### Simple self-consistent prediction methods for the phase space of dark matter: from galactic dynamics to phenomenology

Thomas Lacroix

Collaborators: M. Boudaud, J. Lavalle, E. Nezri,

A. Núñez Castiñeyra, P. Salati, M. Stref

**TeVPA 2018** 

August 29, 2018







# Phase-space distribution of dark matter: source of theoretical uncertainties

#### Direct searches

$$\frac{\mathrm{d}R}{\mathrm{d}E} \propto \rho_{\odot} \int_{v_{\min} \leq |\vec{v}| \leq v_{\mathrm{esc}}} \frac{f_{\odot}(\vec{v})}{|\vec{v}|} \, \mathrm{d}^{3}v$$

Impact at low masses  $v_{\min} \sim v_{\rm esc}$ 

### Speed-dependent annihilation

•  $\langle \sigma v \rangle(r) \propto \langle v_{\rm r}^2 \rangle$  p-wave  $\sigma v = s_1 v_{\rm r}^2 \Rightarrow \langle \sigma v \rangle \equiv \langle \sigma v \rangle(r)$ 

$$\langle \sigma v \rangle(r) = s_1 \int d^3 v_1 d^3 v_2 f_r(\vec{v}_1) f_r(\vec{v}_2) v_r^2$$

•  $\langle \sigma v \rangle (r) \propto \langle 1/v_{\rm r} \rangle$  or  $\langle 1/v_{\rm r}^2 \rangle$ Sommerfeld

### Primordial black holes

• Gravitational microlensing event rates (EROS, MACHO)

$$\frac{\mathrm{d}\Gamma}{\mathrm{d}t} \propto \rho(r) \int v f(\vec{v}, \vec{r}) \, \mathrm{d}^3 v$$

- Merger rates (gravitational waves)
- + DM substructures (test masses), disruption of stellar binaries

### Standard approaches 1: "Standard halo model"

#### Standard halo model (SHM)

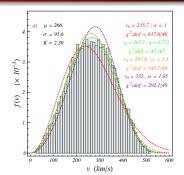
Maxwell-Boltzmann distribution

$$f(\vec{v}) = \frac{1}{v_c^3 \pi^{3/2}} e^{-\left(\frac{\vec{v}}{v_c}\right)^2}$$

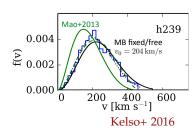
#### Oversimplification

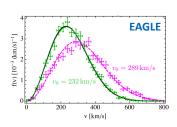
- Isothermal sphere
- Infinite system
- Ad hoc truncation at  $v_{\rm esc}$

### Standard approaches 2: direct fits to simulations

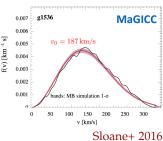


Ling+ 2010, Mollitor+ 2014





Bozorgnia+ 2016



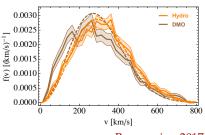
### Standard approaches 2: direct fits to simulations

### General insight

Generic features found in simulations (e.g., cusp/cores)

#### But insufficient approach

- Extrapolations based on fits at 8 kpc
- MW constrained systems (e.g., Gaia)
- Subgrid physics (recipes for star formation, ...)



Bozorgnia+ 2017

### Self-consistent approach required

Eddington-like methods: next-to-minimal approach

### Phase space of dark matter from first principles

### Phase-space distribution $f(\vec{v}, \vec{r})$ : closed system

• Collisionless Boltzmann equation, steady state

$$\{f, H\} = \vec{v} \cdot \frac{\partial f}{\partial \vec{r}} - \frac{\partial \Phi}{\partial \vec{r}} \cdot \frac{\partial f}{\partial \vec{v}} = 0$$

$$\longrightarrow$$
 Jeans' theorem:  $f \equiv f(I_1, ..., I_N)$  where  $\{I_i, H\} = 0$ 

Poisson equation

$$\Delta \Phi = 4\pi G \rho$$
 with  $\rho = \int f(\vec{v}, \vec{r}) d^3 v$ 

### Eddington's inversion formula (Eddington 1916)

Isotropic system + spherical symmetry:  $f(\vec{v}, \vec{r}) \equiv f(\mathcal{E})$ 

with 
$$\mathcal{E} = \Psi(r) - \frac{v^2}{2}$$
 and  $\Psi(r) = \Phi(R_{\text{max}}) - \Phi(r)$ 

$$f(\mathcal{E}) = \frac{1}{\sqrt{8}\pi^2} \left[ \frac{1}{\sqrt{\mathcal{E}}} \left( \frac{\mathrm{d}\rho}{\mathrm{d}\Psi} \right)_{\Psi=0} + \int_0^{\mathcal{E}} \frac{\mathrm{d}^2\rho}{\mathrm{d}\Psi^2} \, \frac{\mathrm{d}\Psi}{\sqrt{\mathcal{E}-\Psi}} \right]$$

### Phase space of dark matter from first principles

### Phase-space distribution $f(\vec{v}, \vec{r})$ : closed system

• Collisionless Boltzmann equation, steady state

$$\{f,H\} = \vec{v} \cdot \frac{\partial f}{\partial \vec{r}} - \frac{\partial \Phi}{\partial \vec{r}} \cdot \frac{\partial f}{\partial \vec{v}} = 0$$

$$\longrightarrow$$
 Jeans' theorem:  $f \equiv f(I_1, ..., I_N)$  where  $\{I_i, H\} = 0$ 

Poisson equation

$$\Delta \Phi = 4\pi G \rho$$
 with  $\rho = \int f(\vec{v}, \vec{r}) d^3 v$ 

### Eddington's inversion formula (Eddington 1916)

Isotropic system + spherical symmetry:  $f(\vec{v}, \vec{r}) \equiv f(\mathcal{E})$ 

with 
$$\mathcal{E} = \Psi(r) - \frac{v^2}{2}$$
 and  $\Psi(r) = \Phi(R_{\text{max}}) - \Phi(r)$ 

$$f(\mathcal{E}) = \frac{1}{\sqrt{8}\pi^2} \left[ \frac{1}{\sqrt{\mathcal{E}}} \left( \frac{\mathrm{d}\rho}{\mathrm{d}\Psi} \right)_{\Psi=0} + \int_0^{\mathcal{E}} \frac{\mathrm{d}^2\rho}{\mathrm{d}\Psi^2} \frac{\mathrm{d}\Psi}{\sqrt{\mathcal{E}-\Psi}} \right]$$

### **Extensions**

#### **Anisotropic systems**

E.g. constant anisotropy, Osipkov-Merritt (radial anisotropy)

### Going beyond spherical spherical symmetry

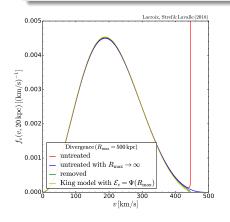
- Angle-action coordinates more suitable coordinate system if no spherical symmetry
  - Binney & Tremaine 1987
- Best way to account for complexity revealed by Gaia
- Level of refinement not necessarily required for DM searches
  - → Evaluate astrophysical uncertainties
- Eddington: lower level of technicalities to account for dynamical constraints
  - $\rightarrow$  Study validity range in detail

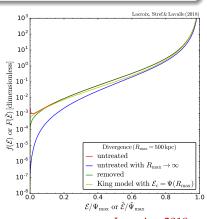
Eddington: method applied blindly to direct searches so far

### Theoretical consistency and radial boundary

### Imposing a radial boundary

- Finite system  $(R_{\text{max}}) \Rightarrow \text{divergence of } f(\vec{r}, \vec{v}) \text{ at } v_{\text{esc}} \text{ (from } 1/\sqrt{\mathcal{E}})$
- Phase-space compression
- $\bullet$   $v_{\rm esc}$  crucial (direct DM searches at low masses, stellar surveys)





### **Proper treatment critical for self-consistency**

#### Removing divergence by hand?

- Approach no longer self-consistent (Poisson equation)
- Poor reconstruction of density and mass in outer regions

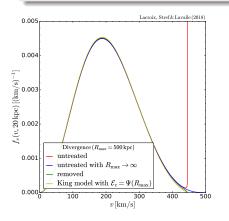
#### Even more critical for radial anisotropy

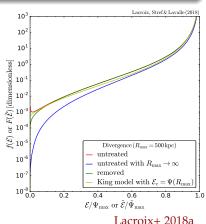
Very large discrepancy between original and reconstructed mass

### Theoretical consistency and radial boundary

### Regularization

- Modified profile, flat at R<sub>max</sub>
- Energy cutoff (King)





Not possible for radial anisotropy

### Theoretical consistency: instabilities

### Validity range of the method

• Standard criterion:

$$f \geqslant 0$$

- Antonov instabilities for some DM-baryon configurations
- Stable solution if  $\frac{\mathrm{d}f}{\mathrm{d}\mathcal{E}} > 0 \Leftrightarrow \frac{\mathrm{d}^2\rho}{\mathrm{d}\Psi^2} > 0$ Doremus+ 1971, Kandrup & Sygnet 1985
- Select mass models

Lacroix+ 2018a

For anisotropic systems criteria against radial perturbations only Doremus+ 1973

### **Impact on predictions for DM searches**

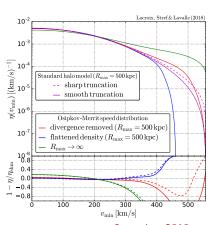
#### Event rate proportional to

$$\eta(v_{\min}) = \int_{v_{\min} \leqslant v \leqslant v_{\oplus} + v_{\mathrm{esc}}} \frac{f_{\vec{v}, \oplus}(\vec{v})}{v} \, \mathrm{d}^3 v$$

### Isotropic

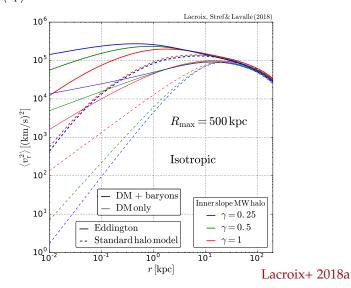
#### Lacroix, Stref & Lavalle (2018) 10-1 Standard halo model ( $R_{max} = 500 \, \text{kpc}$ ) 10-2 -- sharp truncation smooth truncation $[10^{-3})^{\prime\prime} (10^{-4})^{\prime\prime} = 10^{-6}$ Isotropic speed distribution $(R_{\text{max}} = 500 \text{ kpc})$ divergence removed flattened density $R_{\mathrm{max}} \rightarrow \infty$ 10<sup>-7</sup> King model with $\mathcal{E}_c = \Psi(R_{max})$ King model with $\tilde{\Psi}(R_{\text{max}}) = 0$ 10.8 0.0 200 300 400 500 $v_{\rm min} \, [{\rm km/s}]$

### Osipkov-Merritt



### **Impact on predictions for DM searches**

Prototypical case: p-wave annihilation  $\langle \sigma v \rangle(r) \propto \langle v_{\rm r}^2 \rangle$ 



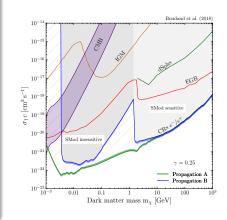
### **Impact on predictions for DM searches**

### Application: p-wave annihilation

$$\sigma v = \sigma_1 v_r^2 \Rightarrow \langle \sigma v \rangle \equiv \langle \sigma v \rangle(r)$$
$$\langle \sigma v \rangle(r) = \sigma_1 \int d^3 v_1 d^3 v_2 f_r(\vec{v}_1) f_r(\vec{v}_2) v_r^2$$

$$\Rightarrow \psi_{\rm e} \neq \langle \sigma v \rangle \int \rho^2(r) \, {\rm d}^3 r$$

- Very strong  $e^+$  constraints (Voyager, AMS-02)
- Justifies focusing on Eddington's methods
- Treatment of theoretical uncertainties



Boudaud+ 2018, in prep.

### Summary and outlook

### Eddington's inversion method

- A few physical assumptions
- Moderate level of technicalities
- Mass model direct input
- Better control astrophysical uncertainties for DM searches
- Dramatic changes wrt Maxwell-Boltzmann

### Self-consistency: theoretical validity range

- Radial boundary (direct searches)
- Positive DF + stability

#### Actual predictivity?

- Testing the method against cosmological simulations
- Not direct fits!
- Eddington globally better than SHM Lacroix+ 2018b, in prep.
- See A. Núñez's talk on Thursday

## Thank you for your attention!