

From stringy vacua with particle physics spectrum to the effective action - exemplified by a non-factorisable orientifold -

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Motivation: String Phenomenology & Cosmology

Status quo: *string theory* \equiv *unification of all interactions*

▶ **string pheno:** (*very roughly*)

- ▶ many compact CY and T^6/Γ models with **SM-like spectrum topology** \leftrightarrow **chiral matter** \rightsquigarrow *machine learning techniques*

see Vaudrevange's talk

- ▶ **rudimentary** results on effective field theory from **dim.**

reduction of SUGRA & DBI $\rightsquigarrow S_{EH}$ & $1/g_{U(1)/SU(N)}^2$

- ▶ depends on size & shape moduli of compact space, M_{string}
- ▶ already **Yukawas challenging** - only computable with **CFT techniques** on T^6/Γ

see Liyanage's talk

▶ **string cosmo:** (*very roughly*)

- ▶ assume that **particle physics** sector is **localized** & OK ?
- ▶ choose 'global' CY geometry to explain **inflation** ?
- ▶ 'dial' terms in the action for moduli stabilisation ?

see Valenzuela's, Grimm's, Westphal's talks

moduli dep. of action relevant to particle physics & cosmology

- ▶ Motivation ✓
- ▶ Stringy vacua of particle physics
 - ▶ SM & new physics matter
 - ▶ effective action
 - ▶ example: non-factorisable orbifold
- ▶ Interplay with cosmology
 - ▶ constraints on moduli from $1/g^2$, 1-loop
- ▶ Conclusions

Stringy vacua of particle physics

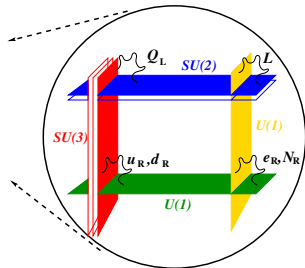
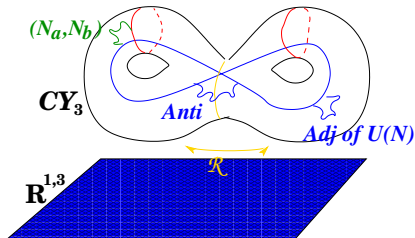
Stringy Vacua of Particle Physics: Matter Sector

Type II string theories: - and dual het. $SO(32)$ -

- ▶ **closed** strings \rightsquigarrow gravity, moduli

& some $U(1)$'s related to internal symmetries

- ▶ **open** strings \rightsquigarrow $\prod U(N)$ groups & matter in **Bifund** & **Adj**
orientifolding: $U(N) \rightarrow SO(2N)$ or $USp(2N)$ & **(Anti)Sym**



- ▶ **restricts** ways to realise the **SM** & **new physics**
 - ▶ two endpoints \rightsquigarrow constraints on reps.
 - ▶ at least 3 kinds of D-branes
 - ▶ all couplings determined by D-brane configuration

Stringy Vacua of Particle Physics cont'd

Stringy consistency conditions:

- ▶ fix D-brane configuration - dependent on CY or T^6/Γ background
 - ▶ complete (massless) matter spectrum determined
- ▶ ensure absence of all gauge & gravitational anomalies
 - ▶ some $U(1)$'s acquire mass
 - ▶ some \mathbb{Z}_n 's $\subset U(1)$'s can survive \rightsquigarrow selection rules on couplings
- ▶ SUSY ensures stability of vacuum @ M_{string}
 - ▶ requires SUSY below M_{string}

above relies on:

- ▶ **topology** of compact space

but more info on **geometry** required for:

- ▶ Yukawas $\propto e^{-\text{Area}_{xyz}}$
- ▶ $1/g_{x,\text{tree}}^2 \propto \text{Vol}(\Pi_x)$ - D_x -brane extends along Π_x in compact space
- ▶ ... higher loop-orders and n -point couplings \rightsquigarrow all kinds of moduli involved!

Stringy Vacua of Particle Physics: Effective Action

formal dimensional reduction of SUGRA & DBI:

- ▶ for D6-branes: $\frac{1}{g_{\text{tree}}^2} \propto \frac{\text{Vol}(\Pi_3)}{\sqrt{\text{Vol}(\text{CY}_3)}} \frac{M_{\text{Planck}}}{M_{\text{string}}} \text{ (compl. str. moduli)}$

beyond this reduction:

- ▶ only possible by CFT techniques $\leftrightarrow T^6/\Gamma$ or 'non-geometric'
- ▶ Yukawas from disks w/ 3 vertex op. @ boundary (Kähler moduli)
- ▶ $\frac{1}{g_{1\text{-loop}}^2}$ also from magnetically gauged vacuum partition fct.
 - ▶ for D6-branes at angles ϕ : $\sim \ln\left(\frac{\Gamma(\phi)}{\Gamma(1-\phi)}\right)$ (compl. str. moduli)
 - ▶ for parallel D6-branes: KK & winding sums (Kähler moduli)

$$\Lambda_{0,0}(\nu) \equiv -\frac{1}{4\pi} \ln(\eta(i\nu)) \xrightarrow{\nu \gg 1} \frac{\nu}{48}$$

$$\Lambda_{\tau,\sigma}(\nu) \equiv -\frac{1}{4\pi} \ln\left(e^{-\frac{\pi\sigma^2\nu}{4}} \frac{|\vartheta_1\left(\frac{\tau-i\sigma\nu}{2}, i\nu\right)|}{\eta(i\nu)}\right) \xrightarrow{\nu \gg 1} \frac{[3(1-\sigma)^2 - 1]\nu}{48} - \delta_{\sigma,0} \frac{\ln[2\sin(\frac{\pi\tau}{2})]}{4\pi}$$

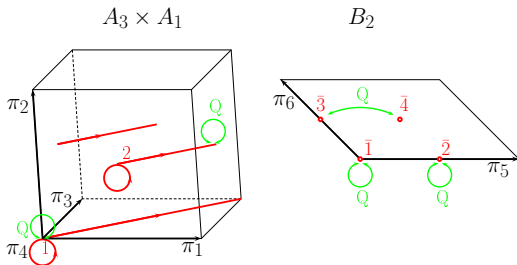
linear in Kähler modulus ν

- ▶ formulas *a priori* for factorisations $T^6 = (T^2)^3$

Example: Non-Factorizable Orientifold T^6/\mathbb{Z}_4

Why \mathbb{Z}_4 on $A_3 \times A_1 \times B_2$?

- ▶ access to compact backgrounds beyond $(T^2)^3/\Gamma$
- ▶ less twisted moduli require stabilisation: $(h_{11}^{\mathbb{Z}_2}, h_{21}^{\mathbb{Z}_2}) = (6, 2)$
vs. $(10, 6)$ for $B_2 \times (A_1)^2 \times B_2$
- ▶ \mathbb{Z}_4 action given by Coxeter element on Lie algebra lattices



- ▶ construction of 3-cycles, SUSY & stringy consistency✓
 - ▶ only 2 or 4 generation models
 - ▶ but good playground
- \rightsquigarrow non-factorisable \mathbb{Z}'_6 on $D_4 \times A_2$ has 3 generation models

(preliminary) Seifert '18

Example: Non-Factorizable Orientifold cont'd

Example: 4-generation model

► $SU(4)_a \times USp(2)_b \times USp(2)_c \times USp(4)_h$

► $[C] = 4 \times \left[(4, \bar{2}, 1, 1) + (\bar{4}, 1, 2, 1) \right]$

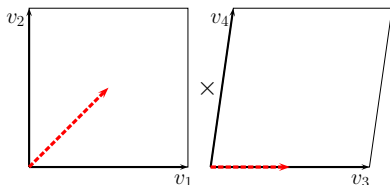
↖ from topological intersection #

↙ from CFT

$$[V] = (15, 1, 1, 1) + (1, 3_s, 1, 1) + (1, 1, 3_s, 1) + 2 \times (1, 1, 1, 10_s) + (1, 1, 1, 6_A) \\ + (5 \times (6_A, 1, 1, 1) + (10_s, 1, 1, 1) + c.c.) + 3 \times (1, 2, 2, 1) + (1, 2, 1, 4) + (1, 1, 2, 4)$$

► CFT methods known for factorisation $(T^2)^3$

► $(A_3 \times A_1) \times B_2 \simeq (T^2 \times T^2) \times T^2$ by shift symmetry



Example: Non-Factorizable Orientifold cont'd

► **@ tree-level:** $\frac{16\pi^2}{g_{x,\text{tree}}^2} = \frac{M_{\text{Planck}}}{M_{\text{string}}} \times \begin{cases} 2\sqrt{u_2} & \text{for } x = a \\ \frac{1}{2\sqrt{u_2}} & \text{for } x = b, c, h \end{cases}$

► **1-loop threshold corrections:**

$$\begin{aligned} \Delta_{SU(4)_a} &= 4(\tilde{\Delta}_{aa}^{\text{total}} + \tilde{\Delta}_{aa'}^{\text{total}}) + 2\tilde{\Delta}_{ab}^{\text{total}} + 2\tilde{\Delta}_{ac}^{\text{total}} + 4\tilde{\Delta}_{ah}^{\text{total}} + \Delta_{a,\Omega\mathcal{R}}^{\text{total}} \\ &= -4\Lambda(0, 0, \mathbf{v}_1, 2) - 4\Lambda(1, 0, \mathbf{v}_1, 2) + 4\hat{\Lambda}(1, 0, \mathbf{v}_1, 4) \\ &\quad - 16\Lambda(0, 0, \mathbf{v}_2, u_2) - 16\Lambda(1, 0, \mathbf{v}_2, u_2) + 8\hat{\Lambda}(0, 0, \mathbf{v}_2, 2u_2) + 8\hat{\Lambda}(1, 0, \mathbf{v}_2, 2u_2) \\ &\quad - 4\Lambda(0, 0, \mathbf{v}_3, 2) - 4\Lambda(1, 1, \mathbf{v}_3, 2) + 4\hat{\Lambda}(1, 1, \mathbf{v}_3, 4) \end{aligned}$$

$$\Delta_{USp(2)_{b/c}} = -4\Lambda(0, 0, \mathbf{v}_1, 2) - 3\Lambda(0, 0, \mathbf{v}_2, 1/u_2) + 2\hat{\Lambda}(0, 0, 2\mathbf{v}_2, 2/u_2) - 4\Lambda(0, 0, \mathbf{v}_3, 2) + \frac{\ln 2}{2}$$

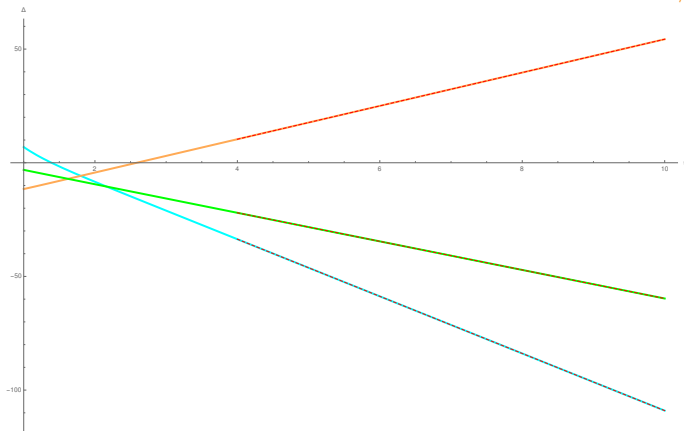
$$\Delta_{USp(4)_h} = -4\Lambda(1, 0, \mathbf{v}_1, 2) - 2\Lambda(0, 0, \mathbf{v}_2, 1/u_2) + 2\hat{\Lambda}(0, 0, 2\mathbf{v}_2, 2/u_2) - 4\Lambda(1, 1, \mathbf{v}_3, 2) + \frac{\ln 2}{2}$$

$\hat{\Lambda}(\sigma, \tau, \dots)$ from Möbius strip - exact shape **unknown**

Example: Pati-Salam w/ 4 Gen. cont'd: Gauge Thresholds

- ▶ (unrealistic) assumption $\hat{\Lambda}(\sigma, \tau, \dots) = \Lambda(0, 0, \dots)$ for $SU(4)_a$
- ▶ universal Kähler modulus $v \equiv v_{i=1,2,3}$

$USp(2)_{b/c}$, $USp(4)_h$



\rightsquigarrow enhancement/reduction of $\frac{1}{g_x^2}$ possible: fct.(Kähler modulus v)

Interplay with cosmology

Interplay with cosmology

Just to mention...

- ▶ QCD axion from open strings \rightsquigarrow EFT determined
trans-Planckian regime **not** in SUGRA approximation

Here more on hierarchies of scales:

- ▶ **tree-level:**
$$\frac{16\pi^2}{g_{x,\text{tree}}^2} = \frac{M_{\text{Planck}}}{M_{\text{string}}} \times \begin{cases} 2\sqrt{u_2} & \text{for } x = a \\ \frac{1}{2\sqrt{u_2}} & \text{for } x = b, c, h \end{cases}$$

with $u_2 \propto \frac{R_1}{R_3}$ - e.g. hierarchy $M_{\text{Planck}} \gg M_{\text{string}}, R_1 \ll R_3$ (\checkmark)

- ▶ **1-loop threshold correction:**

$$v_2 \propto R_1 R_3 > 1$$

$$\Delta_{SU(4)_a} \rightarrow \frac{2\pi}{3} \left((1 - 2c_{1,0}^{(1)}) v_1 - 4c_{1,0}^{(2)} v_2 + (1 - 2c_{1,1}^{(3)}) v_3 \right) + \text{const.}$$

$$\Delta_{USp(2)_{b/c}} \rightarrow \frac{4\pi}{3} \left(v_1 - \frac{v_2}{4} + v_3 \right) + \text{const.}$$

$$\Delta_{USp(4)_h} \rightarrow -\frac{2\pi}{3} (v_1 + v_2 + v_3) + \text{const.}$$

\rightsquigarrow upper bound on **LARGE volume scenario**: $\frac{1}{g_{b,c,h,1\text{-loop}}^2} \stackrel{v_2 \gg 1}{\leq} 0 \lll$

Conclusions & Outlook

Conclusions:

- ▶ stringy **particle physics**: chiral states ✓
vector-like states (also massive) only with CFT techniques
 \rightsquigarrow loop corrections computable
- ▶ **moduli** can be
 - ▶ absent by construction (here: twisted ones)
 - ▶ stabilised by D-branes \rightsquigarrow FI terms
- ▶ **no separation** of **particle** \leftrightarrow **cosmology** properties of string vacua due to loop effects

Outlook:

- ▶ rigorous derivation of stringy EFT needed
- ▶ popular simple cosmological scenarios need to be reconsidered