### Exact results in QFT

Alessandro Pini

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Some information about myself





## • I was born in Milan (Italy)

ack Sea

## • I did my bachelor and my master in Milan (Italy)

BICK

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LANTIC

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Exact results in QFT

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12 December, 2017 3 / 15

# • I was born in Milan (Italy)

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## I did my PhD in Oviedo

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This is my first postdoc

ack

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### Free time activities



#### Running



### Free time activities



I'm also learning how to play the classic guitar (but I'm still at the beginning).

### Contents of the talk





- Seiberg-Witten theory. [N. Seiberg, E. Witten (1995)]
- The AdS/CFT correspondence. [J. Maldacena (1997)]
- Supersymmetric localization. [V.Pestun (2012)]

### The Seiberg-Witten theory in a nutshell

[N. Seiberg, E. Witten (1995)]

We want to solve 
$$\mathcal{N} = 0$$
  $\textcircled{\otimes}$   
(UV) High energy perurbative QCD  $\Rightarrow$  (IR) Low-energy confinement  
 $\mathcal{N} = 1 \quad B \leftrightarrow F$   $\mathcal{N} = 2 \quad B \leftrightarrow F \leftrightarrow B$   
Let's consider  $\mathcal{N} = 2$   $\textcircled{\otimes}$   
(UV)  $4d \quad \mathcal{N} = 2 \quad QFT$   $\Rightarrow$  (IR) Exact low energy effective action

# The AdS/CFT correspondence-holography



Dictionary

Observable in QFT  $\Leftrightarrow$  Observable in the gravity theory

• energy scale in QFT ⇔ radial coordinate in the gravity dual

Exact results in QFT

Supersymmetric QFT on a compact manifold  $\mathcal{M}$  (e.g.  $\mathbb{S}^d$ ) with a supercharge  $\mathcal{Q}$ 

Partition function

$$Z = \int \mathcal{D}[\Phi] e^{-S[\Phi]}$$

Infinite dimensional path integral over all the fields configuration  $Z=\int da f(a)$ 

normal finite dimensional integral

 $\Rightarrow$  exact observables computation in a supersymmetric QFTs

 $\Rightarrow$ 

Exact results in QFT

### PhD research activity

- 5d N = 1 QFT are not renormalizable → fixed point theory, for a proper choice of the gauge group and of the matter content.
   [N. Seiberg (1995)], [K.A. Intriligator, D.R. Morrison and N. Seiberg (1997)]
- In some cases: enhancement of the global symmetry group at UV fixed point.

[N. Seiberg (1995)]

$$(V) \quad G_{UV} \rightarrow \text{RG flow} \rightarrow (IR) \quad G_{IR} \qquad G_{UV} > G_{IR}$$

• Gravity dual for  $5d \mathcal{N} = 1 \text{ QFT} \Rightarrow study of holographic RG flows.$ 

[A. Brandhuber and Y. Oz (1999)],[O. Bergman and D. Rodríguez-Gómez (2012)]

### PhD research activity

- The study of gauge theories on curved background → reaction to curvature measures different things: partition functions, index ...
   [V.Pestun (2012)]...
- Procedure that allows to define a SQFT on a curved background. [G.Festuccia and N. Seiberg (2011)]
   ⇒ study of 5d SQFT on a Riemann manifold.
- A particularly useful background: S<sup>1</sup> × S<sup>4</sup> → computation of the Superconformal Index (SCI) of the theory.

[J.Kinney, J.M.Maldacena, S.Minwalla and S.Raju (2007)],[J.Bhattacharya, S.Bhattacharyya, S.Minwalla and

S.Raju (2008)],[H.-C. Kim,S.-S.Kim, K. Lee (2012)]

• Limits of the SCI  $\leftrightarrow$  subset of operators of the theory

[A.Gadde, L. Rastelli and collaborators (2012)]  $\Rightarrow$  particular limit limit of the 5*d* SCI.

### PhD research activity

• The partition function takes the form

$$Z = \int Z_{pert} Z_{inst}$$

in the case of pure gauge theories (with 8 supercharges)  $Z_{inst}$  coincides with the Hilbert series of the instanton moduli space.

[D. Rodríguez-Gómez and G.Zafrir (2014)],[ C.A.Keller,N. Mekareeya, J.Song and Y.Tachikawa (2012)]

• We can use string theory to study instantons on  $\mathbb{C}^2$ . Higgs branch of a system of  $D_p$  and  $D_{p+4}$  branes (e.g. p = 3)



For the moment I will mostly focused on 4d QFTs.

- $4d \mathcal{N} = 2 \rightarrow 4d \mathcal{N} = 1$ [D.Gaiotto, S. Razamat (2015)]
- Discovery of 4*d* QFTs with N = 3 (Non Lagrangian theories!). [II. Garcia Etxabarria, D. Regalado (2015)
  - $\rightarrow$  a lot of things to explore
    - Study of the SCI.
    - Seiberg-Witten curve.

#### THANKS FOR YOUR ATTENTION

### 4d $\mathcal{N}$ = 2 prepotential and low energy effective action

#### The prepotential $\mathcal{F}$

$$\mathcal{F}(A) = \frac{1}{2}\tau_0 A^2 + \frac{i}{\pi} A^2 \log(\frac{A^2}{\Lambda^2}) + \frac{1}{2\pi i} A^2 \sum_{l=1}^{+\infty} c_l (\frac{\Lambda}{A})^{4l}$$

#### The low energy effective action

$$\frac{1}{4\pi} \operatorname{Imm}\left[\int d^4\theta \frac{\partial \mathcal{F}}{\partial A} \overline{A} + \int d^2\theta \frac{1}{2} \frac{\partial^2 \mathcal{F}}{\partial A^2} W_{\alpha} W^{\alpha}\right]$$

- QFT with a fermionic odd charge Q, such that  $Q^2 = B$
- BPS observable  $\mathcal{QO}_{BPS} = 0$
- we want to evaluate

$$\langle \mathcal{O}_{BPS} \rangle = \int_{F} \mathcal{D}[X] \mathcal{O}_{BPS} e^{-S[X]}$$

it holds that

$$\langle \mathcal{O}_{BPS} \rangle = \langle \mathcal{O}_{BPS} + \mathcal{Q}O \rangle$$

### Supersymmetric localization 2

Let's consider the deformed path-integral

$$\langle \mathcal{O}_{BPS} \rangle = \int_{F} \mathcal{D}[X] \mathcal{O}_{BPS} e^{-S[X] - t\mathcal{Q}P_{F}[X]}$$

It holds that

$$\frac{d}{dt} \langle \mathcal{O}_{BPS} \rangle = 0 \quad \Rightarrow \quad t \to +\infty$$

the integral is dominated by the saddle points of the localising action

$$X = X_0 + \frac{1}{\sqrt{t}}\delta X \quad \Rightarrow \quad S[X_0] + \frac{1}{2}\int \int \frac{\delta^2 S_{loc}[X_0]}{\delta X^2} \mid_{X = X_0} (\delta X)^2$$