

# CALCULATIONS WITH OFF-SHELL MATRIX ELEMENTS, TMD PARTON DENSITIES AND TMD PARTON SHOWERS

Marcin Bury

Institute of Nuclear Physics, Polish Academy of Sciences, Kraków

xFitter workshop  
4-7 March 2018



Based on arXiv:1712.05932 (MB, A. van Hameren, H. Jung, K. Kutak, S. Sapeta and M. Serino)

- Introduction
- High Energy Factorization
- KaTie: parton-level event generator
- TMD densities: Kimber-Martin-Ryskin-Watt procedure
- CASCADE: parton showers
- Example: predictions for high  $p_t$  dijets
- Conclusions and outlook

# Introduction

- Theoretical predictions for strong interactions have often larger uncertainties than experimental ones
- Loop calculations in HEF too demanding so far (work in progress by A. van Hameren)
- Parton showers are used to simulate higher order corrections
- Collinear IPS changes the kinematics of the hard process by changing transverse momentum
- TMD parton distributions allows us to keep the exact kinematics of the initial state

# High Energy Factorization ( $k_T$ -factorization)

Catani, Ciafaloni, Hautmann

Collins, Ellis

$$\sigma_{AB \rightarrow q\bar{q}} = \int d^2k_{TA} \frac{dx_A}{x_A} \mathcal{F}(x_A, k_{TA}) d^2k_{TB} \frac{dx_B}{x_B} \mathcal{F}(x_B, k_{TB}) \hat{\sigma}_{g^*g^*} \left( \frac{m^2}{x_A x_B s}, \frac{k_{TA}}{m}, \frac{k_{TB}}{m} \right)$$

- reduces to collinear factorization for  $s \gg m^2 \gg k_T^2$ , but holds also for  $s \gg m^2 \sim k_T^2$
- allows for higher-order kinematical effects at leading order

# High Energy Factorization ( $k_T$ -factorization)

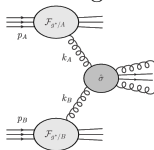
Catani, Ciafaloni, Hautmann

Collins, Ellis

$$\sigma_{AB \rightarrow q\bar{q}} = \int d^2k_{TA} \frac{dx_A}{x_A} \mathcal{F}(x_A, k_{TA}) d^2k_{TB} \frac{dx_B}{x_B} \mathcal{F}(x_B, k_{TB}) \hat{\sigma}_{g^*g^*} \left( \frac{m^2}{x_A x_B s}, \frac{k_{TA}}{m}, \frac{k_{TB}}{m} \right)$$

- reduces to collinear factorization for  $s \gg m^2 \gg k_T^2$ , but holds also for  $s \gg m^2 \sim k_T^2$
- allows for higher-order kinematical effects at leading order

- requires matrix elements with *off-shell* initial-state partons with  $k_i^2 = k_{iT}^2 < 0$



$$k_A^\mu = x_A p_A^\mu + k_{TA}^\mu$$

$$k_B^\mu = x_B p_B^\mu + k_{TB}^\mu$$

- $k_T$ -dependent  $\mathcal{F}$  may satisfy BFKL, CCFM, BK, KGBJS evolution equations
- typically associated with small- $x$  and forward physics, saturation, heavy-ions

# KaTie: parton-level event generator

KATIE (A. van Hameren) [arXiv:1611.00680](https://arxiv.org/abs/1611.00680)

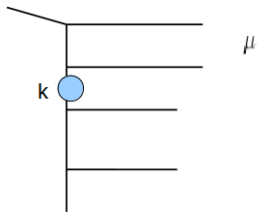
- Complete Monte Carlo program for tree-level calculations
- Any process within the Standard Model
- On-shell and/or off-shell initial-state partons
- Employs numerical Dyson-Schwinger recursion to calculate helicity amplitudes
- Automatic phase space optimization
- Produces weighted parton-level event files in the Les Houches format
- Single- and multi-parton scattering

# TMD densities: Kimber-Martin-Ryskin-Watt procedure

Survival probability without emissions

$$T_s(\mu^2, k^2) = \exp\left(-\int_{\mu^2}^{k^2} \frac{dk'^2}{k'^2} \frac{\alpha_s(k'^2)}{2\pi} \times \sum_{a'} \int_0^{1-\Delta} dz' P_{aa'}(z')\right)$$

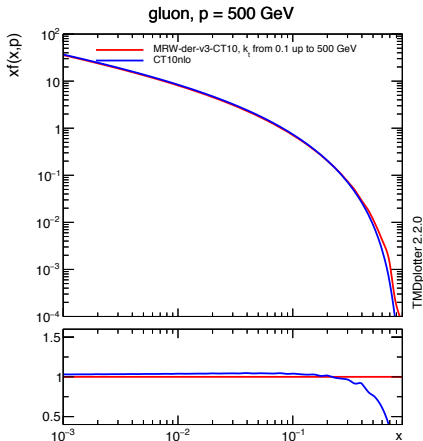
$$\mathcal{F}(x, k^2, \mu^2) \sim \partial_{\lambda^2} (T_s(\lambda^2, \mu^2) x g(x, \lambda^2)) \Big|_{\lambda^2=k^2}$$



$$\Delta = \frac{\mu}{\mu + k}$$

$\mu = \text{hard scale}$

# TMD densities: Kimber-Martin-Ryskin-Watt procedure



Integrating the updf we get back collinear PDF

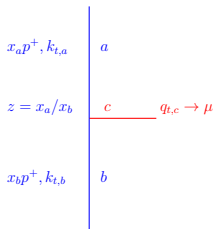
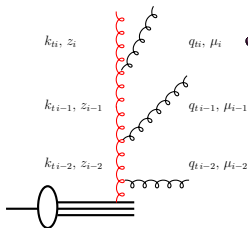
$$\int_0^{\mu^2} dk_t^2 \mathcal{F}_a(x, k_t^2, \mu^2) = xf(x, \mu^2)$$



CASCADE (H. Jung) [Comput.Phys.Commun. 143 \(2002\) 100](#)

- $k_t$ -factorization based Monte Carlo program doing ISR and FSR
- Backward evolution of the initial state
- New features
  - full TMD evolution of all flavour PS
  - interface to LHE files, for input from KATIE and POWHEG
- Hadronization performed via PYTHIA 6
- Analysis is done with RIVET package
- Jets reconstructed via the anti- $k_t$  algorithm with FASTJET

# Initial-state parton shower

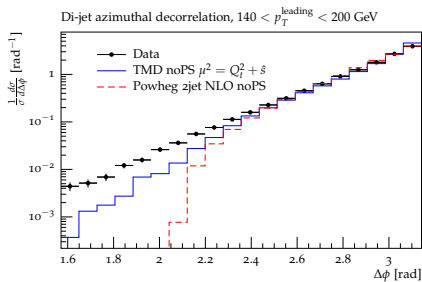
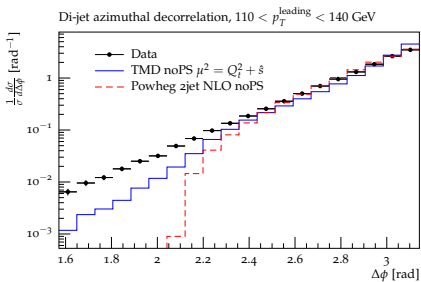


- Starting with the hard scale  $\mu_i$ , the parton shower algorithm searches for the next scale  $\mu_{i-1}$ , at which a resolvable branching occurs, which is selected from the Sudakov form factor

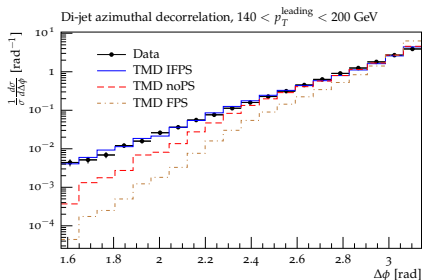
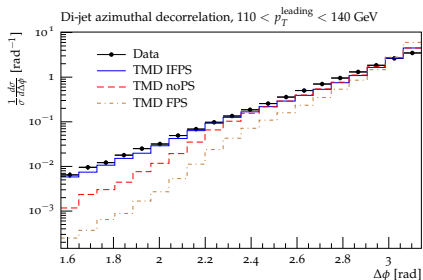
$$\Delta_S(x, \mu_i, \mu_{i-1}) = \exp \left[ - \int_{\mu_{i-1}}^{\mu_i} \frac{d\mu'}{\mu'} \frac{\alpha_S(\mu')}{2\pi} \sum_a \int dz P_{a \rightarrow bc}(z) \frac{x' \mathcal{A}_a(x', k'_t, \mu')}{x \mathcal{A}_b(x, k_t, \mu')} \right]$$

- The transverse momentum can be obtained by giving a physical interpretation to the evolution scale, which is associated with the angle of the emission
- Once the transverse momentum of the emitted parton is known, the  $k_t$  of the propagating parton can be calculated
- The whole procedure continues until cutoff (lowest scale) is reached

# Azimuthal angle correlations - no shower



# Azimuthal angle correlations - showers



# Conclusions and outlook

## Conclusions

- A framework for evaluation of cross section in HEF supplemented with initial- and final-state parton showers
- Full set of TMD parton densities using the KMRW approach
- TMD parton shower including all flavours and following the TMD distribution, without the need for adjusting further parameters

## Outlook

- More comparisons to data for  $p + p$  and  $p + A$  using KATIE + CASCADE
- Refinement of used TMDs i.e. fits to latest data
- Corrections of higher orders to ME