



Identification of hadronically decaying Tau Leptons at the LHC



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- Physics potential with tau leptons
- General tau lepton properties
- ATLAS/CMS Tau Reconstruction and Identification Algorithms
- ATLAS/CMS Tau Trigger
- Remarks on Tau ID performance in early data
- Physics with tau leptons at the LHC with early data
- Concluding Remarks



Physics with Tau Leptons





"WHAT IF WE SPEND ALL THESE BILLIONS AND THERE JUST AREN'T ANY MORE PARTICLES TO FIND ?" Finding new particles at the LHC often involve T-leptons in the final state

- Standard Model Higgs boson
- Supersymmetry searches
- MSSM Neutral Higgs boson
- MSSM Charged Higgs boson
- W '/ Z ' bosons with enhanced 3rd generation coupling
- other scenarios possible

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di-tau (had-had) discovery reach in the m_0-m_{1/2} plane
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m<sub>0</sub>: universal scalar mass (GUT scale)
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m_{1/2} : universal gaugino mass (GUT scale)



Best of both worlds - SUSY Higgs



The Higgs Sector gets much more complex with SUSY: 5 Higgs bosons!

H⁰, h⁰: CP Even neutral A⁰: CP Odd neutral H⁺, H⁻: charged Higgs (shown)



Note high branching ratio to τv especially for high tan β and low mass

Tau final states also important for MSSM neutral Higgs boson searches

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Exotic Gauge Boson Production



New Heavy Gauge Bosons such as W' and Z' bosons are predicted for some models





Basic Tau Properties



TAU BRANCHING RATIOS

- Leptonic Decay Modes $e/\mu\nu_{e\mu}\nu_{\tau}$ 35%
- 1Prong Hadronic Decay Modes $\pi^- \nu_{\tau}$ 11%

 $\begin{array}{ll} \pi^{-} \pi^{0} \nu_{\tau} & 25\% \\ \pi^{-} \pi^{0} \pi^{0} \nu_{\tau} & 9\% \\ \pi^{-} \pi^{0} \pi^{0} \pi^{0} \nu_{\tau} & 1\% \\ K^{-} + \text{Neutrals} & 1.5\% \end{array}$

- 3Prong Hadronic Decay Modes $\pi^{-}\pi^{+}\pi^{-}\nu_{\tau}$ 9% $\pi^{-}\pi^{+}\pi^{-}\pi^{0}\nu_{\tau}$ 4.5% $K^{-}\pi^{+}\pi^{-}\nu_{\tau}$ 0.4% 15%
- Other Modes (\sim 3%)

TAU CHARACTERISTICS

- $m_{ au} \sim 1.7 {
 m GeV}$
- $c\tau = 87 \mu m$
- Hadronic Decays are well Collimated Collection of Charged and Neutral

47% Pions/Kaons

- Most have 1 or 3 Charged Tracks
- Leading Pion Direction Reproduces τ Direction Well.



Three Generations of Matter

observed first in 1977 by Martin Perl et al. (SLAC-LBL)

Because leptonic tau decays cannot really be separated from electrons and muons, we say "tau reconstruction and ID" to refer to hadronically decaying tau leptons

We try to distinguish between 1-prong and 3-prong tau decays (neglect 5-prong) We may in some cases try to reconstruct the number of π^0 in the tau decay

Taus and Jets



Both hadronically decaying taus and jets deposit energy in the EM and Had calorimeters

Challenge: separate out a clean sample of taus from the overwhelming QCD jet rate!



Reconstructing the $\tau\tau$ invariant mass



Use the collinear approximation to reconstruct the $\tau \tau$ mass

- · A good approximation when the parent particle is heavily boosted
- · Assume that the tau decay products are collinear to the tau direction
- Note: approximation breaks down when the decay daughters are back-to-back

$$\vec{p}_{\tau} = \frac{\vec{p}_{l}}{x_{\tau}}$$
Fraction of the tau momentum carried by the visible daughter
$$x_{\tau_{h}} = \frac{p_{x,h}p_{y,l}-p_{y,h}p_{x,l}}{p_{x,h}p_{y,l}+p_{x,miss}p_{y,l}-p_{y,h}p_{x,l}-p_{y,miss}p_{x,l}}$$

$$x_{\tau_{l}} = \frac{p_{x,h}p_{y,l}-p_{y,h}p_{x,l}}{p_{x,h}p_{y,l}-p_{y,h}p_{x,l}-p_{y,miss}p_{x,l}}$$

$$p_{\tau_l} = \frac{p_{x,h}p_{y,l} - p_{x,miss}p_{y,h} - p_{y,h}p_{x,l} + p_{y,miss}p_{x,h}}{p_{x,h}p_{y,l} - p_{x,miss}p_{y,h} - p_{y,h}p_{x,l} + p_{y,miss}p_{x,h}}$$

$$M_{\tau\tau} \approx \frac{M_{lh}}{\sqrt{x_{\tau_l} x_{\tau_h}}}$$

 $p_{T,\tau_l} + p_{T,\tau_h} = p_{T,l} + p_{T,h} + p_{T,\text{miss}}$

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The ATLAS Tau Reconstruction Algorithm



The ATLAS Tau Reconstruction Algorithm can be organized in four steps

1) Seed Building

- track seed finding
- calorimeter seed finding

2) Base Reconstruction

- track association
- calculation of ID variables
- energy/energy flow calculation

3) Tau Substructure Reconstruction

- piO cluster reconstruction
- conversion track removal

4) Identification Algorithms

- electron and muon vetoes
- cut-based ID
- multi-variate ID: projective likelihood, NN, boosted decision trees, probability density range searches

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ATLAS properties:

ATLAS has pixel/SCT/TRT with B = 2 T

tracker provides coverage to $|\eta|$ < 2.5

ECAL (LAr calorimeter) covers to $|\eta|$ < 3.2 with high granularity

HCAL is 25 times coarser than the ECAL: proper coverage to $|\eta| < 3.2$ forward coverage 3.2 < $|\eta| < 4.9$

tracker $X/X_0 = 0.5 - 1.8$ lengths

Tau Seeds at ATLAS



Seed Cell



eta



Base Reconstruction at ATLAS

Direction (η, ϕ) of tau determined by leading track (unless no track seed - then determined by jet barycentre)

Quality tracks associated in $\triangle R < 0.2$ of seed direction Energy calculated by:

- cell weighting (default): $E_{calib} = w(\eta, \phi, i) \times E_{raw}$
- energy flow: neutral energy + charged track momenta terms (corrected for energy leakage)



ID variables calculated: isolation, width, ratios of calo energy and track momenta







Reconstruction of π^0 showers in tau decays



Just as 1-prong and 3-prong taus have different calo properties, so do $\tau^- \rightarrow \pi^- \upsilon$ and $\tau^- \rightarrow \rho^- \upsilon \rightarrow \pi^- \pi^0 \upsilon$

If possible, it makes sense to require different cuts/pdfs to identify $\tau^- \rightarrow \pi^- \upsilon$ and $\tau^- \rightarrow \rho^- \upsilon$ decay if we can separate them

ATLAS reconstructs π^0 subshowers by reconstructing EM topoclusters - but need to eliminate clusters from charged particles:

parametrized shower subtraction for clusters matched to tracks

Subtraction in EMB2



DUITEITEQIUN				
decay	0 π ⁰	1 π ⁰	\geq 2 π^0	
τ→πυ	61%	23%	16%	
τ→ρυ	18%	50%	32%	
τ→a ₁ υ	9%	36%	55%	

 π^0

p

Barrel region

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Conversion Track Tagging in tau decays



In approximately 40.5% of $\tau\text{-decays}$ a π^0 is present

The process $\pi^0 \rightarrow \gamma\gamma$ and $\gamma \rightarrow e^+e^-$ distorts the track spectrum of τ -leptons - can we tag such tracks as conversion tracks?

There are studies at ATLAS to identify two-track vertex candidates and then calculate electron probability that track originated from an electron

Main property: high threshold TRT (Transition Radiation Tracker) hits • electrons radiate photons which also can produce hits in the TRT gas hence providing more hits at "high" threshold



if we tag tracks from conversions, we can reduce the bias introduced on the τ track multiplicity spectrum





Electron Veto at ATLAS



Tau candidates are divided into categories based on:

- E^{HCAL} (energy deposited in HCAL)
- E^{strip}max (energy in strip layer that is not associated to leading track)
- Then (category-dependent) cuts to veto electrons are placed on:
- E^{ECAL}_T/p_T^{lead trk} (ratio of transverse energy in ECAL to transverse momentum of leading track)
- $N^{TRT}_{HT}/N^{TRT}_{LT}$ (ratio of high threshold to low threshold hits in the TRT)



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Tau ID algorithms at ATLAS



The ID at ATLAS is needed to separate $\tau\text{-leptons}$ from jets

ID algorithms are based on the variables such as isolation fraction, EM radius calculated in the base reconstruction

ATLAS has studied the following:

- cut-based
- Bayesian Projective Likelihood
- Probability Density Range Searches
- Neural Networks
- Boosted Decision Trees

Note that while multi-variate techniques can give very high performance on MC - for first data, it may be that "simpler is better"

cut-based approaches have preference in the early data taking period (more on this a bit later)



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The CMS Tau Reconstruction Algorithm



The CMS Tau Reconstruction Algorithm is organized in three steps

1) General Particle Flow Reconstruction

- iterative tracking algorithm
- calorimeter clustering using topological clusters
- links from tracks to clusters
- particle ID and energy calibration

2) Base Tau Reconstruction and Identification

- Cone algorithm to make jets of particles
- match to charged track
- isolation from charged hadrons and photons

3) Higher Level Identification

- other variables to suppress QCD jets
- special criteria to reject electrons and muons

CMS properties:

CMS has silicon tracker with B = 3.8 T

tracker provides coverage to $|\eta|$ < 2.6

ECAL (crystal calorimeter) covers to $|\eta|$ < 3.0 with high granularity

HCAL is 25 times coarser than the ECAL: proper coverage to $|\eta| < 3.0$ forward coverage 3.0 < $|\eta| < 5.0$

tracker $X/X_0 = 0.4 - 1.6$ lengths



Particle Flow Algorithm at CMS



Particle Flow: Reconstruction of all particles in the event using all detector subsystems

Iterative tracking:

- start with very tight criteria on reconstructed tracks (seed tracks)
- remove hits associated to selected tracks from consideration
- loosen tracking criteria to reconstruct additonal tracks without removed hits
- iterative procedure gives high efficiency, while reducing fakes due to hit removal of already selected tracks
- \bullet finally loosen constraints on primary vertex to pick up tracks from K0_S and Λ tracks

Calo clustering (topoclusters):

- cluster seed local maxima exceeding threshold
- add iteratively neighbouring cells exceeding threshold
- separate clustering in different calorimeter components



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Particle Flow Algorithm at CMS

comparison of tracks to calorimeter information



Track extrapolation to calorimeter systems at the expected depth and linked to clusters if the extrapolation is within cluster boundaries

determines which hypothesis is best

- Clusters are linked to other clusters similarly (using cluster boundaries)
- Particle ID performed to assign particle flow objects in the following order:
- muons
- electrons
- charged hadrons
- photons
- neutral hadrons

Energy Calibration: $E_{calib} = a + b(E,\eta)E_{ECAL} + c(E,\eta)E_{HCAL}$ where coefficients are tuned to MC for now

Cone algorithm of 0.5 to make "particle-flow jets" - start of tau reco



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Base Tau Reconstruction and ID at CMS



Take "particle flow jets" with $p_T > 15$ GeV and require charged hadron candidate with $p_T > 5$ GeV within $\Delta R<0.1$ of jet direction (forms the leading track)

Define a "signal cone" and "isolation annulus" around the leading track • fixed cones: 0-0.07 for signal, and 0.07-0.45 for isolation in (η , ϕ) space • shrinking cones: signal cone limit at 5/E_T within 0.07-0.15 Note: p_T cut should be tuned to analysis

No charged hadron candidates or photon candidates allowed in isolation annulus

In development: use particle-flow algorithm to reconstruct secondary tracks from photon conversions (to further reject photons) τ -jet axis



Energy and Direction Resolutions at CMS





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Higher Level Tau ID at CMS

Higher level tools to further suppress QCD jets constantly being developed/refined But many fake taus can come from electrons - electron vetoes are necessary! Use multi-variate techniques of tracking+calorimeter information to pre-ID electrons

The ATLAS Tau Trigger

At Level 1, hardware trigger uses calorimeter towers of $\Delta\eta x \Delta \phi$ =0.1x0.1 in 2x1 group in ECAL and 2x2 group in HCAL

EM/Hadronic isolation also used in 0.1 wide ring around the seed tower

HLT L2 uses ID variables such as: track multiplicity, isolation, narrowness of candidate

HLT EventFilter uses "pseudo-offline" algorithms to reconstruct candidate and filter out fakes

Detector Understanding with First LHC data - 01. July 2009 80 100 12 True Visible E_T (GeV)

120

• L1

△ L2

• EF

60

The CMS Tau Trigger - Level 1

At L1, the CMS calorimeter trigger towers are about 0.087x0.087 in ($_{\eta,\phi})$

4x4 regions are examined and matched to predetermined patterns

Then the area is expanded to 3x3 regions of 4x4 towers (12x12 in total) If central region has more energy than its eight neighbours

Level 1 $_{\tau}$ -leptons with lower transverse energy need to be prescaled to control the trigger rates

Multi τ -lepton triggers are also possible

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1.290 m

The CMS Tau Trigger - HLT

t→1,3h+X)

Level-1 τ -jets provide a seed to HLT

Calo HLT τ -candidates must pass isolation criteria: $P_{isol} = \Sigma E_T^{0.4} - \Sigma E_T^{0.13} < P_{threshold}$

regional jet finding (in region of L1 τ -jets) speeds up processing

Pixel HLT τ -candidates are reconstructed around L1 τ -jets with the fast track finding algorithm

Requires no tracks found in isolation cone

L=2x10³³cm⁻²s⁻¹

Calo Tau Trigger on first Calo jet

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LvI-2 τ -jet axis

signal cone R_s

р

Needs	Measures for tau ID performance	Analyses affected	
jet/electron suppression	efficiency/rejection	all	
same sign background determination	good charge ID	$Z \rightarrow \tau \tau$, H/A $\rightarrow \tau \tau$, Z' $\rightarrow \tau \tau$, some SUSY searches	
additional evidence of tau excess (in early data)	unbiased track multiplicity	$W \rightarrow \tau v, Z \rightarrow \tau \tau,$ inclusive SUSY searches	
trigger on taus	trigger efficiency/rejection and prescale	$W \rightarrow \tau \nu, H^{\pm} \rightarrow \tau \nu$	
invariant/visible mass reconstruction	resolutions on tau energy and direction (Ε, η, φ)	Z \rightarrow ττ, H/A \rightarrow ττ, Z' \rightarrow ττ, SUSY stau mass measurement	
label tau decay modes	π^0 reconstruction	exclusive SUSY searches, $Z' \rightarrow \tau \tau$ polarization studies	

Rejection vs Efficiency

With QCD processes having a rate of 10⁶ higher than Electroweak processes - suppression of multi-jet backgrounds is crucial

Algorithms also need to suppress electron/muon fakes as well - but not at such a high level of rejection

Tight cuts: high purity, low signal yield
Loose cuts: high efficiency, background contamination

Optimal working point depends on analysis and the luminosity available:

- W/Z analyses tighter cuts?
- Higgs/SUSY searches loose cuts?
- low luminosity looser cuts?
- high luminosity tighter cuts?

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Note that the efficiency of tau ID and rejection on jets 0.04
need to be measured in data!
(more on this later)
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Charge Determination & Track Multiplicity

Charge determination & track multiplicity heavily depend on the track selection • rejecting fake tracks and efficient ID of pion/kaon tracks is crucial • tagging tracks from conversion can help correct the track multiplicity

Analyses with first data will face skepticism that the signal is "from τ -leptons" • one good cross check of this is the track multiplicity spectrum for signal region τ -candidates - dominated by 1-prong and 3-prong candidates

Note that QCD fakes tend to have large track multiplicity:

• higher fraction of fakes in 3-prong bin

Advanced analyses (i.e. VBF $H \rightarrow \tau\tau$) may want to bias the track spectrum towards lower track multiplicity to kill multijet background.

(other ways to confirm signal excess - i.e. invariant mass peak above background distributions)

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CMS Preliminary

Energy & Direction Resolutions

Analyses looking to reconstruct the $\tau\tau$ invariant or visible mass will want good resolution of the τ -lepton kinematics (E_T , η , ϕ) CMS Preliminary

Energy/particle flow determinations of such kinematics take into account track info which has better resolution at low $p_{\rm T}$

At high p_T , there may be difficulties because the narrowing of τ -candidates makes individual components harder to distinguish (calo-determination may be more stable)

Energy/particle flow usually underestimate QCD jet energy, while being accurate for τ -leptons because of prominence of EM calorimeter info to HAD calorimeter (this reduces the effective QCD cross-section)

Tau ID in early data

Optimizing the rejection/efficiency of τ -ID in Monte Carlo is interesting but data will be a much different story!

• proper MC modeling of variables is going to be inaccurate

- many variables require precision alignment/calibration before they are effective (secondary vertexing, track impact parameters, hadronic calibration)
- many variables will have different behaviour due to underlying event and pile-up (track and calorimeter isolation with wide cones)

For early data - can try to use only variables suspected to be under control with tens/hundreds of pb⁻¹

Try to avoid using multi-variate techniques, and use rectangular cuts instead

ATLAS looked at optimization of rectangular cuts on a set of *"safe variables*"

i.e. EM radius, calorimeter isolation with narrow cone, width of track system, ratio of E_T^{EM} to E_T^{HAD} at the EM scale, ratio of EM/HAD energy to Σp_T^{tracks}

Optimization for different p_{T} bins, different efficiencies is being studied

The $Z \rightarrow \tau \tau$ Channel

The Z $\rightarrow \tau\tau$ channel (one τ_{lep} decays leptonically, one τ_{had} decays hadronically) will be the most important channel to make data-based studies of τ_{had} -lepton properties

- main backgrounds: top pairs, W+jets, $Z \rightarrow \mu\mu$, $Z \rightarrow ee$, multi-jets
- typical cuts may include: electron/muon ID and isolation, $M_T(lep, E_T^{MISS})$, τ_{had} -ID ΣE_T^{jets} , dilepton veto, b-tag veto, E_T^{MISS}
- identifying one leg as electron/muon (tag-side) gives an unbiased sample of τ_{had} -leptons on the other leg (probe-side) that can be used for study

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Understanding τ -ID with $Z \rightarrow \tau\tau$

The $Z \rightarrow \tau \tau$ sample will be a key control sample for understanding τ -ID with data

- Tag (electron/muon) and probe (τ -lepton) methods can be used to understand:
- the τ -ID efficiency (inclusive / 1-prong / 3-prong)
- the τ -trigger efficiency (for different thresholds/trigger streams)

Using the known $\sigma(pp \rightarrow Z) \times BR(Z \rightarrow \tau \tau)$, efficiencies can be measured in data

- \cdot global scale factor to correct for such differences in MC and data
- eventually PDFs for ID variables can also be corrected with data input

The visible mass distribution of the $Z \rightarrow \tau \tau$ products can be used to evaluate the energy scale for the τ -leptons (cell weighting/energy flow)

The invariant mass of $Z \rightarrow \tau \tau$ products can be reconstructed using collinear approximation

- \cdot should peak at well measured Z mass (91.2 GeV)
- $\boldsymbol{\cdot}$ uncertainty dominated by E_{T}^{MISS}
- peak of distribution gives us a handle on the $\mathsf{E}_\mathsf{T}^{\mathsf{MISS}}$ scale

Measuring τ fake rates in jet samples This can be done with very early data - and studies can be achieved before the $Z \rightarrow \tau \tau$ sample is well-established Fake Rate = (# probe jets passing τ -ID) / (# probe jets) Tag jet probe jet Real τ -production is ~negligible in comparison with jet production • can be corrected by MC subtraction of real τ -lepton contribution) (MC more reliable for W/Z and top events than for QCD jet events) Can further ensure a <u>"jet pure"</u> sample by: requiring back-to-back jets

- p_T balance between jets: $\Delta p_T(jets) < p_T^{lead jet} / 2$
- very "jet-like" tag jet: $N_{track}^{tag-jet} > 4$ (or increase track cut for higher p_T)

No other requirements on probe jet - need to keep unbiased sample of jets

Fake rate should be computed for different E_T , η bins (maybe even as function of ϕ) Statistics will not be a problem (so will do better than MC studies!) Can generate background PDFs from data

The W $\rightarrow \tau v$ Channel

Despite ~10 greater signal yield for $W \rightarrow \tau v$ compared to $Z \rightarrow \tau \tau$, this is more challenging due to lack of additional τ -lepton leg

- heavily dependent on τ -lepton triggers and E_T^{MISS}
- in fact τ + E_T^{MISS} triggers used rates will have to be controlled (prescales for higher luminosity)

Backgrounds: $W \rightarrow ev$, $W \rightarrow \mu v$, $Z \rightarrow \tau \tau$, $Z \rightarrow ee$, top pairs, multi-jets

 $W \rightarrow \tau v$ channel will depend on the τ -ID fake rate measurement for background estimation (no same sign control sample)

QCD control samples can be defined with τ -candidates failing isolation criteria

Taus with Tops

τ-leptons coming from top pair production give a τ-sample with higher p_T than $Z \rightarrow \tau \tau$ Typical topology is $t\bar{t} \rightarrow WbW\bar{b} \rightarrow qq'b\tau \upsilon \bar{b}$ or $t\bar{t} \rightarrow WbW\bar{b} \rightarrow l\nu b\tau \upsilon \bar{b}$

Trigger paths: τ trigger, τ + E_T^{MISS} trigger, multi-jet trigger, lepton trigger

Can use b-tagging to help select top events since W+jets background is large

For lepton- τ channel, same sign events can be used as background control sample

Nice cross check to $Z \rightarrow \tau\tau$ for τ -ID efficiency and τ -trigger efficiency when using events selected with unbiased trigger

Could be the best study sample of $\tau\text{-leptons}$ with higher p_T

Perhaps even τ -energy scale determination is possible by comparing p_T^{top} in tau decay channels vs lepton decay channels

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Physics with τ -lepton final states offers many exciting possibilities

CMS/ATLAS have well developed algorithms and detailed MC studies for hadronically decaying $_{\tau}$ -lepton reconstruction and identification

CMS/ATLAS have infrastructure/plans for studying τ -ID with first data

A variety of trigger paths at both experiments involve τ -signatures, whether alone, multi- τ triggers or in combination with leptons objects/ E_{τ}^{MISS}

W/Z boson and top production are key channels in the early data era at LHC to understand τ -ID and measure: offline/trigger efficiency, energy scale, fake rates, other characteristics

Lots of work/opportunities for young researchers to make their impact on this important signature and exciting physics over the next few years!

> Many thanks to the Workshop Organizers for the invitation and the privelige to participate!

Further Reading

Further reading CMS:

"CMS TDR: Physics Performance" J. Phys. G34: 995 (2007)

"Particle-Flow Event Reconstruction in CMS and Performance for Jets, Taus, and $E_{\rm T}^{\rm MISS"}$ CMS PAS PFT-09/001

"CMS Strategies for Tau Reconstruction and Identification using Particle-Flow Techniques" CMS PAS PFT-08/001

Further reading ATLAS:

"Expected Performance of the ATLAS Experiment: Detector, Trigger, Physics" CERN-OPEN-2008-020

"The ATLAS Experiment at the CERN Large Hadron Collider" JINST 3 (2008) S08003

Detector Depth: Radiation Lengths

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Track Quality Criteria at ATLAS

	Leading Track	Default Track	Loose Track
р _т >	6 GeV	1 GeV	1 GeV
d ₀ <	2 mm	1 mm	2 mm
$ \Delta z $ *sin(θ) <	10 mm	1.5 mm	10 mm
N _{Si} ^{Hit} ≥	7	7	7
N _{Pix} ^{Hit} ≥	0	2	0
N _{Blay} ^{Hit} ≥	0	1	0

H1 Weighting Method (pedagogical example)

$$E' = w E$$

$$w = [c_1 \exp(-c_2 E/V) + c_3]$$

• $w \rightarrow 1$ for large E/V:

- $c_3 \approx 1$
- weighting does not change electromagnetic clusters
- small energy density dominated by hadronic activity: w > 1:
 - $c_{1,2} > 0$
 - exact values depend on total cluster energy, choice of weighted unit (cell or cluster), ...
- plot shows 30 GeV pions from 2002 EMEC–HEC test beam as a simple cluster weight example
 - restrict sample to pions fully contained in the EMEC
 - plot E_{beam} / E vs. E / V with E, V: cluster energy and volume, respectively
 - extract weight function
 - compare resolution for weighted and unweighted sample

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Energy: energy flow details at ATLAS

$$E_T^{eflow} = E_T^{emcl} + E_T^{neuEM} + \sum p_T^{track} + \sum resE_T^{chrgEMTrk} + resE_T^{neuEM}$$

E_T^{emcl}: EM clustered energy not matched to track

 $\Sigma p_T^{\text{track}}$: sum of track transverse momenta replaces E_T^{chrgEM} and E_T^{chrgHAD} E_T^{neuEM} : EM energy left over collected by cells in 0.2 of leading track

 $\Sigma resE_T^{chrgEMTrk}$: max(0, E_T^{chrgEM} -0.7xp $_T^{track}$) corrects for leakage of photon showers to cells which matched to charged hadrons

 $\label{eq:reserved} res E_T^{neuEM}: -0.1 \times p_T^{track} \\ corrects for double counting of of EM leakage from charged hadrons$

$\pi^{\rm O}$ reconstruction - shower subtraction

The reconstruction of π^0 clusters requires the removal of energy deposited by π^{\pm} mesons: cell energy subtraction using parametrized shower development

$$w_{i} = c_{L} \int f_{L}(\eta - \eta_{trk}, \varphi - \varphi_{trk}) \cosh(\eta) d\eta d\varphi$$

cellvolume

c_L: longitudnal weights (based on which layer in ECAL)

 f_L : lateral parametrization in (η, ϕ) space

Can test/tune on sample of $\tau^{\pm} \rightarrow \pi^{\pm} \upsilon$ true decays - should get zero energy in calorimeter after the subtraction

Stau Mass Measurement

Most important tau related SUSY parameter: $\tilde{\tau_1}$ mass measurement

Obtained from di-tau invariant mass spectrum in decay chain

$$ilde{\chi}^0_2 o ilde{ au}_1 au o ilde{\chi}^0_1 au^\pm au^\mp$$

Endpoint of undecayed tau mass spectrum depends on $\tilde{\tau_1}$ mass.

Escaping neutrino energy in tau decays makes this endpoint hardly directly measurable

Choose e. g. inflection point of trailing edge as endpoint sensitive observable (needs calibration: inflection point vs. endpoint)

Tau Polarization in SUSY

Polarisation of taus from stau decay $\tilde{\tau}_1^{\pm} \rightarrow \tau^{\pm} \tilde{\chi}_1^0$ is a probe of the $\tilde{\chi}_1^0$ composition and can be used to discriminate between different models of SUSY breaking:

• Universal SUGRA models:

• For most non-universal SUGRA models: $P_{ au}\simeq\cos^2 heta_{ au}-\sin^2 heta_{ au}$

- AMSB models:
- For many GMSB models:

 $P_{\tau} \simeq +1$

 $P \sim -1$

$$P_{\tau} = \sin^2 \theta_{\tau} - \cos^2 \theta_{\tau}$$

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