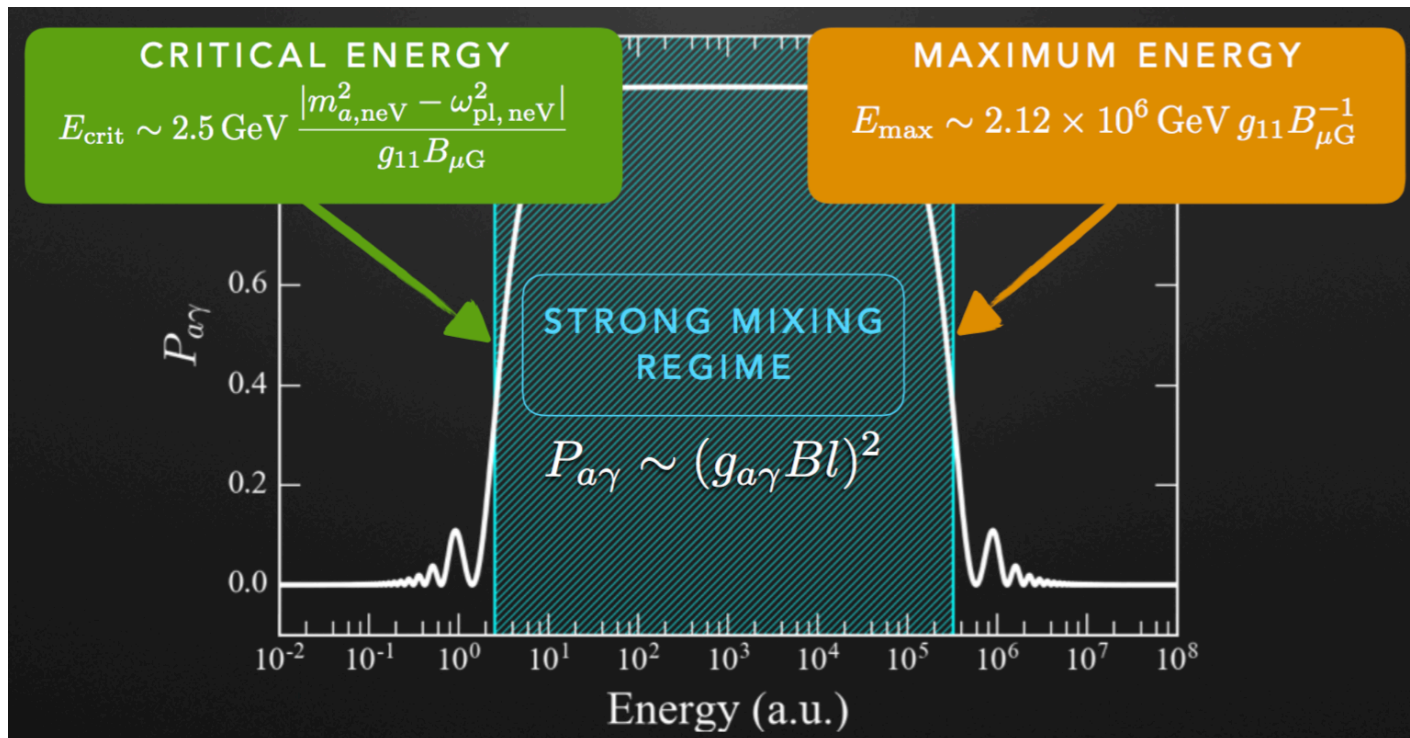
▶ $\Delta_{\text{pl}} \equiv -\frac{\omega_{\text{pl}}^2}{2E} \simeq -1.1 \times 10^{-2} \left(\frac{E}{\text{keV}} \right)^{-1} \left(\frac{n_e}{10^{-7} \text{cm}^{-3}} \right) \text{Mpc}^{-1}$

▶ $\Delta_{\text{CMB}} \simeq 0.80 \times 10^{-10} \left(\frac{E}{\text{keV}} \right) \text{Mpc}^{-1}$



Photon – ALP Conversion in Cosmic Magnetic Fields?

- Photons propagating in magnetic field experience photon-ALP oscillations:

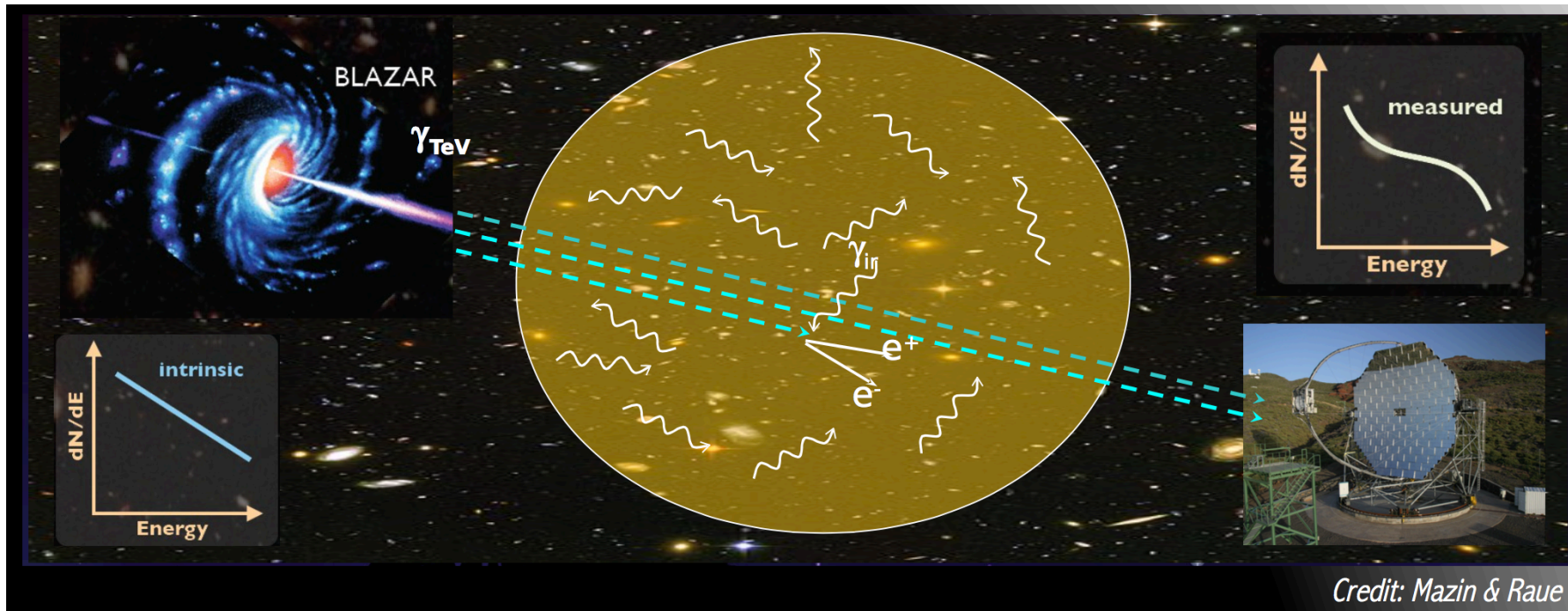


- Signatures for ALPs in AGN spectra:

- Apparent boost in photon flux (reduced opacity of Universe)
- Apparent spectral irregularities

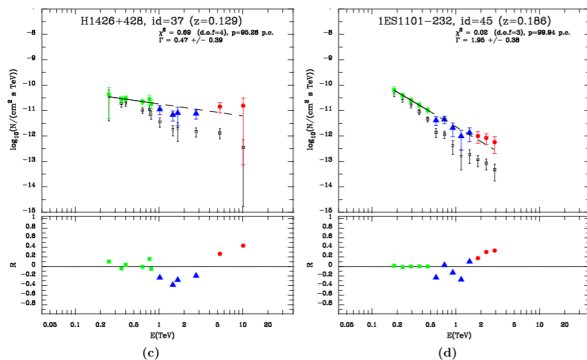
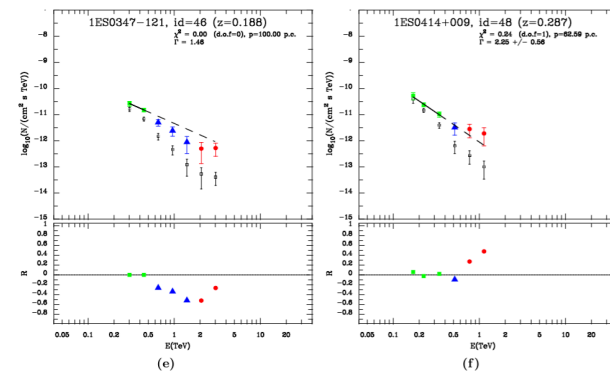
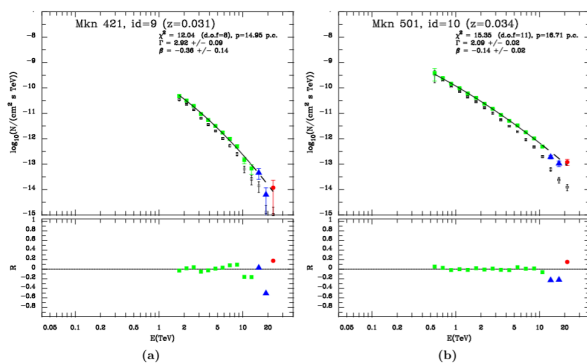
Photon – ALP Conversion in Cosmic Magnetic Fields?

- Gamma ray spectra from distant AGNs should show an energy and redshift dependent exponential attenuation, due to pair production at Extragalactic Background Light (EBL)

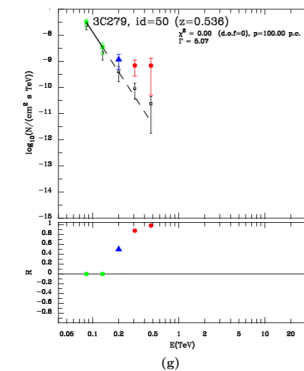


Photon – ALP Conversion in Cosmic Magnetic Fields?

- Gamma ray spectra from distant AGNs should show an energy and redshift dependent exponential attenuation, due to pair production at Extragalactic Background Light (EBL)
- Indication of anomalous gamma transparency: attenuation observed by **IACT** and **Fermi-LAT** too small [Aharonian et al. 07; de Angelis, Roncadelli et al. 07; ...; Horns, Meyer 12; Meyer, Horns, Raue 13; Rubtsov, Troitsky 14; Kohri, Kodama 17]



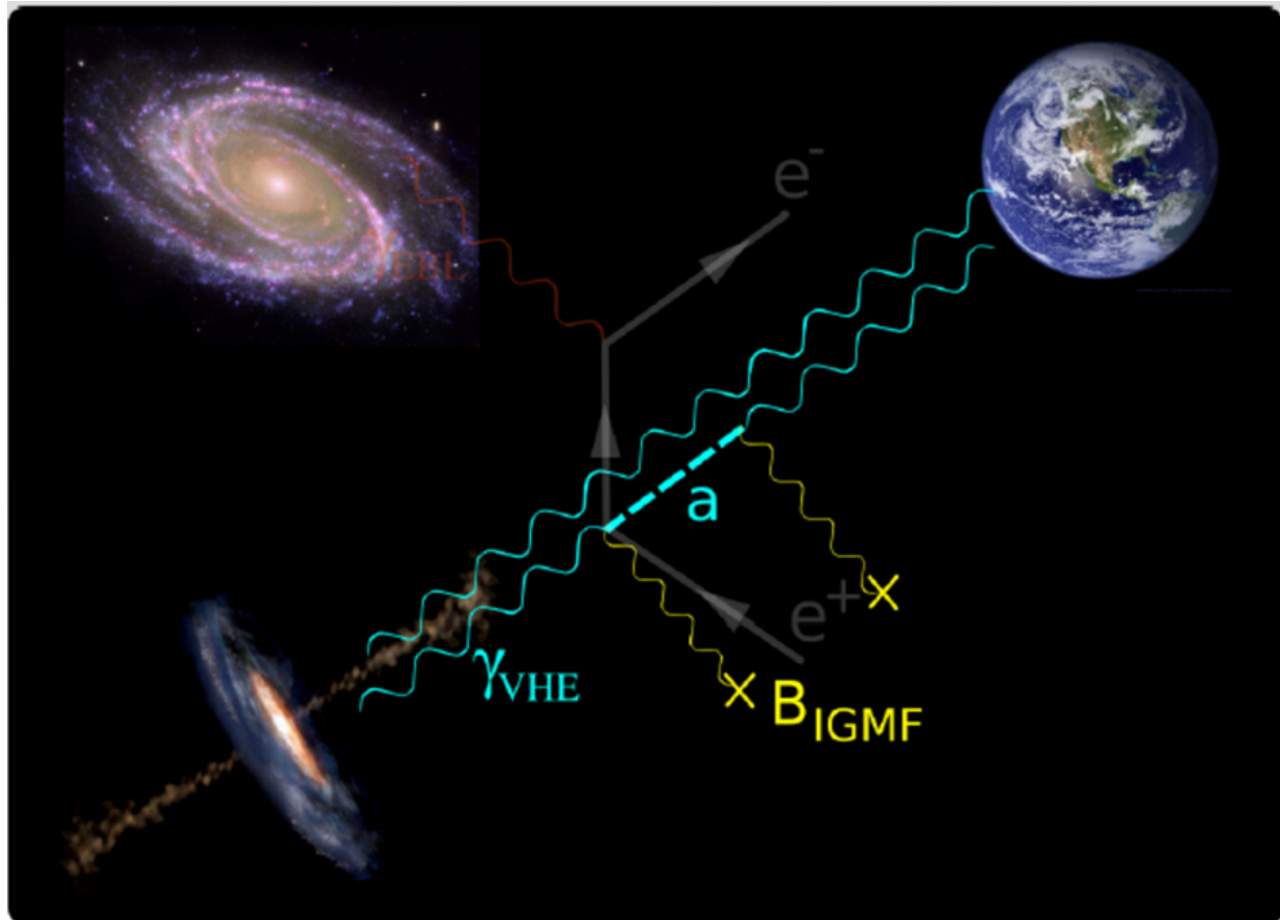
[Horns, Meyer 12]



Photon – ALP Conversion in Cosmic Magnetic Fields?

- > Possible explanation: photon \leftrightarrow ALP conversions in magnetic fields

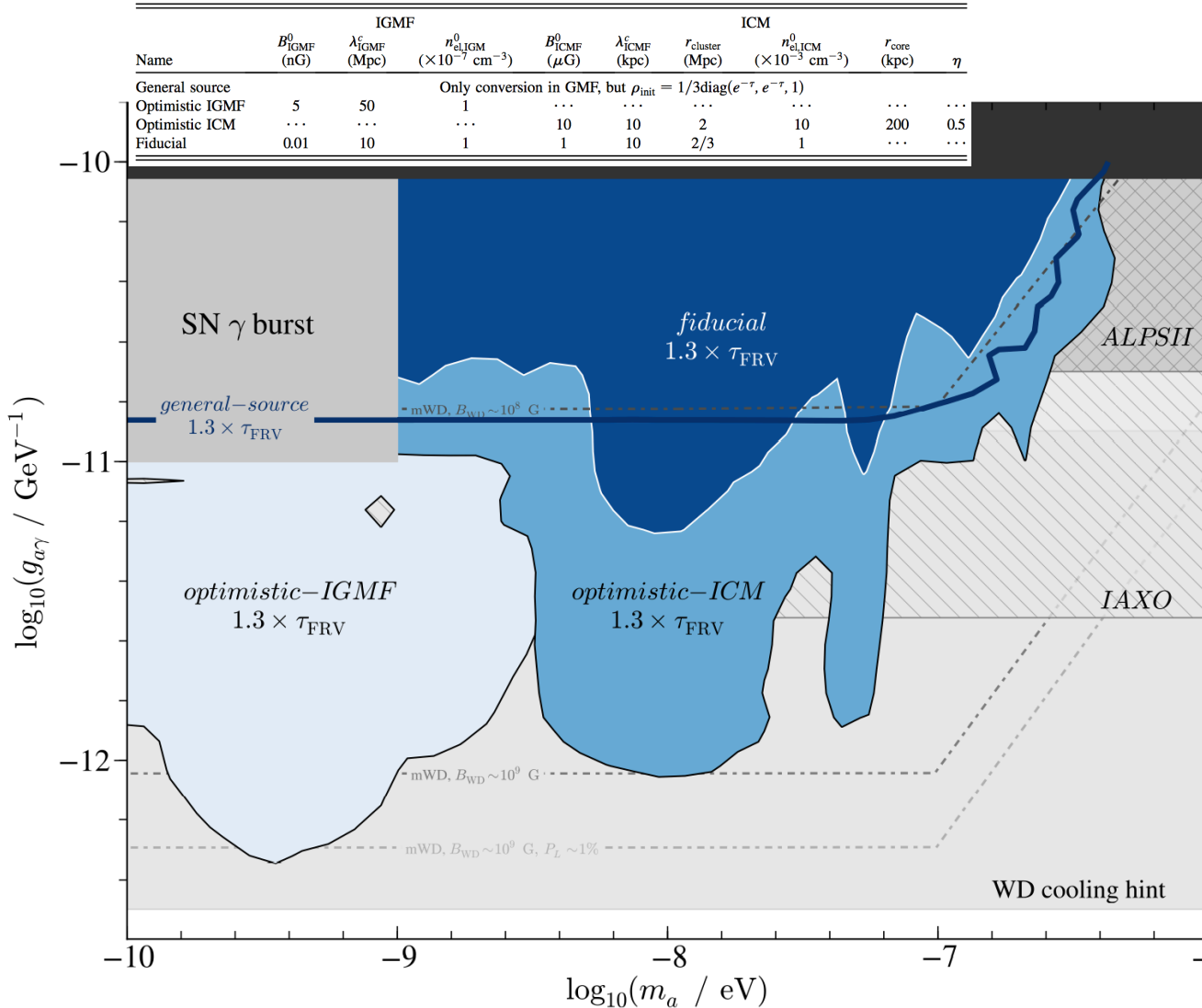
[De Angelis et al 07; Simet et al 08; Sanchez-Conde et al 09; Meyer,Horns,Raue 13; Rubtsov,Troitsky 14; Kohri,Kodama 17]



Photon – ALP Conversion in Cosmic Magnetic Fields?

➤ Possible explanation: photon \leftrightarrow ALP conversions in magnetic fields

[De Angelis et al 07; Simet et al 08; Sanchez-Conde et al 09; Meyer,Horns,Raue 13; Rubtsov,Troitsky 14; Kohri,Kodama 17]



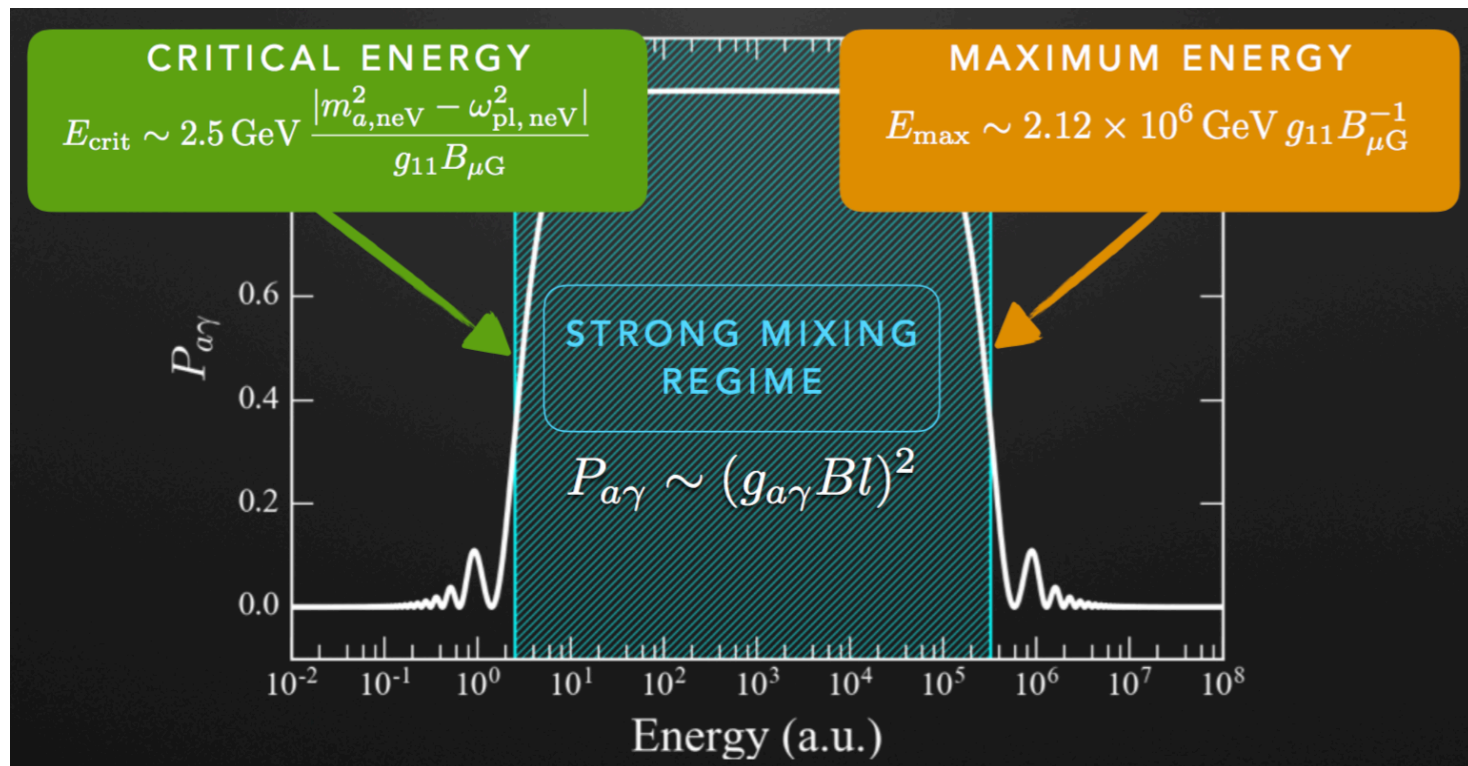
[Meyer,Horns,Raue 13]



Photon – ALP Conversion in Cosmic Magnetic Fields?

> Part of parameter space of interest ruled out by

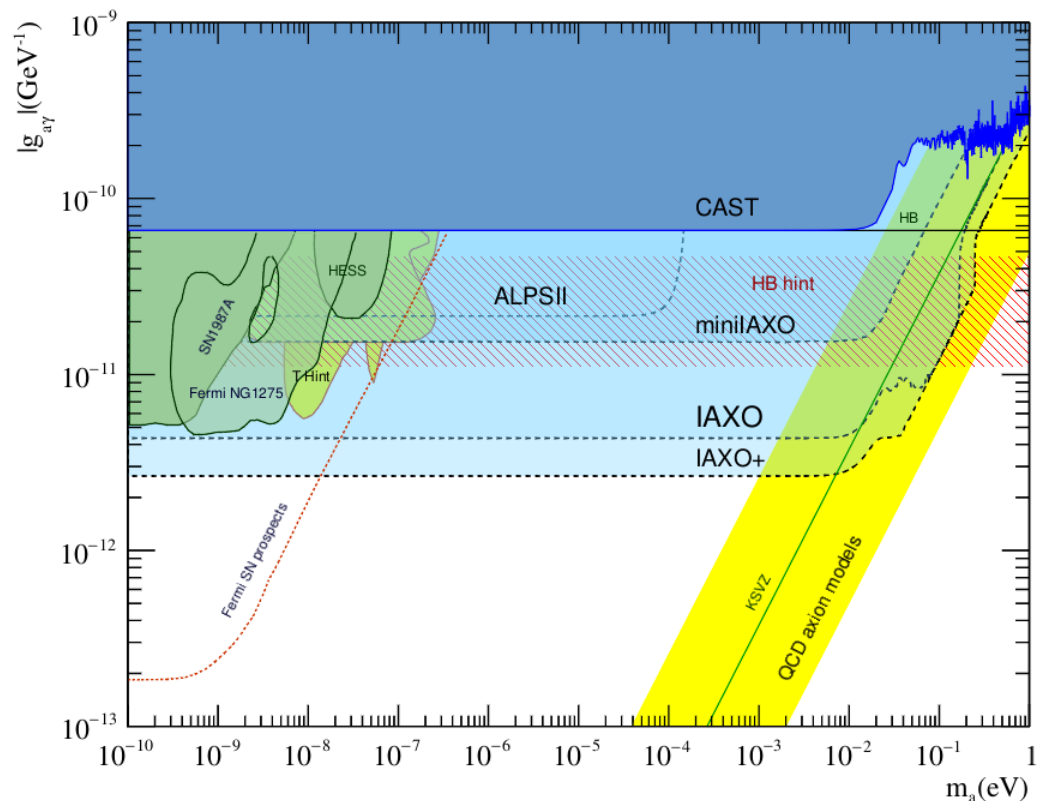
- non-observation of gamma-ray signal from SN 1987 A due to ALP-photon conversion in galactic magnetic field [Brockway et al. 96; Grifols et al. 96; Payez et al. 15]
- non-observation of spectral irregularities in AGN, due to photon-ALP conversions [Abramowski 13; Ajello et al. 16; Berg et al. 16]



Photon – ALP Conversion in Cosmic Magnetic Fields?

> Part of parameter space of interest ruled out by

- non-observation of gamma-ray signal from SN 1987 A due to ALP-photon conversion in galactic magnetic field [Brockway et al. 96; Grifols et al. 96; Payez et al. 15]
- non-observation of spectral irregularities in AGN, due to photon-ALP conversions [Abramowski 13; Ajello et al. 16; Berg et al. 16]



[Irastorza 17]

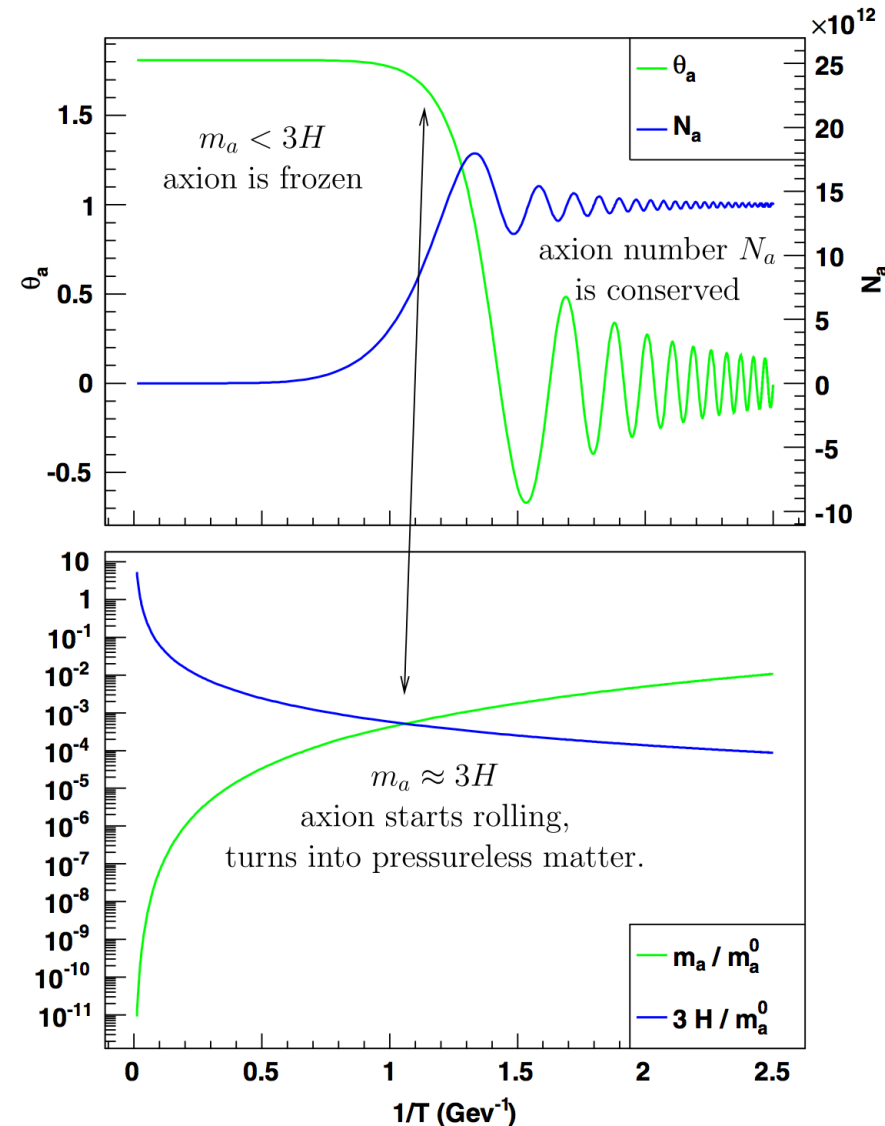


Axion Cold Dark Matter

> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]

- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM



[Wantz,Shellard '09]



Axion Cold Dark Matter

> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]

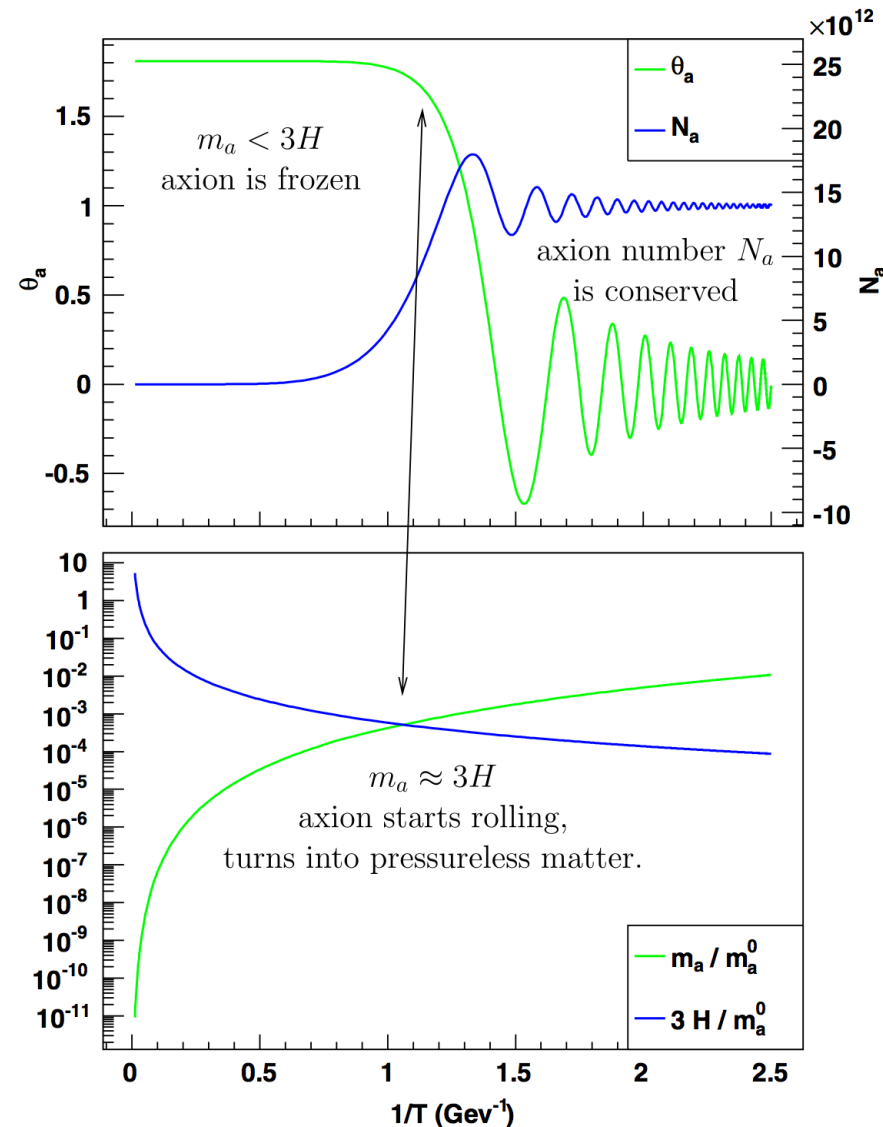
- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM

> Crucial QCD input for prediction of axion DM abundance:

- Equation of state at temperatures around 1 GeV: determines $H(T)$
- Topological susceptibility:

$$\chi(T) \equiv \int d^4x \langle q(x)q(0) \rangle_T$$

determines $m_A^2(T) = \chi(T)/f_A^2$



[Wantz,Shellard '09]



Axion Cold Dark Matter

> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,...]

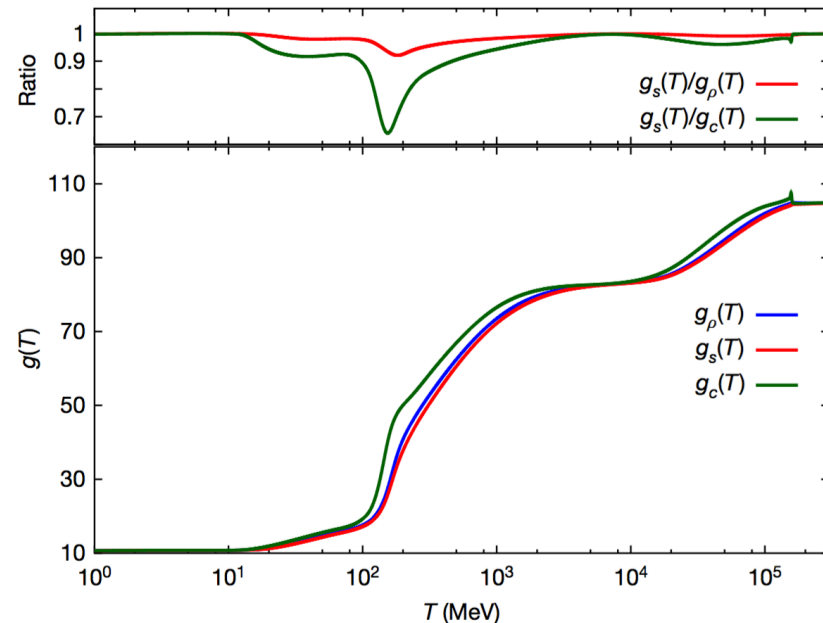
- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM

> Crucial QCD input for prediction of axion DM abundance:

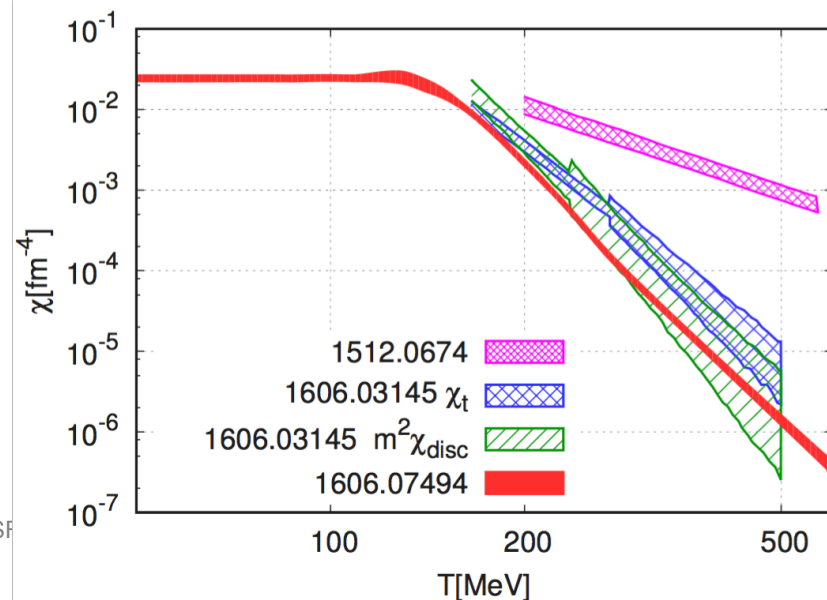
- Equation of state at temperatures around 1 GeV: determines $H(T)$
- Topological susceptibility:

$$\chi(T) \equiv \int d^4x \langle q(x)q(0) \rangle_T$$

determines $m_A^2(T) = \chi(T)/f_A^2$



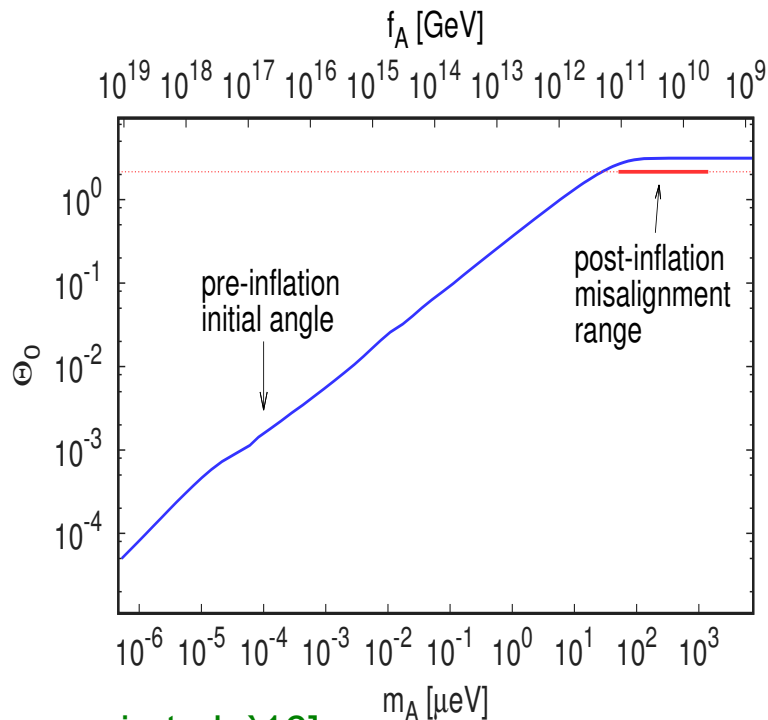
[Borsanyi et al., Nature '16 [1606.0794]]



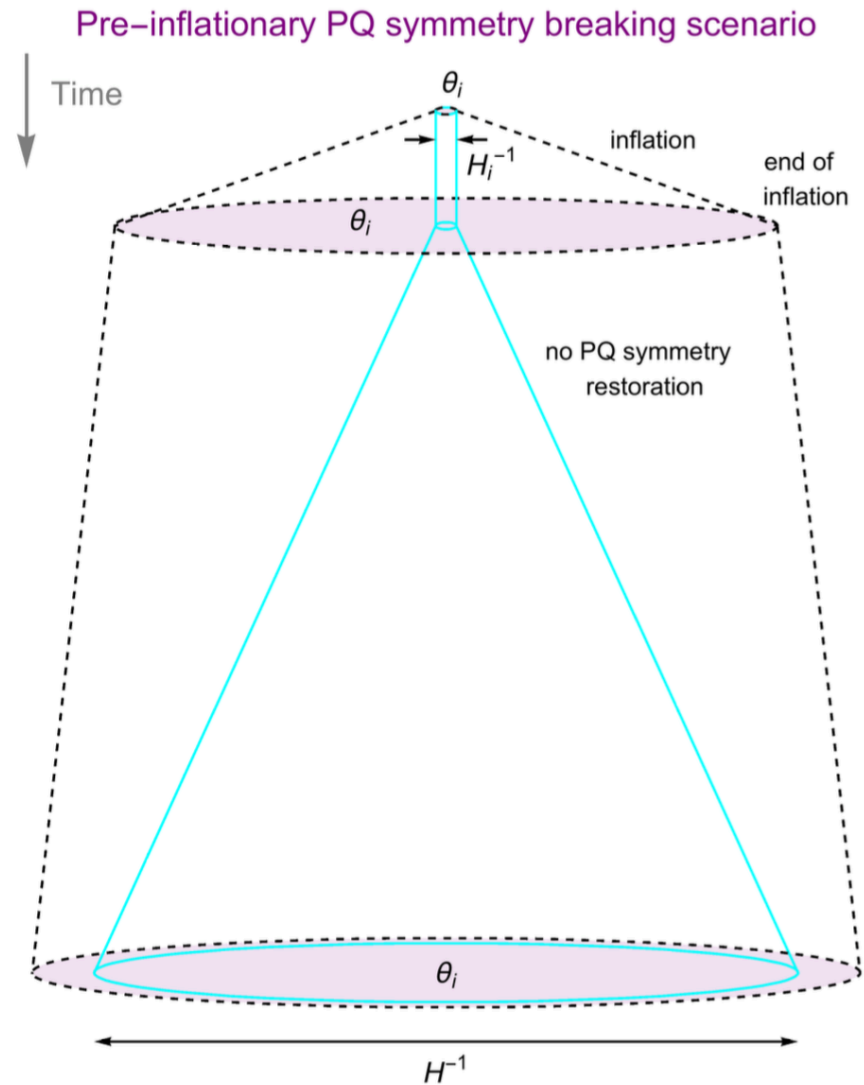
Axion Cold Dark Matter

> If PQ symmetry broken during inflation and not restored afterwards (pre-inflationary PQ breaking scenario)

- Axion CDM density depends on single initial angle during inflation and f_A



[Borsanyi et al. '16]



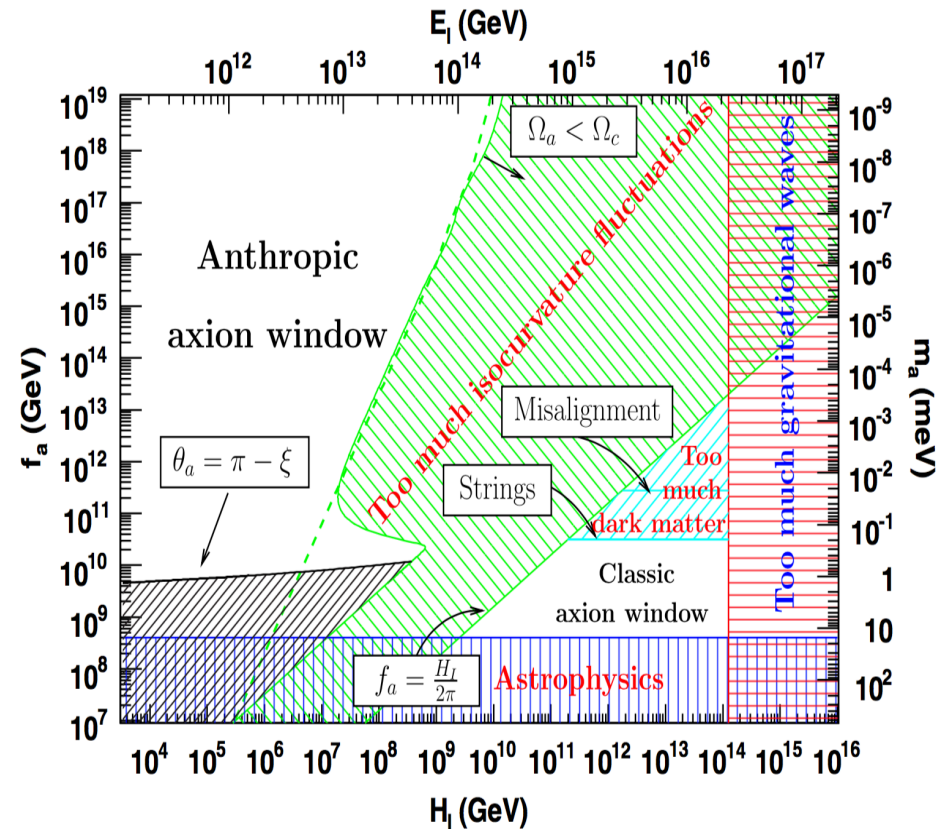
[Saikawa]



Axion Cold Dark Matter

➤ If PQ symmetry broken during inflation and not restored afterwards (pre-inflationary PQ breaking scenario)

- Axion CDM density depends on single initial angle during inflation and f_A
- Axion is present during inflation and creates isocurvature fluctuations which are not erased after inflation



[Wilczek, Turner '91; Beltran et al. 06;
Hertzberg, Tegmark, Wilczek 08; Visinelli, Gondolo 09;
Hamann et al. 09; **Wantz, Shellard 09**]



Axion Cold Dark Matter

➤ If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)

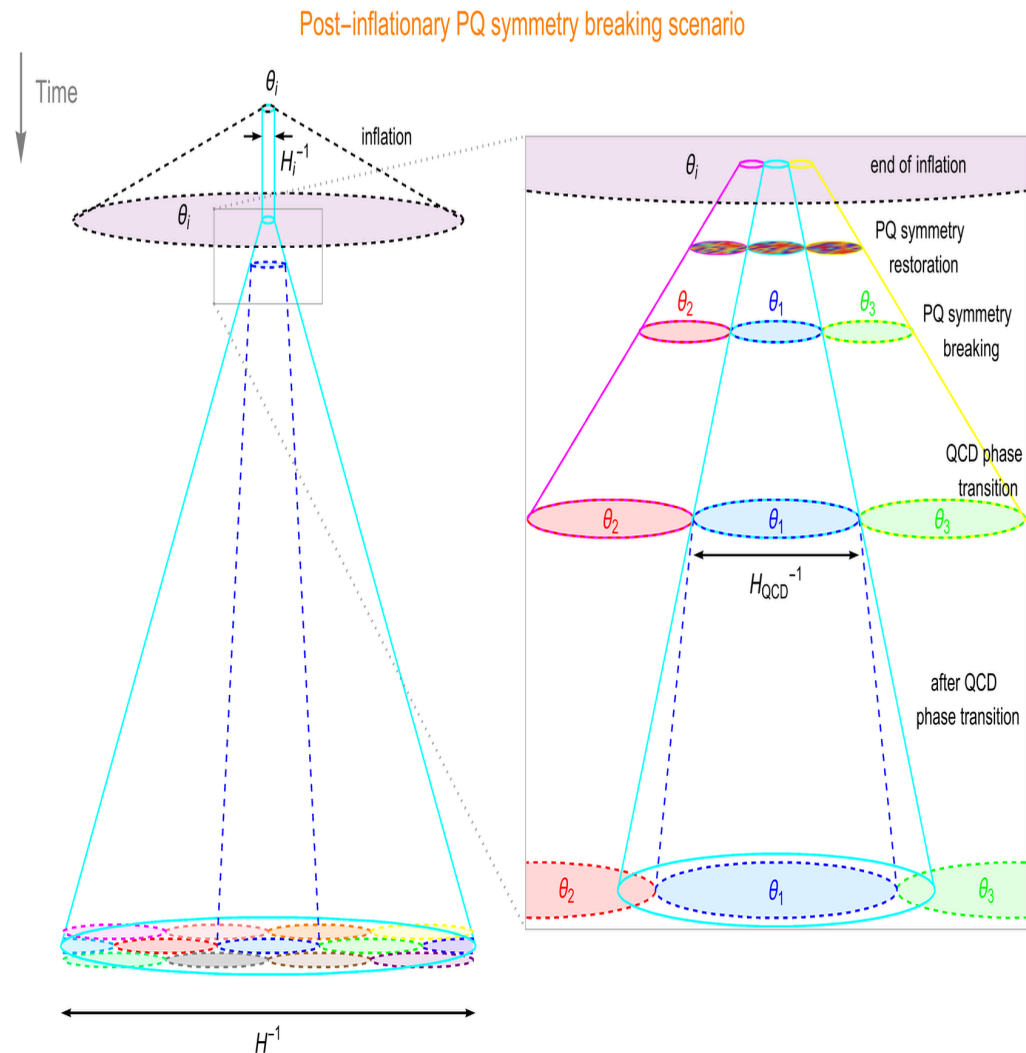
- Vacuum realignment contribution depends on spatially averaged initial misalignment angle and f_A

$$\Omega_{A,\text{real}} h^2 = (3.8 \pm 0.6) \times 10^{-3} \times \left(\frac{f_A}{10^{10} \text{ GeV}} \right)^{1.165}$$

- Upper limit on f_A from requirement that realignment contribution should not exceed DM abundance gives lower limit on axion mass:

$$m_A > 28(2) \mu\text{eV}$$

[Borsanyi et al. `16]



Axion Cold Dark Matter

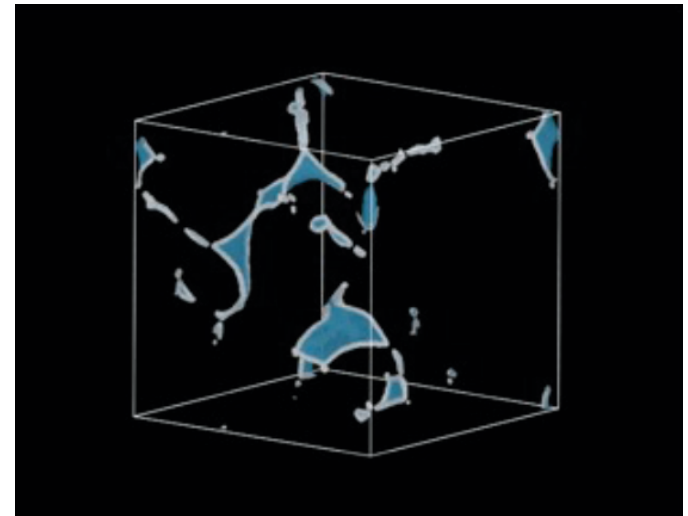
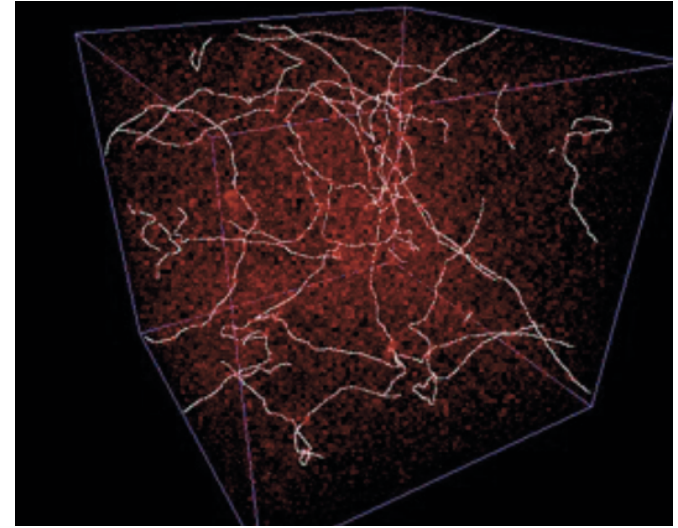
➤ If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)

- Additional contributions arise from decay of topological defects (axion strings and axion domain walls)

$$\Omega_{A,\text{tot}} h^2 = \Omega_{A,\text{real}} h^2 + \Omega_{A,\text{string}} h^2 + \Omega_{A,\text{wall}} h^2$$

- Latter determined by lattice simulations of cosmic string and wall networks

[Hiramatsu et al. 12; Kawasaki et al. 15]



[Hiramatsu et al. 12]

Axion Cold Dark Matter

➤ If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)

- For $N_{\text{DW}} = 1$, string-wall systems short lived, leading to

$$\Omega_{A,\text{tot}} h^2 \approx 1.6_{-0.7}^{+1.0} \times 10^{-2} \times \left(\frac{f_A}{10^{10} \text{ GeV}} \right)^{1.165}$$

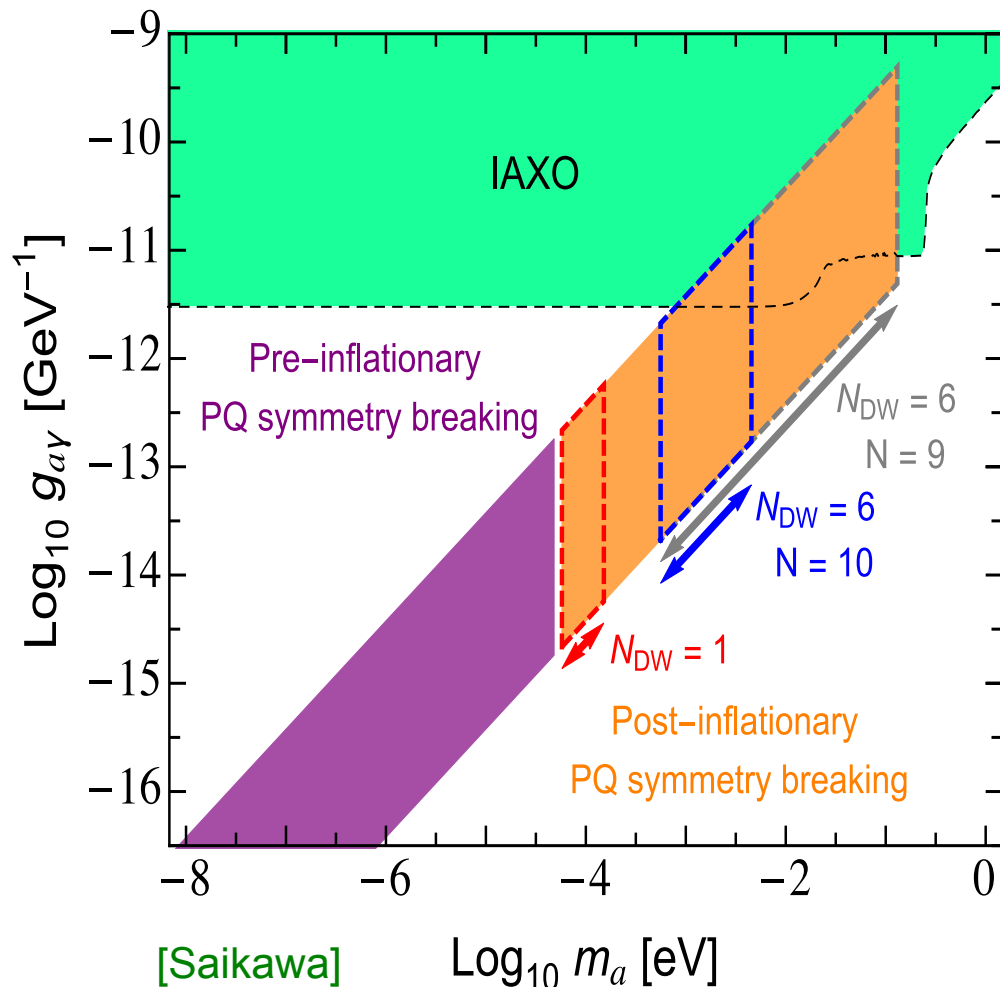
CDM explained for

$$m_A \approx (50-200) \mu\text{eV}$$

- For $N_{\text{DW}} > 1$, string-wall systems absolutely stable, eventually overclosing Universe, unless PQ symmetry explicitly broken e.g. by Planck suppressed operators

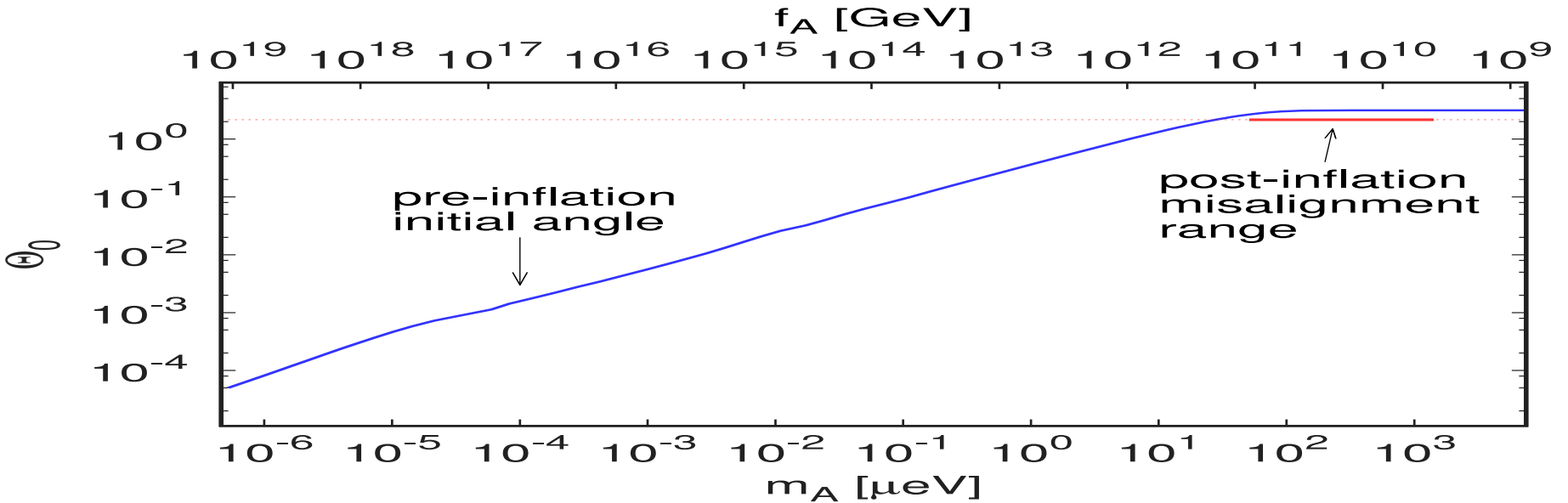
$$\Delta V = g \sigma^N / M_{\text{P}}^{N-4}$$

[AR, Saikawa '16]



Axion Dark Matter Direct Detection Experiments

- Upcoming generation of axion dark matter direct detection experiments can probe entire mass range:



CASPEr

ABRACADABRA

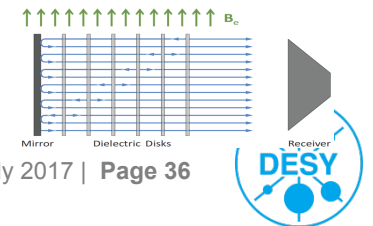
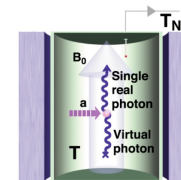
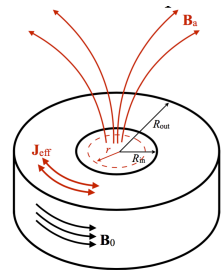
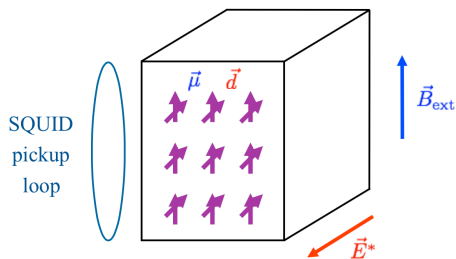
ADMX

MADMAX BRASS

HAYSTACK ORPHEUS

CULTASK ORGAN

QUAX

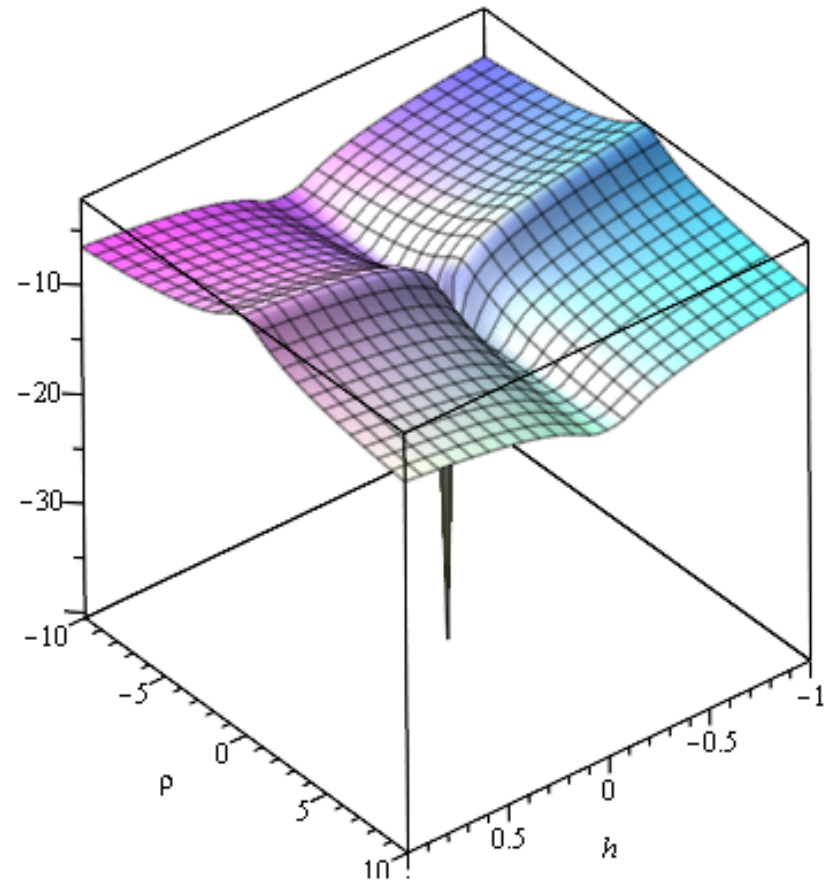


Unifying Inflation and Dark Matter with PQ Field

- $|\sigma| = \rho/\sqrt{2}$ or mixture with Higgs modulus may play role of inflaton, if it has non-minimal coupling to gravity,

$$S \supset - \int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_\sigma \sigma^* \sigma \right] R$$

[Fairbairn,Hogan,Marsh `14]



[Ballesteros,Redondo, AR,Tamarit `16]

Unifying Inflation and Dark Matter with PQ Field

- > $|\sigma| = \rho/\sqrt{2}$ or mixture with Higgs modulus may play role of inflaton, if it has non-minimal coupling to gravity,

$$S \supset - \int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_\sigma \sigma^* \sigma \right] R$$

[Fairbairn,Hogan,Marsh `14]

- > CMB observables

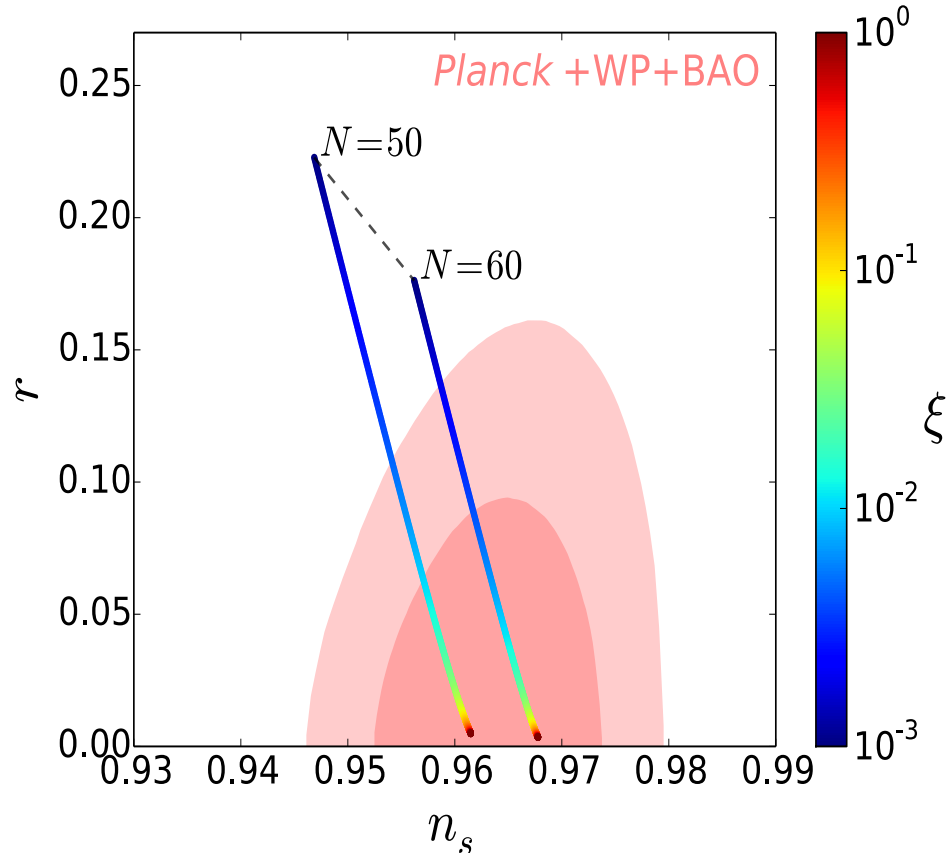
$$A_s = (2.20 \pm 0.08) \times 10^{-9},$$

$$n_s = 0.967 \pm 0.004,$$

$$r < 0.07$$

fit by

$$\xi \simeq 2 \times 10^5 \sqrt{\lambda} \gtrsim 10^{-3}$$



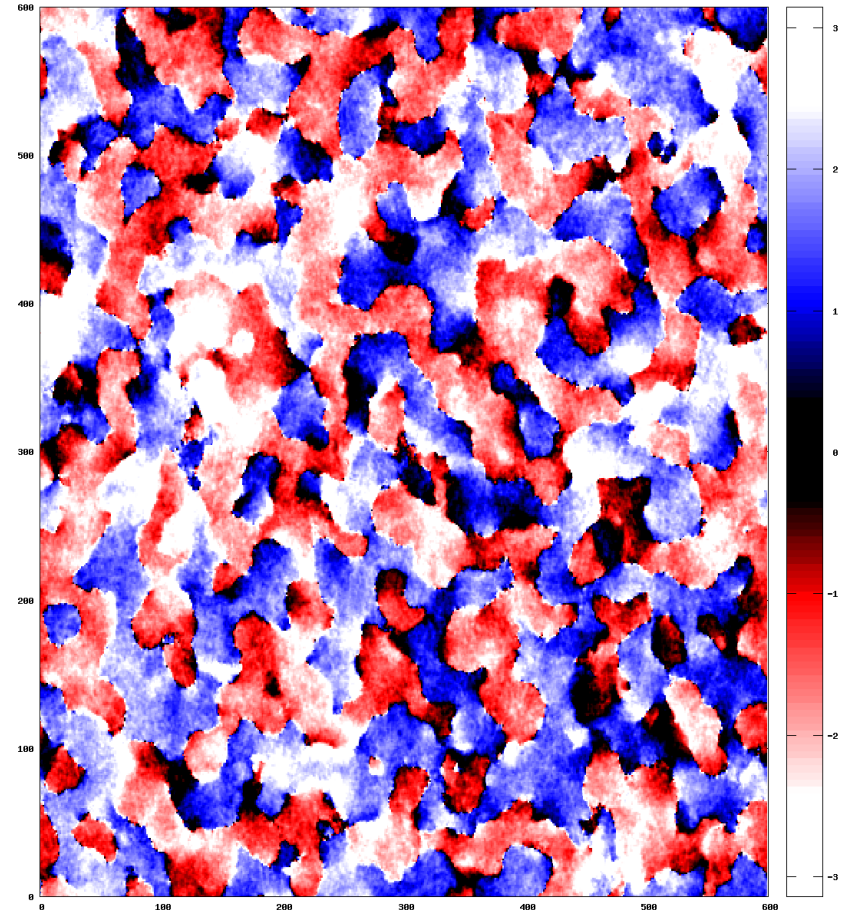
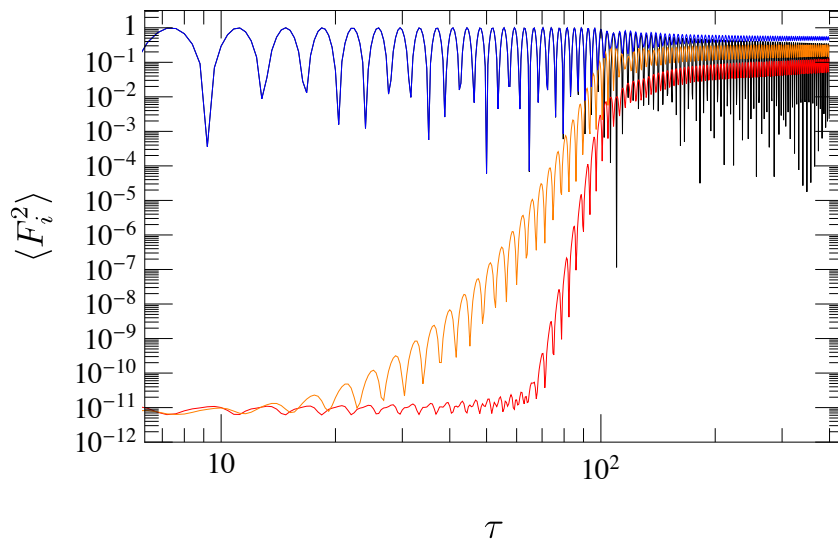
[Fairbairn,Hogan,Marsh `14]



Unifying Inflation and Dark Matter with PQ Field

- PQ symmetry restored after inflation already in preheating stage when PQ field undergoes Hubble damped oscillations in quartic potential

[Ballesteros, Redondo, AR, Tamarit '16]



Unifying Inflation and Dark Matter with PQ Field

- PQ symmetry restored after inflation already in preheating stage when PQ field undergoes Hubble damped oscillations in quartic potential

[Ballesteros, Redondo, AR, Tamarit '16]

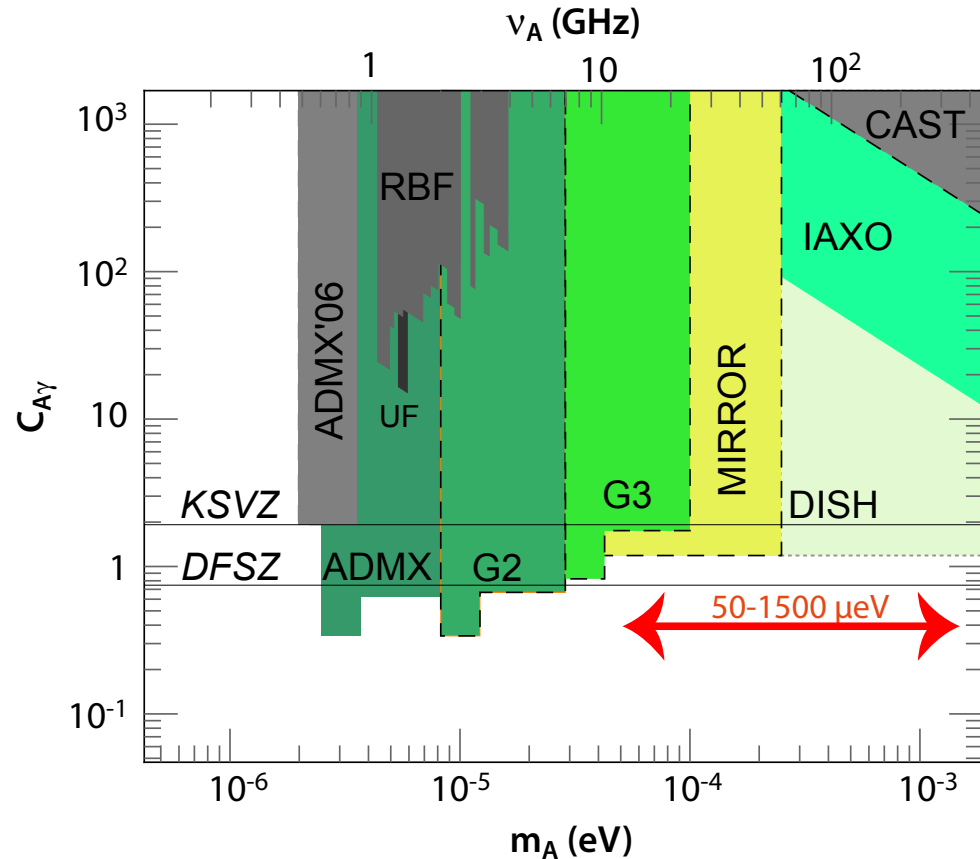
- Axion mass predicted

- For $N_{\text{DW}} = 1$:

$$50 \mu\text{eV} \lesssim m_A \lesssim 200 \mu\text{eV}$$

- For $N_{\text{DW}} > 1$:

$$100 \mu\text{eV} \lesssim m_A \lesssim 10 \text{meV}$$



[Borsanyi et al. '16]



Unifying Inflation and Dark Matter with PQ Field

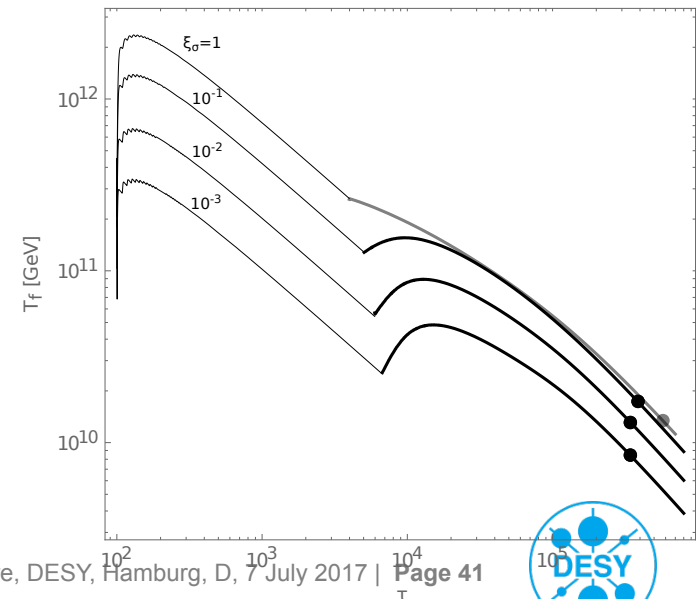
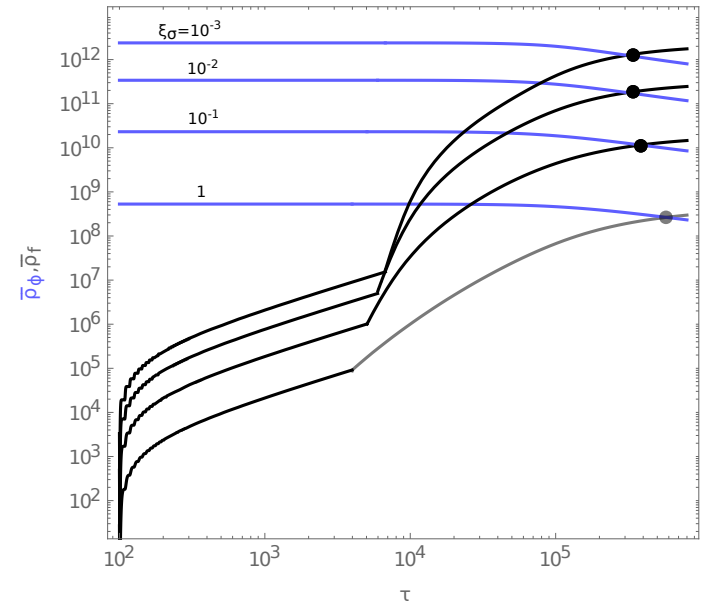
- PQ symmetry restored after inflation already in preheating stage when PQ field undergoes Hubble damped oscillations in quartic potential

[Ballesteros, Redondo, AR, Tamarit '16]

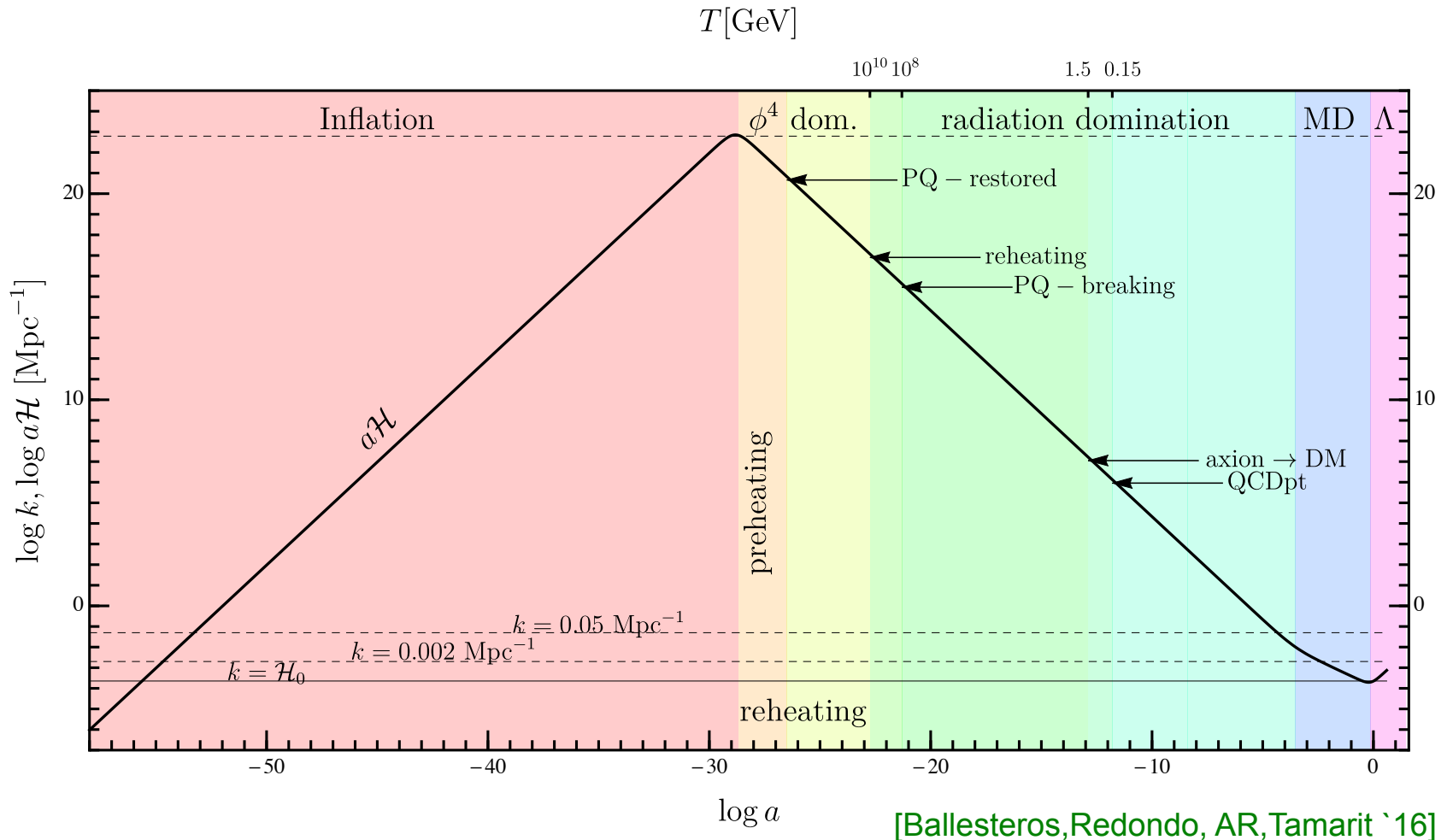
- Axion mass predicted

- Large reheating temperature

- 10^7 GeV for pure PQ scalar inflation ($\lambda_{H\sigma} > 0$)
- 10^{10} GeV for mixed PQ scalar/Higgs inflation ($\lambda_{H\sigma} < 0$)



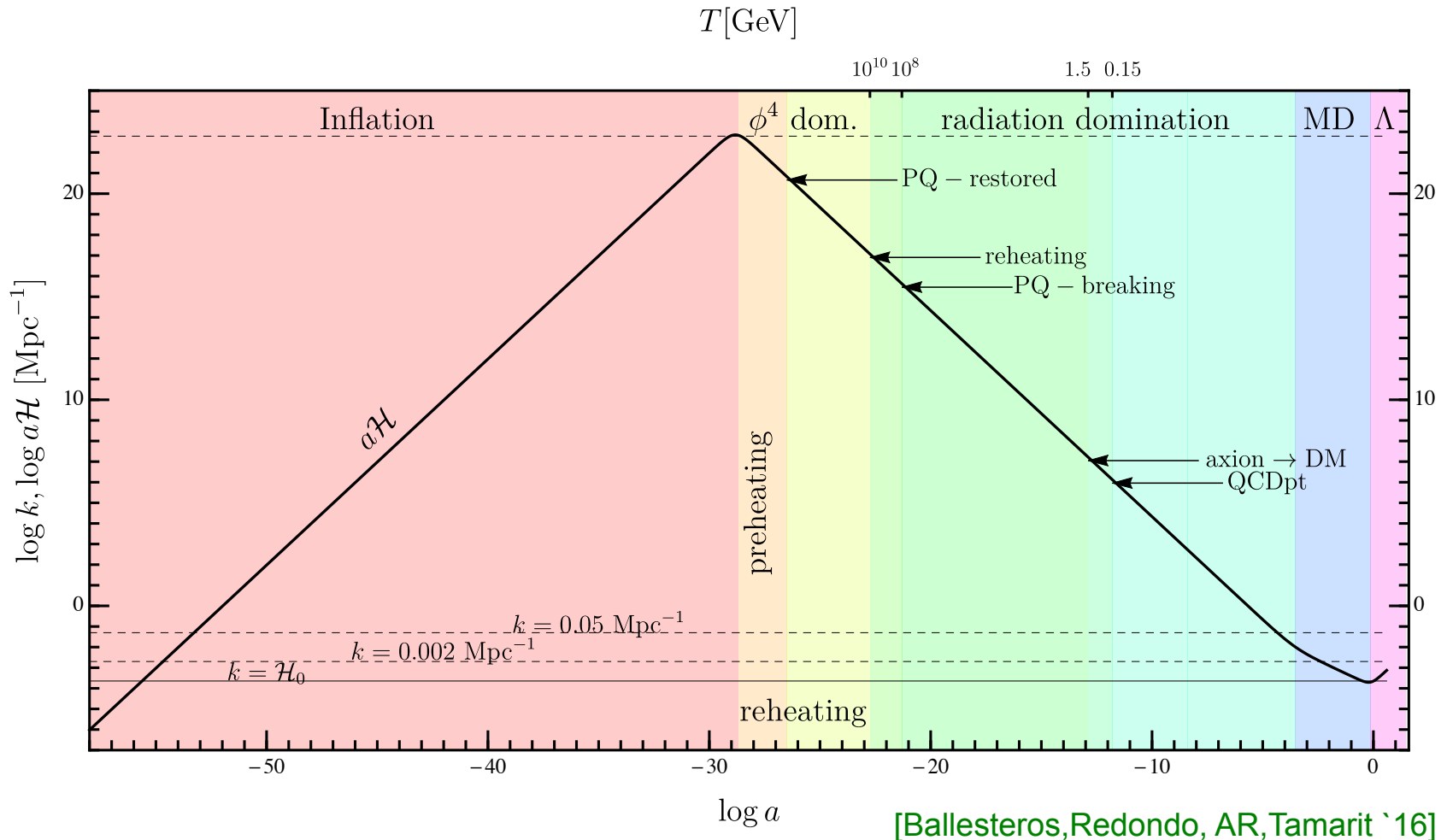
Unifying Inflation and Dark Matter with PQ Field



- Universe expands as in a radiation-dominated era ($w = 1/3$) from the end of inflation until matter-radiation equality



Unifying Inflation and Dark Matter with PQ Field

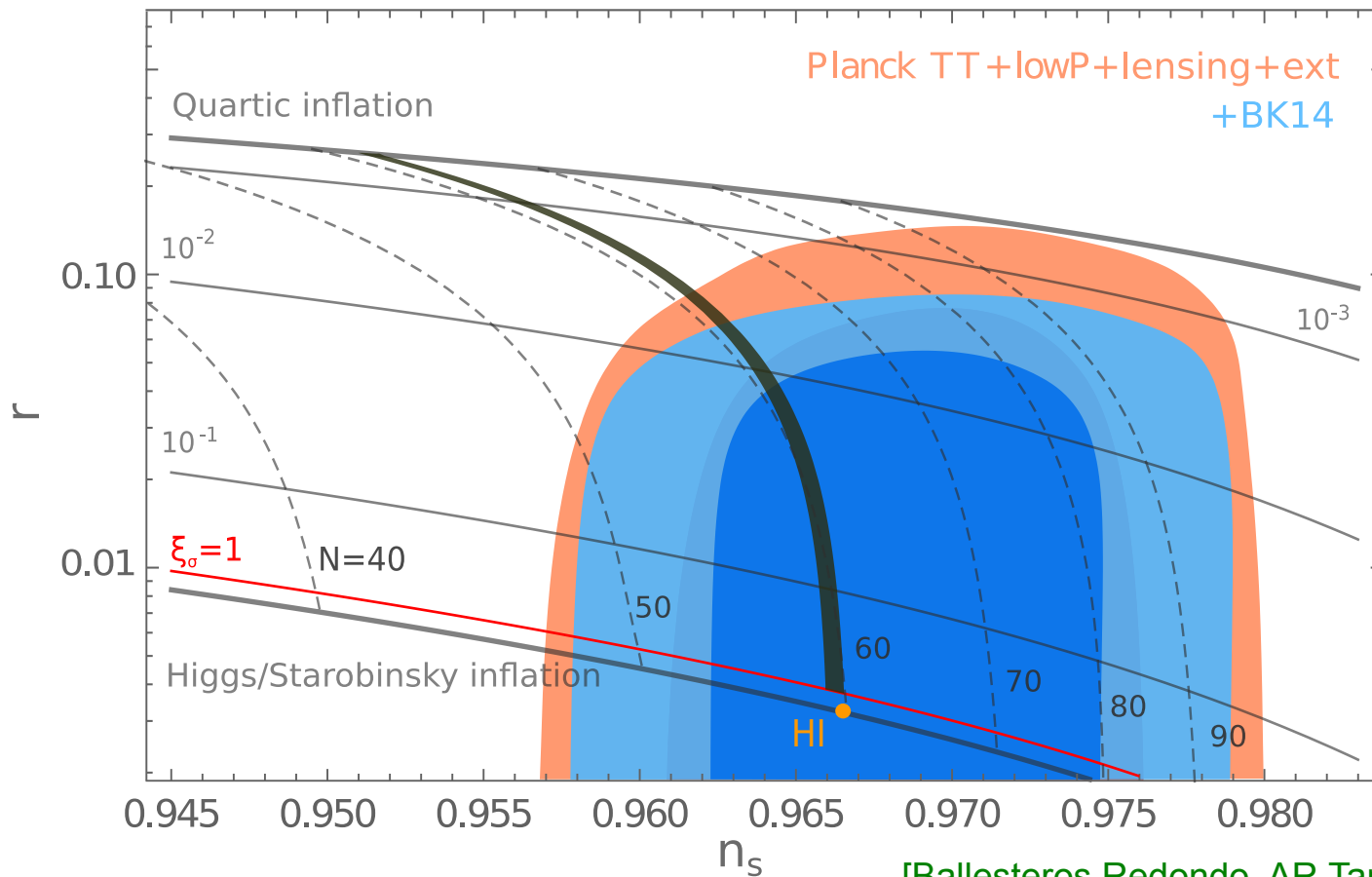


- Allows to calculate number of e-folds $N(k) = \ln(a_{\text{end}}/a(k))$ from the time a given comoving scale k leaves horizon until end of inflation

[Liddle, Leach 03]



Unifying Inflation and Dark Matter with PQ Field



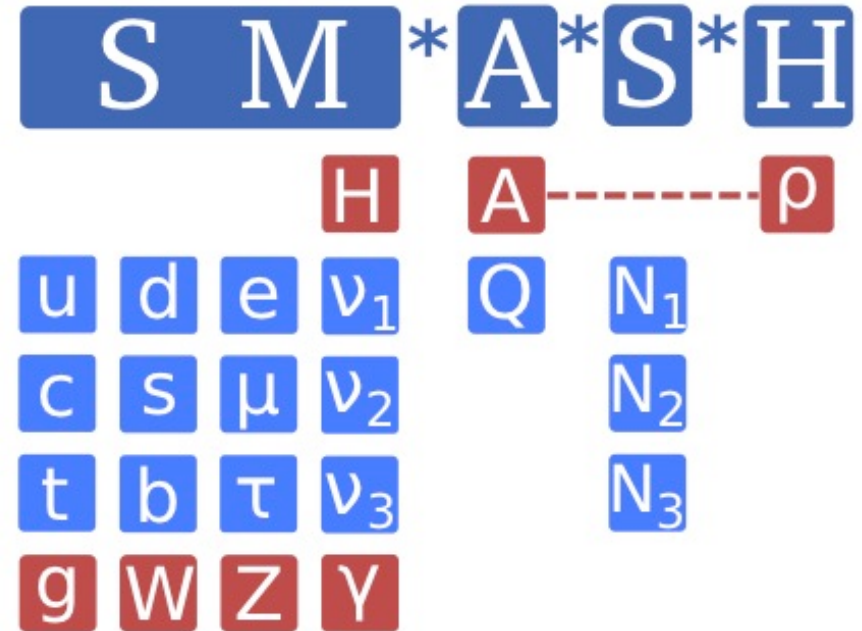
- > Sharp prediction of r vs n_s for fixed pivot scale, e.g. $k_0 = 0.002 \text{ Mpc}^{-1}$
- > Can be probed decisively by next generation CMB experiments (e.g. **LiteBIRD**, **PRISM**)

Unifying Inflation, Dark Matter, and Seesaw with PQ Field

> Augmenting axion models with three SM singlet neutrinos, getting their Majorana masses also through the vev $v_\sigma = N_{\text{DW}} f_A$

- no strong CP problem
- dark matter
- inflation
- neutrino masses and mixing
- baryogenesis via leptogenesis

[Dias et al. '14; Ballesteros et al. '16]



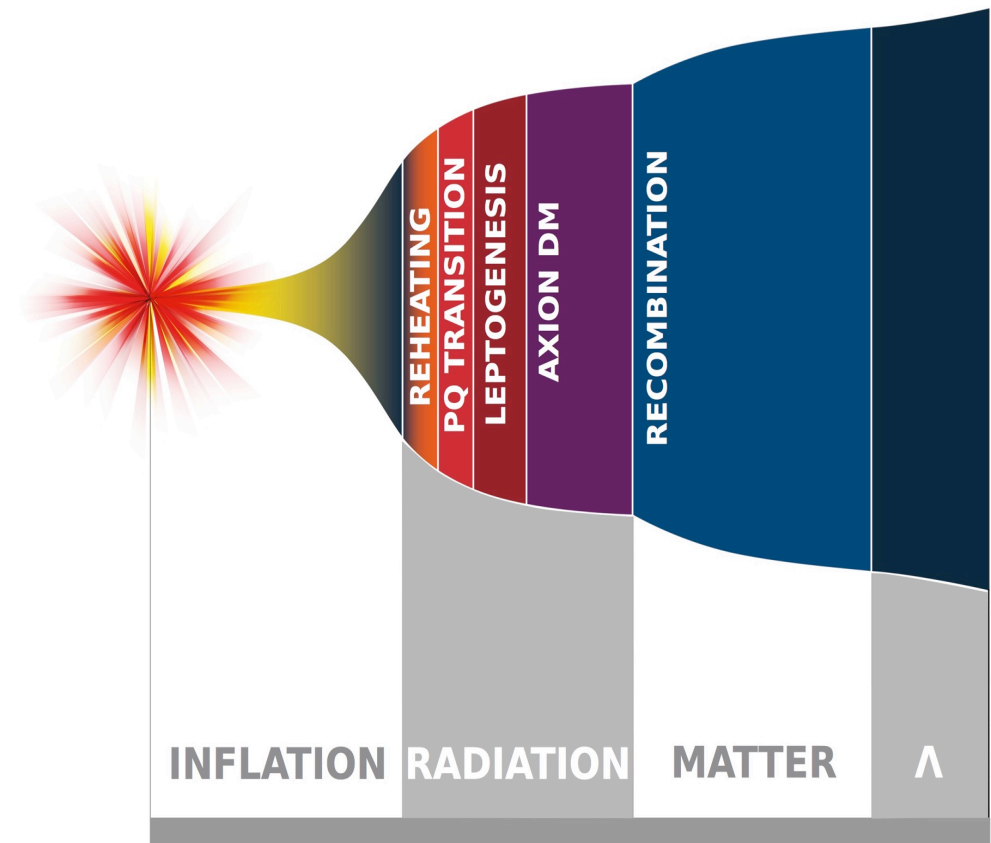
Unifying Inflation, Dark Matter, and Seesaw with PQ Field

- Augmenting axion models with three SM singlet neutrinos, getting their Majorana masses also through the vev $v_\sigma = N_{\text{DW}} f_A$

- no strong CP problem
- dark matter
- inflation
- neutrino masses and mixing
- baryogenesis via leptogenesis

[Dias et al. '14; Ballesteros et al. '16]

- Complete and consistent history of the universe from inflation to now



[desy.de]

Summary

- Axion/ALPs may explain puzzles from astrophysics and cosmology:
 - Dark matter
 - Stellar energy losses
 - Anomalous transparency of the universe for gamma rays
- Relevant parameter range can will be probed decisively by next generation experiments

