

Pseudo-Goldstone scalars in the minimal $SO(10)$ Higgs model

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[Murayama, Pierce: 2001]
[Bajc, Lavignac, Mede: 2015]
- **Minimal SUSY $SO(10)$ (without 120_S)**
 - ▶ Fails to reproduce the absolute neutrino mass scale
[Bertolini, Schwetz, Malinský: 2006]
[Aulakh, Garg: 2006]

The minimal renormalizable $SO(10)$ GUT

Matter fields

$$16_F = (1, 2, -\tfrac{1}{2}) \oplus (\bar{3}, 2, \tfrac{1}{3}) \oplus (3, 2, \tfrac{1}{6}) \oplus (\bar{3}, 1, -\tfrac{2}{3}) \oplus (1, 1, 1) \oplus (1, 1, 0)$$

$\downarrow$	$\downarrow$	$\downarrow$	$\downarrow$	$\downarrow$	$\downarrow$
L_L	d_L^c	Q_L	u_L^c	e_L^c	N_L^c
$\underbrace{\hspace{10em}}$		$\underbrace{\hspace{10em}}$			$\underbrace{\hspace{5em}}$
$\bar{5}$		10			1

- Automatic anomaly cancellation

The minimal renormalizable $SO(10)$ Higgs model

Gauge fields

$$\begin{array}{ccccccc} 45_G = & (8, 1, 0) & \oplus & (1, 3, 0) & \oplus & (1, 1, 0) & \oplus & (1, 1, 0) & \oplus & (3, 1, \frac{2}{3}) & \oplus & (3, 2, -\frac{5}{6}) & \oplus \\ & \downarrow & & \downarrow & & \downarrow & & & & & & & & \\ & G_\mu^b & & A_\mu^a & & B_\mu, Y_\mu & & & & (3, 2, \frac{1}{6}) & \oplus & (1, 1, 1) & + h.c. \\ & & & & & & & & & \downarrow & & & \\ & & & & & & & & & X_\mu & & & \end{array}$$

Scalar sector

- GUT scale SSB by 45_S
 - ▶ Minimal, preserves rank
 - ▶ Robust GUT scale estimates with respect to gravitational effects
 - ▶ Two real SM singlets:

$$\langle (1, 1, 0)_{15} \rangle = \sqrt{3}\omega_{BL} \quad \langle (1, 1, 0)_1 \rangle = \sqrt{2}\omega_R$$

- Rank reduction by 126_S
 - ▶ Seesaw scale
 - ▶ Renormalizable Yukawa interaction $\leftarrow 16 \times 16 \supset 126, 16$ & ~~Witten's loop~~
 - ▶ One complex SM singlet

$$\langle (1, 1, 0)_{\overline{10}} \rangle = \sqrt{2}\sigma$$

- ▶ Phenomenology $\Rightarrow |\sigma| \ll \max(\omega_{BL}, \omega_R) \equiv \omega_{max}$

Tree level scalar spectrum

Scalar fields $45_S \oplus 126_S \leftrightarrow \Phi = (\phi_{ij}, \Sigma_{ijklm}, \Sigma_{ijklm}^*)$ form the most general scalar potential

$$V_0(\phi, \Sigma, \Sigma^*) = V_\phi(\phi) + V_\Sigma(\Sigma, \Sigma^*) + V_{\phi, \Sigma}(\phi, \Sigma, \Sigma^*)$$

$$V_\phi = -\frac{\mu^2}{4} (\phi\phi)_0 + \frac{a_0}{4} (\phi\phi)_0 (\phi\phi)_0 + \frac{a_2}{4} (\phi\phi)_2 (\phi\phi)_2,$$

$$\begin{aligned} V_\Sigma = & -\frac{\nu^2}{5!} (\Sigma\Sigma^*)_0 + \frac{\lambda_0}{(5!)^2} (\Sigma\Sigma^*)_0 (\Sigma\Sigma^*)_0 + \frac{\lambda_2}{(4!)^2} (\Sigma\Sigma^*)_2 (\Sigma\Sigma^*)_2 + \\ & + \frac{\lambda_4}{(3!)^2 (2!)^2} (\Sigma\Sigma^*)_4 (\Sigma\Sigma^*)_4 + \frac{\lambda'_4}{(3!)^2} (\Sigma\Sigma^*)_{4'} (\Sigma\Sigma^*)_{4'} + \\ & + \frac{\eta_2}{(4!)^2} (\Sigma\Sigma)_2 (\Sigma\Sigma)_2 + \frac{\eta_2^*}{(4!)^2} (\Sigma^*\Sigma^*)_2 (\Sigma^*\Sigma^*)_2, \end{aligned}$$

$$\begin{aligned} V_{\phi, \Sigma} = & \frac{i\tau}{4!} (\phi)_2 (\Sigma\Sigma^*)_2 + \frac{\alpha}{2 \cdot 5!} (\phi\phi)_0 (\Sigma\Sigma^*)_0 + \frac{\beta_4}{4 \cdot 3!} (\phi\phi)_4 (\Sigma\Sigma^*)_4 + \\ & + \frac{\beta'_4}{3!} (\phi\phi)_{4'} (\Sigma\Sigma^*)_{4'} + \frac{\gamma_2}{4!} (\phi\phi)_2 (\Sigma\Sigma)_2 + \frac{\gamma_2^*}{4!} (\phi\phi)_2 (\Sigma^*\Sigma^*)_2. \end{aligned}$$

Tree level mass matrix contains tachyons

$$\begin{aligned} M_S^2[(8, 1, 0)] &= 2a_2(\omega_{BL} - \omega_R)(\omega_R + 2\omega_{BL}), \\ \underbrace{M_S^2[(1, 3, 0)] &= 2a_2(\omega_R - \omega_{BL})(2\omega_R + \omega_{BL})}_{\text{nontachyonic iff near } SU(5)' \times U(1)_{Z'} \text{ SSB chain}} \end{aligned}$$

- SM multiplets $(8, 1, 0), (1, 3, 0) \leftrightarrow$ pseudo-Goldstone bosons of the broken global symmetry $O(45)$ in the limit $\sigma \rightarrow 0, a_2 \rightarrow 0$
- If $a_2 \ll 1$, loop corrections dominant [Bertolini, Di Luzio, Malinský: 2010]
[Gráf, Malinský, Mede, Susič: 2017]
- Makes sense at one-loop level [Kolešová, Malinský: 2014]

One more pseudo-Goldstone

SM singlet 2×2 mass submatrix (only 45_S) has an eigenvalue

$$m_{PG}^2 = a_2 \left(-\frac{45\omega_{BL}^4}{3\omega_{BL}^2 + 2\omega_R^2} + 13\omega_{BL}^2 - 2\omega_{BL}\omega_R - 2\omega_R^2 \right) + O(a_2^2) + O\left(\frac{\sigma^2}{\omega_{max}^2}\right)$$

- $m_{PG}^2 \propto a_2$ up to $O\left(\frac{\sigma^2}{\omega_{max}^2}\right)$ corrections
- Significant one-loop corrections to m_{PG}^2
 - ▶ One-loop effective potential calculation

$$\begin{aligned}
\mathcal{M}_{11}^2 = & \underbrace{-2a_2 \omega^2[-2, 1, 1] + 24a_0\omega_{BL}^2 - 12\sigma^2\beta'_4}_{\text{Tree level}} + \frac{g^4}{16\pi^2} (\omega^2[40, 1, 13] + \\
& + 13\omega_R^2 - 6\sigma^2) + \frac{3g^4}{32\pi^2\omega[1, -1]} \left\{ \omega[1, -1]^3 \log \left[\frac{g^2\omega[1, -1]^2}{2\mu_R^2} \right] + \right. \\
& + (4\sigma^2\omega[3, -1] + \omega[7, -5]\omega[1, 1]^2) \log \left[\frac{g^2(4\sigma^2 + \omega[1, 1]^2)}{2\mu_R^2} \right] - \\
& - 16(\sigma^2\omega_R - 2\omega_{BL}^3 + 3\omega_{BL}^2\omega_R) \log \left[\frac{2g^2(\sigma^2 + \omega_{BL}^2)}{\mu_R^2} \right] + \\
& \left. + 8\omega_R(\sigma^2 + \omega_R^2) \log \left[\frac{2g^2(\sigma^2 + \omega_R^2)}{\mu_R^2} \right] \right\} + \frac{35\tau^2}{8\pi^2} + \\
& + \frac{\omega^5[-5, 19, -170, 1222, -5, 19]\omega_R}{16R^2\pi^2} \beta_4^2 + \frac{15\omega[1, 1]\omega_R}{4\pi^2} \beta_4\beta'_4 + \\
& + \frac{5\omega_R\omega[-5, 11]}{4\pi^2} \beta_4'^2 +
\end{aligned}$$

$$\begin{aligned}
& + \frac{1}{16\omega[1, -1]\pi^2} F \left[-9\omega[1, 1], -36\omega_{BL}\omega[1, -1], \omega^2[-18, 72, 24], 36\omega_{BL}\omega_R, \right. \\
& \quad 108\omega_{BL}^2\omega[1, -1], 144\omega_{BL}^2\omega[1, -1], 144\omega_{BL}^2\omega[1, -1], 24\omega_{BL}^2\omega[2, -3], \\
& \quad \left. 12\omega^3[8, -15, -12, -4], 24\omega^3[2, 0, 3, 1] \right] \log \left[\frac{2(\beta_4\omega_{BL}^2 + \omega[2, 1](2\beta_4'\omega[1, 1] - \tau))}{\mu_R^2} \right] \\
& + \frac{1}{16\omega[1, -1]\pi^2} F \left[-3\omega[3, 5], -36\omega_{BL}\omega[1, -1], 0, 24\omega^2[3, -6, -1], 36\omega_{BL}^2\omega[1, -1], \right. \\
& \quad \left. 0, -144\omega_{BL}^2\omega[1, -1], 0, 0, 72\omega^3[2, -2, 3, 1] \right] \log \left[\frac{2\omega[3, 1](2\omega_R\beta_4' - \tau)}{\mu_R^2} \right] \\
& + \frac{1}{8\omega[1, 1]\pi^2} F \left[3\omega[3, 1], 36\omega_{BL}\omega[1, -1], 0, -24\omega_R\omega[3, 1], 108\omega_{BL}^2\omega[1, -1], 0, \right. \\
& \quad \left. 144\omega_{BL}^2\omega[1, -1], 0, 0, 24\omega^3[2, 0, 3, 1] \right] \log \left[\frac{2\omega[1, 1](2\beta_4'\omega[2, 1] - \tau)}{\mu_R^2} \right] + \dots
\end{aligned}$$

Consistency checks

Particular symmetry breaking limits ($\sigma \rightarrow 0$):

- $\mathbf{SU(5)} \times \mathbf{U(1)_Z}$ ($\omega_R \rightarrow \pm\omega_{BL}$):

$$(1, 1, 0)_{PG}, (8, 1, 0), (1, 3, 0) \subset (24, 0),$$

- $\mathbf{SU(4)_c} \times \mathbf{SU(2)_L} \times \mathbf{U(1)_R}$ ($\omega_{BL} \rightarrow 0$):

$$(1, 1, 0)_{PG}, (8, 1, 0) \subset (15, 1, 0),$$

- $\mathbf{SU(3)_c} \times \mathbf{SU(2)_L} \times \mathbf{SU(2)_R} \times \mathbf{U(1)_{BL}}$ ($\omega_R \rightarrow 0$):

$$(1, 1, 0)_{PG} \sim (1, 3, 0) \longleftarrow \text{D-parity.}$$

Summary and Outlook

- The minimal renormalizable $SO(10)$ Higgs model:
 - ▶ gauge bosons 45_G ,
 - ▶ scalar sector $45_S \oplus 126_S$.
- Potentially realistic scenarios involve tachyonic instabilities.
- Existence of tachyons is relic of the tree-level calculations.
- Scalar spectrum contains pseudo-Goldstones - $(8, 1, 0)$, $(1, 3, 0)$ and newly calculated $(1, 1, 0)$.
- One-loop corrections to the pseudo-Goldstone masses improve the proton lifetime estimates.