



# Anomaly Cancellation in Effective Supergravity Theories from the Heterotic String: Two Simple Examples

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## Supergravity from the $E_8 \times E_8$ Heterotic String

• Goal was to study anomalies arising from the massless spectrum of the  $\rm Z_3$  and  $\rm Z_7$  toroidal orbifold compactifications of the  $\rm E_8~x~E_8$  Heterotic String

$$\mathbb{R}^{9,1} \to \mathbb{R}^{3,1} \times X_6 \qquad E_8 \times E_8 \to E_8 \times E_6 \times G$$
$$X_6 = T^6/\mathbb{Z}_3 \qquad \mathbb{Z}_3 : G = SU(3)$$
$$\mathbb{Z}_7 : G = U(1) \times U(1)$$

• The 4D spectrum is N=1 supergravity with vector and chiral matter supermultiplets.

• The decomposition from compactificaton

$$\{g_{MN}, B_{MN}\} \rightarrow \{g_{\mu\nu}, B_{\mu\nu}, T^{ij}, \dots\}$$

• Gives the so-called diagonal Kahler moduli:

$$T_i = G_{i\bar{i}} + iB_{i\bar{i}}$$

• With Kahler potential

$$K = e^{g^{a}} \bar{\Phi}^{a} \Phi^{a} + g(T^{i}, \bar{T}^{i})$$
$$g(T^{i}, \bar{T}^{i}) = \sum_{i=1}^{3} g^{i}(T^{i}, \bar{T}^{i}) = -\sum_{i=1}^{3} \ln(T^{i} + \bar{T}^{i})$$
$$g^{a}(T^{i}, \bar{T}^{i}) = \sum_{i} q_{i}^{a} g^{i}(T^{i}, \bar{T}^{i})$$

• Anomalies arise from the SL(2,R) modular transformation of the three diagonal Kahler moduli:

$$T^i \to \frac{a_i T_i - ib_i}{ic_i T_i + d_i} \qquad \qquad a_i b_i - c_i d_i = 1$$

• Chiral matter supermultiplets must also transform

$$\Phi^a \to \exp\left(-\sum_i q_i^a F^i\right) \Phi^a$$
  $F_i = ic_i T_i + d_i$ 

• These field redefinitions in turn force a Kahler transformation:

$$K \to K + \sum_{i} (F^i + \bar{F}^i) = K + F + \bar{F}$$

- Thus the ``modular'' anomalies are a combination of Kahler + reparametrization anomalies and we get modular/gauge/gauge, modular/gravity/gravity, and pure modular anomalies.
- Our project showed that the anomalies in these supergravity theories has the correct universal form such that they can be cancelled via the 4D Green-Schwarz mechanism

### Motivation

- Modular transformations correspond to T-Duality transformations of the underlying string theory – consistency of the low energy QFT and its relation to string theory relies on the cancellation of anomalies
- On the string side, showing anomaly cancellation is easy, but one would want to ensure that the cancellation holds up purely from a QFT point of view.
- However, the QFT calculation is subtle. Previous investigations studied only subsets of the total anomaly or found a non-universal total anomaly due to difficulties with calculating the pure modular anomaly.
- Also phenomenological reasons for wanting a preserved symmetry residual R-parity

## Anomaly Cancellation Outline

- Calculate the one-loop effective action
- Introduce Pauli-Villars fields with non-invariant masses to cancel divergences\*
- Anomaly arises from variation of terms involving PV mass terms in the effective action that are finite as the masses are taken to infinity
- Cancel the result via the GS mechanism
- Example:  $\mathcal{L}_{YM} = -\sqrt{g} \frac{s}{8} \sum_{a} \left( F^{a}_{\mu\nu} - i\tilde{F}^{a}_{\mu\nu} \right) F^{\mu\nu}_{a} + \text{h.c.}$   $\Delta \mathcal{L}_{YM1} = -\frac{\sqrt{g}}{64\pi^{2}} \left( 8\pi^{2}b \sum_{i} F^{i} \right) \left( F^{a}_{\mu\nu} - i\tilde{F}^{\mu\nu}_{a} \right) F^{\mu\nu}_{a} + \text{h.c.}$
- Cancellation by imposing:

#### One-Loop Lagrangian

$$\mathcal{L}_{eff} = \mathcal{L}(g, K_R) + \sqrt{g} \frac{\ln(\Lambda^2)}{32\pi^2} (L_D + L_F) + \mathcal{L}_Q$$

• F-Term:

$$\Phi = \Phi_1 + \Phi_2 + \frac{3}{2}C_a\Phi_{YM}^a + \frac{1}{36}\Phi_X(N - 9N_G - 79) + \frac{1}{6}\Phi_w(41 + N - 3N_G) + N_G\Phi_g'$$

$$\Phi_{1} = -\frac{1}{2} \left( C_{a}^{M} \Phi_{YM}^{a} + \Gamma_{j}^{i\alpha} \left( \Gamma_{i\alpha}^{j} + 2(T^{a})_{i}^{j} W_{\alpha}^{a} \right) \right) \qquad \qquad \Gamma_{j\alpha}^{i} = -\frac{1}{8} (\bar{\mathcal{D}}^{2} - 8R) \mathcal{D}_{\alpha} Z^{k} \Gamma_{jk}^{i}$$

$$\Phi_{2} = \frac{1}{3} X^{\alpha} \left( \Gamma_{\alpha} + 2(T^{a})_{i}^{i} W_{\alpha}^{a} \right) \qquad \qquad X_{\alpha} = -\frac{1}{8} (\bar{\mathcal{D}}^{2} - 8R) \mathcal{D}_{\alpha} K$$

• Quadratic:

$$\mathcal{L}_Q = -\sqrt{g} \frac{\Lambda^2}{32\pi^2} \left( \frac{1}{2} (3 + N_G - N) \mathcal{D}^{\alpha} X_{\alpha} | + (\hat{V} + M^2) (7 + 3N_G - N) + N_G \mathcal{D}^{\alpha} k_{\alpha} + \mathcal{D}^{\alpha} \Gamma_{\alpha} | + \frac{2}{x} \mathcal{D}_a \operatorname{Tr}(T^a) \right)$$

#### PV Fields

• A variety of fields are needed to cancel divergences and produce the correct anomaly. Many have Kahler potential & superpotential terms based on the couplings of the matter fields. As an example, a subset of chiral PV fields with non-invariant masses have

$$W(\Phi^P, \Phi'^P) = \mu_P \Phi^P \Phi'^P$$
$$K(\Phi^P, \Phi'^P) = e^{f^P} |\Phi^P| + e^{f'^P} |\Phi'^P|$$
$$f^P = \alpha^P K(Z, \bar{Z}) + \beta^P g(T, \bar{T}) + \delta^P k(S, \bar{S}) + \sum_n q_n^C g^n(T^n, \bar{T}^n)$$

# Cancellation of Divergences

- To cancel divergences we require
  - Quadratic Divergences:

$$\operatorname{Tr}(\eta\Gamma_{\alpha}) = 0$$
  $\Gamma_{j\alpha}^{i} = -\frac{1}{8}(\bar{\mathcal{D}}^{2} - 8R)\mathcal{D}_{\alpha}Z^{k}\Gamma_{jk}^{i}$ 

• Logarithmic Divergences:

$$\operatorname{Tr}(\eta\Gamma_{\alpha}\Gamma_{\beta}) = \operatorname{Tr}(\eta\Gamma_{\alpha}T^{a}) = \operatorname{Tr}(\eta(T^{a})^{2}) = 0$$

• Linear Divergences:

$$\operatorname{Tr}(\eta \phi G \cdot \tilde{G}) = 0$$

- Where  $G_{\mu\nu}$  is the field strength associated with the fermion connections  $G_{\mu\nu} = [D_{\mu}, D_{\nu}] = -\Gamma^C_{D\mu\nu} + iF^a_{\mu\nu}(T^a)^C_D + \frac{1}{2}X_{\mu\nu}\delta^C_D$
- And is the parameter for a chiral rotation corresponding to a T-Duality transformation:  $\phi^{C} = \left(\frac{1}{2} - \alpha^{C} - \beta^{C}\right)F - \sum_{i}F^{i}q_{i}^{C}$ 9 of 14

#### Some Sum Rules

$$\sum_{C} \eta^{C} = -N - 29$$

$$\sum_{C} \eta^{C} f^{C} = -10K - \sum_{n} q_{n}^{p} g^{n}$$

$$\sum_{C} \eta^{C} \alpha^{C} \alpha^{C} = -4$$

$$\sum_{\gamma} \eta_{\gamma}^{V} = -12 - N_{G}$$

$$30\delta_{nm} = \sum_C \eta^C q_n^C q_m^C (1 - 2\bar{\gamma}^C)$$

10 of 14

### Anomaly Result

• After satisfying the regularization conditions & anomaly matching conditions, we find

$$\begin{split} \delta \mathcal{L}_{anom} &= b \int d^4 \theta E (F + \bar{F}) \Omega \\ &= \frac{b \sqrt{g} \mathrm{Im}(F)}{4} \left( F_a \cdot \tilde{F}^a + R \cdot \tilde{R} + \frac{1}{2} \sum_i g^i \cdot \tilde{g}^i + \cdots \right) \\ &8 \pi^2 b = \frac{1}{24} \left( 2 \sum_p q_n^p - N + N_G - 21 \right) \\ &= C_a - C_a^M + 2 \sum_b C_a^b q_n^b \\ &= 30 \end{split}$$

## Anomaly Cancellation

 In the chiral superfield description of the dilaton, we have a tree level coupling between the dilaton and the Chern-Simons anomaly superfield of the form

$$\mathcal{L} = \int d^4\theta E(S + \bar{S})\Omega = -\frac{1}{8} \int d^4\theta \frac{E}{R} S\Phi + \text{h.c.}$$

• So that the anomaly can be cancelled by allowing the dilaton to transform under a modular transformation as

$$\Delta S = -bF$$

## Conclusions & Future Work

- We have shown that, for at least two examples, string-derived supergravities can be regulated in such a way that the anomalies take a universal form and can be eliminated with the GS mechanism
- Next step is to apply the procedure to more realistic string compactifications to make statements about the phenomenology of string models :
  - FIQS Z<sub>3</sub> orbifold with two Wilson lines, with gauge group

 $SO(10) \times SU(3) \times SU(2) \times [U(1)]^7 \times U(1)_X$ 

- $U(1)_X$  is anomalous, but can be included in the GS mechanism
- Hypercharge defined by a linear combination of the non-anomalous U(1)'s
- Quark and lepton fields charged under the hypercharge and  $SU(3) \times SU(2)$
- Hidden sector with SO(10)

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