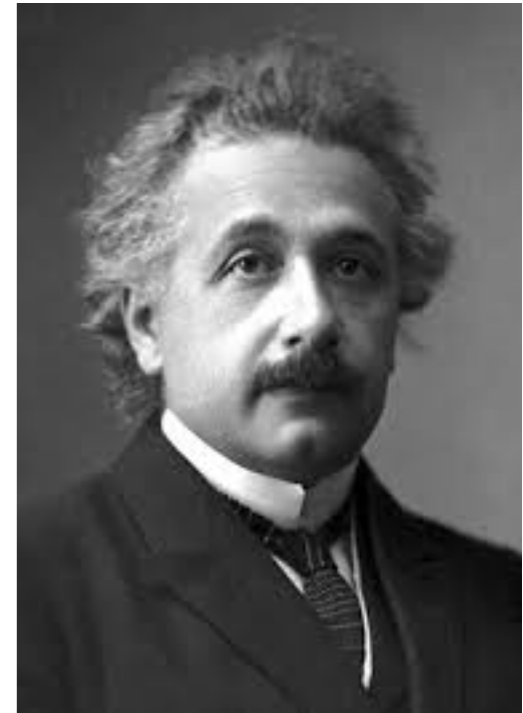


Naturalness and weak scale SUSY

Howard Baer
University of Oklahoma

DESY naturalness forum
April 26, 2018

twin pillars of guidance:
naturalness & simplicity



“The appearance of fine-tuning in a scientific theory is like a cry of distress from nature, complaining that something needs to be better explained”

S. Weinberg

“Everything should be made as simple as possible, but not simpler”

A. Einstein

Lessons from yesterday

- Multi-parameter effective SUSY theories are useful because independent soft terms parametrize our ignorance of mechanics of SUSY breaking
- However, the fine-tuning calculation should be applied to the underlying more fundamental theory where soft terms are surely correlated
- With correlated soft terms, then both UV measures reduce to IR measure at tree level: $\Delta_{HS}, \Delta_{BG} \rightarrow \Delta_{EW}$
- This is good news: there is only one answer for amount of fine-tuning and it is independent of scale choice at which spectra is generated: given some spectra, no matter where it comes from, all should agree on whether or not it is fine-tuned
- Re-evaluation of sparticle mass upper bounds shows low $\mu < 200-300$ GeV large A -term models with $m(h) \sim 125$ GeV are still natural with $m_{gl} < 6$ TeV, $m_{t1} < 3$ TeV

Some topics lightly or not discussed

- Only required weak scale sparticles are higgsinos: hard to see at LHC but huge motivation for e^+e^- collider with roots $> 2m(\text{higgsino})$
- HL-LHC not enough to falsify natural SUSY: will need HE-LHC to probe $m(\text{gl}) \sim 6 \text{ TeV}$ and $m(\text{t1}) \sim 3 \text{ TeV}$
- Requiring naturalness in both EW and QCD sector: DM = higgsino-like WIMP + axion admixture: usually mostly axions, but with suppressed axion/WIMP detection rates
- Non-holonomic soft terms may soften light higgsino argument, but these seem highly suppressed in more UV complete theories such as string-based models
- Some natural models all have light higgsinos but: NUHM2 (unified gauginos), natural generalized mirage mediation (compressed gauginos), natural anomaly mediation ($m_0 < M_2 < M_1 < M_3$)
- Precision higgsino/collider measurements can reveal which of these cases or other would be realized in nature

Recall yesterday talk conclusion:

First order question:

why is the weak scale $m(W,Z,h) \sim 100 \text{ GeV}$?

Because $\mu(\text{weak})$, $m_{H_u}(\text{weak}) \sim 100\text{--}200 \text{ GeV}$
and top squarks \sim few TeV but highly mixed

Second order questions:

1. Why might $\mu \ll m(\text{soft})$

2. Why might soft terms be at multi-TeV

scale but with $m(H_u)$ driven

radiatively to $m_{H_u}^2(\text{weak}) \sim -(100\text{--}200 \text{ GeV})^2$?

SUSY μ problem: μ term is SUSY, not SUSY breaking:
expect $\mu \sim M(\text{Pl})$ but phenomenology requires $\mu \sim m(\text{Z})$

● NMSSM: $\mu \sim m(3/2)$; but beware singlets!

● Giudice-Masiero: μ forbidden by some symmetry:
generate via Higgs coupling to hidden sector

● **Kim-Nilles**: invoke SUSY version of DFSZ axion

solution to strong CP:

$$W \ni \lambda_\mu S^2 H_u H_d / m_P$$

KN: PQ symmetry forbids μ term,
but then it is generated via PQ breaking

$$\mu \sim \lambda_\mu f_a^2 / m_P$$

Little Hierarchy due to mismatch between
PQ breaking and SUSY breaking scales?

$$m_{3/2} \sim m_{hid}^2 / M_P$$

$$f_a \ll m_{hid}$$

**Higgs mass tells us where
to look for axion!**

$$m_a \sim 6.2 \mu\text{eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$$

Little Hierarchy from radiative PQ breaking? exhibited within context of MSY/CCK model

Murayama, Suzuki, Yanagida (1992);
Gherghetta, Kane (1995)

Choi, Chun, Kim (1996)

Bae, HB, Serce, PRD91 (2015) 015003

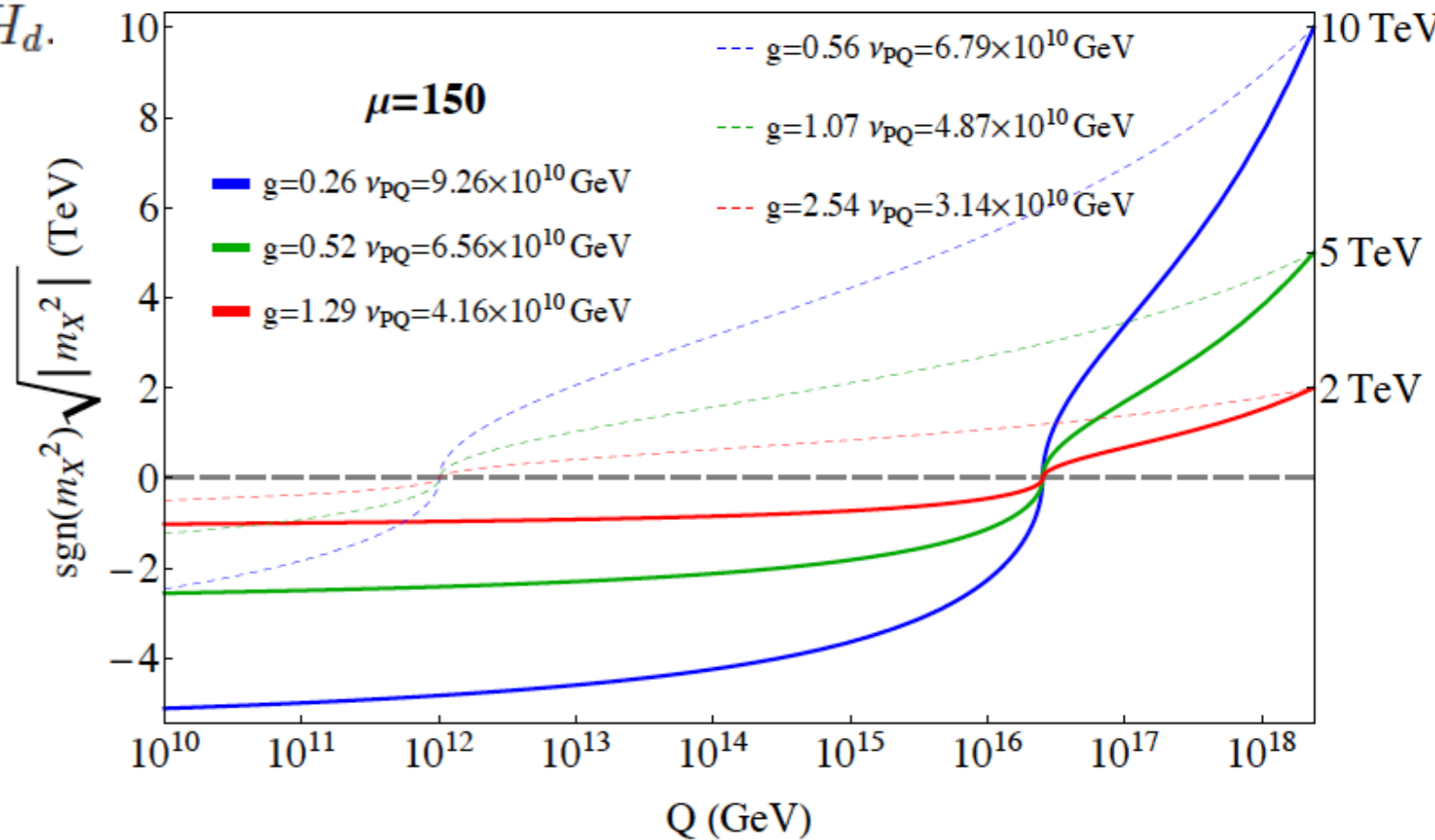
augment MSSM with PQ charges/fields:

$$\hat{f}' = \frac{1}{2} h_{ij} \hat{X} \hat{N}_i^c \hat{N}_j^c + \frac{f}{M_P} \hat{X}^3 \hat{Y} + \frac{g}{M_P} \hat{X} \hat{Y} \hat{H}_u \hat{H}_d.$$

SUSY breaking triggers
PQ breaking:
generate f_a and M_N

$$M_{N_i^c} = v_X h_i |_{Q=v_X}$$

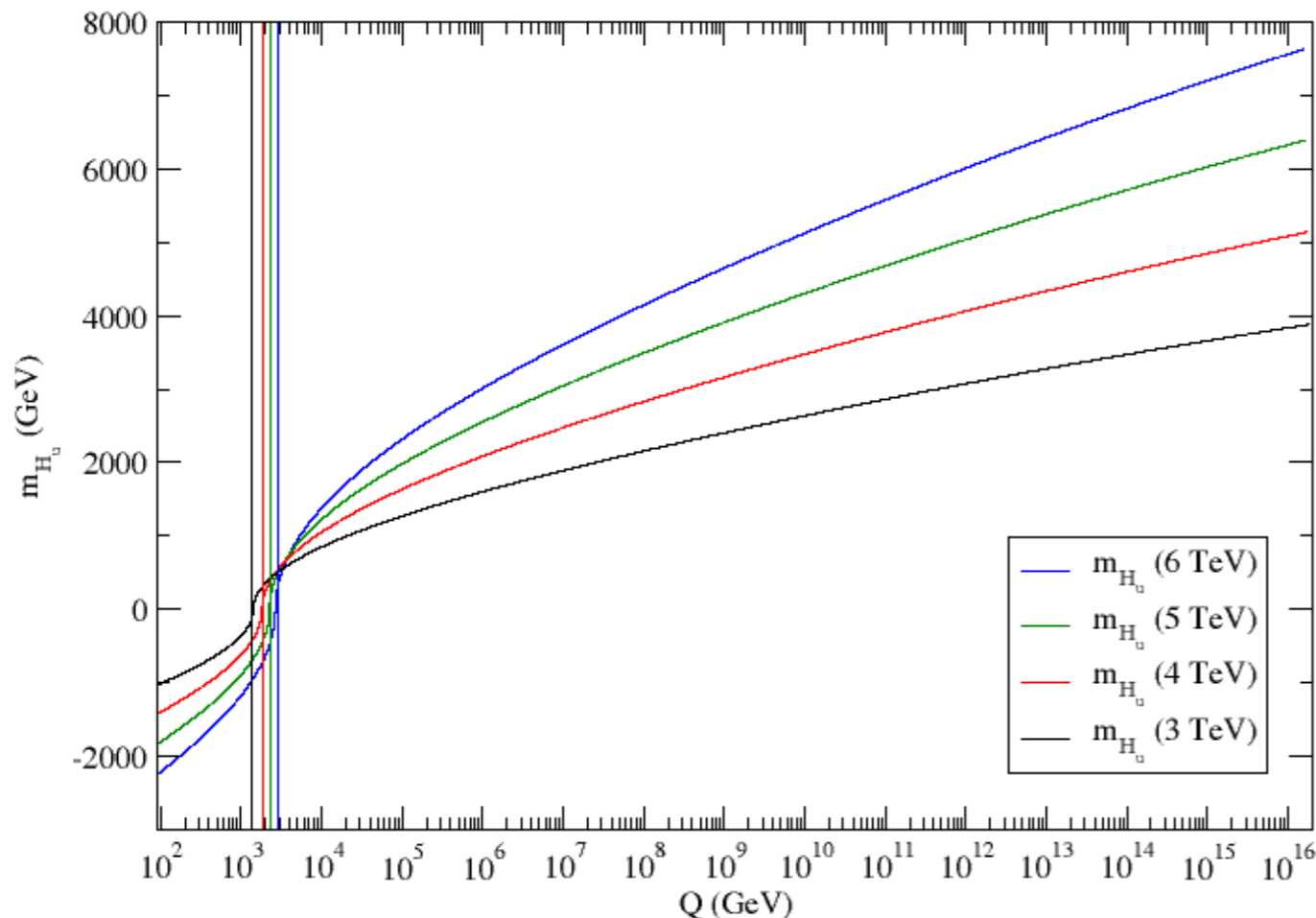
$$\mu = g \frac{v_X v_Y}{M_P}.$$



Large $m_{3/2}$ generates small $\mu \sim 100 - 200$ GeV!

Why might m_{H_u} have the value needed to give naturalness at weak scale?

1. For right correlations amongst soft terms, get “generalized focus point”



e.g.

For $\mu = 150$ GeV, $\tan \beta = 10$ and

$$m_0^2 = m_{3/2}^2$$

$$A_0 = -1.6 m_{3/2}$$

$$m_{1/2} = m_{3/2}/5$$

$$m_{H_d}^2 = m_{3/2}^2/2.$$

$$m_{H_u}^2(GUT) = 1.8 m_{3/2}^2 - (212.52 \text{ GeV})^2.$$

$m_{H_u}^2(weak) \sim a \cdot m_{3/2}^2$ with correlated soft terms such that a is small: generalized focus point behavior

Statistical analysis of SUSY breaking scale: M. Douglas, hep-th/0405279

some reasonable assumptions

- string theory landscape contains vast ensemble of $N=1, d=4$ SUGRA EFTs at high scales
- the EFTs contain the SM as weak scale EFT
- the EFTs contain visible sector +potentially large hidden sector
- visible sector contains MSSM plus extra gauge singlets (e.g. a PQ sector, RN neutrinos,...)
- SUGRA is broken spontaneously via superHiggs mechanism via either F- or D- terms or in general a combination

Why do soft terms take on values needed for natural (barely-broken) EWSB?

2. string theory landscape?

- assume model like MSY/CCK where $\mu \sim 100$ GeV
- then $m(\text{weak})^2 \sim |m_{H_u}^2|$
- If all values of SUSY breaking field $\langle F_X \rangle$ equally likely, then mild (linear) statistical draw towards large soft terms
- This is balanced by anthropic requirement of weak scale $m_{\text{weak}} \sim 100$ GeV

Anthropic selection of $m_{\text{weak}} \sim 100$ GeV:

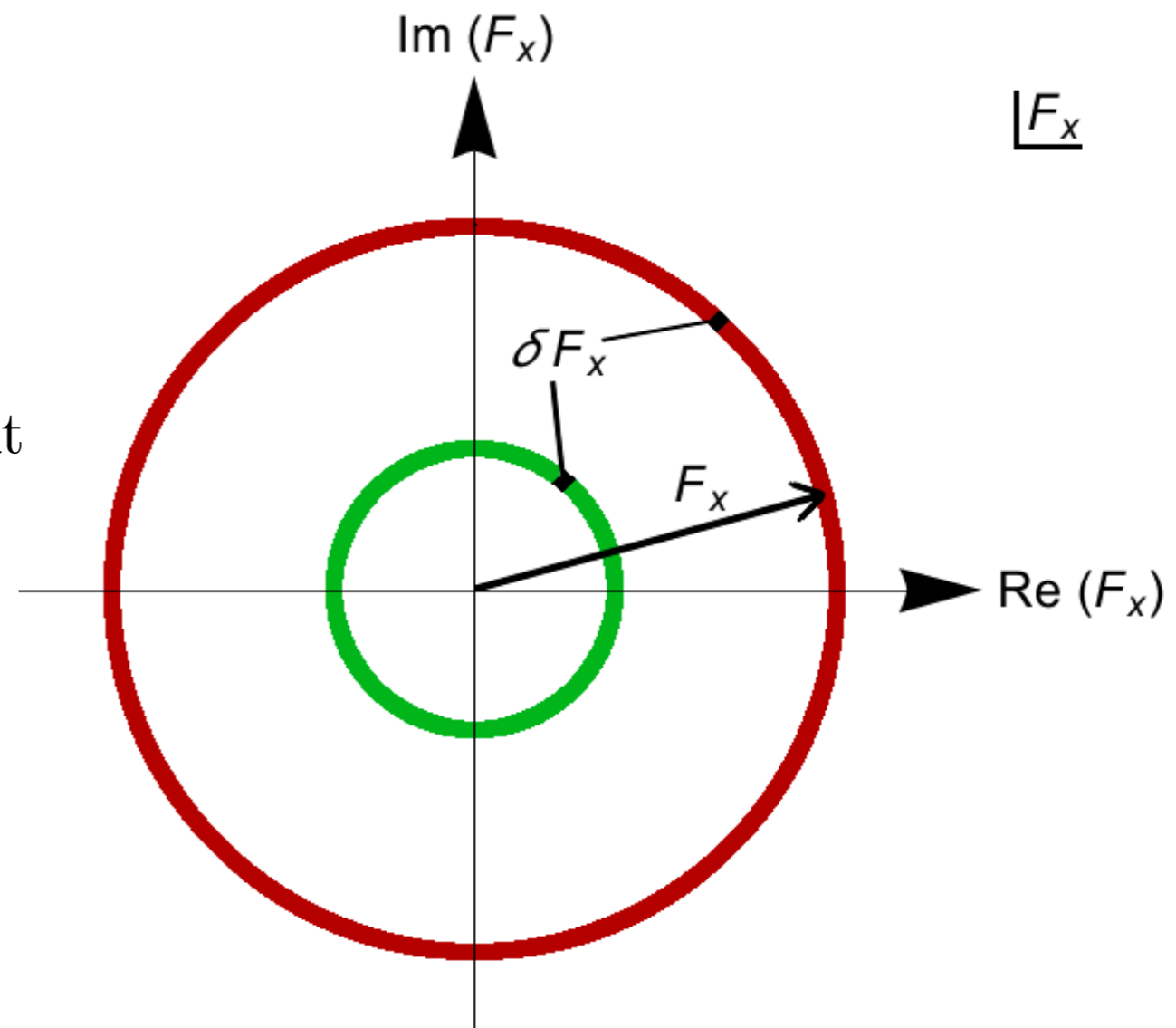
If m_W too large, then weak interactions

$\sim (1/m_W^4)$ too weak

weak decays, fusion reactions suppressed

elements not as we know them

$$m(\text{weak}) < \sim 400 \text{ GeV (Agrawal et al.)}$$



Scalar potential is given by usual SUGRA form:

$$\begin{aligned}
 V &= e^{K/m_P^2} \left(g^{i\bar{j}} D_i W D_{\bar{j}} W^* - \frac{3}{m_P^2} |W|^2 \right) + \frac{1}{2} \sum_{\alpha} D_{\alpha}^2 \\
 &= e^{K/m_P^2} \left(\sum_i |F_i|^2 - 3 \frac{|W|^2}{m_P^2} \right) + \frac{1}{2} \sum_{\alpha} D_{\alpha}^2
 \end{aligned}$$

- W = holomorphic superpotential
- K = real Kähler function
- $F_i = D_i W = DW/D\phi^i \equiv \partial W/\partial\phi^i + (1/m_P^2)(\partial K/\partial\phi^i)W$ are F -terms
- $D_{\alpha} \sim \sum \phi^{\dagger} g t_{\alpha} \phi$ are D -terms
- ϕ^i are chiral superfields

minimize V :

- $\partial V/\partial\phi^i = 0$
- $\partial^2 V/\partial\phi^i \partial\phi^j > 0$
- $\Lambda_{cc} = m_{hidden}^4 - 3e^{K/m_P^2} |W|^2/m_P^2$ with
- $m_{hidden}^4 = \sum_i |F_i|^2 + \frac{1}{2} \sum_{\alpha} D_{\alpha}^2$ is hidden sector mass scale

gravitino mass $m_{3/2} = e^{K/2m_P^2} m_P \sim m_{hidden}^2/m_P$ with $m_{hidden} \sim 10^{12}$ GeV

Denef&Douglas: statistics of SUSY breaking in landscape

DD observation: W_0 distributed uniformly as complex variable allows dynamical neutralization of Λ while not influencing SUSY breaking

Then, number of flux vacua containing spontaneously broken SUGRA with SUSY breaking scale m_{hidden}^2 is:

$$dN_{vac}[m_{hidden}^2, m_{weak}, \Lambda] = f_{SUSY}(m_{hidden}^2) \cdot f_{EWFT} \cdot f_{cc} dm_{hidden}^2$$

- $f_{cc} \sim \Lambda/m^4$ where DD maintain $m \sim m_{string}$ and not m_{hidden}
- $f_{SUSY}(m_{hidden}^2) \sim (m_{hidden}^2)^{2n_F+n_D-1}$ for uniformly distributed values of F and D breaking fields
- $f_{EWFT} \sim m_{weak}^2/m_{soft}^2$ (?) where $m_{soft} \sim m_{3/2} \sim m_{hidden}^2/m_P$

$$n = 2n_F + n_D - 1$$

$$f_{SUSY} \sim m_{soft}^n$$

n_F	n_D	n
0	1	0
1	0	1
0	2	1
1	1	2
0	3	2
2	0	3
2	1	4

landscape favors high scale SUSY breaking
tempered by $f(EWFT)$ anthropic penalty!

What about DD/AD anthropic penalty $f_{EWFT} \sim m_{weak}^2/m_{soft}^2$?

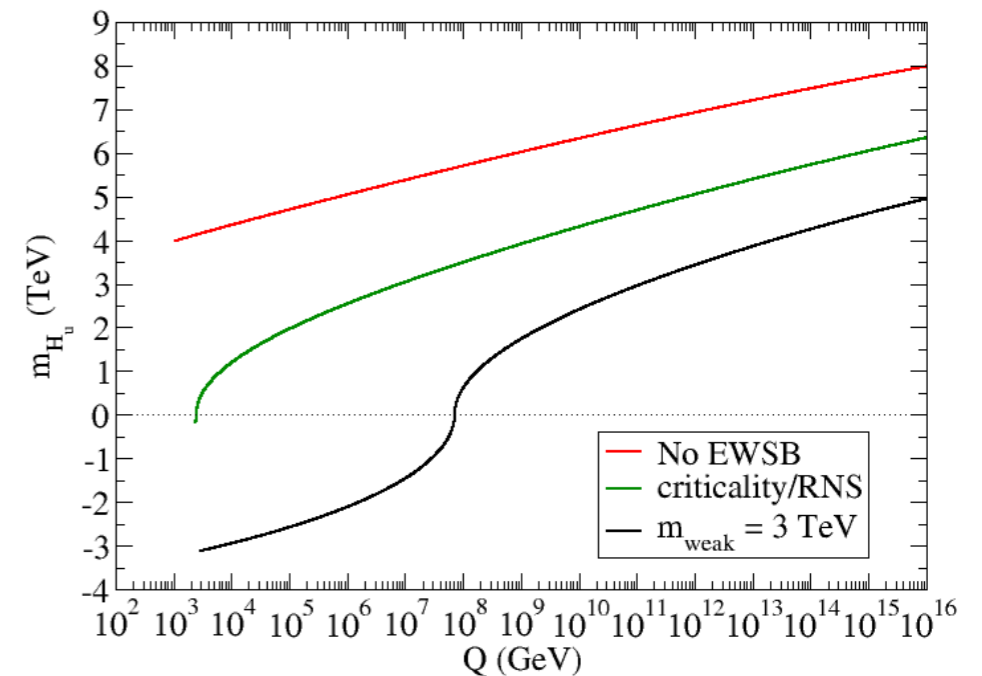
This fails in a variety of *practical* cases:

- A -terms get large: \Rightarrow CCB minima
- $m_{H_u}^2$ too large: fail to break EW symmetry

Must require proper EWSB!

Even if EWS properly broken, then

- large A_t reduces EWFT in the $\Sigma_u^u(\tilde{t}_{1,2})$
- large $m_{H_u}^2(m_{GUT})$ needed to radiatively drive $m_{H_u}^2$ to natural value at weak scale



Better proposal: $f_{EWFT} \Rightarrow \Theta(30 - \Delta_{EW})$

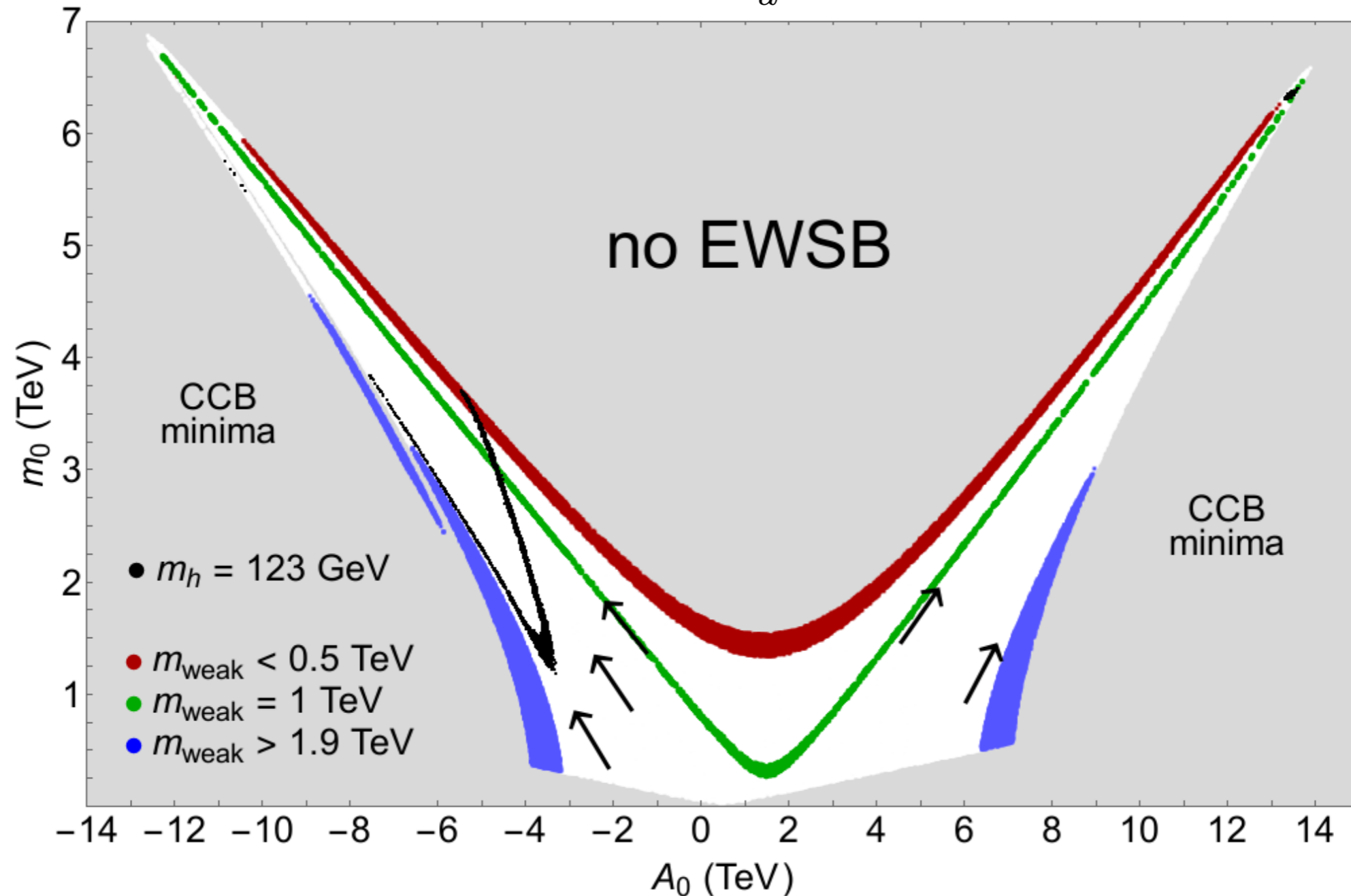
keeps calculated weak scale within factor ~ 4 of measured weak scale

$$m_{weak} \equiv m_{W,Z,h} \sim 100 \text{ GeV}$$

Assume $\mu \sim 100 - 200 \text{ GeV}$ via *e.g.* rad PW breaking: then m_Z variable and may be large depending on soft terms $m_{H_{u,d}}^2$ and $\Sigma_{u,d}^{u,d}(i)$

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2$$

$$m_{H_u} = 1.3m_0$$



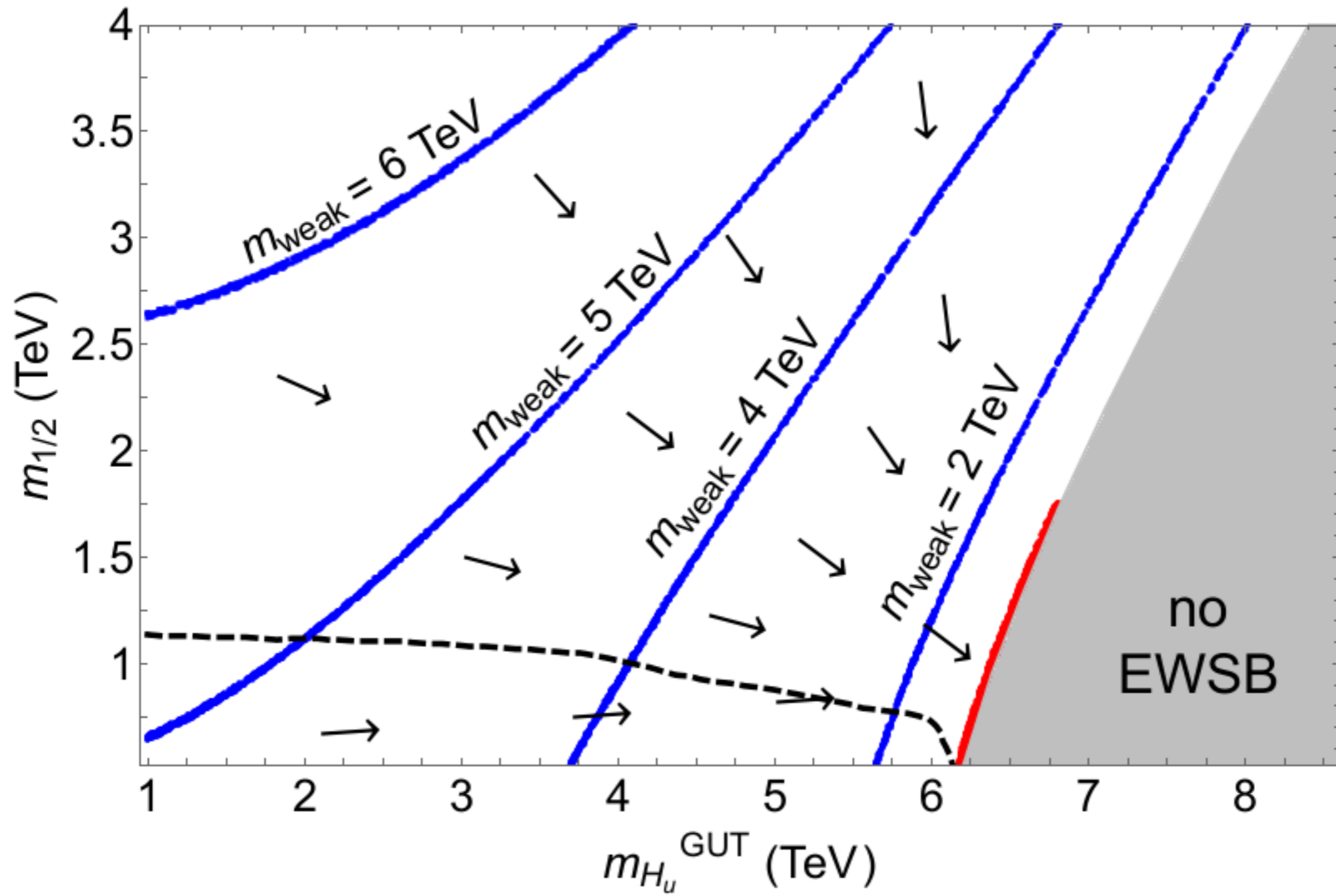
statistical draw to large soft terms balanced by anthropic draw toward red ($m(\text{weak}) \sim 100$ GeV): then $m(\text{Higgs}) \sim 125$ GeV and natural SUSY spectrum!

Denef, Douglas, JHEP0405 (2004) 072

Giudice, Rattazzi, NPB757 (2006) 19;

HB, Barger, Savoy, Serce, PLB758 (2016) 113

$$m_0 = 5 \text{ TeV}$$



statistical/anthropic draw toward FP-like region

For practical calculations, adopt NUHM3 SUGRA model:

- $m_0(1, 2) = \text{gen}(1,2)$ common soft mass
- $m_0(3) = \text{gen}(3)$ common soft mass
- $m_{H_u}^2$ up-Higgs soft mass
- $m_{H_d}^2$ down-Higgs soft mass
- $m_{1/2} =$ unified gaugino mass
- $A_0 =$ unified trilinear soft term
- $\tan \beta$

Trade $m_{H_u}^2, m_{H_d}^2 \Leftrightarrow \mu, m_A$

$m_0(1, 2), m_0(3), m_{1/2}, A_0, \tan \beta, \mu, m_A$ (NUHM3)

Recent work: place on more quantitative footing:
scan soft SUSY breaking parameters as $m(\text{soft})^n$
along with $f(\text{EWFT})$ penalty

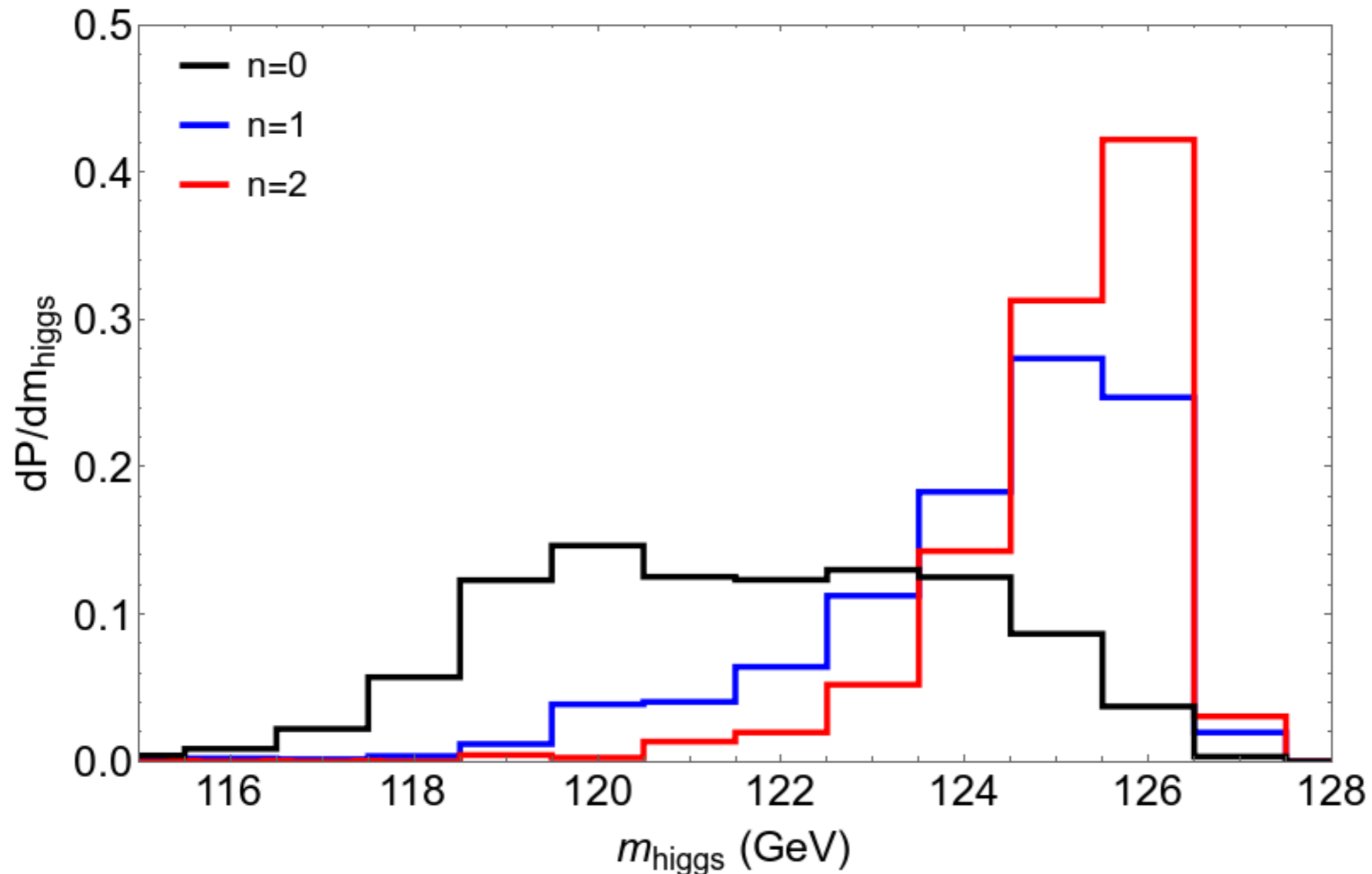
We scan according to m_{soft}^n over:

- $m_0(1, 2) : 0.1 - 40 \text{ TeV},$
 - $m_0(3) : 0.1 - 20 \text{ TeV},$
 - $m_{1/2} : 0.5 - 10 \text{ TeV},$
 - $A_0 : 0 - -60 \text{ TeV},$
 - $m_A : 0.3 - 10 \text{ TeV},$
- $\tan \beta : 3 - 60 \quad (\text{flat})$

$\mu=150 \text{ GeV (fixed)}$

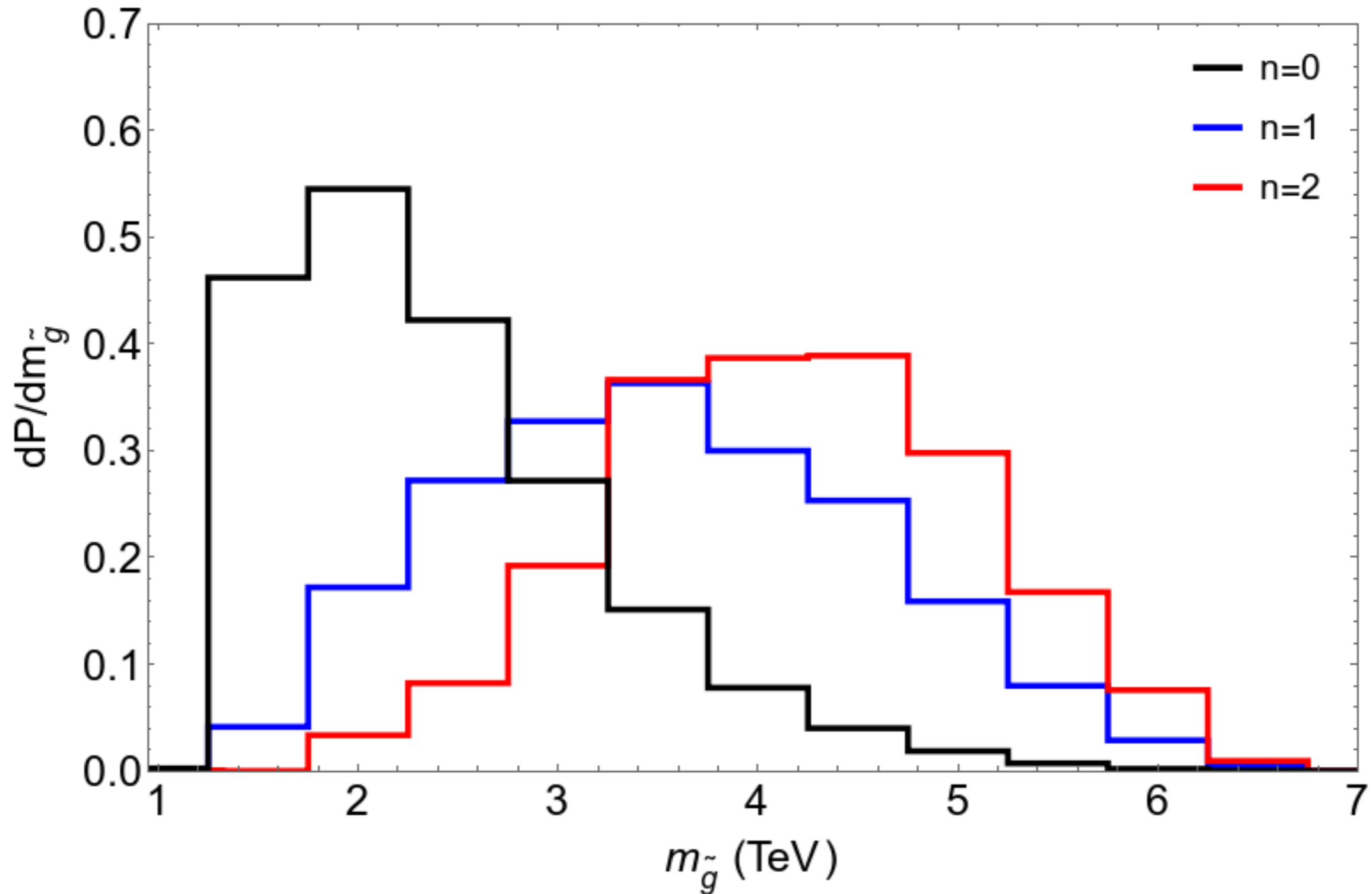
Making the picture more quantitative:

$$dN_{vac}[m_{hidden}^2, m_{weak}, \Lambda] = f_{SUSY}(m_{hidden}^2) \cdot f_{EWFT} \cdot f_{cc} dm_{hidden}^2$$



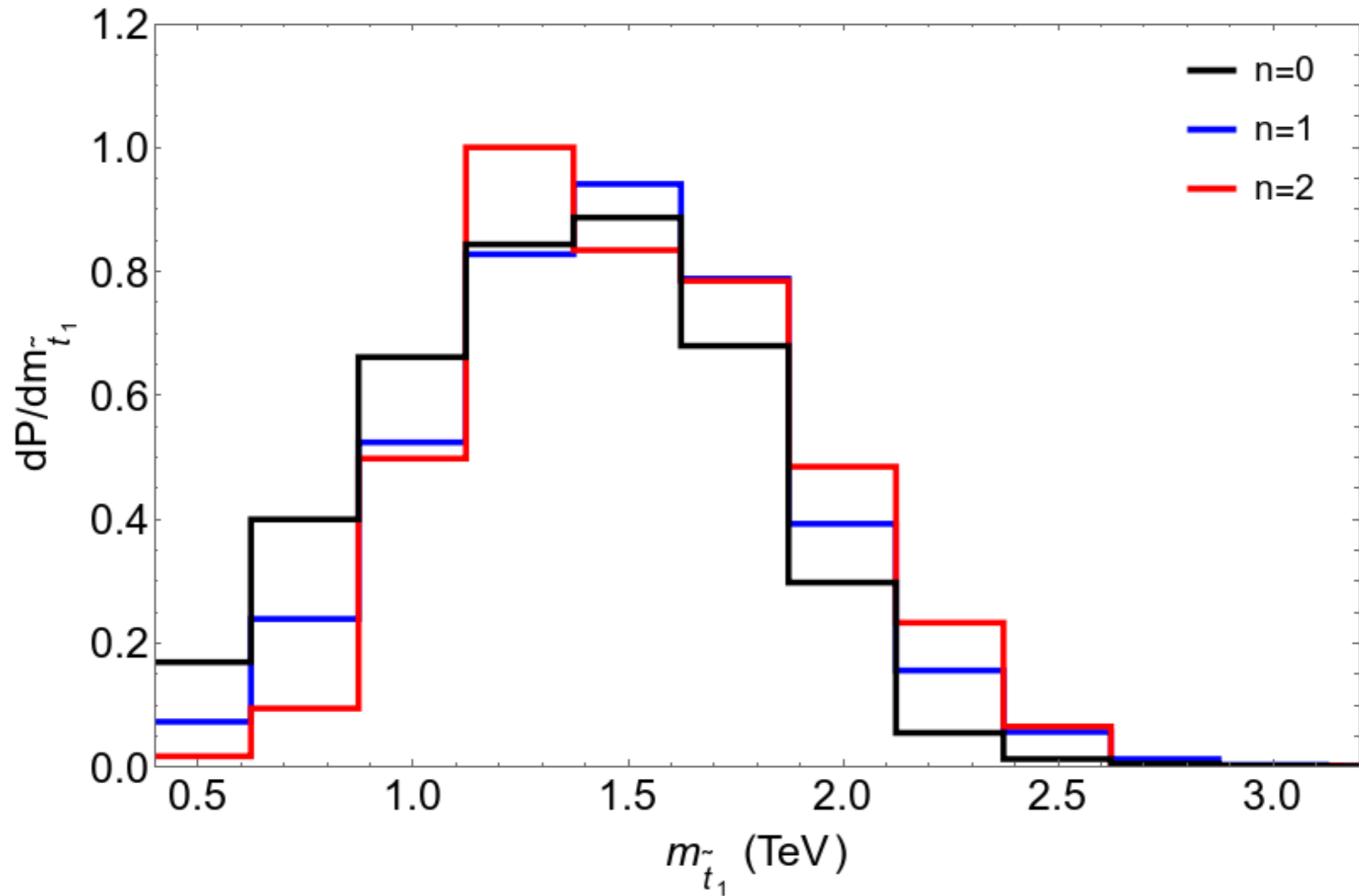
$m(h) \sim 125$ most favored for $n=1,2$

What is corresponding distribution for gluino mass?



typically beyond LHC 14 reach (may need HE-LHC)

and m_{t_1} ?



Summary $n=1,2$:

- $m_h \sim 125 \pm 2$ GeV
- $m_{\tilde{g}} \sim 4 \pm 2$ TeV,
- $m_{\tilde{t}_1} \sim 1.5 \pm 0.5$ TeV,
- $m_A \sim 3 \pm 2$ TeV,
- $\tan \beta \sim 13 \pm 7$,
- $m_{\tilde{W}_1, \tilde{Z}_{1,2}} \sim 200 \pm 100$ GeV and
- $m_{\tilde{Z}_2} - m_{\tilde{Z}_1} \sim 7 \pm 3$ GeV with
- $m_0(1, 2) \sim 20 \pm 10$ TeV (for first/second generation matter scalars)

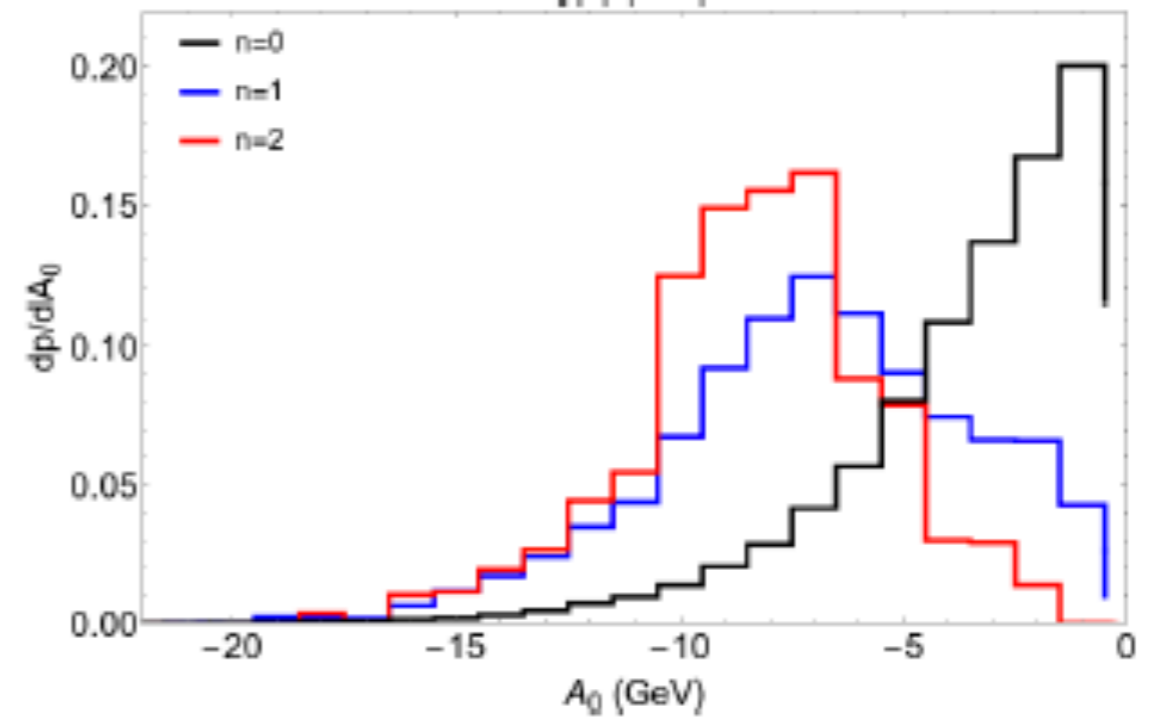
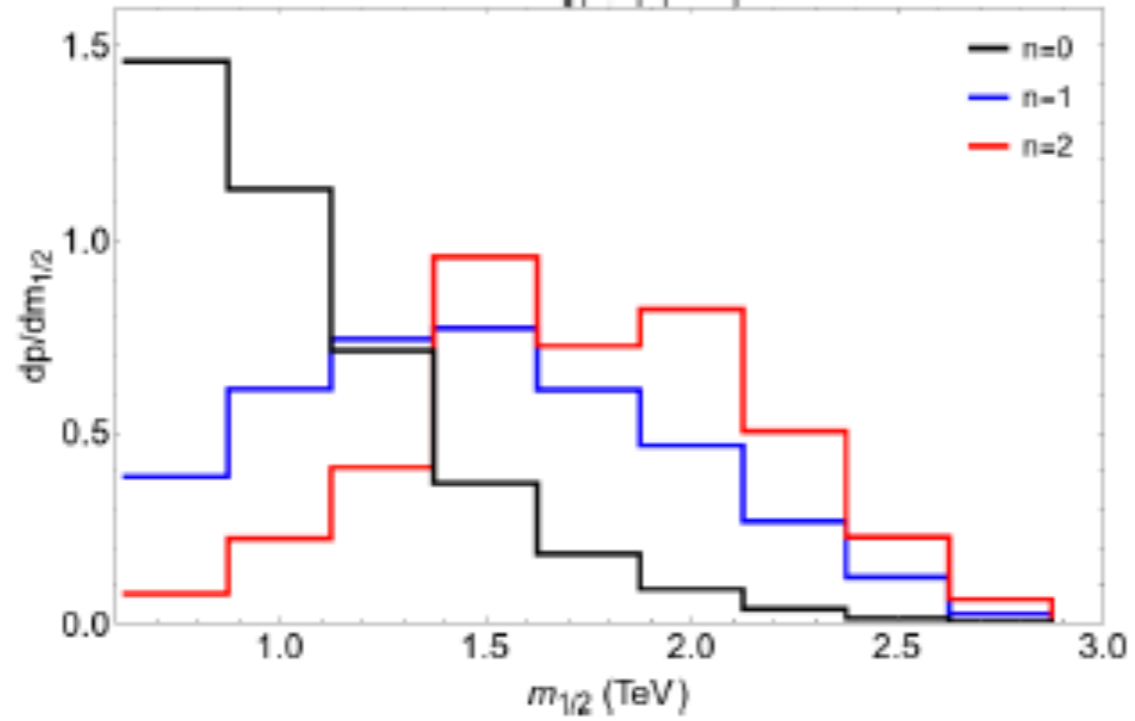
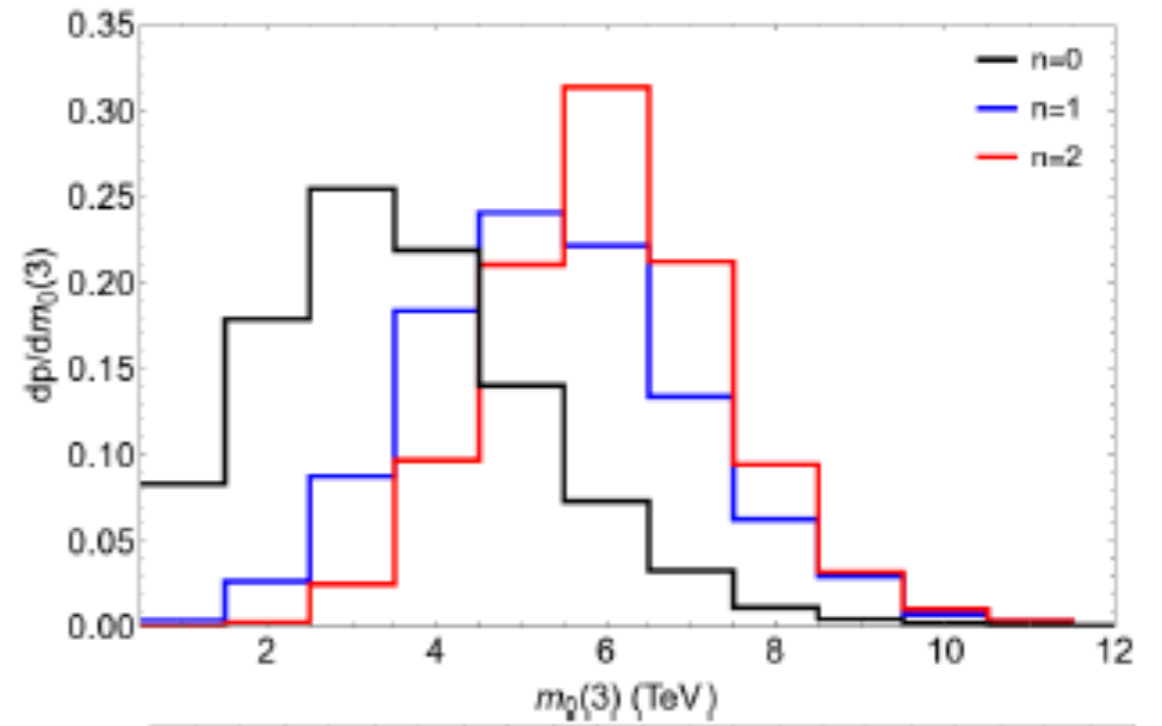
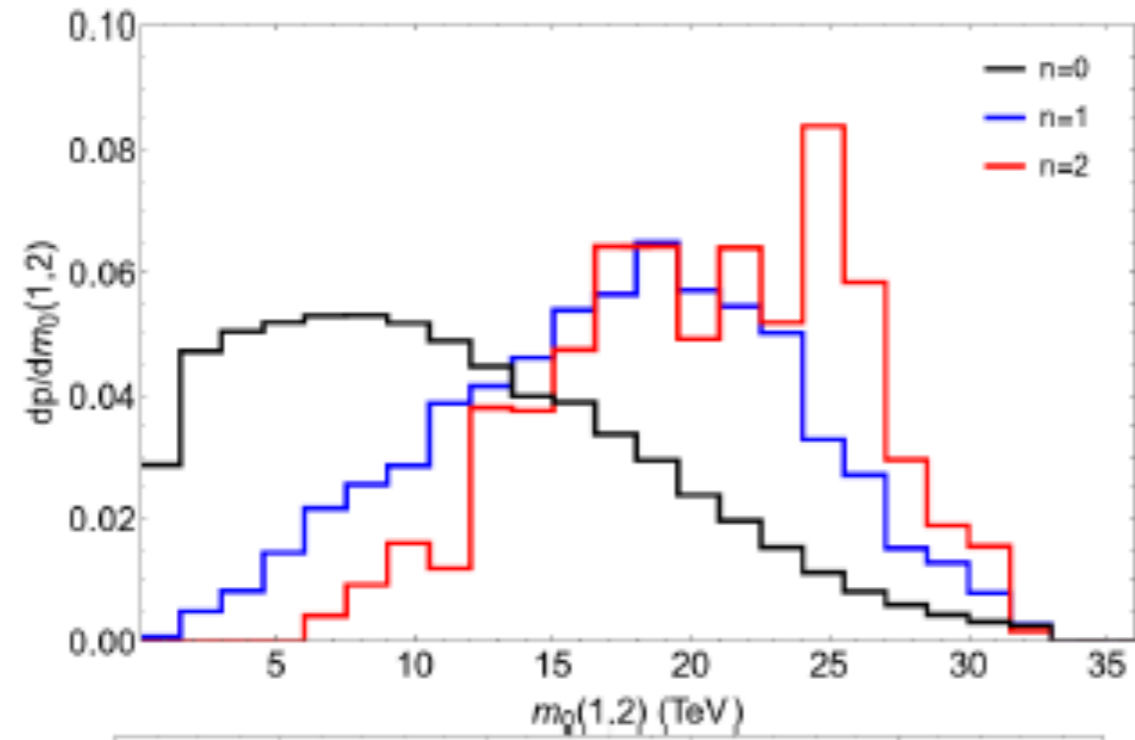
$n \geq 3$ case: soft terms pulled so hard usually gives CCB or no EWSB minima in scalar potential

some conclusions

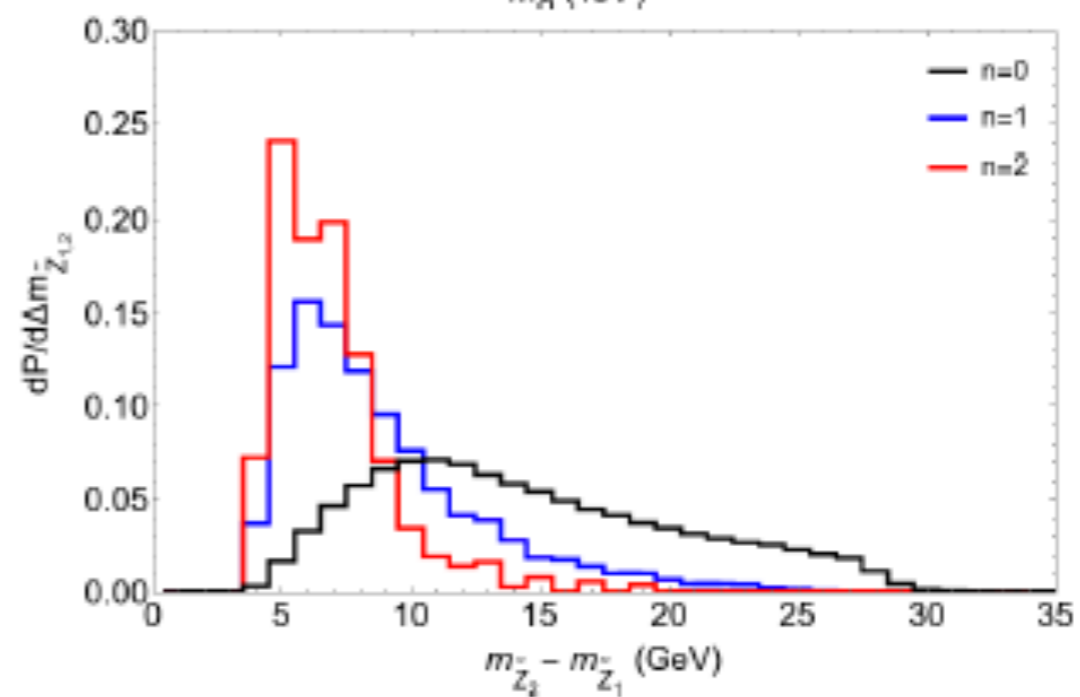
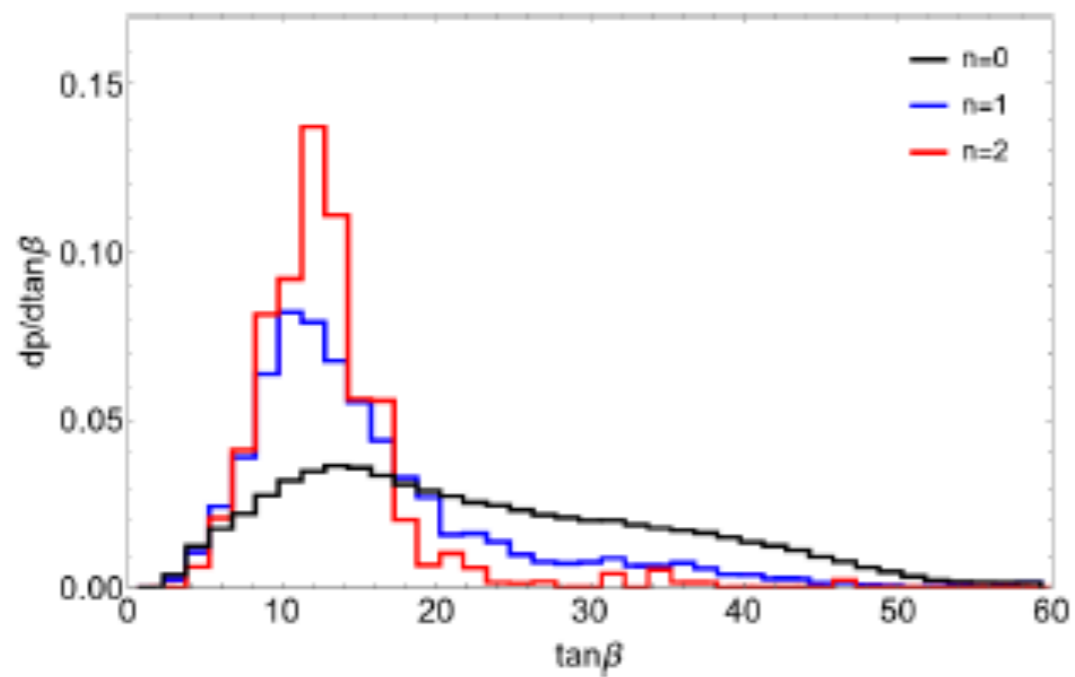
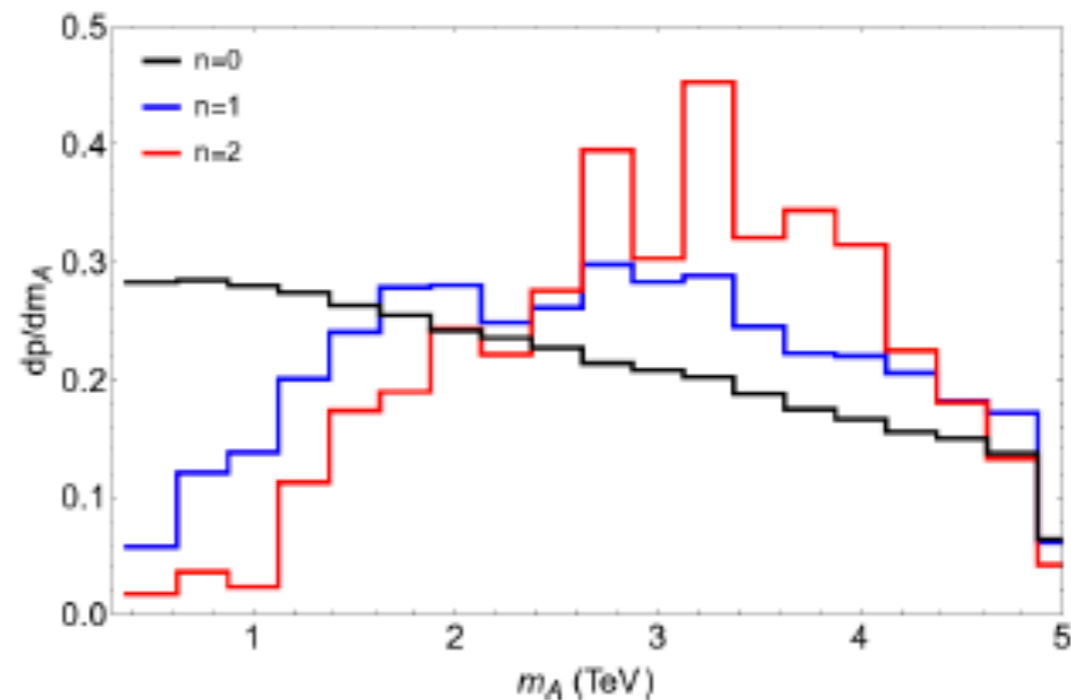
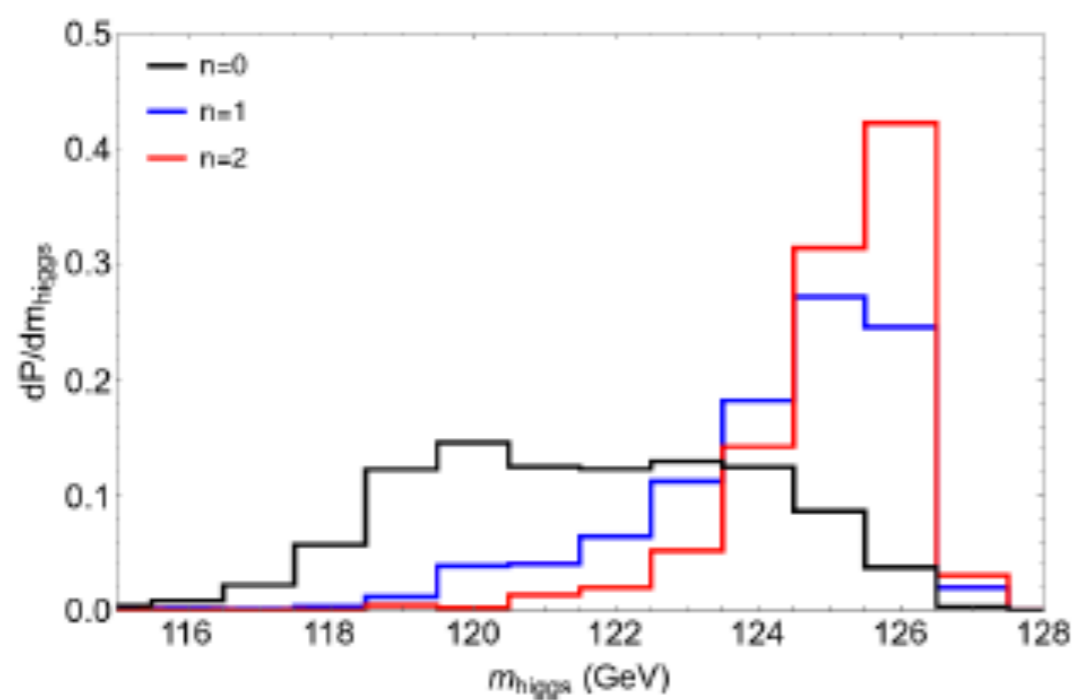
- mu problem modified: $m(\text{weak}) \sim 100\text{--}200 \text{ GeV} \ll m(\text{soft}) \sim \text{multi-TeV}$
- mu term linked to axion physics: Kim-Nilles/SUSY DFSZ
- PQ symmetry radiatively broken as consequence of SUSY breaking: unifies 3 intermediate mass scales: SUSY-breaking, PQ, Majorana ν
- A mild statistical draw on soft terms from the string landscape coupled with anthropic pull of weak scale to $\sim 100 \text{ GeV} \rightarrow m(h) \sim 125 \text{ GeV}$
- The same draw provides a decoupling solution to SUSY flavor, CP, gravitino problem (and cosmological moduli problem) and expect $m(3/2) \sim 10\text{--}30 \text{ TeV}$
- Explains why LHC has so far seen no sign of SUSY
- most robust LHC search channel may be $z_1 z_2 \rightarrow llj + \text{MET}$
- HE-LHC ($\sqrt{s} > 27 \text{ TeV}$) may be needed for gluino/stop discovery
- dark matter a wimp/axion admixture?
- At ILC250, expect Higgs couplings very SM-like; need $\sqrt{s} > 2m(\text{higgsino})$ to establish SUSY discovery/BSM physics

Backup

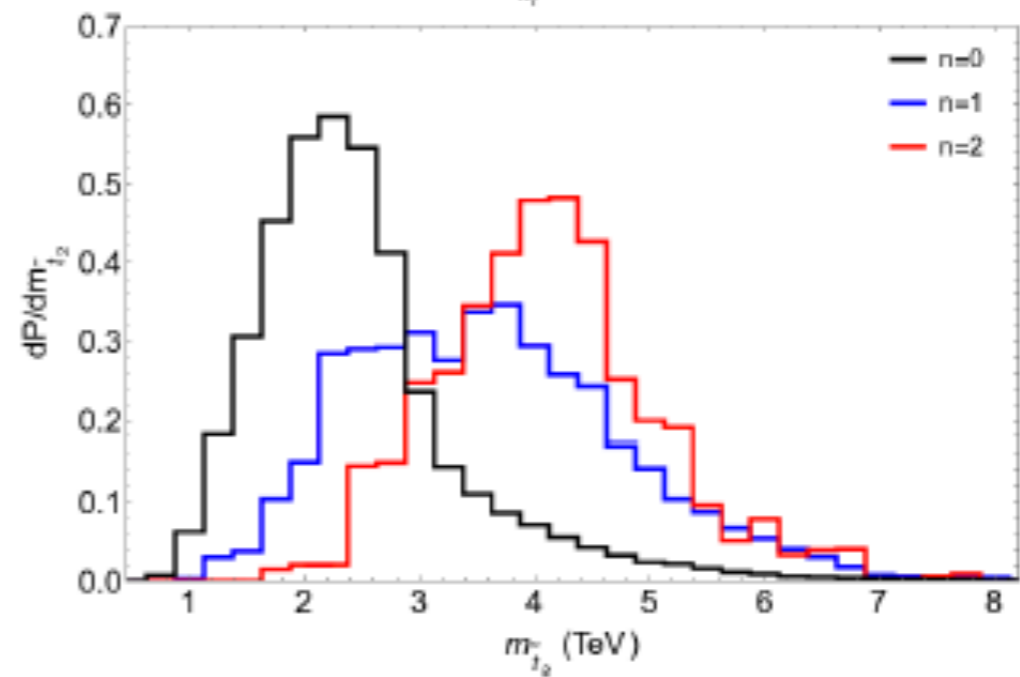
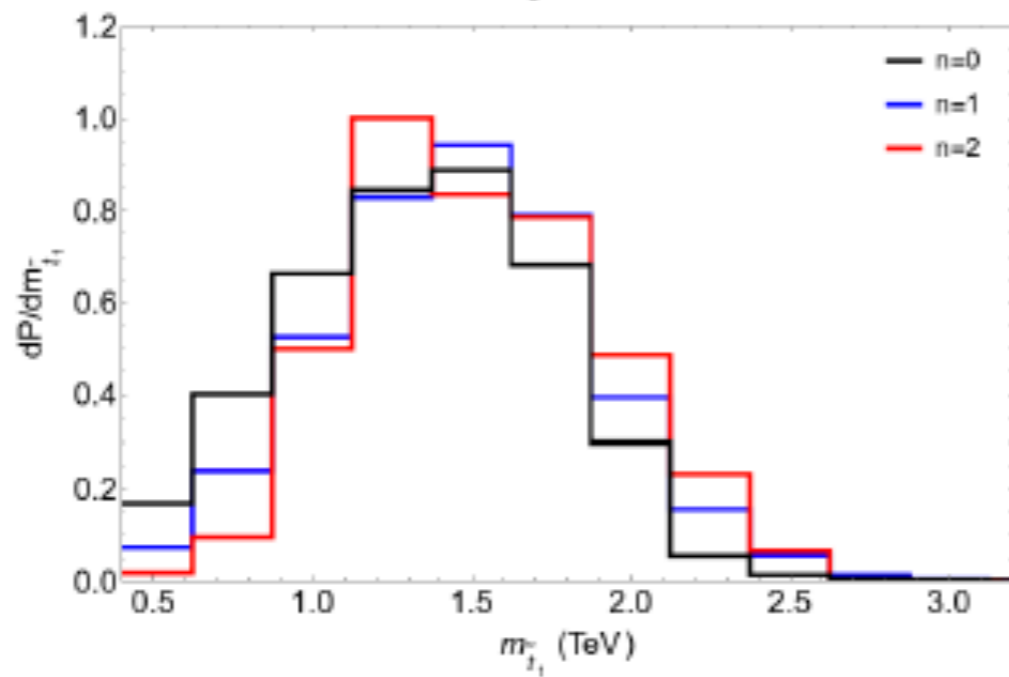
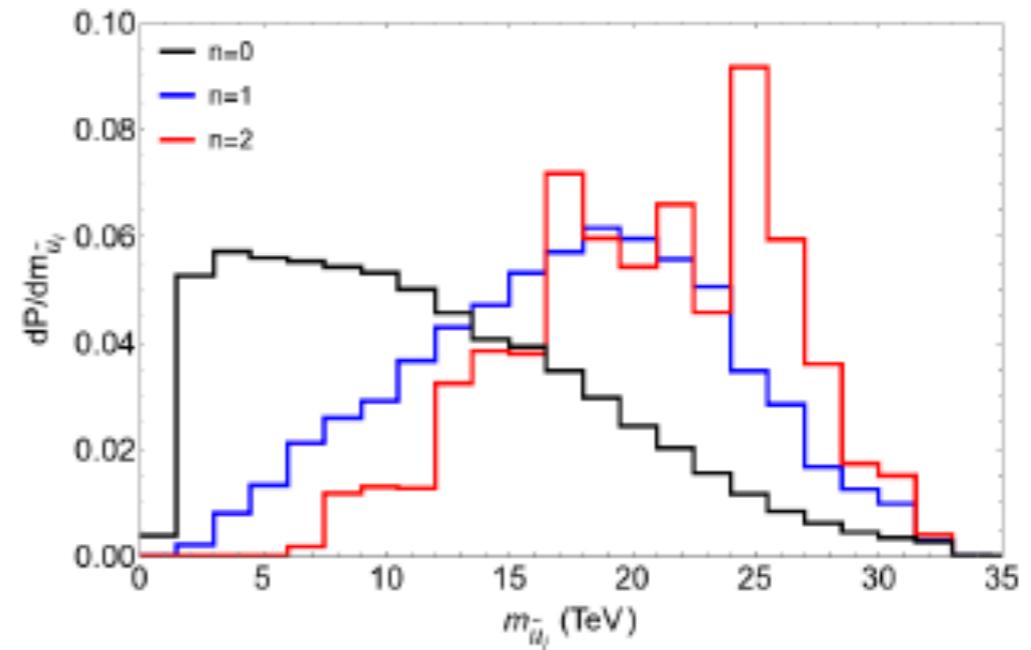
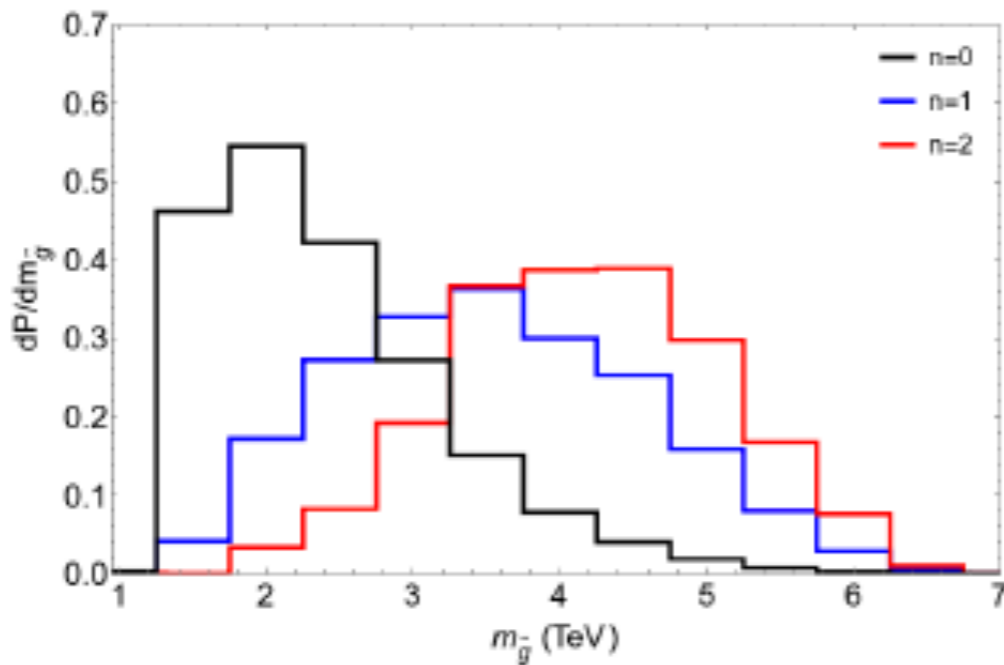
Probability distributions for input soft terms



Higgs and -ino sector



Some key strongly interacting sparticle probability distributions



Prospects for SUSY at LHC:

signature list for radiatively-driven natural SUSY:

- $\tilde{g}\tilde{g}$
- $\tilde{t}_1\tilde{t}_1^*$
- $\tilde{Z}_1\tilde{Z}_2$ (higgsino pair production)
- $\tilde{W}_2^\pm\tilde{Z}_4$ (wino pair production)

Summary of collider searches

- In light of recent LHC bounds ($m(\text{glino}) > 2 \text{ TeV}$, $m(\text{t1}) > 1 \text{ TeV}$) and $m(\text{h})$ requiring TeV-scale highly mixed top squarks, concern has arisen about an emerging Little Hierarchy problem characterized by $m(\text{weak}) \sim 100 \text{ GeV} \ll m(\text{SUSY}) \sim \text{multi-TeV}$ rendering perhaps SUSY as “unnatural”
- We propose an improved naturalness measure based upon scalar potential minimization condition

$$m_Z^2/2 = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \Sigma_u^u(\tilde{t}_{1,2}) - \mu^2$$

This leads to upper bounds from naturalness:

- $m(\text{higgsinos}) \sim 100\text{--}300 \text{ GeV}$ (the lighter the better)
- $m(\text{t1}) < \sim 3 \text{ TeV}$
- $m(\text{glino}) < \sim 6 \text{ TeV}$

DM=WIMP/axion mix?

Conclusions:

1. SUSY still natural;
2. hunt for nSUSY has only begun;
3. HL-LHC handle most SUSY with ino-mass unification;
4. other (e.g. mirage) may require HE-LHC to complete search

process	current	HL-LHC	HE-LHC
glino-glino	$m(\text{glino}) > 2 \text{ TeV}$	$\sim 2.8 \text{ TeV}$	5.5 TeV
t1-t1	$m(\text{t1}) > 1 \text{ TeV}$	1.3 TeV	3.5 TeV
SSdB (winos)	x	$m(\text{W2}) \sim 1 \text{ TeV}$?
z1z2j- >l+l+j+MET	barely	$\mu \sim 250 \text{ GeV}$?

HB, Barger, Gainer, Huang, Tata
Savoy, Mustafayev
Sengupta, Serce

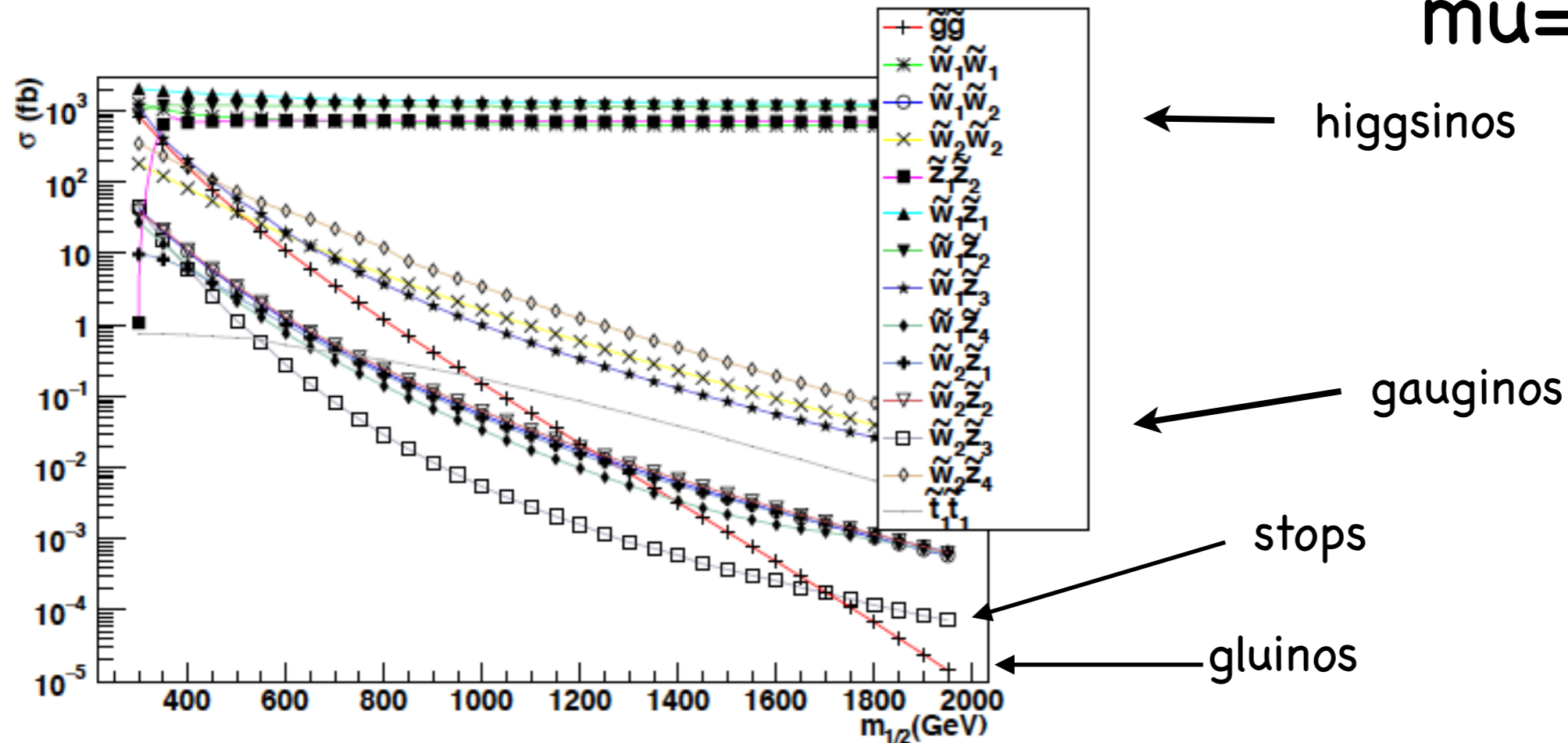
Conclusion: SUSY IS alive and well!

- many calculations of naturalness over-estimate fine-tuning
- naturalness: Little Hierarchy $\mu \ll m(\text{SUSY})$ allowed
- radiatively-driven naturalness: $\mu \sim 100\text{--}200$ GeV, $m(t_1) < 3$ TeV, $m(\text{gluino}) < 5\text{--}6$ TeV
- SUSY DFSZ axion: solve strong CP, solve SUSY μ problem; generate $\mu \ll m(\text{SUSY})$
- landscape pull on soft terms towards RNS, $m(h) \sim 125$ GeV
- natural mirage-mediation/mini-landscape
- natural NUHM2: HL-LHC can cover via $SSdB+Z1Z2j$ channels
- natural mirage/mini-landscape may escape detection at HL-LHC; need LHC33!
- expect ILC as higgsino factory
- DM= axion+higgsino-like WIMP admixture: detect both?

Backup

Sparticle prod'n along RNS model-line at LHC14:

$\mu=150$ GeV



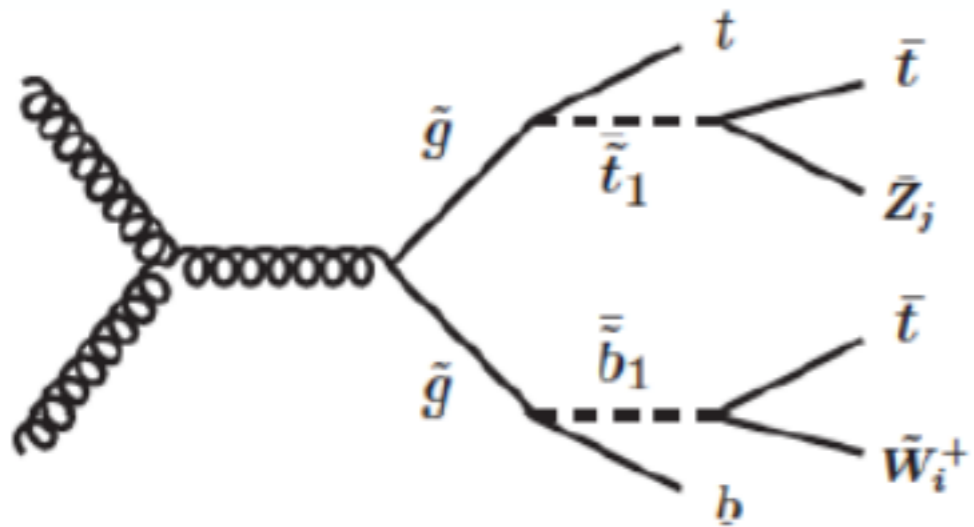
higgsino pair production dominant-but only soft visible energy release from higgsino decays

largest visible cross section: wino pairs

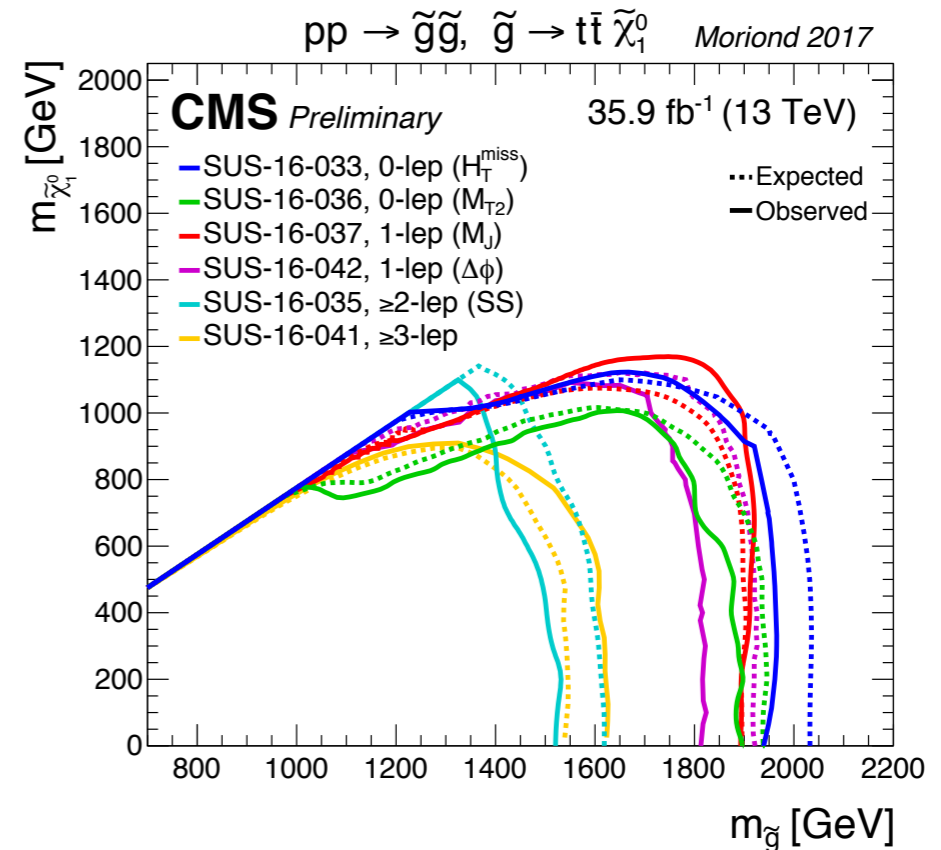
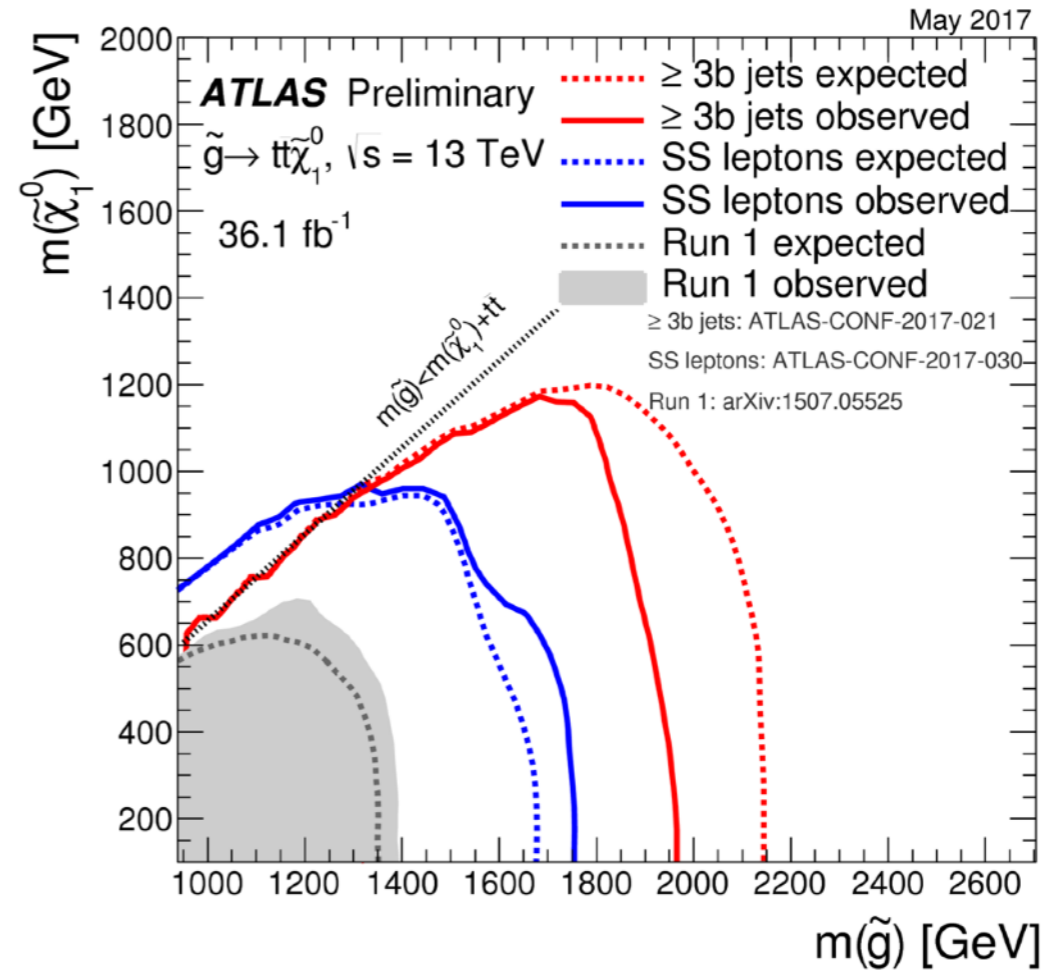
gluino pairs sharply dropping

stops at bottom

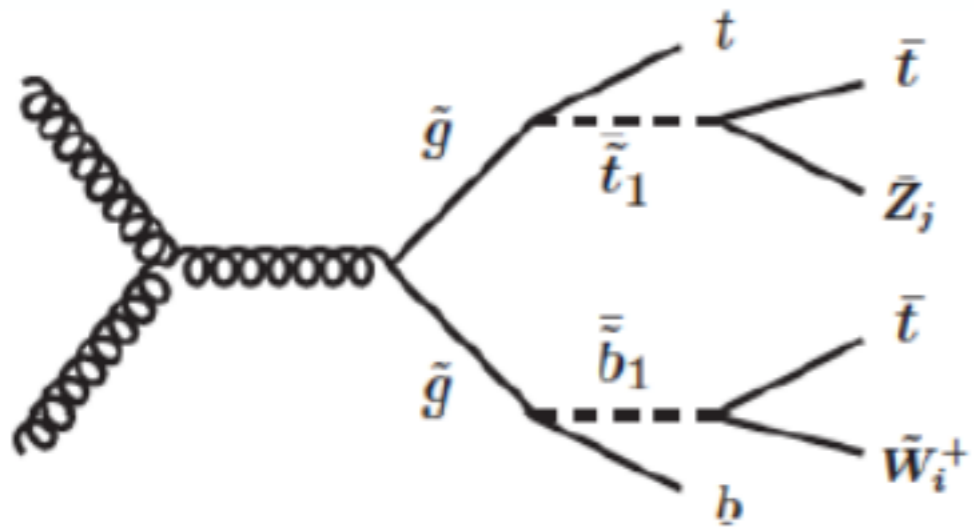
gluino pair cascade decay signatures



Current limits for $m(Z_1) \sim 150$ GeV:
 $m(\text{gluino}) > \sim 2$ TeV

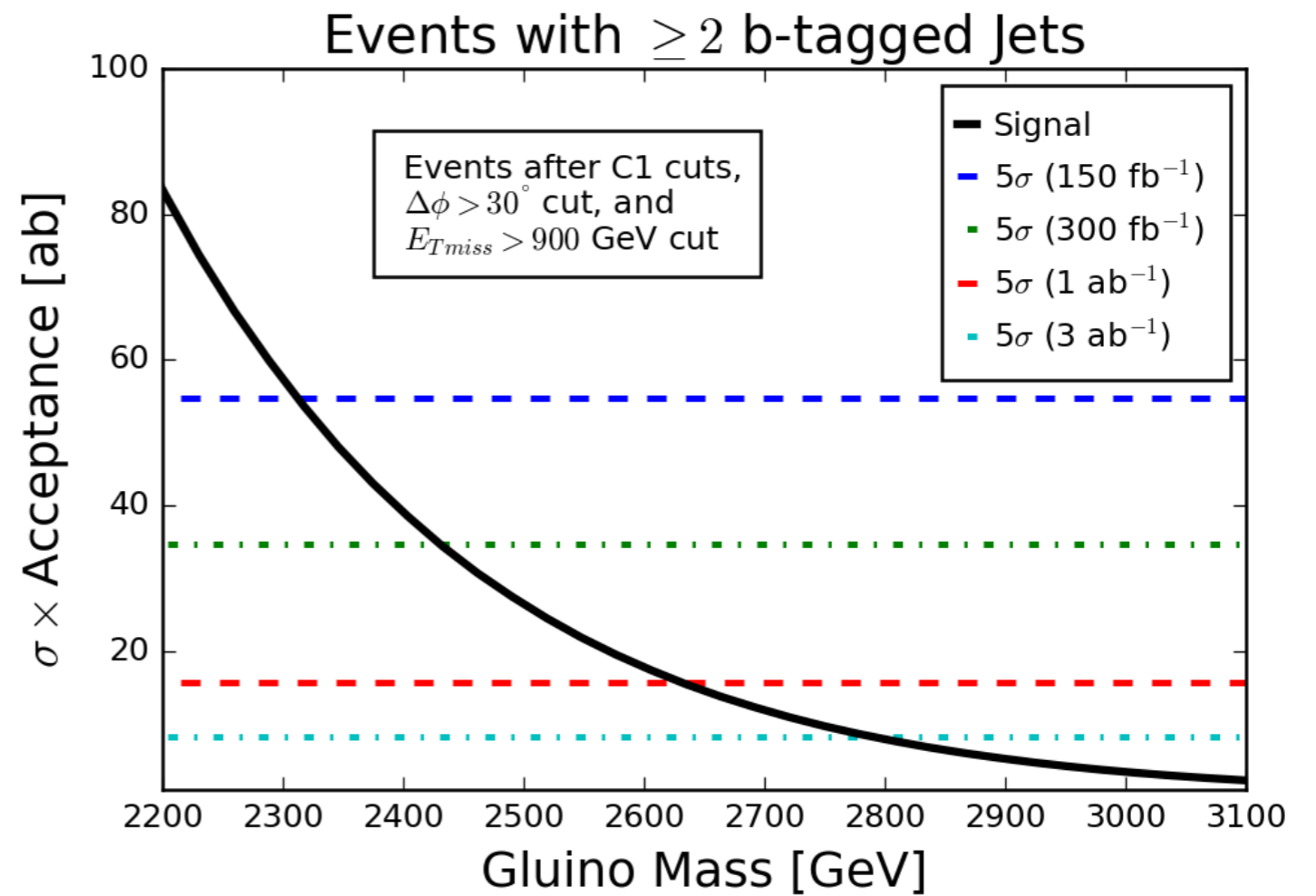


gluino pair cascade decay signatures

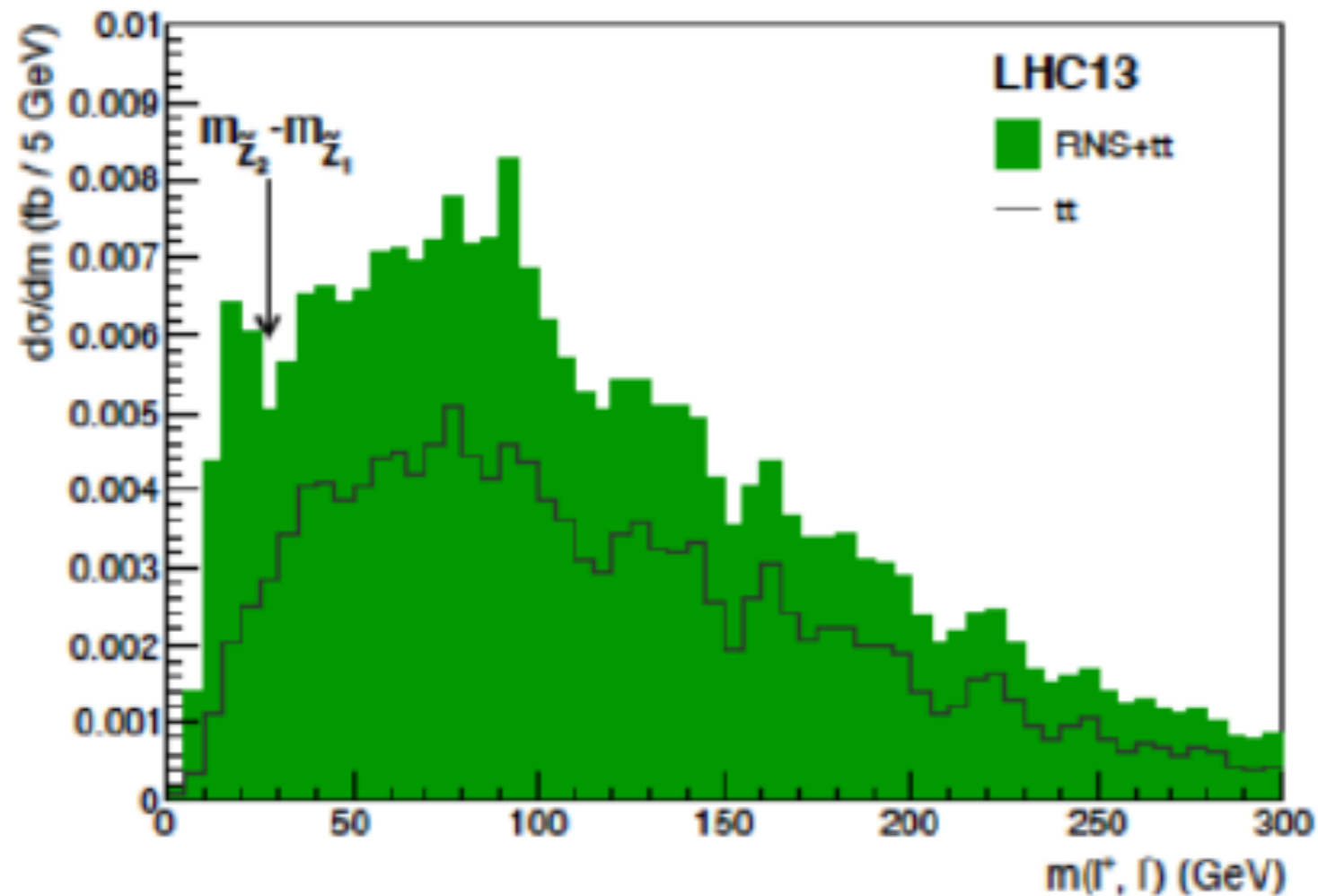


Estimated HL-LHC reach for gluinos

HL-LHC reach to
 $m(\text{gluino}) \sim 2.8 \text{ TeV}$;
 RNS: $m(\text{gluino}) < \sim 5 \text{ TeV}$

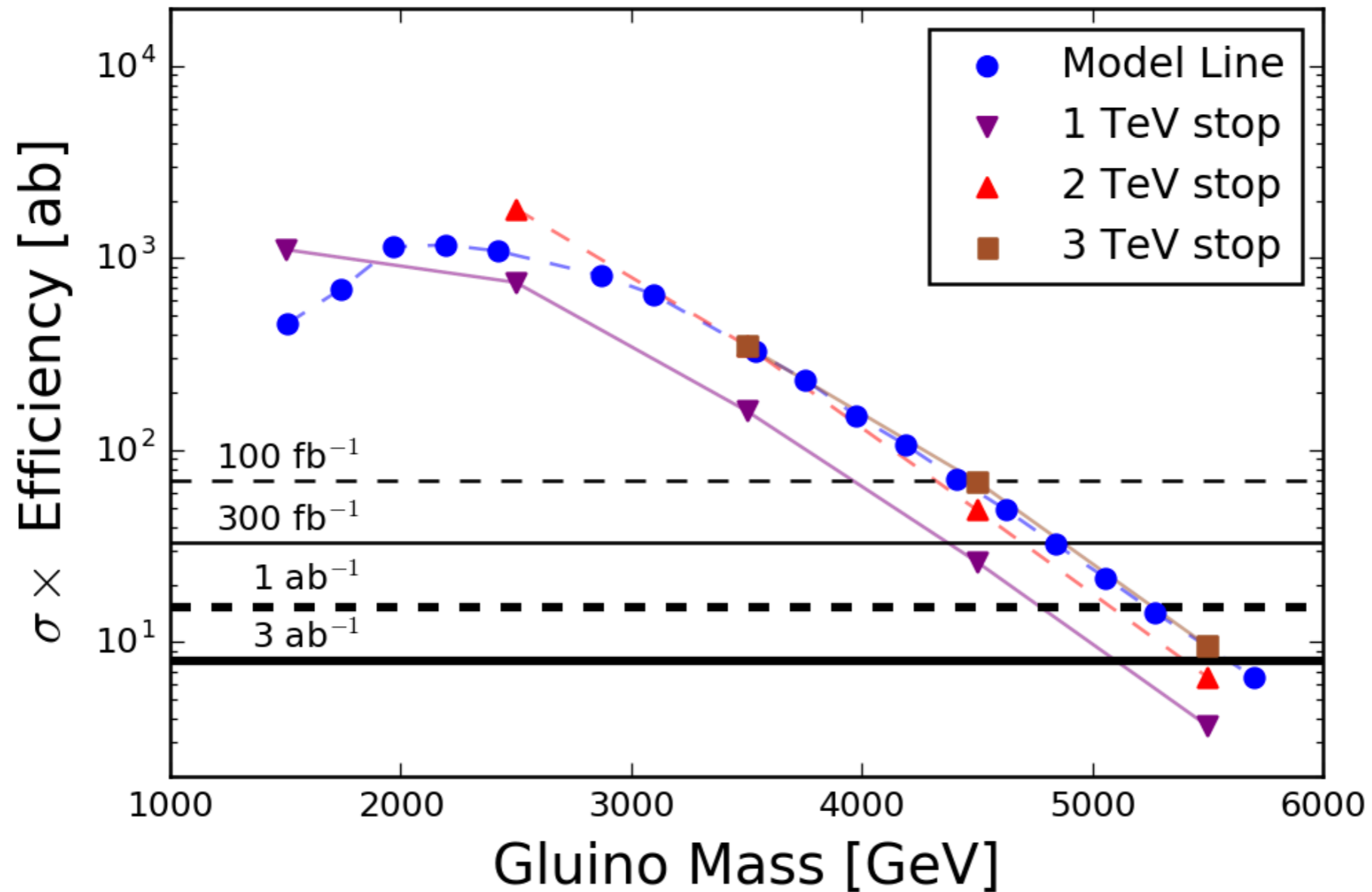


LHC14 has some reach for
gluino pair production in RNS;
if a signal is seen,
should be distinctive



OS/SF dilepton mass
edge apparent from
cascade decays
with $z_2 \rightarrow z_1 + l + l^{\text{bar}}$

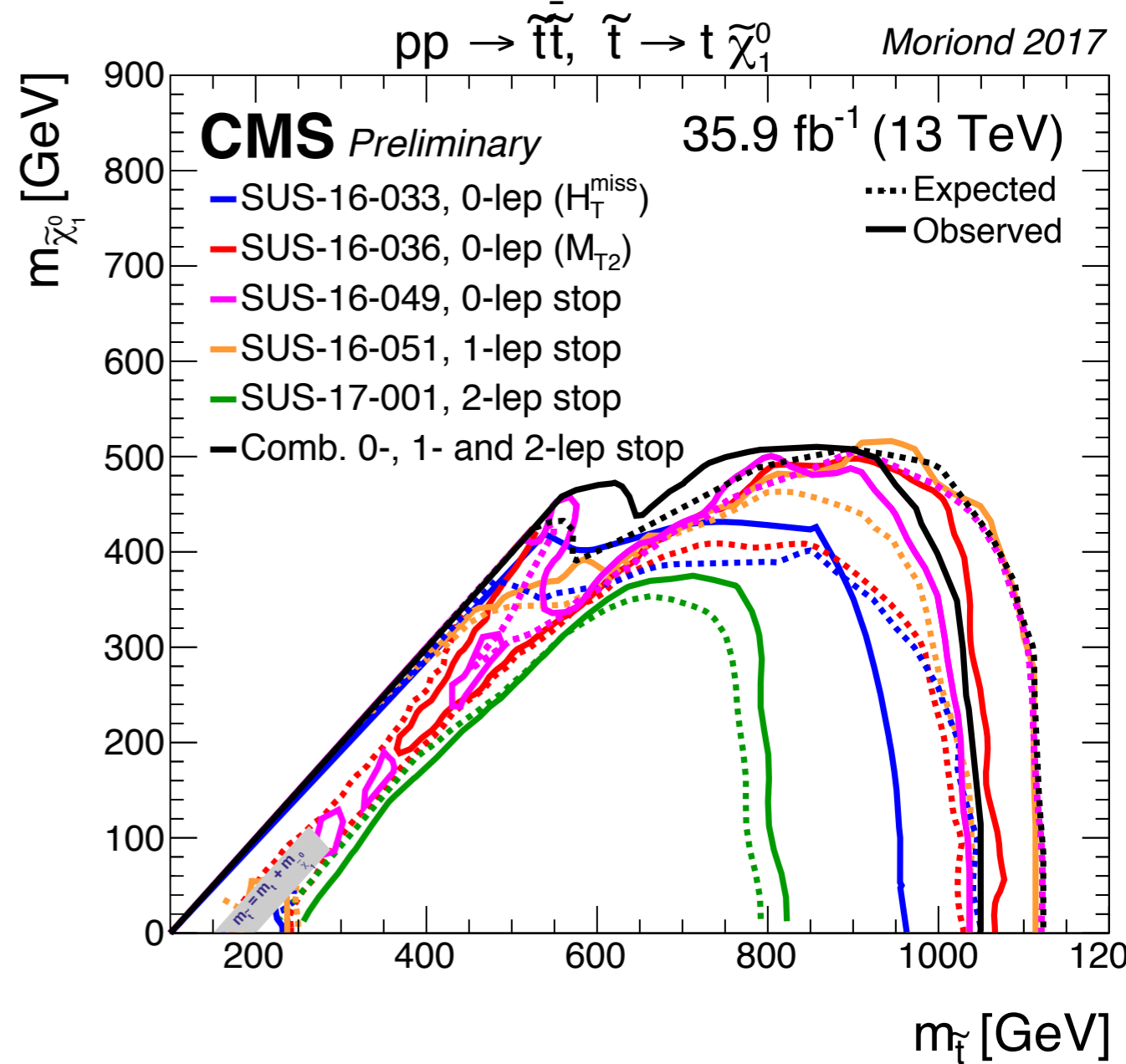
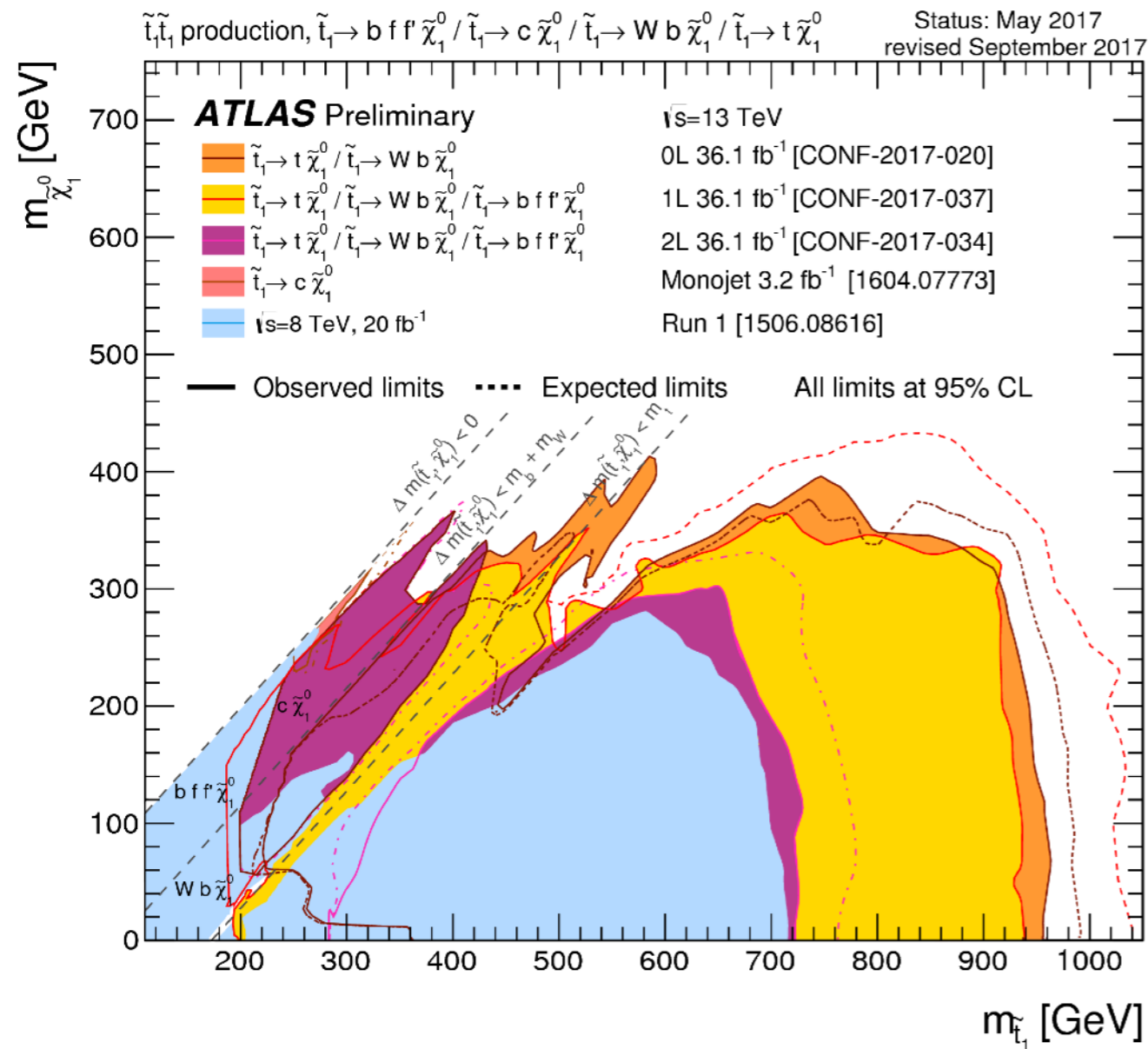
Glauino 5-sigma reach at LHC33: to about $m(\text{gluino}) \sim 5\text{--}5.5 \text{ TeV}$



≥ 4 jets; ≥ 2 -b-jets; $\text{MET} > 1500 \text{ GeV}$

HB, Barger, Gainer, Huang, Savoy, Serce, Tata

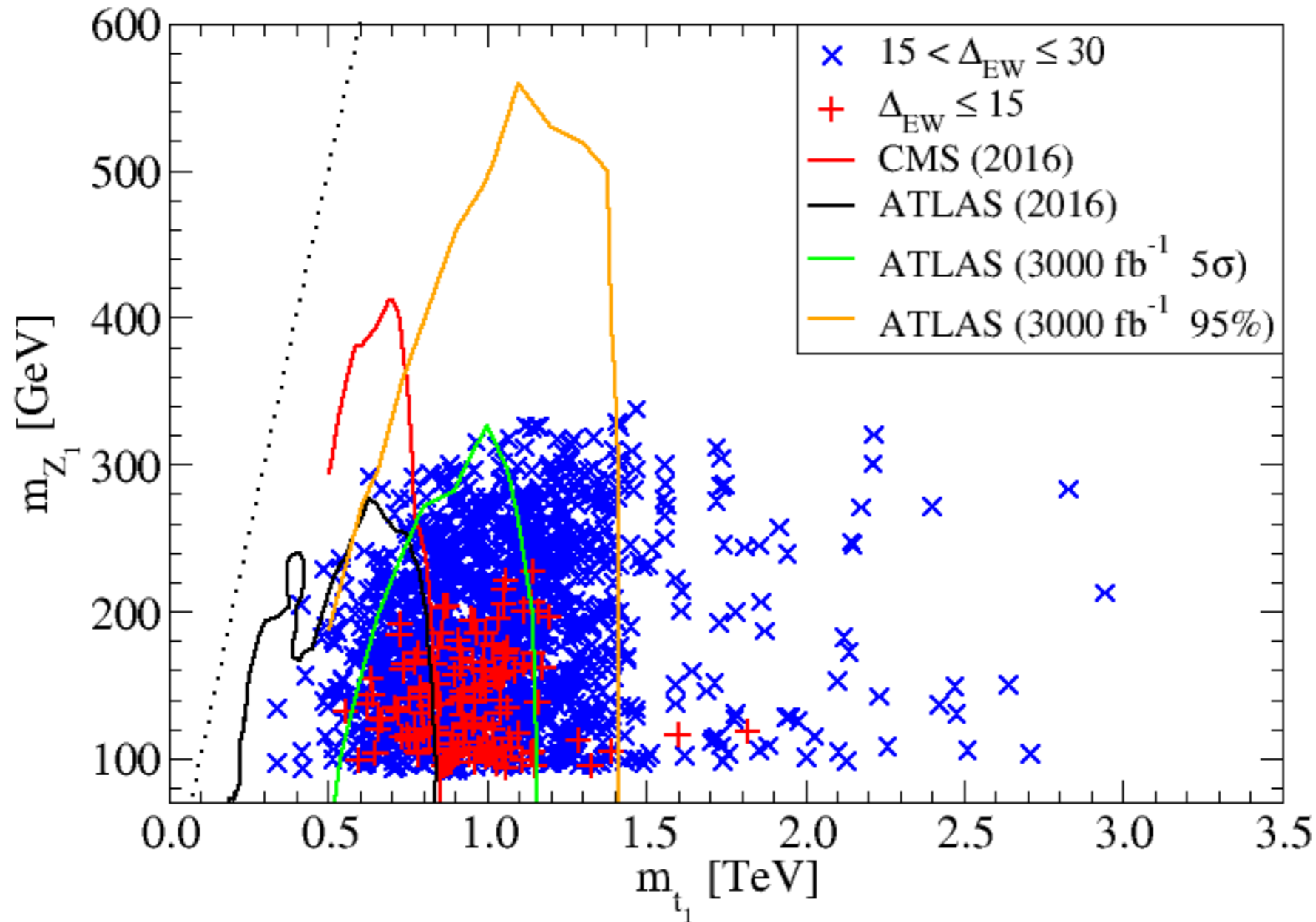
Present limits on top squarks from LHC



Evidently $m(t_1) > \sim 1 \text{ TeV}$ for $m(\text{LSP}) \sim 150 \text{ GeV}$

- * TeV-scale top squark needed for $m(h) \sim 125 \text{ GeV}$
- * Also needed for $b \rightarrow s \text{ gamma}$

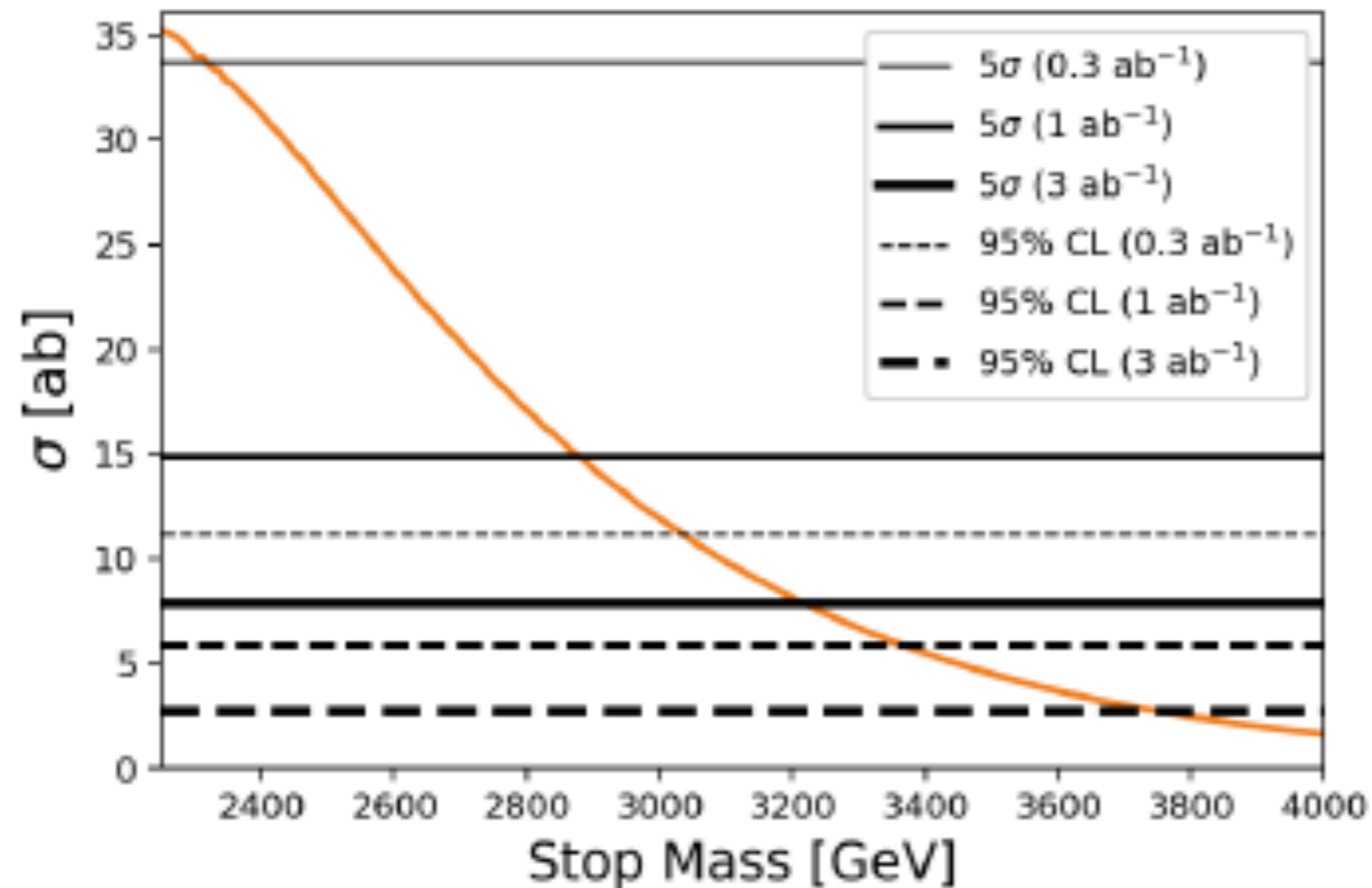
Prospects for top squarks in natural SUSY



$m(t_1)$ can range up to 3 TeV with little cost to naturalness;
the hunt for stops has only begun!

HL-LHC reach extends to $m(t_1) \sim 1.2-1.4$ TeV

Reach of LHC33 for top squarks



- $\tilde{t}_1 \rightarrow b\tilde{W}_1; \sim 50\%$

- $\tilde{t}_1 \rightarrow t\tilde{Z}_1; \sim 25\%$

- $\tilde{t}_1 \rightarrow t\tilde{Z}_2; \sim 25\%$

- A. $\tilde{t}_1\tilde{t}_1^* \rightarrow b\bar{b} + E_T^{\text{miss}} \sim 25\%$,

- B. $\tilde{t}_1\tilde{t}_1^* \rightarrow b\bar{t}, \bar{b}t + E_T^{\text{miss}} \sim 50\%$,

- C. $\tilde{t}_1\tilde{t}_1^* \rightarrow t\bar{t} + E_T^{\text{miss}} \sim 25\%$.

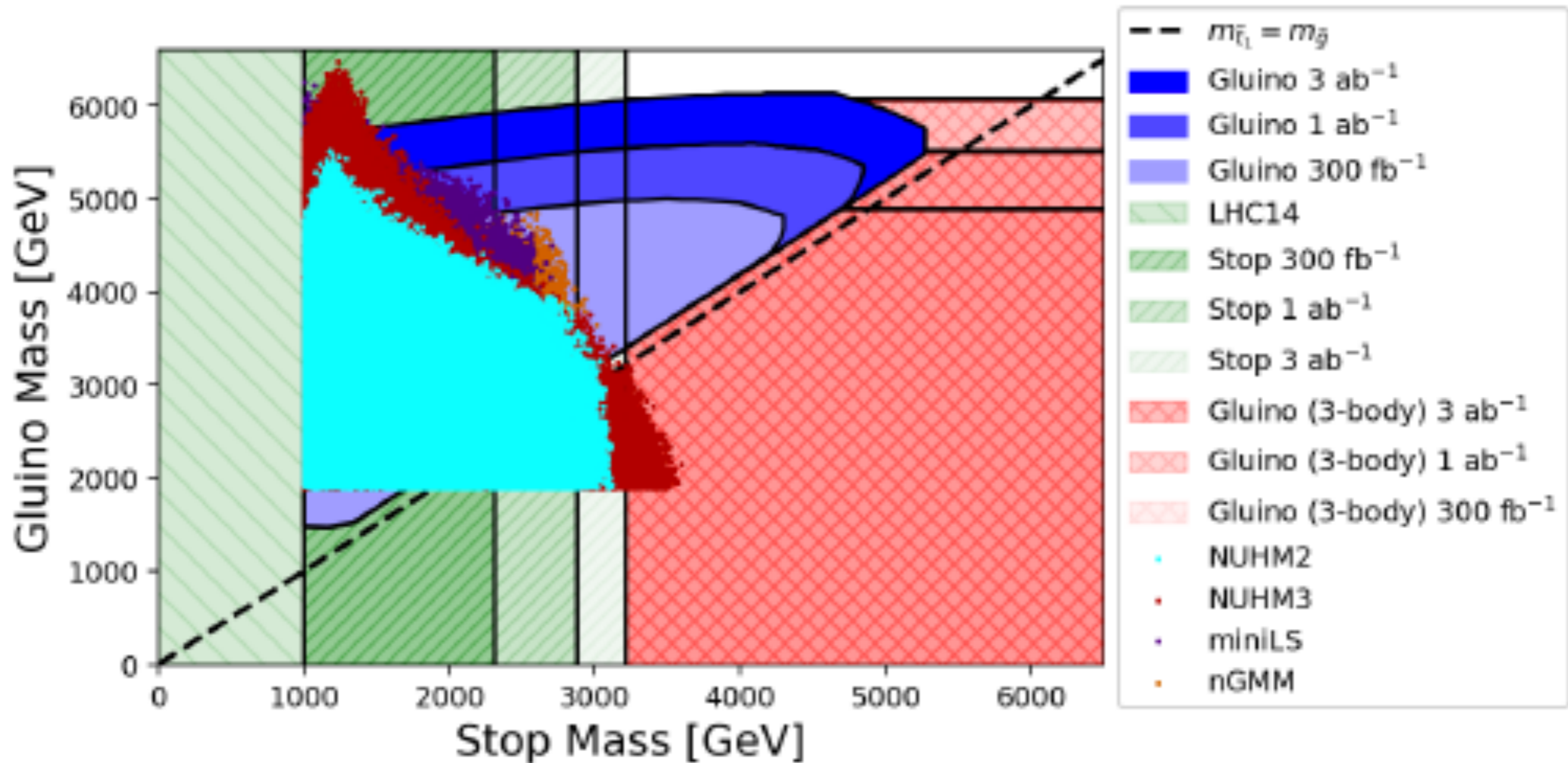
LHC33 reach extends to $m(t_1) \sim 3-3.8$ TeV

$n(b\text{-jets}) \geq 2; \text{MET} > 750$ GeV

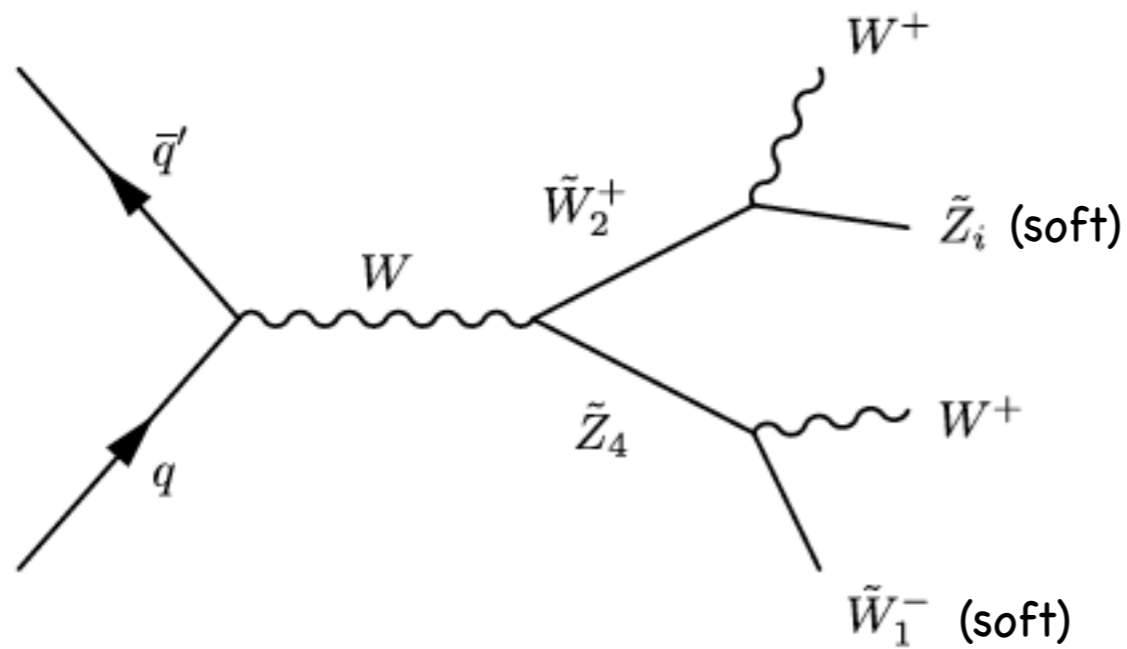
HB, Barger, Gainer, Serce, Tata

Combined LHC33 reach for $t1$ and $g1no$
covers all natural SUSY p-space!

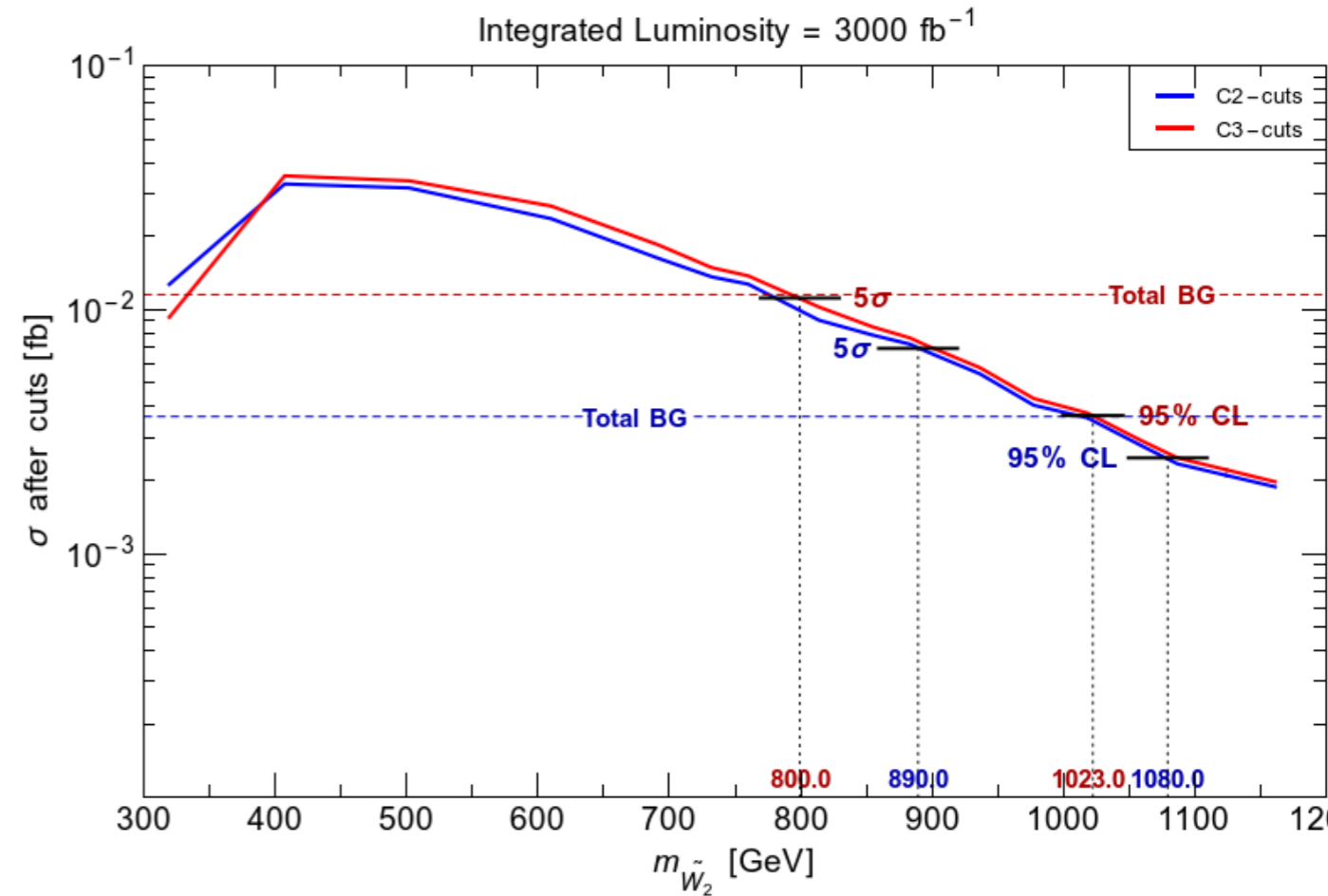
(need to re-do for LHC27)



Distinctive same-sign diboson (SSdB) signature from SUSY models with light higgsinos!

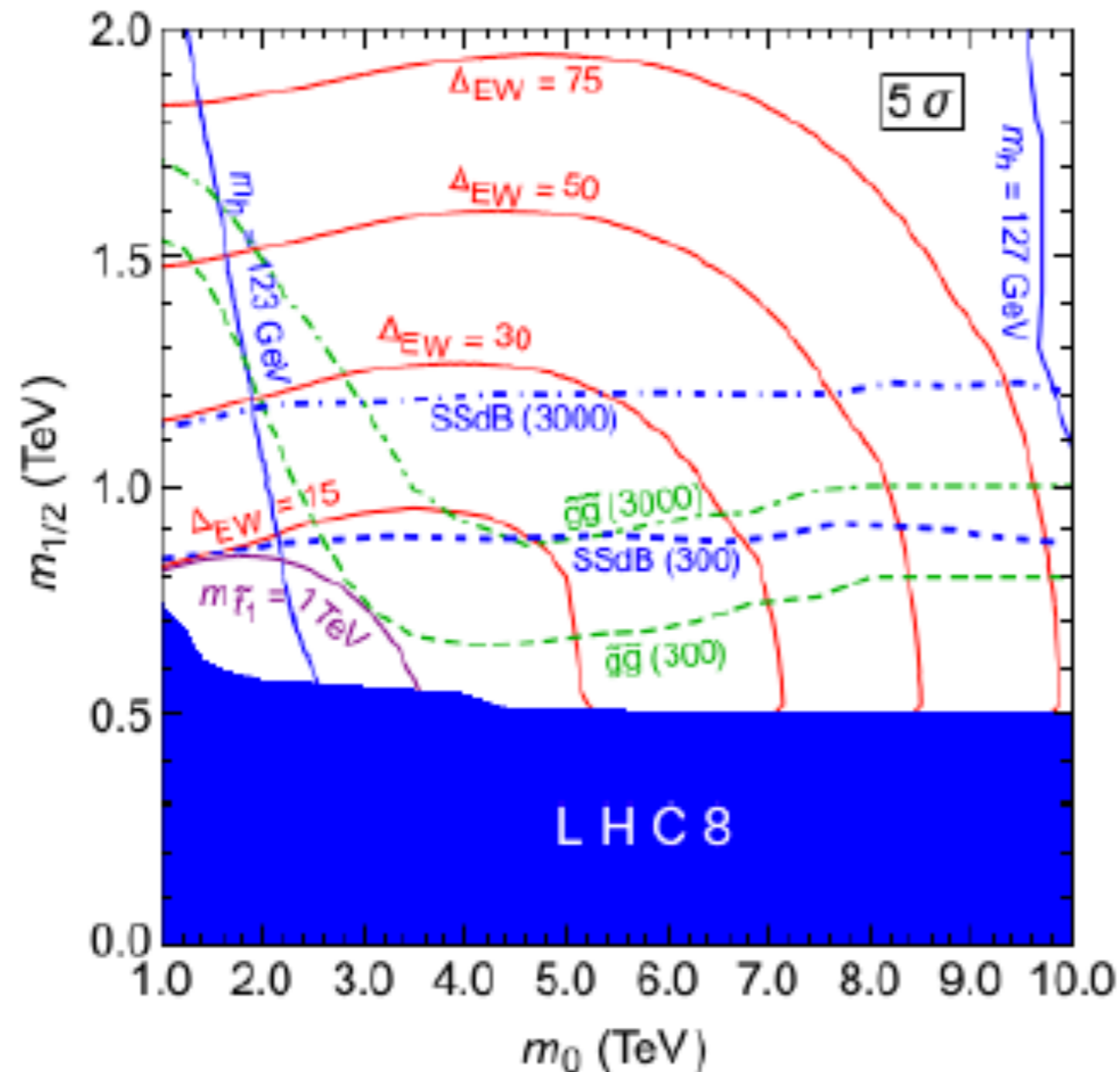


wino pair production



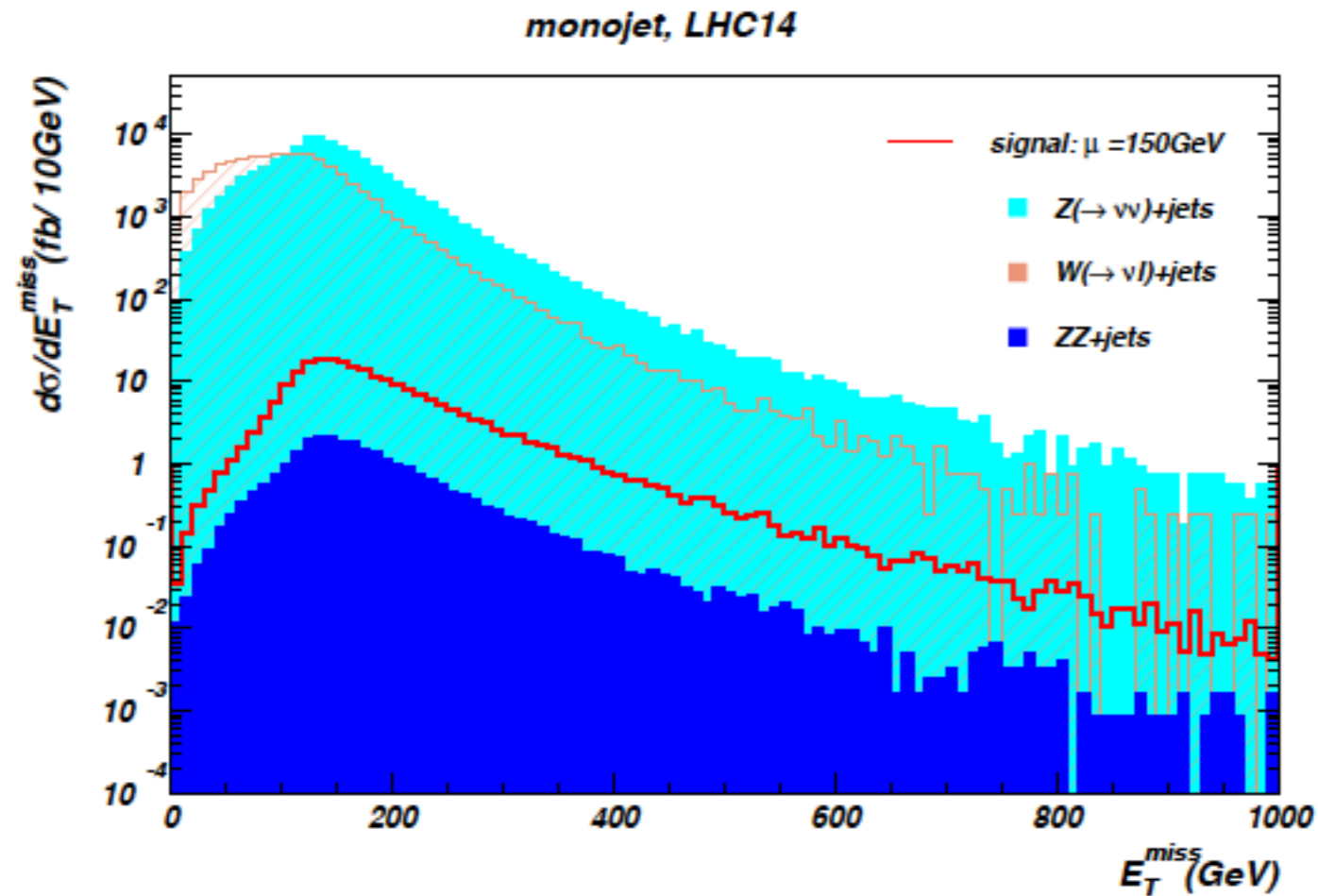
This channel offers good reach of LHC14 for RNS;
it is also indicative of wino-pair prod'n
followed by decay to higgsinos

Good old m_0 vs. $m_{1/2}$ plane still viable, but needs $\mu \sim 100\text{--}200$ GeV as possible in NUHM2 instead of CMSSM/mSUGRA



For models with no mass unif'n, reach via SSdB may exceed $g\bar{g}$ pairs for high luminosity

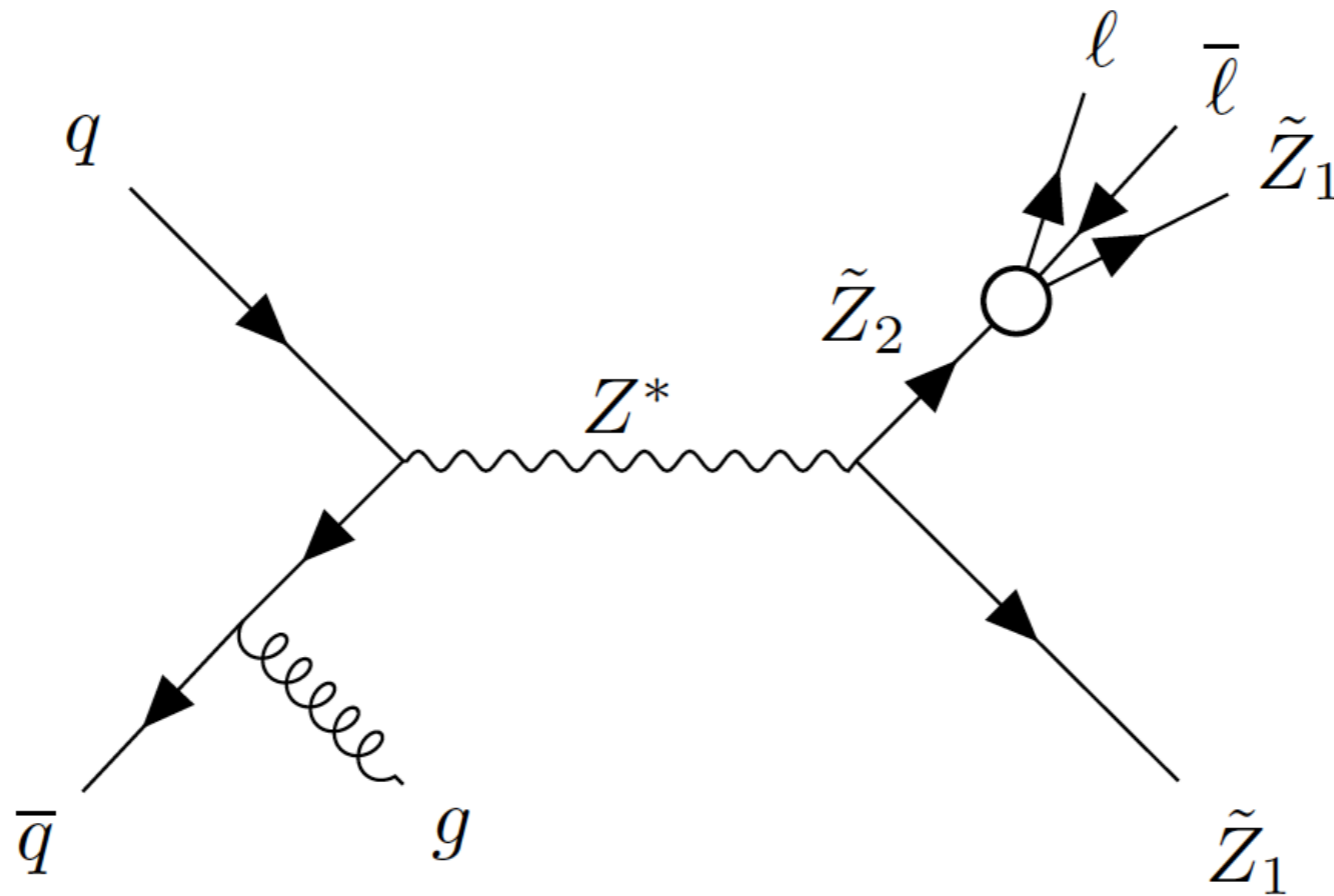
See direct higgsino pair production recoiling from ISR (monojet signal)?



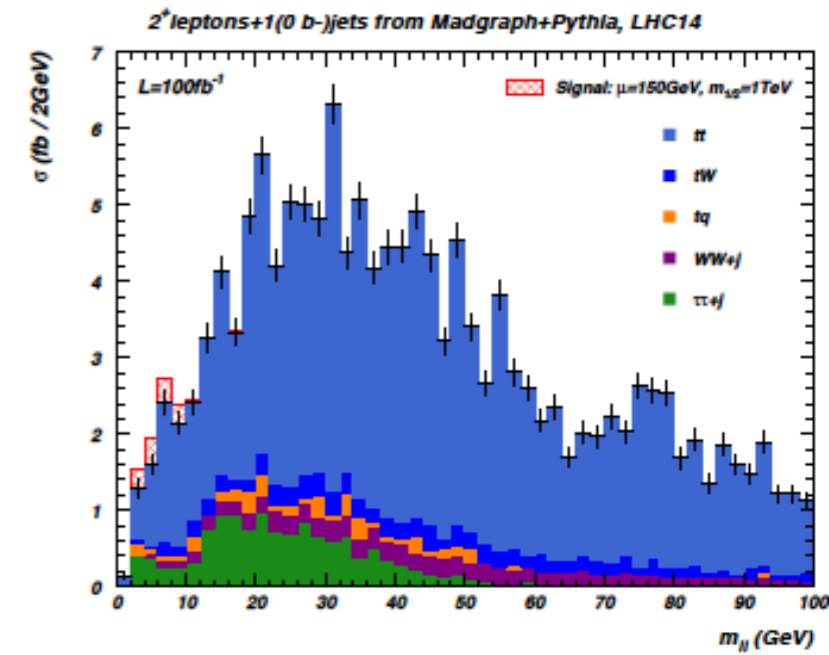
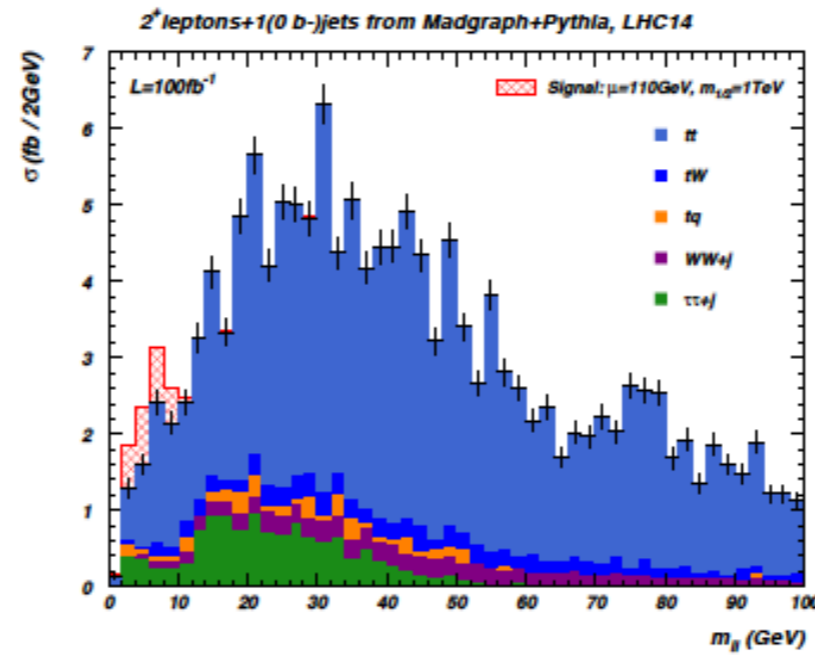
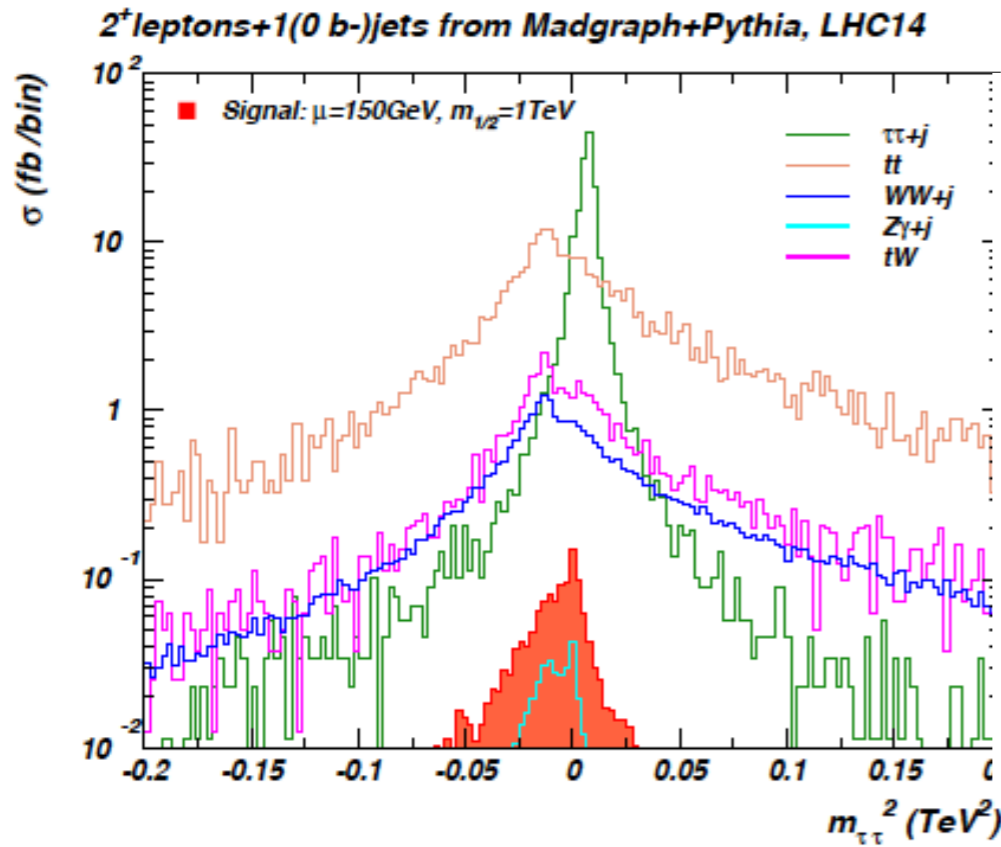
typically 1% S/BG after cuts:
very tough to do!

What about $pp \rightarrow \tilde{Z}_1 \tilde{Z}_2 j$ with $\tilde{Z}_2 \rightarrow \tilde{Z}_1 \ell^+ \ell^-$?

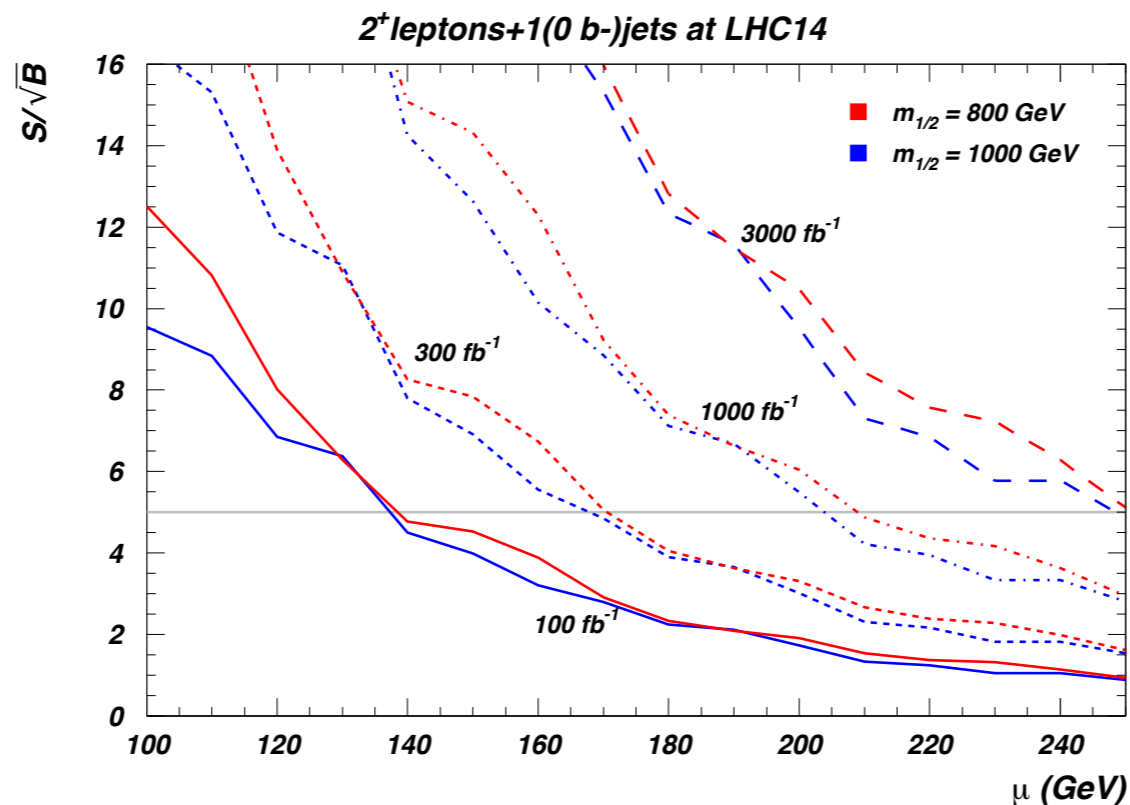
Han, Kribs, Martin, Menon, PRD89 (2014) 075007;
HB, Mustafayev, Tata, PRD90 (2014) 115007;



use MET to construct $m^2(\text{tau-tau})$



cut $m(\text{ditau})^2 < 0$



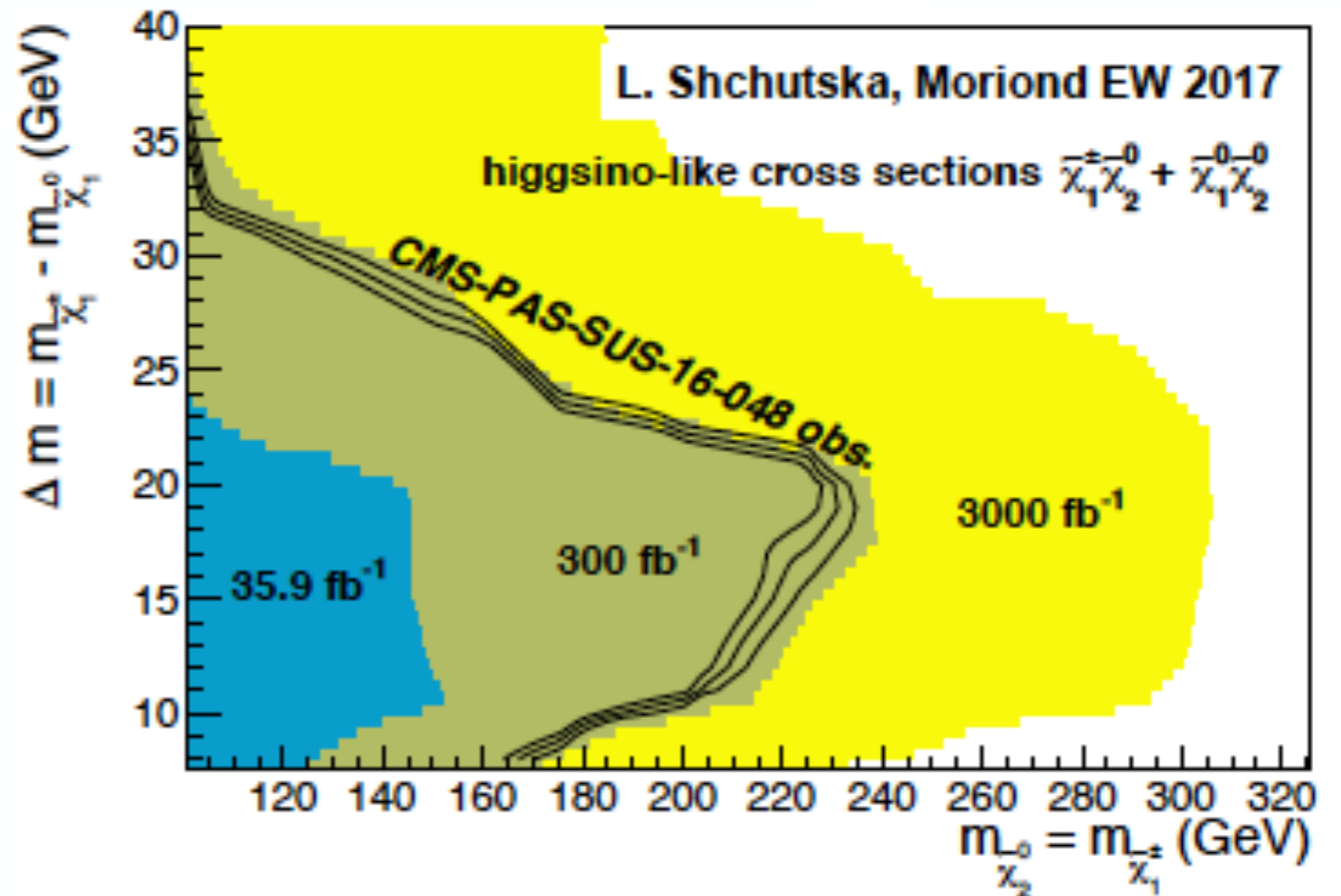
HL-LHC 5-sigma reach
to $\mu \sim 250\text{ GeV}$!

HB, Mustafayev, Tata

CMS analysis: this may be **the most important SUSY discovery channel at LHC** since it directly probes higgsinos which can't be too far from $m(W,Z,h)$

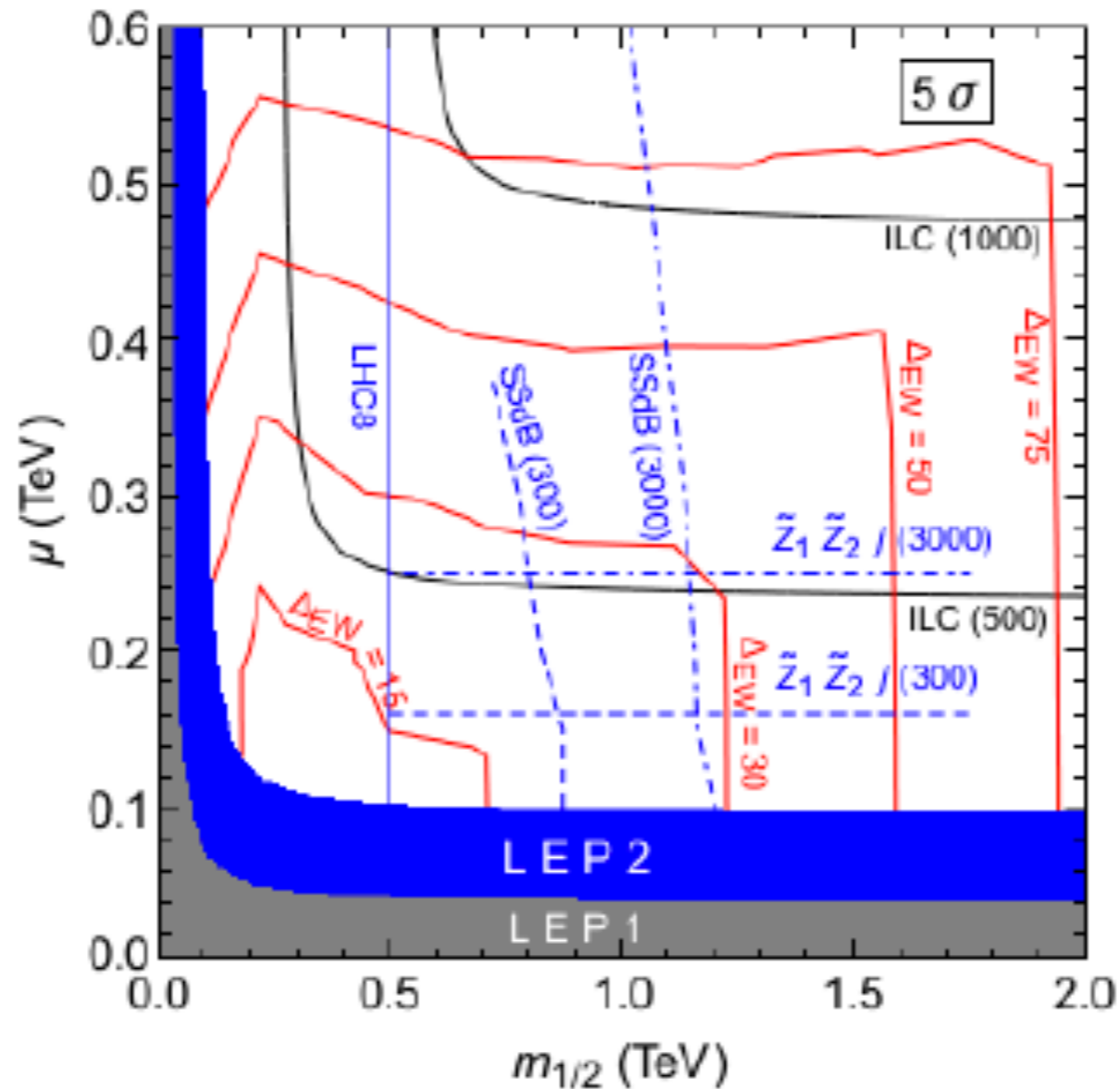
Atlas study underway- results soon?

Higgsino cross section (projection only)



NatSUSY z2-z1 mass gap may range down to 3 GeV so need to ID very soft, low $m(l\bar{l})$ leptons

panoramic view of reach of HL-LHC for natural SUSY



Combined SSdB/ljMET searches may cover all Nat SUSY p-space at HL-LHC for models with no mass unification; in mirage scenario, z_2-z_1 mass gap can be reduced and M_2 can be much higher than in NUHM2

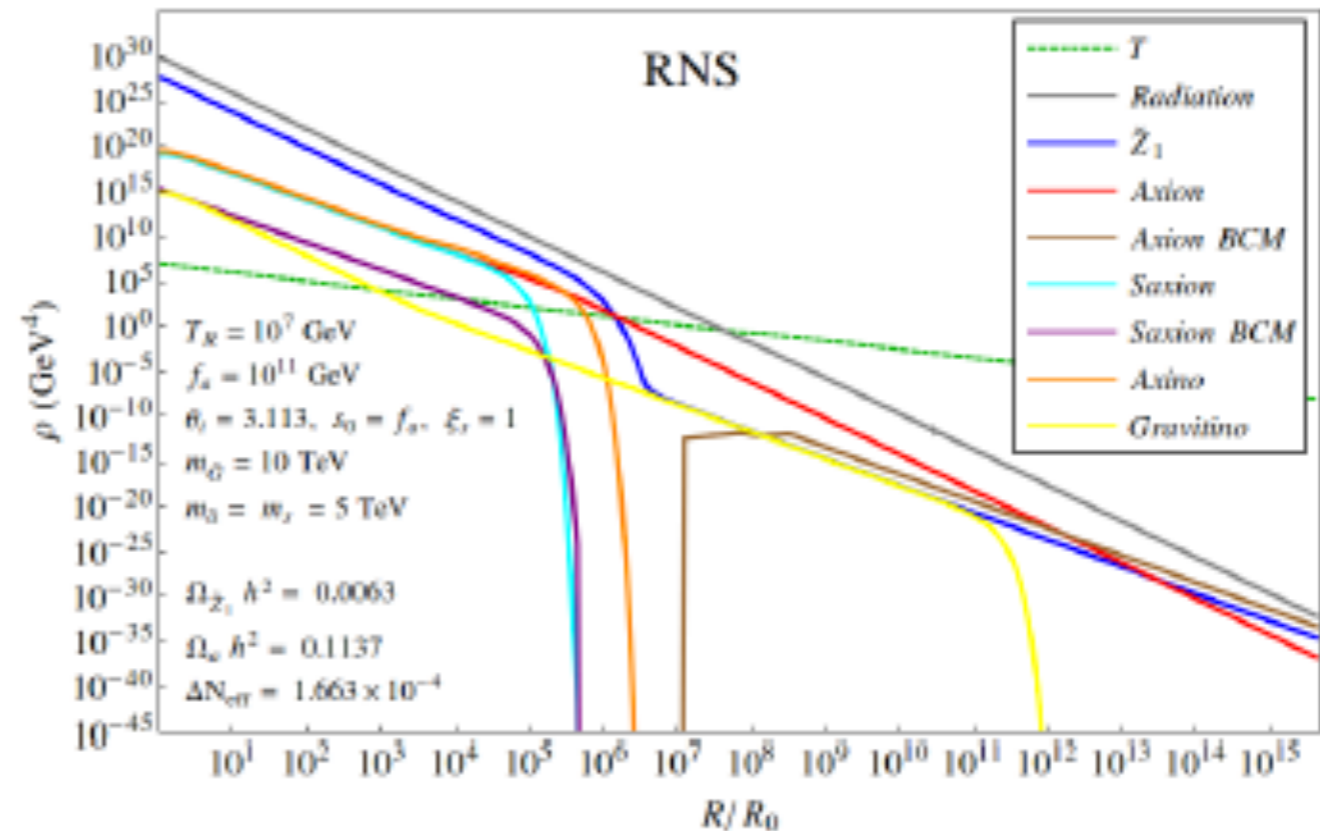
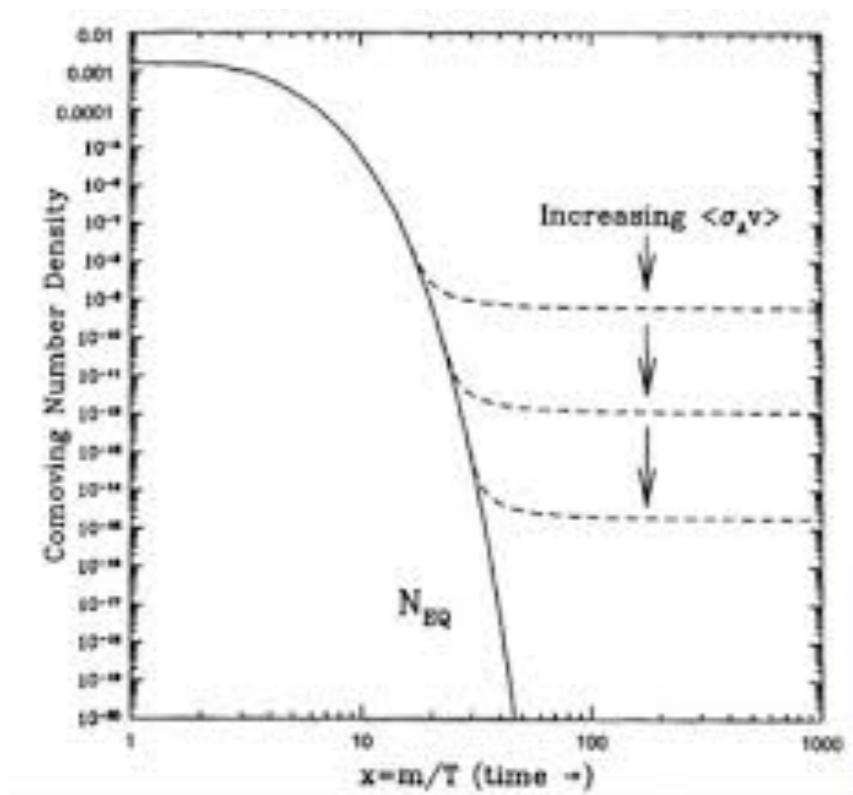
What happens to SUSY WIMP dark matter?

- higgsino-like WIMPs thermally underproduced
- 3 not four light pions \Rightarrow QCD theta vacuum
- $F\tilde{F}$ term should be present but neutron(EDM) \Rightarrow it is tiny
- strong CP problem \Rightarrow axions: no fine-tuning in QCD sector
- SUSY context: axion superfield, axinos and saxions
- DM = axion+higgsino-like WIMP admixture
- DFSZ SUSY axion: solves mu problem with $\mu \ll m_{3/2}$!
- ultimately detect both WIMP and axion?

usual picture

=>

mixed axion/WIMP



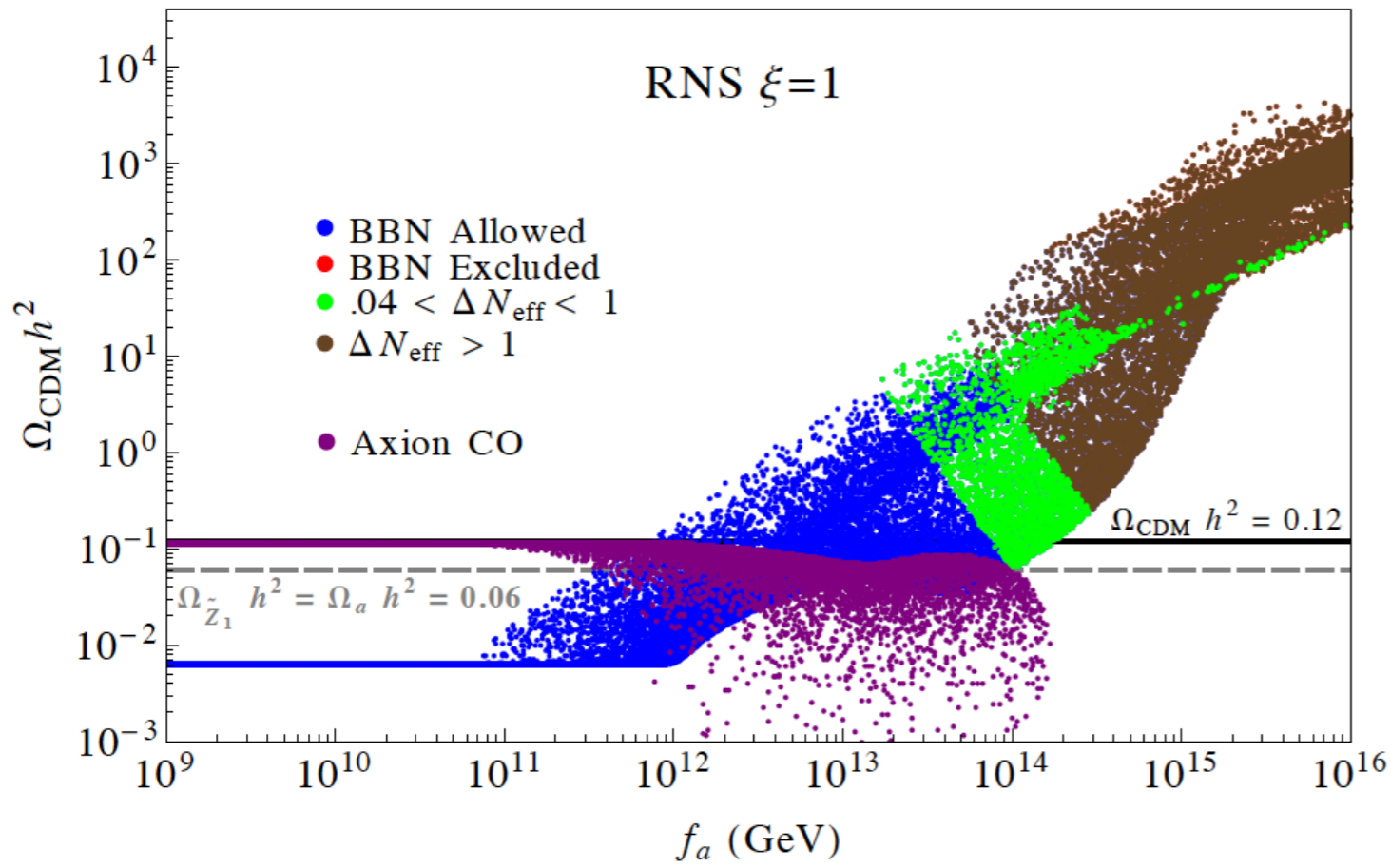
KJ Bae, HB, Lessa, Serce

much of parameter space is axion-dominated
with 10-15% WIMPs



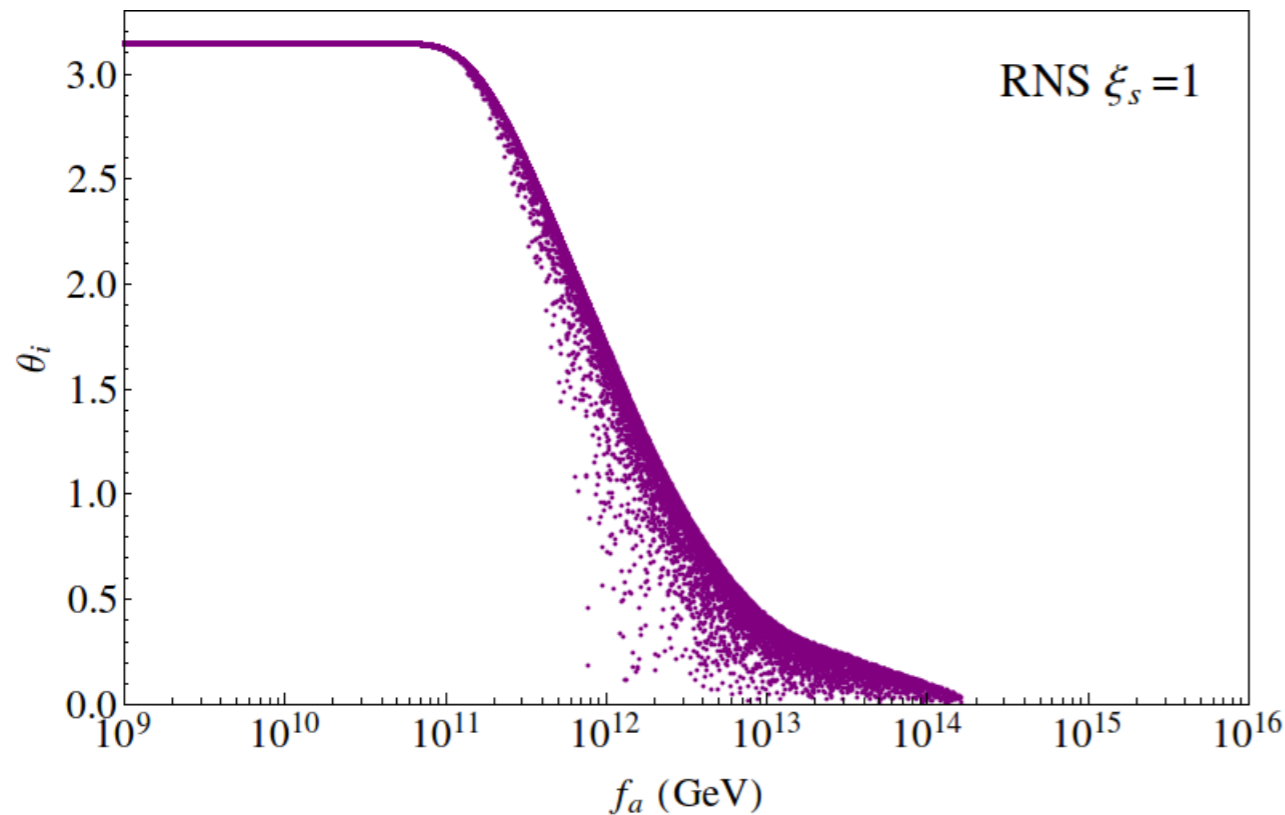
\Rightarrow





higgsino abundance

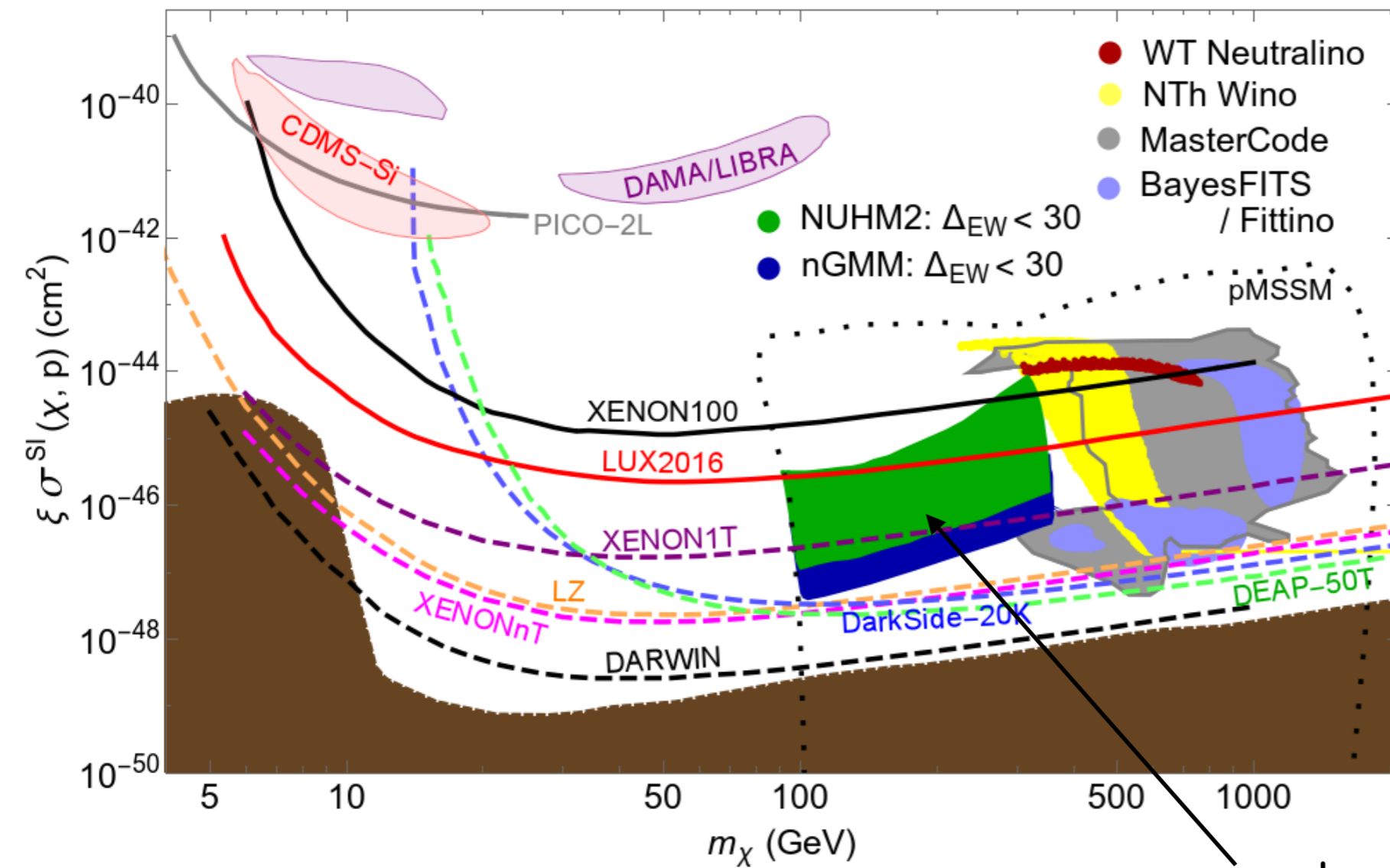
axion abundance



mainly axion CDM
 for $f_a < \sim 10^{12}$ GeV;
 for higher f_a , then
 get increasing wimp
 abundance

Direct higgsino detection rescaled

for minimal local abundance $\xi \equiv \Omega_{\chi}^{TP} h^2 / 0.12$



Bae, HB, Barger, Savoy, Serce

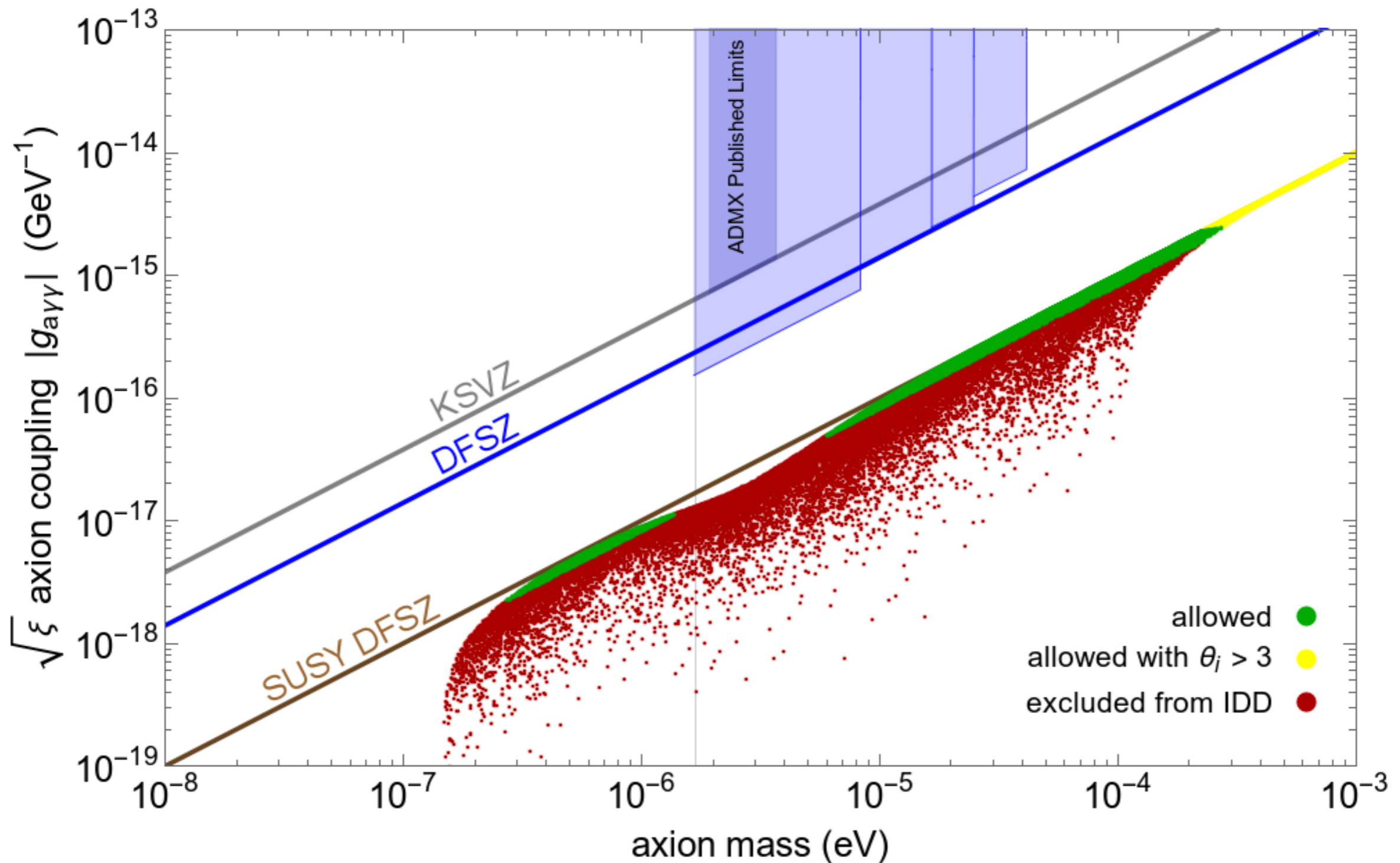
$$\mathcal{L} \ni -X_{11}^h \bar{\tilde{Z}}_1 \tilde{Z}_1 h$$

$$X_{11}^h = -\frac{1}{2} (v_2^{(1)} \sin \alpha - v_1^{(1)} \cos \alpha) (g v_3^{(1)} - g' v_4^{(1)})$$

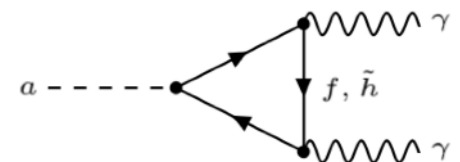
Xe-1-ton
now operating!

natural SUSY

Can test completely with ton scale detector
or equivalent (subject to minor caveats)

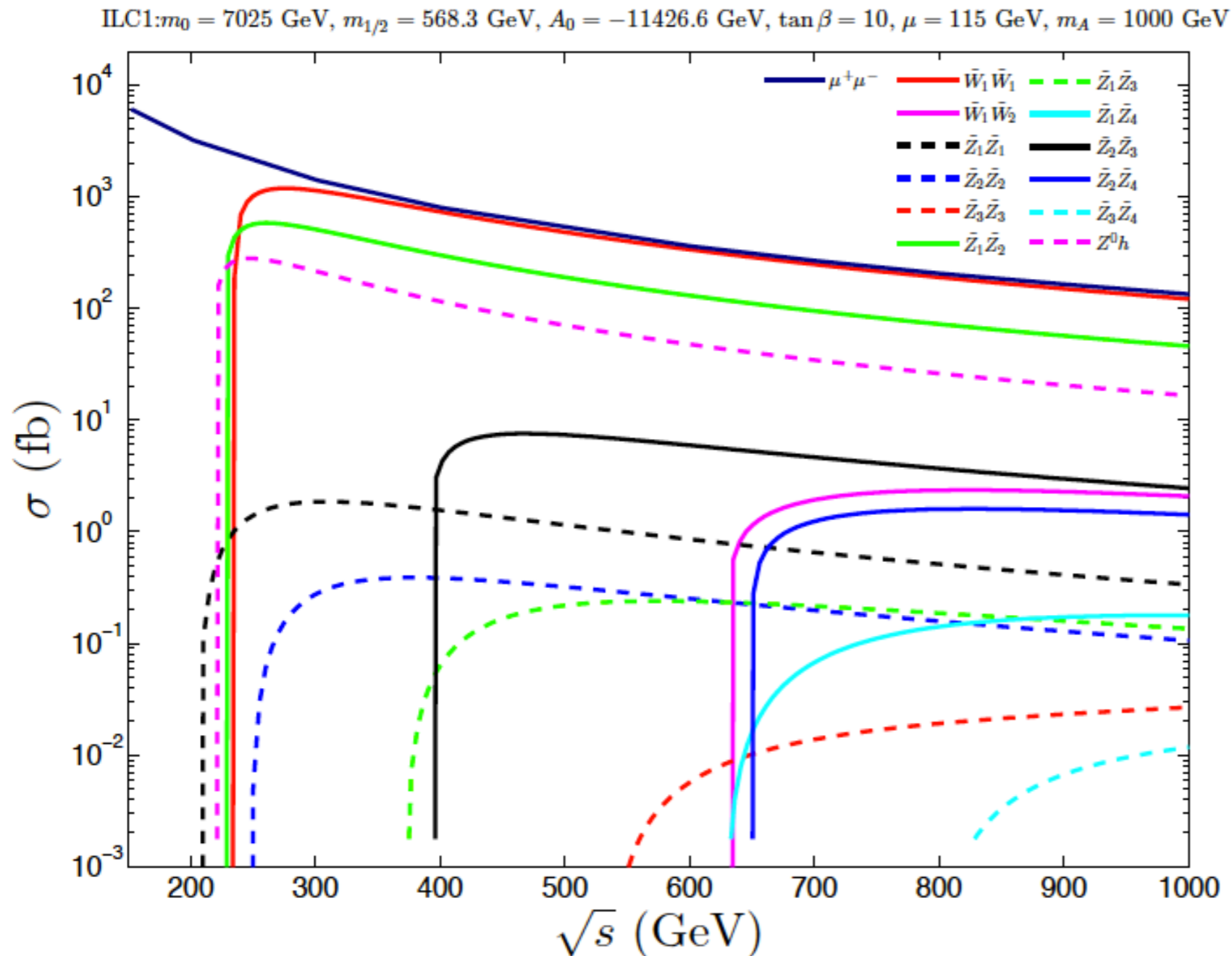


SUSY DFSZ axion: large range in $m(a)$ but coupling reduced
 may need to probe broader and deeper!



Smoking gun signature: light higgsinos at ILC:

ILC is Higgs/higgsino factory!



$$\sigma(\text{higgsino}) \gg \sigma(Zh)$$

10–15 GeV higgsino mass
gaps no problem
in clean ILC environment

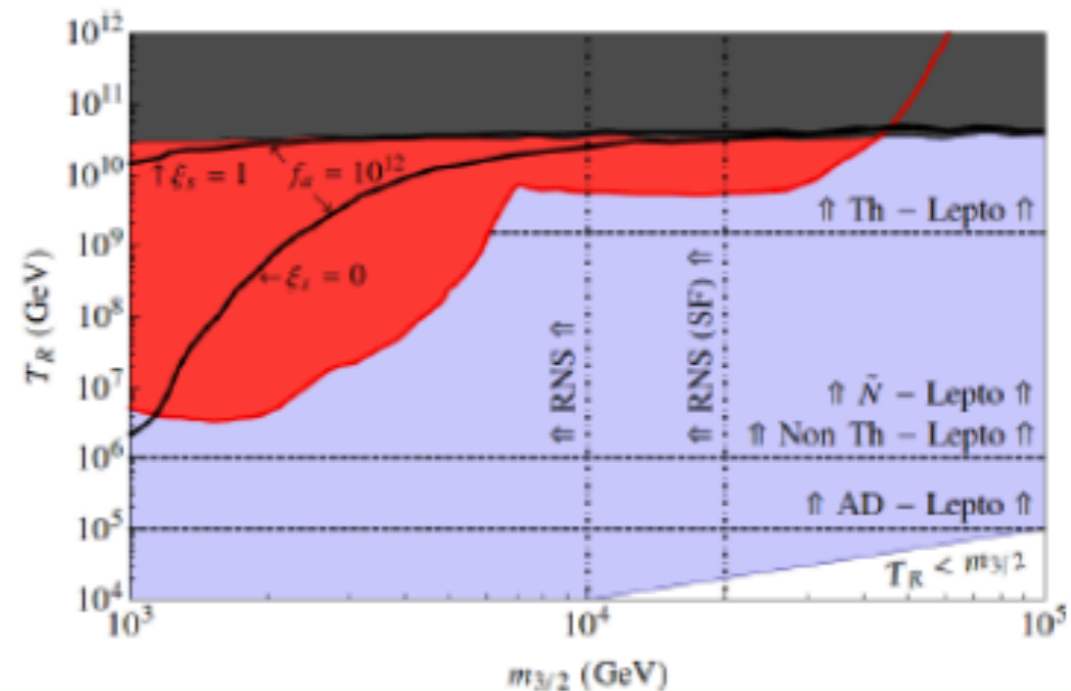
HB, Barger, Mickelson, Mustafayev,
Tata
arXiv:1404:7510

ILC either sees light higgsinos or MSSM dead

Baryogenesis scenarios for radiative natural SUSY

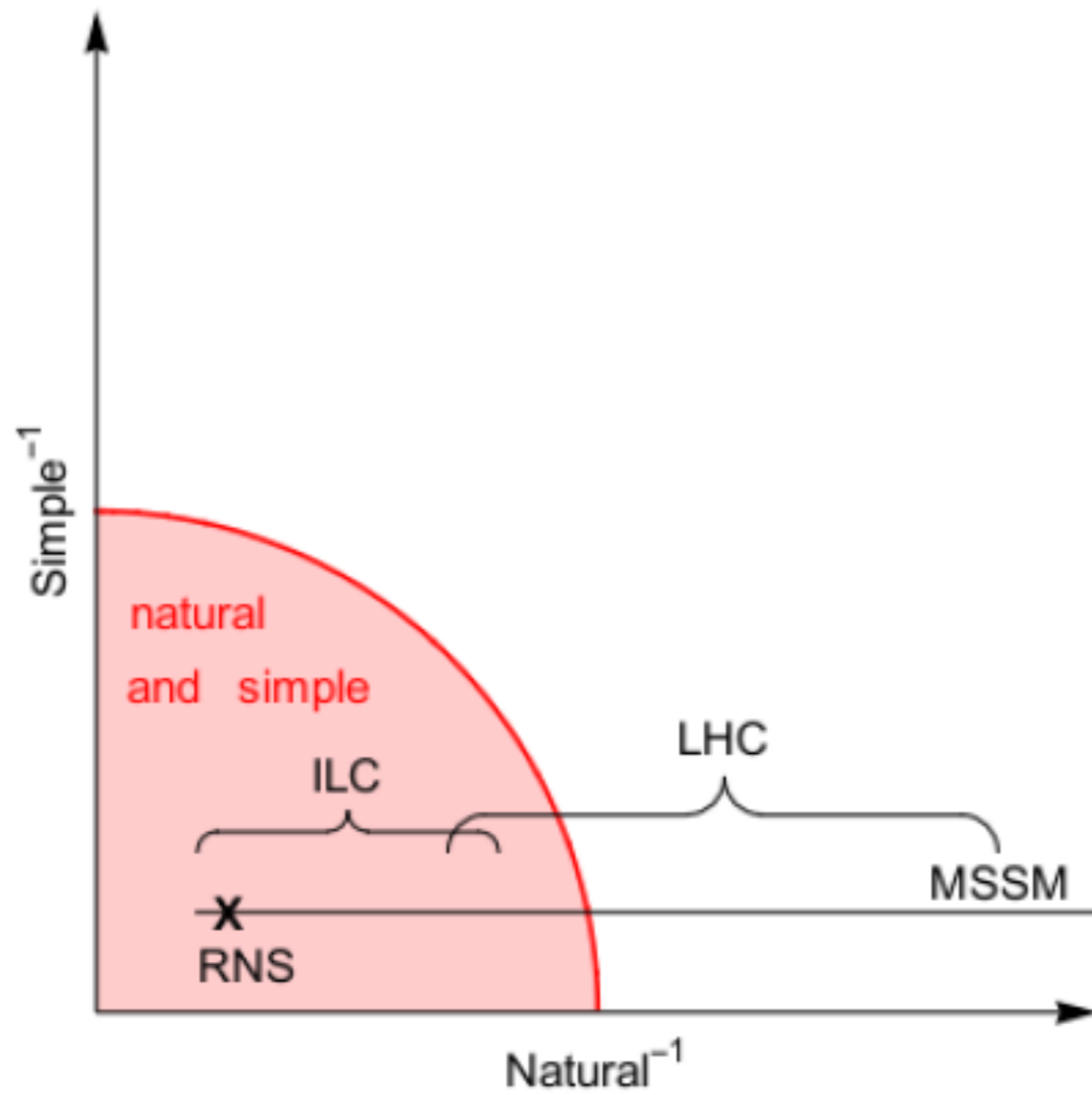
- thermal leptogenesis
- non-thermal (inflaton decay)
- oscillating sneutrino
- Affleck-Dine (AD)

gravitino problem plus
axino/saxion problem:
still plenty room



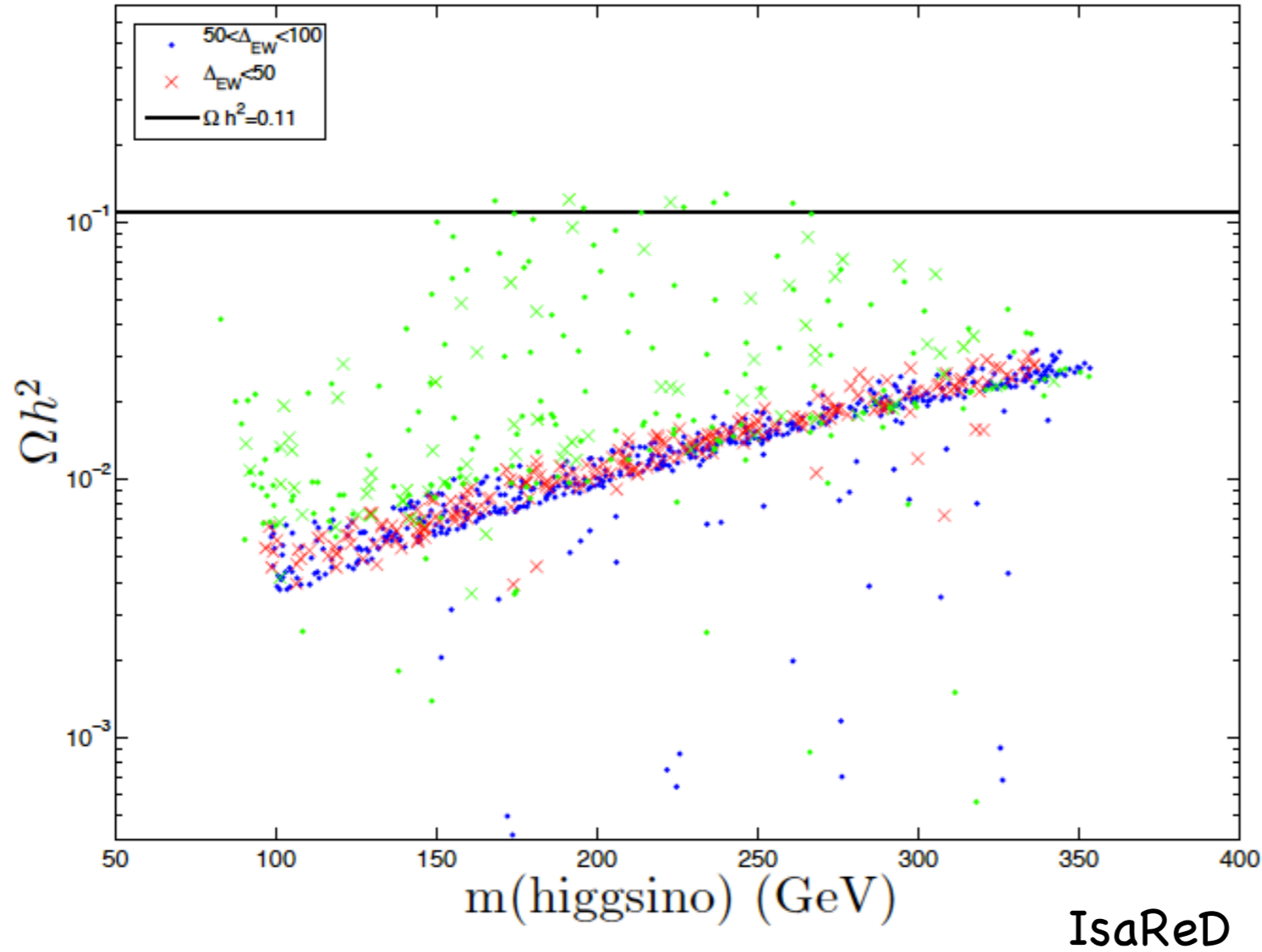
$$f_a = 10^{11}, 10^{12} \text{ GeV}$$

Bae, HB, Serce, Zhang, arXiv:1510.00724



Dark matter in RNS

Mainly higgsino-like WIMPs thermally underproduce DM



green: excluded;
red/blue: allowed

HB, Barger, Mickelson

Factor of 10–15 too low

But so far we have addressed only **Part 1**
of fine-tuning problem:

In QCD sector, the term $\frac{\bar{\theta}}{32\pi^2} F_{A\mu\nu} \tilde{F}_A^{\mu\nu}$ must occur

But neutron EDM says it is not there: strong CP problem

(frequently ignored by SUSY types)

Best solution after 35 years:

PQWW/KSVZ/DFSZ **invisible axion**

In SUSY, axion accompanied by axino and saxion

Changes DM calculus:

expect mixed WIMP/axion DM (**2 particles**)

Axion cosmology

★ Axion field eq'n of motion: $\theta = a(x)/f_a$

$$- \ddot{\theta} + 3H(T)\dot{\theta} + \frac{1}{f_a^2} \frac{\partial V(\theta)}{\partial \theta} = 0$$

$$- V(\theta) = m_a^2(T) f_a^2 (1 - \cos \theta)$$

– Solution for T large, $m_a(T) \sim 0$:

$$\theta = \text{const.}$$

– $m_a(T)$ turn-on ~ 1 GeV

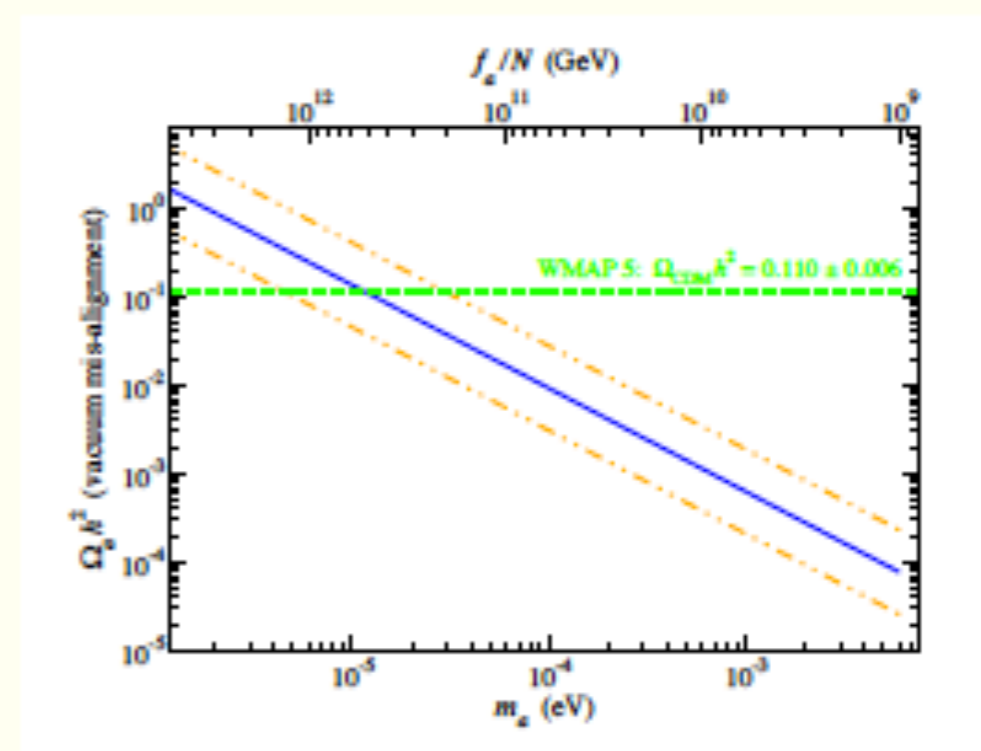
★ $a(x)$ oscillates,

creates axions with $\vec{p} \sim 0$:

production via vacuum mis-alignment

$$\star \Omega_a h^2 \sim \frac{1}{2} \left[\frac{6 \times 10^{-6} \text{ eV}}{m_a} \right]^{7/6} \theta_i^2 h^2$$

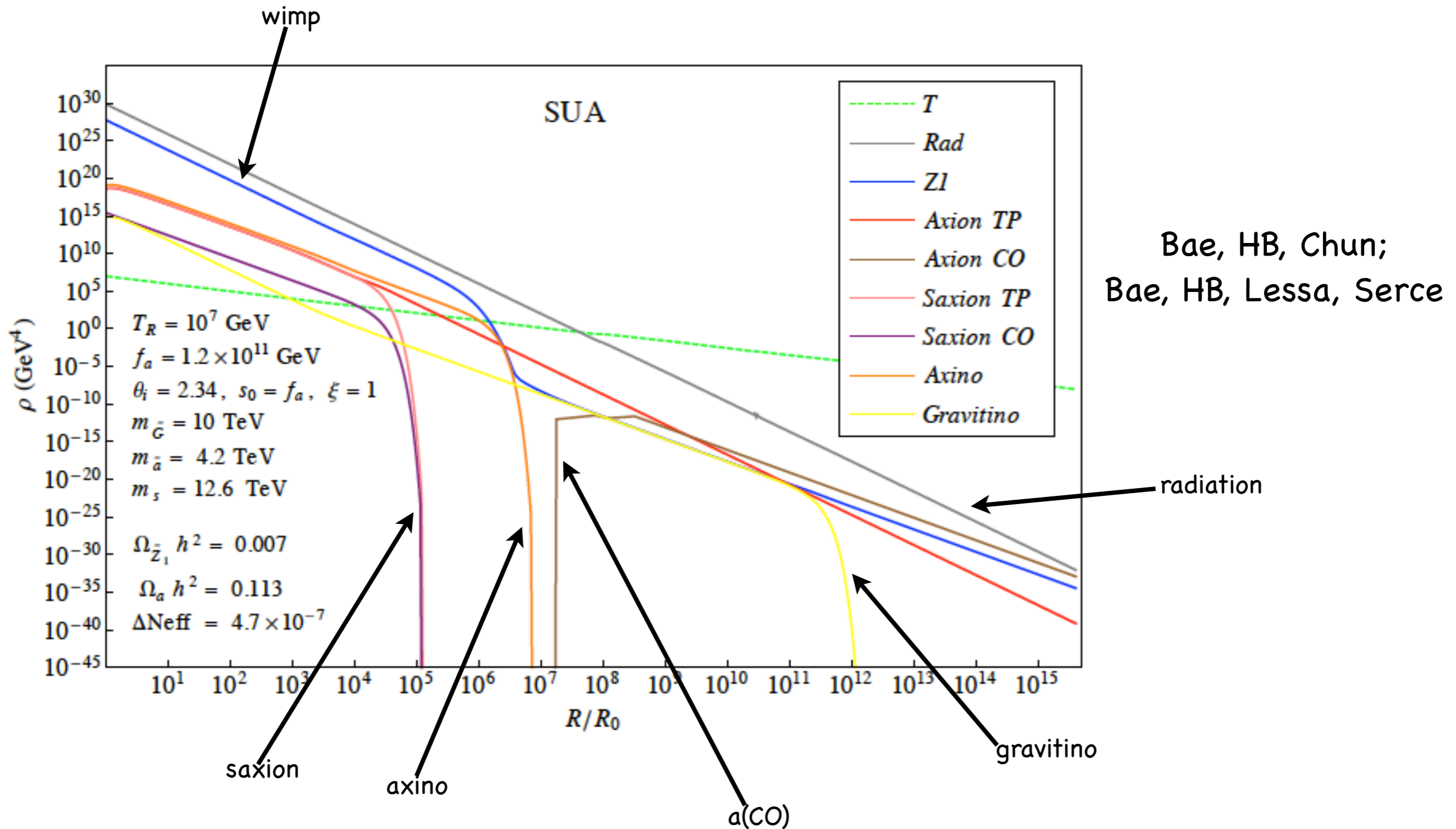
★ astro bound: stellar cooling $\Rightarrow f_a \gtrsim 10^9 \text{ GeV}$

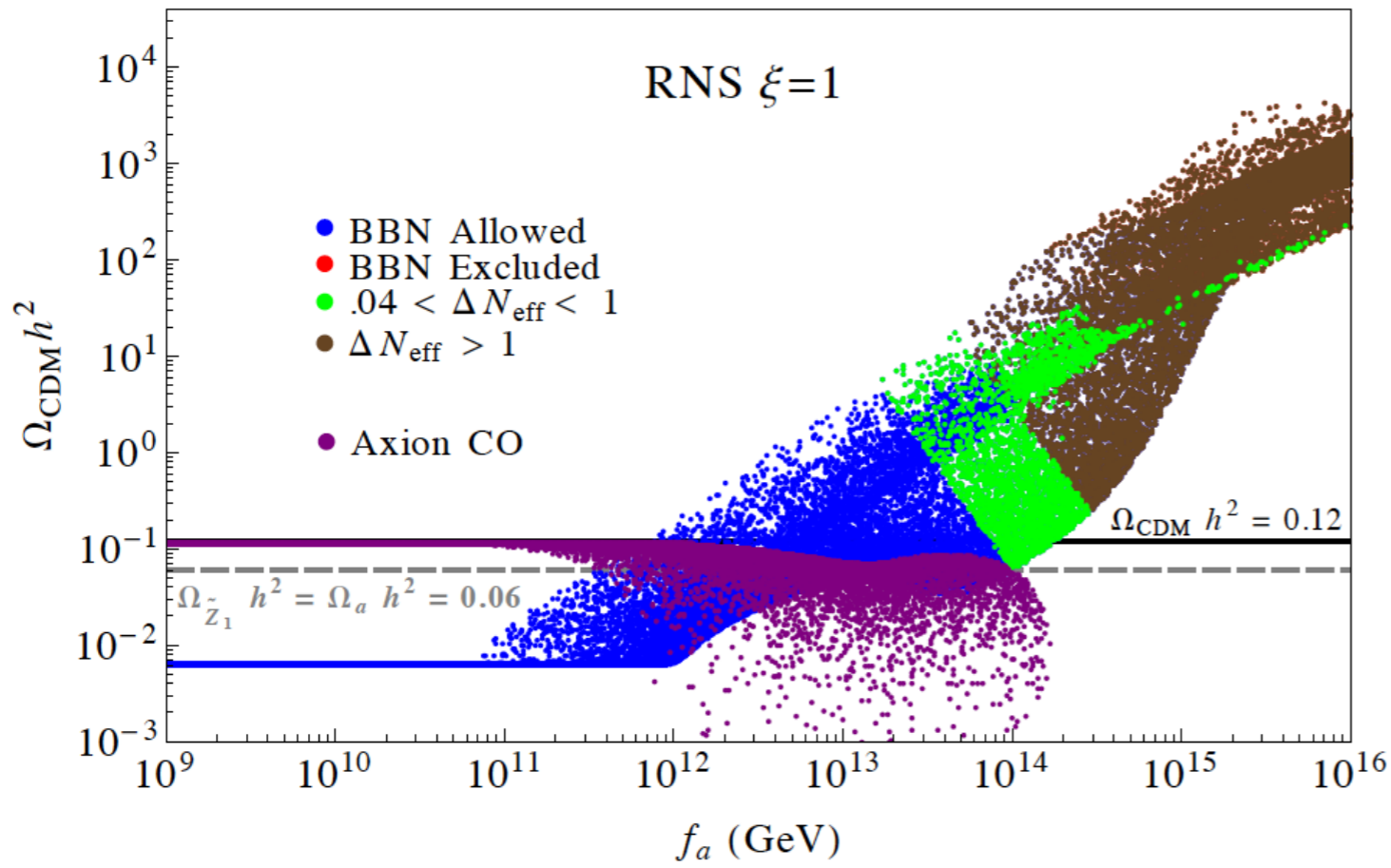


mixed axion-neutralino production in early universe

- neutralinos: thermally produced (TP) or NTP via \tilde{a} , s or \tilde{G} decays
 - re-annihilation at $T_D^{s,\tilde{a}}$
- axions: TP, NTP via $s \rightarrow aa$, bose coherent motion (BCM)
- saxions: TP or via BCM
 - $s \rightarrow gg$: entropy dilution
 - $s \rightarrow SUSY$: augment neutralinos
 - $s \rightarrow aa$: dark radiation ($\Delta N_{eff} < 1.6$)
- axinos: TP
 - $\tilde{a} \rightarrow SUSY$ augments neutralinos
- gravitinos: TP, decay to SUSY

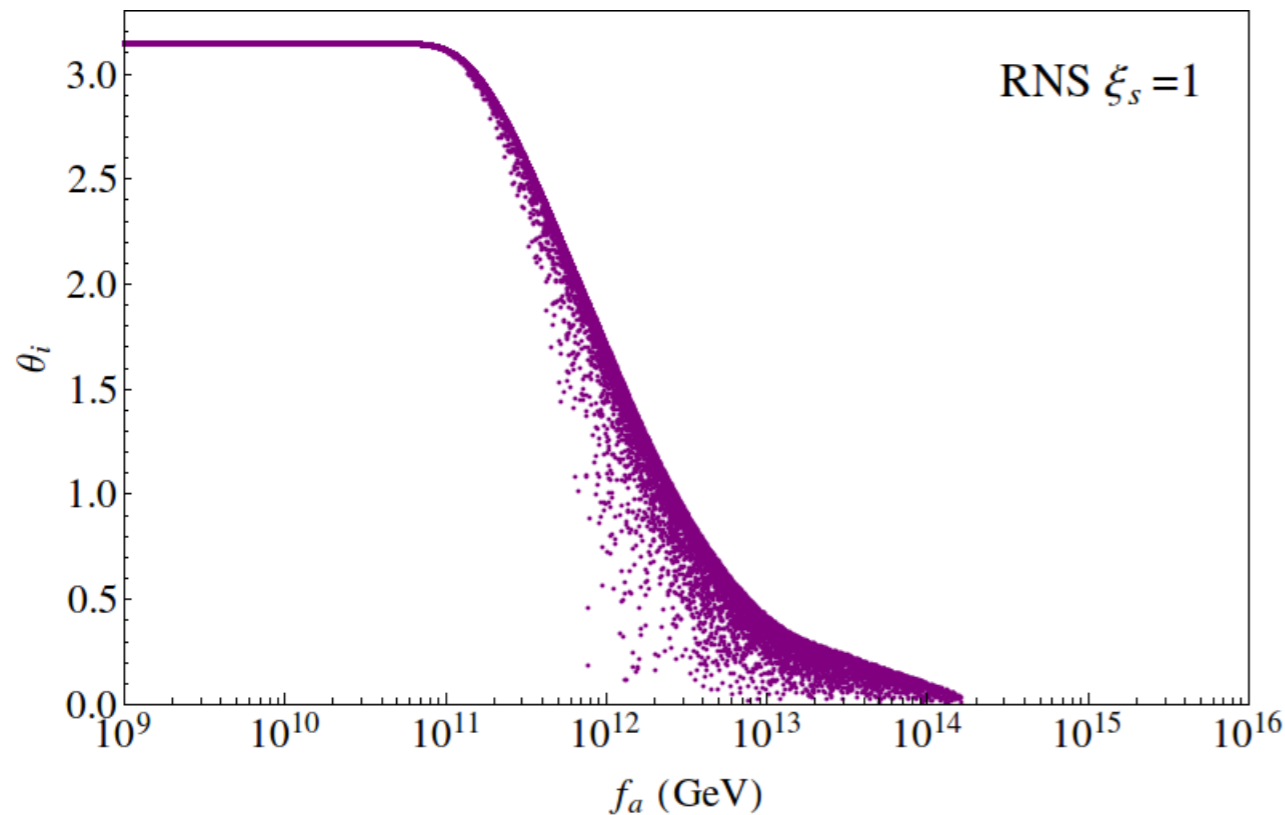
DM production in SUSY DFSZ: solve eight coupled Boltzmann equations





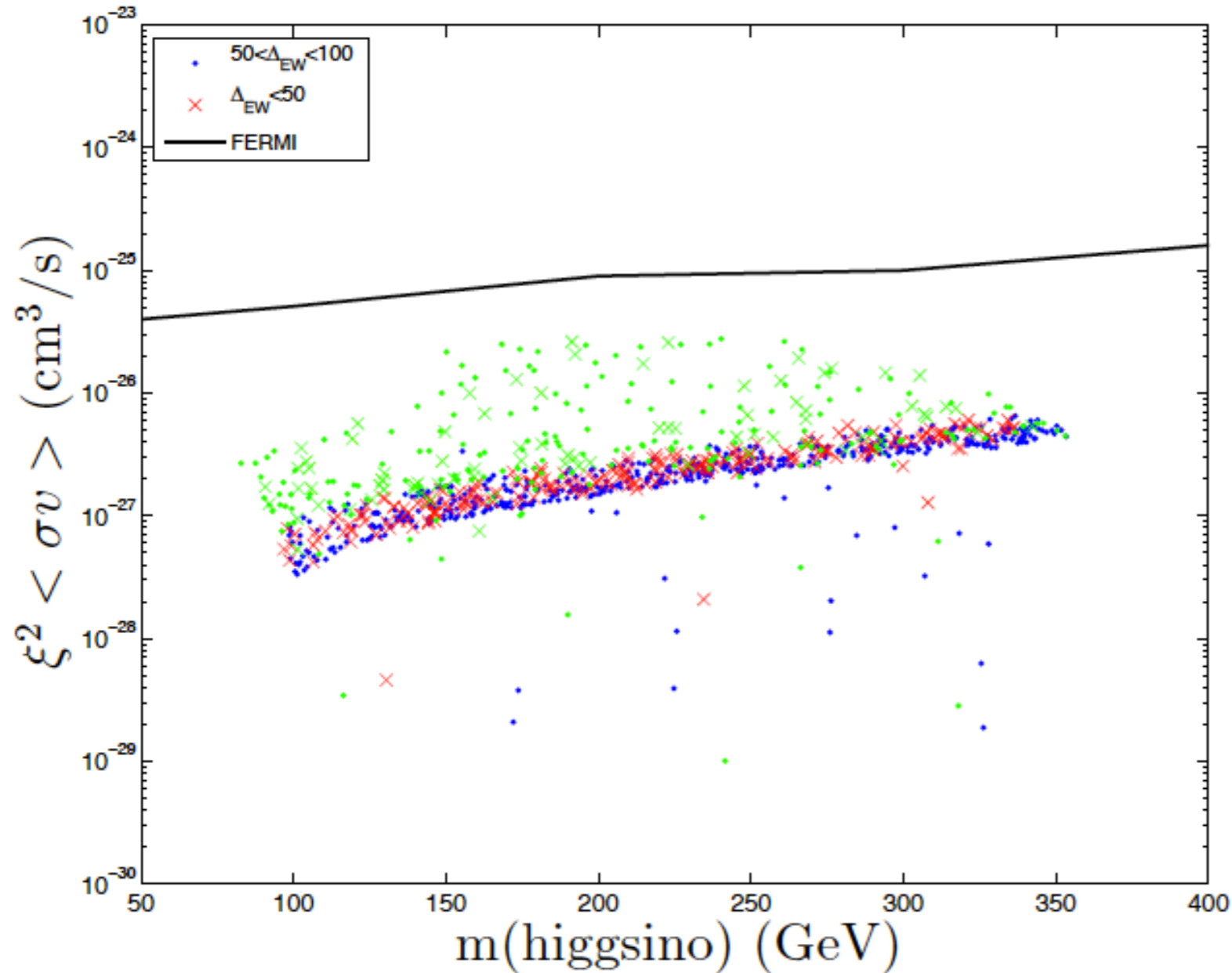
higgsino abundance

axion abundance



mainly axion CDM
 for $f_a < \sim 10^{12}$ GeV;
 for higher f_a , then
 get increasing wimp
 abundance

Higgsino detection via halo annihilations:



green: excluded by Xe-100

annihilation rate is high but rescaling is **squared**

Gamma-ray sky signal is factor 10-20 below current limits

Recommendation: put this horse out to pasture

$$\delta m_{H_u}^2 \sim -\frac{3f_t^2}{8\pi^2} (m_{Q_3}^2 + m_{U_3}^2 + A_t^2) \ln(\Lambda/m_{SUSY})$$

R.I.P.

sub-TeV 3rd generation squarks **not** required for naturalness