Entanglement entropy and related observables in gauge/gravity duality and in tensor networks

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- Quantum gravity
- Quantum information
- Gauge/gravity duality (AdS/CFT correspondence)

# Duality



# Duality



Gauge/gravity duality:

A quantum theory without gravity is related to a gravity theory

- Conjecture which follows from a low-energy limit of string theory
- Duality:

Quantum field theory at strong coupling

⇔ Theory of gravitation at weak coupling

Holography:

Quantum field theory in *d* dimensions

 $\Leftrightarrow$  Gravitational theory in d + 1 dimensions

Quantum field theory defined on the boundary of the d+1-dimensional space



Example of gauge/gravity duality with huge amount of symmetry

AdS: Anti-de Sitter space: Hyperbolic space with constant negative curvature

$$ds^2 = \frac{L^2}{z^2} \left( \eta_{\mu\nu} dx^\mu dx^\nu + dz^2 \right)$$



CFT: Conformal field theory Example: QFT at RG fixed point

Quelle: Institute of Physics, Copyright: C. Escher

Quantum observables at the boundary of the curved space may be calculated from propagation through curved space



Quantum theory at finite temperature:

Dual to gravity theory with black hole



Hawking temperature identified with temperature in the dual field theory

# **Retarded Green's Functions in Strongly Coupled Systems**



subject to infalling boundary condition at horizon

# Book on gauge/gravity duality



# Gauge/Gravity Duality

Foundations and Applications

Martin Ammon Johanna Erdmenger Density matrix  $ho = \sum\limits_n |\Psi_n
angle \langle \Psi_n|$ 

Von Neumann entropy  $S_{vN} = -\text{Tr}(\rho \ln \rho)$ 

Maximised when  $\rho$  diagonal with equal entries, vanishes for pure states where  $\rho^2=\rho$ 

Density matrix  $ho = \sum_{n} |\Psi_n\rangle \langle \Psi_n|$ 

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Maximised when  $\rho$  diagonal with equal entries, vanishes for pure states where  $\rho^2=\rho$ 

Consider product Hilbert space  $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$ Reduced density matrix

$$\rho_A = \mathrm{Tr}_B \rho_{\mathrm{tot}}$$

Entanglement entropy

$$S_A = -\mathrm{Tr}_A \rho_A \ln \rho_A$$

Analogy to black hole entropy ('Lost information' hidden in *B*)

# Entanglement entropy: Gauge/gravity duality



Ryu-Takayanagi 2006:

$$S_A = \frac{\operatorname{Area}\gamma_A}{4G_N}$$

 $\gamma_A$ : Minimal area bulk surface with  $\partial A = \partial \gamma_A$ 

Satisfies strong subadditivity

 $S_A + S_B \ge S_{A \cup B} + S_{A \cap B}$ 

Conformal field theory in 1+1 dimensions (Cardy, Calabrese):

 $S = \frac{c}{3}\ln(\ell\Lambda)$ 

Reproduced by Ryu-Takayanagi result

 $\Lambda \propto 1/\epsilon, \epsilon$  boundary cut-off in radial direction  $c=3L/(2G_3)$ 

Finite temperature (at small  $\ell$ ):

$$S(\ell) = \frac{c}{3} \ln \left( \frac{1}{\pi \epsilon T} \sinh(2\pi \ell T) \right)$$

# J.E., Miekley 1709.07016

#### Analytic expression in closed form for strip region:

$$S_{EE} = \frac{L^{d-1} \left(\tilde{\ell}/\epsilon\right)^{d-2}}{2(d-2)G_N} + \frac{\sqrt{\pi}L^{d-1}}{4(d-1)G_N} \frac{\tilde{\ell}^{d-2}}{z_{\star}^{d-2}} \sum_{\Delta m=0}^{\frac{2(d-1)}{\chi}-1} \frac{(1/2)_{\Delta m}}{\Delta m!} \frac{\Gamma\left(\frac{d}{\chi}a_{-1/2}^{\text{EE}}\right)}{\Gamma\left(\frac{d}{\chi}a_0^{\text{EE}}\right)} \left(\frac{z_{\star}}{z_h}\right)^{\Delta m d}$$
(5.5b)  
$$\times_{\frac{3d-2}{\chi}+1} F_{\frac{3d-2}{\chi}} \left(1, a_{-\frac{1}{2}}^{\text{EE}}, \dots, a_{\frac{d}{\chi}-\frac{3}{2}}^{\text{EE}}, b_{\frac{1}{2}}^{\text{EE}}, \dots, b_{\frac{2(d-1)}{\chi}-\frac{1}{2}}^{\text{EE}}; a_0^{\text{EE}}, \dots, a_{\frac{d}{\chi}-1}^{\text{EE}}, b_1^{\text{EE}}, \dots, b_{\frac{2(d-1)}{\chi}}^{\text{EE}}; \left(\frac{z_{\star}}{z_h}\right)^{\frac{2(d-1)d}{\chi}}\right)$$

 $z_*$ : Turning point of minimal surface

Given implicity in terms of strip width  $\ell$ 

J.E., Miekley 1709.07016

#### Entanglement density



Non-monotonic behaviour signals violation of area theorem

## Holographic entanglement entropy: Finite temperature

Hubeny, Rangamani, Takayanagi 0705.0016



figure by Raimond Abt

# Entanglement entropy: Tensor networks



MERA networks: CFT ground states Implement RG idea

Networks defined on discretizations of hyperbolic space

cf. AdS/CFT: Extra dimension corresponds to RG scale

MERA Network: Entanglement entropy bounded from above by Ryu-Takayanagi formula (Swingle 0905.1317) Susskind, Stanford

A fixed spatial volume is associated to each tensor

Subnetwork connecting RT surface to boundary provides map  $\mathcal{H}_{\mathrm{RT}} \rightarrow \mathcal{H}_{A}$ 

**RT surface:** Smallest Hilbert space for any cut through the network bounded by  $\partial A$  Susskind, Stanford

A fixed spatial volume is associated to each tensor

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**RT surface:** Smallest Hilbert space for any cut through the network bounded by  $\partial A$ 

Natural realization of complexity = volume conjecture Number of tensors measures complexity Quantum error-correcting codes satisfy discretized Ryu-Takayanagi formula (Pastawski, Yoshida, Harlow, Preskill 1503.06237)

Building block: Tensor with maximal entanglement along any bipartition

 $\Rightarrow$  Isometry from the bulk Hilbert space to the boundary Hilbert space

Random tensor network:

Observables obtained by averaging over tensor network states built from random tensors living on a fixed graph

Random tensor networks may be mapped to an associated Ising model Hayden et al 1601.01694

Average value of second Renyi entropy related to partition function

 $\overline{\mathrm{Tr}(\rho_A^2)} \sim Z_A$ 

New paper (complexity): Abt, J.E., Hinrichsen, Melby-Thompson, Meyer, Northe, Reyes

Numerical simulation of entanglement entropy in black hole backgroundMap to associated Ising modelH. Hinrichsen, Würzburg University, 1710.01327



FIG. 6. BTZ black hole (left) with radial coordinate  $\arctan(r/L)$  mapped to conformal coordinates  $\phi, \eta$  (right) where an Ising model on a square lattice is embedded. The top and the bottom row of spins are fixed according to the respective boundary conditions (red= $\uparrow$ , blue= $\downarrow$ ) while the green spins are allowed to fluctuate.

# Gauge/gravity duality: Time evolution of entanglement entropy

J.E., Flory, Fernandez, Megias, Straub, Witkowski 1705.04696



Very large temperature differences:  $\Delta S(t) \propto v_E s_{eq} A t + \dots$ 

# Complexity:

**Susskind:** Complexity = Volume and Complexity = Action proposals

Abt, J.E., Hinrichsen, Melby-Thompson, Meyer, Northe, Reyes: Modify volume proposal to

$${\cal C}({\cal A})=-rac{1}{2}\int\limits_{\Sigma}d^{d+1}xRd\sigma$$



For d = 2, evaluate complexity using Gauss-Bonnet theorem see talk by I. Reyes

For black hole:

$$\mathcal{C} = \frac{x}{\epsilon} - \pi$$
$$\Delta \mathcal{C} = 2\pi$$

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This is reproduced using random tensor networks



FIG. 8. Numerical results on a lattice with  $200 \times 200$  sites for a BTZ black hole with mass M = 0.1. Left: Numerically measured entanglement of the two solutions as functions of the subregion size. As can be seen, the lines cross precisely at the theoretically expected transition point, marked by the vertical green dashed line. Right: Corresponding complexity, reproducing the linear law. The inset shows a magnification where the discontinuous jump occurs.

Kinematic space: Czech et al

Kinematic space of an asymptotically AdS<sub>3</sub> spacetime: Space of its oriented, boundary-anchored geodesics

Use this approach to relate bulk volumes to field-theory entanglement entropy Volume from integration over geodesics that intersect it



Use of geometry widespread in information theory

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Fisher metric in information theory: Metric on space of probability distributions

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Probability distribution  $p(x, \vec{\theta}), x$  a stochastic variable,  $\vec{\theta}$  a set of n external parameters Spectrum  $\gamma(x, \vec{\theta}) \equiv -\ln p(x, \theta)$ 

**Fisher metric** 

$$g_{\mu\nu}(\vec{\theta}) = \int dx \, p(x,\vec{\theta}) \frac{\partial \gamma(x,\theta)}{\partial \theta^{\mu}} \frac{\partial \gamma(x,\theta)}{\partial \theta^{\nu}} = \langle \partial_{\mu} \gamma \partial_{\nu} \gamma \rangle$$

For Gaussian distribution (saddle point approximation)

$$p(x_1,\ldots,x_n) = \frac{1}{(\sqrt{2\pi}\sigma)^n} \exp\left(-\sum_{i=1}^n \frac{(x_i - \bar{x}_i)^2}{2\sigma^2}\right)$$

Fisher metric gives Anti-de Sitter space:

$$ds^2 = rac{1}{\sigma^2} \left( d \bar{x}_i d \bar{x}^i + 2n d \sigma^2 
ight)$$

Question: Understanding the dynamics governing this metric

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- In AdS/CFT, gravity action and dynamics obtained from saddle-point approximation of string theory
- Information theory may help to establish the gravity dynamics more generally
- Further insight using tensor networks?

Banerjee, J.E., Sarkar 1701.02319

Fidelity susceptibility F

Fisher metric

$$G_{mn} = \frac{\partial^2 F}{\partial \lambda^m \partial \lambda^n}$$

Couplings  $\lambda_m$  are dual to to deformations of the AdS metric

**Proposal:**  $F = vol(\lambda) - vol(0)$  (finite expression)

For  $\lambda$  a metric deformation the result matches

- New relations between quantum theory and gravity
- Tensor networks may contribute to a further understanding of gauge/gravity duality
- Conversely, gauge/gravity duality provides predictions for strongly coupled systems