

# **EXTRACTING NUCLEAR FORM FACTORS FROM COHERENT NEUTRINO SCATTERING**



EMILIO CIUFFOLI IMP, CAS, LANZHOU Based on a work in collaboration with J. Evslin, Q. Fu and J. Tang [1]

#### $\mathbf{C}\mathbf{E}\nu\mathbf{N}\mathbf{S}$

Coherent Elastic Neutrino-Nucleus Scattering (CE $\nu$ NS) was observed recently by the COHER-ENT collaboration, using neutrinos created via pion Decay At Rest ( $\pi$ DAR) [2]. The CE $\nu$ NS cross section is given by

$$\frac{d\sigma(E_{\nu}, E_{r})}{dE_{r}} = \frac{G_{F}^{2}[N - (1 - 4\sin^{2}\theta_{w})Z]^{2}F^{2}(Q^{2})M^{2}}{4\pi} \times \frac{1}{M}\left(1 - \frac{E_{r}}{E_{max}}\right)$$

Since at these energies  $(1 - 4 \sin^2 \theta_w) \simeq 0.045$ , the proton contribution is strongly suppressed; in particular the form factor depends almost exclusively on the neutron distribution. Studying  $CE\nu NS$  it is possible to extract 400 important information on the electroweak form factor and the neutron distribution [3]. Facilities: At China Spallation Neutron Source (CSNS) 7 during Phase I a 1.6 GeV, 100 kW pulsed proton beam hits a fixed target, creating neutrinos via  $\pi$ DAR as by-  $\frac{W}{2}$  200 products of the collision; the time structure of the beam  $\overline{2}$ (frequency 25 Hz) will reduce significantly the steady 100 state background. Neutrinos will also be produced at the CIADS facility, currently under construction as part of the China - Accelerator Driven System (C-ADS) project: here the energy will be lower (500 MeV) but, due to the higher power (2.5 MW), the neutrino flux will be around 5 times larger, however the beam will be con- Figure 1: Expected spectrum for 1 ton LAr tinuous, not pulsed, and the background considerably detector, 1 year lifetime higher. In the following calculations, the neutrino beam that will be produced at CSNS is used.



## **MODEL-INDEPENDENT ANALYSIS**

For a model-independent analysis, we considered a Taylor expansion of  $F^2(Q^2)$ ; each term  $Q^{2n}$  is multiplied by a factor proportional to the 2n-th momentum of the radius distribution,  $\langle R^{2n} \rangle$ . We calculated the  $1\sigma$  region in the  $\langle R^2 \rangle$ - $\langle R^4 \rangle$  plane for a 1-ton Argon detector, considering at first only the uncertainty on the total flux normalization, then taking into account also the  $\langle R^6 \rangle$  term and the uncertainty on the QF.

# NEUTRON DISTRIBUTION

We used the Helm model to describe the neutron distribution inside the nucleus; it is a twoparameters model that depends on the distribution radius *R* and the neutron skin thickness *s*. In this energy range the dependence on *s* is negligible: the sensitivity to the form factor is expressed as the 1- $\sigma$  bound on R







# **QUENCHING FACTOR**

We considered a simple linear model to describe the uncertainty on the quenching factor (QF):

$$E_{obs} = E_{real}(1 + \epsilon)$$



0.5 2.0  $\langle R_{\rm fit}^2 \rangle / \langle R_{\rm Ar}^2 \rangle$ **Figure 4:** 1-, 2- and 3- $\sigma$ 's regions in the  $\langle R^2 \rangle - \langle R^4 \rangle$ plane. Upper panel: only total flux uncertainty taken into account. Central panel: expansion up to  $\langle R^6 \rangle$  (treated as a pull parameter). Lower panel: best-fit values for  $\langle R^6 \rangle$  and  $\epsilon$  are used

### REFERENCES

[1] E. Ciuffoli, J. Evslin, Q. Fu, Qiang and J. Tang, arXiv:1801.02166 [physics.ins-det] D. Akimov *et al*. Science, 357(6356):1123-1126, 2017. [2]

[3] M. Cadeddu, C. Giunti, Y. F. Li, and Y. Y. Zhang, Phys. Rev. Lett., 120(7):072501, 2018.