

# Establishing common standards for the EFT interpretation of top-quark measurements at the LHC

Gauthier Durieux  
(DESY)

J. A. Aguilar Saavedra, C. Degrande, G. D.,  
F. Maltoni, E. Vryonidou, C. Zhang (editors),

D. Barducci, I. Brivio, V. Cirigliano, W. Dekens, J. de Vries, C. Englert, M. Fabbrichesi,  
R. Franceschini, C. Grojean, U. Haisch, Y. Jiang, J. Kamenik, M. Mangano,  
D. Marzocca, Y. Kats, E. Mereghetti, K. Mimasu, L. Moore, G. Perez, T. Plehn, F. Riva,  
M. Russell, J. Santiago, M. Schulze, Y. Soreq, A. Tonero, M. Trott, S. Westhoff,  
C. White, A. Wulzer, J. Zupan.

with much support and feedback from  
EFT enthusiasts in ATLAS and CMS,  
and LHC TOP WG conveners

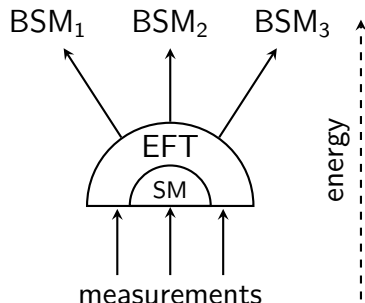


# The standard model effective field theory

systematically parametrizes the theory space  
in direct vicinity of the SM

- ▶ through a proper QFT
- ▶ based on SM fields and symmetries
- ▶ in a low-energy limit
- ▶ systematic when global

*(...) if one writes down the most general possible Lagrangian, including all terms consistent with assumed symmetry principles, (...) the result will simply be the most general possible S-matrix consistent with analyticity, perturbative unitarity, cluster decomposition and the assumed symmetry. [Weinberg '79]*



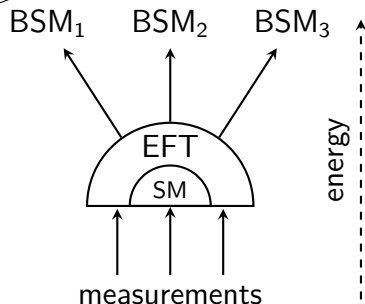
# The standard model effective field theory

systematically parametrizes the theory space  
in direct vicinity of the SM

- ▶ through a proper QFT
- ▶ based on SM fields and symmetries
- ▶ in a low-energy limit
- ▶ systematic when all relevant operators are included

(...) Identify new physics through correlated deviations in precisely measured observables.

(...) the most general possible  $S$ -matrix, including all terms consistent with assumed symmetry principles, (...) the result will simply be the most general possible  $S$ -matrix consistent with analyticity, perturbative unitarity, cluster decomposition and the assumed symmetry. [Weinberg '79]



# Common framework for top physics at the LHC

## First steps

### Delimit an initial scope

- address only processes involving top quarks
- decide which contributions are relevant, in principle
- prioritize the study of flavour structures

### Fix notation

- define d.o.f. natural for top physics at the LHC
- fix notation, normalization, and indicative allowed ranges
- provide simulation tools as TH/EXP interface

### Discuss analysis strategies (one example)

- address the challenges of a global EFT
- highlight useful experimental outputs

Delimit an initial scope

# Relevant operators

Use the *Warsaw* basis of dim-6 operators as reference  
Focus on operators involving a top quark

[Grzadkowski et al '10]

Four-quark operators (11)

$$O_{qq}^{1(ijkl)} = (\bar{q}_i \gamma^\mu q_j)(\bar{q}_k \gamma_\mu q_l),$$

$$O_{qq}^{3(ijkl)} = (\bar{q}_i \gamma^\mu \tau^I q_j)(\bar{q}_k \gamma_\mu \tau^I q_l),$$

$$O_{qu}^{1(ijkl)} = (\bar{q}_i \gamma^\mu q_j)(\bar{u}_k \gamma_\mu u_l),$$

$$O_{qu}^{8(ijkl)} = (\bar{q}_i \gamma^\mu T^A q_j)(\bar{u}_k \gamma_\mu T^A u_l),$$

$$O_{qd}^{1(ijkl)} = (\bar{q}_i \gamma^\mu q_j)(\bar{d}_k \gamma_\mu d_l),$$

$$O_{qd}^{8(ijkl)} = (\bar{q}_i \gamma^\mu T^A q_j)(\bar{d}_k \gamma_\mu T^A d_l),$$

$$O_{uu}^{(ijkl)} = (\bar{u}_i \gamma^\mu u_j)(\bar{u}_k \gamma_\mu u_l),$$

$$O_{ud}^{1(ijkl)} = (\bar{u}_i \gamma^\mu u_j)(\bar{d}_k \gamma_\mu d_l),$$

$$O_{ud}^{8(ijkl)} = (\bar{u}_i \gamma^\mu T^A u_j)(\bar{d}_k \gamma_\mu T^A d_l),$$

$$O_{quqd}^{1(ijkl)} = (\bar{q}_i u_j) \in (\bar{q}_k d_l),$$

$$O_{quqd}^{8(ijkl)} = (\bar{q}_i T^A u_j) \in (\bar{q}_k T^A d_l),$$

Two-quark operators (9)

$$O_{u\varphi}^{(ij)} = \bar{q}_i u_j \tilde{\varphi} (\varphi^\dagger \varphi),$$

$$O_{\varphi q}^{1(ij)} = (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_i \gamma^\mu q_j),$$

$$O_{\varphi q}^{3(ij)} = (\varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_i \gamma^\mu \tau^I q_j),$$

$$O_{\varphi u}^{(ij)} = (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{u}_i \gamma^\mu u_j),$$

$$O_{\varphi ud}^{(ij)} = (\tilde{\varphi}^\dagger i D_\mu \varphi)(\bar{u}_i \gamma^\mu d_j),$$

$$O_{uW}^{(ij)} = (\bar{q}_i \sigma^{\mu\nu} \tau^I u_j) \tilde{\varphi} W_{\mu\nu}^I,$$

$$O_{dW}^{(ij)} = (\bar{q}_i \sigma^{\mu\nu} \tau^I d_j) \varphi W_{\mu\nu}^I,$$

$$O_{uB}^{(ij)} = (\bar{q}_i \sigma^{\mu\nu} u_j) \tilde{\varphi} B_{\mu\nu},$$

$$O_{uG}^{(ij)} = (\bar{q}_i \sigma^{\mu\nu} T^A u_j) \tilde{\varphi} G_{\mu\nu}^A,$$

Two-quark-two-lepton operators (8)

$$O_{lq}^{1(ijkl)} = (\bar{l}_j \gamma^\mu l_j)(\bar{q}_k \gamma_\mu q_l),$$

$$O_{lq}^{3(ijkl)} = (\bar{l}_j \gamma^\mu \tau^I l_j)(\bar{q}_k \gamma_\mu \tau^I q_l),$$

$$O_{lu}^{(ijkl)} = (\bar{l}_j \gamma^\mu l_j)(\bar{u}_k \gamma_\mu u_l),$$

$$O_{eq}^{(ijkl)} = (\bar{e}_j \gamma^\mu e_j)(\bar{q}_k \gamma_\mu q_l),$$

$$O_{eu}^{(ijkl)} = (\bar{e}_j \gamma^\mu e_j)(\bar{u}_k \gamma_\mu u_l),$$

$$O_{lequ}^{1(ijkl)} = (\bar{l}_i e_j) \in (\bar{q}_k u_l),$$

$$O_{lequ}^{3(ijkl)} = (\bar{l}_i \sigma^{\mu\nu} e_j) \in (\bar{q}_k \sigma_{\mu\nu} u_l),$$

$$O_{ledq}^{(ijkl)} = (\bar{l}_i e_j)(\bar{d}_k q_l),$$

+  $\mathcal{B}$  and  $\mathcal{L}$  operators (4 or 5)

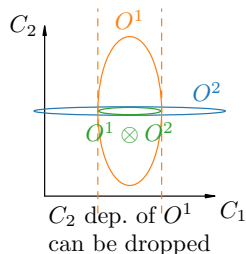
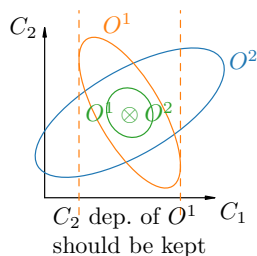
Others should be sufficiently constrained by processes involving no tops.

# Relevant contributions

Use present sensitivities and constraints rather than theoretical prejudices to decide which contributions are relevant.

1. work on an obs.-by-obs. basis ( $O^k$ )
2. evaluate all tree-level contributions
3. discard dependences when irrelevant
4. compute higher orders in SM couplings where necessary

**Note** the relevance of d.o.f.'s in a measurement may change as constraints are collected!



# Flavour assumptions

(FCNCs treated separately)

To prioritize the study of flavour structures

## Lepton sector (not critical)

- rather loose  $[U(1)_{l+e}]^3$  aka flavour diagonality
- could easily be restricted to  $U(3)_{l+e}$ ,  $U(3)_l \times U(3)_e$ , or ...

## Quark sector (baseline and variants)

mostly restrict the large number of four-quark operators

**Baseline**  $U(2)_q \times U(2)_u \times U(2)_d$  among first two generations

$\equiv$  SM flavour symmetry in the limit  $y_{u,d,s,c} \rightarrow 0$ ,  $V_{\text{CKM}} \rightarrow \mathbb{I}$

forces the first two generations to appear as  $\sum_{i=1,2} \bar{q}_i q_i$ ,  $\bar{u}_i u_i$ ,  $\bar{d}_i d_i$

**Extended** to  $U(2)_{q+u+d}$

[sugg. by J.A.Aguilar Saavedra]

- allows light right-handed charged currents  $\sum_{i=1,2} \bar{u}_i d_i$
- allows light chirality flipping currents  $\sum_{i=1,2} \bar{q}_i u_i$ ,  $\bar{q}_i d_i$

**Restricted** to *top-philic* scenario

[sugg. by A.Wulzer]

- assumes NP generates all operators with tops and bosons
- then project that over-complete set on the Warsaw basis with EOM, etc.



Fix notation

# Top-specific d.o.f. definitions

Match interference structures  
and interactions with physical gauge bosons

$$\bullet \begin{pmatrix} O_{\varphi q}^{1(33)} \\ O_{\varphi q}^{3(33)} \end{pmatrix} = \begin{pmatrix} (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_3 \gamma^\mu q_3) \\ (\varphi^\dagger \overleftrightarrow{D}'_\mu \varphi)(\bar{q}_3 \gamma^\mu \tau^I q_3) \end{pmatrix} = \begin{pmatrix} -1 & 1 \\ 1 & 1 \\ 0 & 1 \\ 0 & 1 \end{pmatrix}^T \begin{pmatrix} \frac{+e}{2s_W c_W} (\bar{t} \gamma^\mu P_L t) Z_\mu (v+h)^2 \\ \frac{-e}{2s_W c_W} (\bar{b} \gamma^\mu P_L b) Z_\mu (v+h)^2 \\ \frac{g}{\sqrt{2}} (\bar{t} \gamma^\mu P_L b) W_\mu^+ (v+h)^2 \\ \frac{g}{\sqrt{2}} (\bar{b} \gamma^\mu P_L t) W_\mu^- (v+h)^2 \end{pmatrix}$$

$$c_{\varphi Q}^- \equiv C_{\varphi q}^{1(33)} - C_{\varphi q}^{3(33)} \quad \text{enters in } pp \rightarrow t\bar{t}Z$$

$$c_{\varphi Q}^3 \equiv C_{\varphi q}^{3(33)} \quad \text{enters in } t \rightarrow bW^+$$

$$c_{\varphi Q}^+ \equiv C_{\varphi q}^{1(33)} + C_{\varphi q}^{3(33)} \quad \text{enters in } e^+e^- \rightarrow b\bar{b} \text{ (or } pp \rightarrow b\bar{b}Z)$$

$$\bullet \begin{pmatrix} O_{qq}^{1(ii33)} \\ O_{qq}^{1(i33i)} \\ O_{qq}^{3(ii33)} \\ O_{qq}^{3(i33i)} \end{pmatrix} = \begin{pmatrix} 1 & 1/6 & 0 & 1/2 \\ 0 & 1/6 & 1 & -1/6 \\ 0 & 1 & 0 & 3 \\ 0 & 1 & 0 & -1 \end{pmatrix}^T \begin{pmatrix} (\bar{q}_i \gamma^\mu q_i) (\bar{Q} \gamma_\mu Q) \\ (\bar{q}_i \gamma^\mu \tau^I q_i) (\bar{Q} \gamma_\mu \tau^I Q) \\ (\bar{q}_i \gamma^\mu T^A q_i) (\bar{Q} \gamma_\mu T^A Q) \\ (\bar{q}_i \gamma^\mu \tau^I T^A q_i) (\bar{Q} \gamma_\mu \tau^I T^A Q) \end{pmatrix}$$

$$c_{Qq}^{1,1} \equiv C_{qq}^{1(ii33)} + \frac{1}{6} C_{qq}^{1(i33i)} + \frac{1}{2} C_{qq}^{3(i33i)} \quad \text{interferes with EW NC}$$

$$c_{Qq}^{3,1} \equiv C_{qq}^{3(ii33)} + \frac{1}{6} (C_{qq}^{1(i33i)} - C_{qq}^{3(i33i)}) \quad \text{interferes with EW CC}$$

$$c_{Qq}^{1,8} \equiv C_{qq}^{1(i33i)} + 3 C_{qq}^{3(i33i)}, \quad \text{interferes with QCD}$$

$$c_{Qq}^{3,8} \equiv C_{qq}^{1(i33i)} - C_{qq}^{3(i33i)} \quad \text{doesn't interfere with EW CC}$$

# Counting and constraints

	benchmark	extended	restricted
four heavy quarks	11 + 2 CPV		5
two light and two heavy quarks	14	+10 + 10 CPV	} 5
two heavy quarks and two leptons	(8 + 3 CPV) × 3		
two heavy quarks and bosons	9 + 6 CPV		
			9 + 6 CPV

## Indicative direct constraints:

[many from TopFitter]

Four-heavy (11 + 2 CPV d.o.f.)	Indicative direct limits
$c_{QQ}^1 \equiv 2C_{qq}^{1(3333)} - \frac{2}{3}C_{qq}^{3(3333)}$	
$c_{QQ}^8 \equiv 8C_{qq}^{3(3333)}$	
$!c_{QQ}^+ \equiv C_{qq}^{1(3333)} + C_{qq}^{3(3333)}$	$[-2.92, 2.80] (E_{cut} = 3 \text{ TeV})$ [35]
$c_{Qt}^1 \equiv C_{qu}^{1(3333)}$	$[-4.97, 4.90] (E_{cut} = 3 \text{ TeV})$ [35]
$c_{Qt}^8 \equiv C_{qu}^{8(3333)}$	$[-10.3, 9.33] (E_{cut} = 3 \text{ TeV})$ [35]
$c_{Qb}^1 \equiv C_{qd}^{1(3333)}$	
$c_{Qb}^8 \equiv C_{qd}^{8(3333)}$	
$c_{tt}^1 \equiv C_{uu}^{1(3333)}$	$[-2.92, 2.80] (E_{cut} = 3 \text{ TeV})$ [35]
$c_{tb}^1 \equiv C_{ud}^{1(3333)}$	
$c_{tb}^8 \equiv C_{ud}^{8(3333)}$	
$c_{QtQb}^{1[I]} \equiv \frac{[\text{Im}]}{\text{Re}}\{C_{quqd}^{1(3333)}\}$	
$c_{QtQb}^{8[I]} \equiv \frac{[\text{Im}]}{\text{Re}}\{C_{quqd}^{8(3333)}\}$	
Two-light-two-heavy (14 d.o.f.)	
$c_{Qq}^{3,1} \equiv C_{qq}^{3(i333)} + \frac{1}{6}(C_{qq}^{1(i334)} - C_{qq}^{3(i334)})$	$[-0.66, 1.24]$ [36], $[-3.11, 3.10]$ [35]
$c_{Qq}^{3,8} \equiv C_{qq}^{1(i334)} - C_{qq}^{3(i334)}$	$[-6.06, 6.73]$ [35]
$c_{Qq}^{1,1} \equiv C_{qq}^{1(i333)} + \frac{1}{6}C_{qq}^{1(i334)} + \frac{1}{2}C_{qq}^{3(i334)}$	$[-3.13, 3.15]$ [35]
$c_{Qq}^{1,8} \equiv C_{qq}^{1(i334)} + 3C_{qq}^{3(i334)}$	$[-6.92, 4.93]$ [35]

[ $\Lambda = 1 \text{ TeV}$ ]

# Tree-level UFO implementations

## As TH/EXP interfaces

- ▶ dedicated dim6top
  - $O(90)$  d.o.f.'s of the extended flavour scenario
- ▶ SMEFTsim alternative
  - implementing Warsaw operators
  - providing restriction cards for emulating d.o.f.'s

[Brivio, Jiang, Trott '17]

## Benchmark dependences

(cross checked among the two models)

e.g. linear contributions to total rates:

permil of the SM rate,  $\Lambda = 1 \text{ TeV}$

SM		$pp \rightarrow t\bar{t}$ 5.2 × 10 <sup>2</sup> pb	$pp \rightarrow t\bar{t}b\bar{b}$ 2.3 pb	$pp \rightarrow t\bar{t}t\bar{t}$ 0.0099 pb	$pp \rightarrow t\bar{t}e^+\nu$ 0.02 pb	$pp \rightarrow t\bar{t}e^+e^-$ 0.016 pb	$pp \rightarrow t\bar{t}\gamma$ 1.5 pb	$pp \rightarrow t\bar{t}h$ 0.4 pb
$c_{QQ}^1$	cQQ1	-0.25	-1.5	-1 × 10 <sup>2</sup>		-1.6	-0.66	-0.71
$c_{QQ}^8$	cQQ8	-0.16	-2.5	-32		-0.91	-0.49	-0.28
$c_{Qt}^1$	cQt1	-0.15	-4.3	1 × 10 <sup>2</sup>		-0.77	-0.19	-0.56
$c_{Qt}^8$	cQt8	-0.053	-1.5	-39		-0.18	-0.094	-0.15
$c_{Qb}^1$	cQb1	-0.0055	0.53	-0.051		-0.014	-0.0069	-0.029
$c_{Qb}^8$	cQb8	0.14	3.2	0.12		0.35	0.16	0.57
$c_{tt}^1$	ctt1			-1.6 × 10 <sup>2</sup>				
$c_{tb}^1$	ctb1	-0.0096	0.36	-0.056		-0.02	-0.023	-0.04
$c_{tb}^8$	ctb8	0.14	2.9	0.11		0.26	0.3	0.58
$c_{Qq}^{3,8}$	cQq83	2.6	2	5	-84	-19	10	16
$c_{Qq}^{1,8}$	cQq81	12	20	24	2.6 × 10 <sup>2</sup>	73	36	73
$c_{tq}^8$	ctq8	12	21	27	2.6 × 10 <sup>2</sup>	63	54	73
$c_{Qu}^8$	cQu8	7.2	12	18		21	42	44
$c_{tu}^8$	ctu8	7.5	11	15		14	23	44
$c_{Qd}^8$	cQd8	5	8.3	11		17	6.8	28
$c_{td}^8$	ctd8	4.8	7.2	10		12	14	28
$c_{Qq}^{3,1}$	cQq13	3.3	5.3	5.1	1.1 × 10 <sup>2</sup>	22	11	19
$c_{tq}^{1,1}$	ctq11	0.02	0.10	7.0	6.1	4.8	0.0	0.0



# Discuss analysis strategies

**Warning:** illustrative theorist view!

- to show how the challenges of a global EFT could be addressed
- to fix ideas on what are useful outputs from a TH perspective

# An example of EFT analysis strategy

Choose a (particle-level) fiducial volume close enough to the detector level for unfolding to be very model independent.

to be checked!

—→ facilitates re-interpretations

- in an evolving global EFT picture
- with more sophisticated predictions
- with less restrictive assumptions (about flavour, non-top operators, etc.)
- outside experimental collaboration

—→ facilitates multidimensional EFT analyses

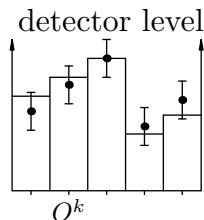
—→ **but** may sometimes be impractical or suboptimal

# An example of EFT analysis strategy

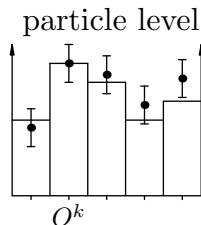
For  $O^k$  observables

total rate, binned  $p_T$ ,  $\eta$ ,  $m_{xy}$ , etc. distributions,  
binned MVA output, ratios, asymmetries, *optimal observables*,...

Unfold



unfold  
the data  
 $\Rightarrow$   
under SM  
hypothesis



Provide

- observable definitions (code if non-standard)
- statistical uncertainties
- systematics breakdown and correlations

( $\rightarrow$  re-interpretable in any model)



## Global EFT interpretation

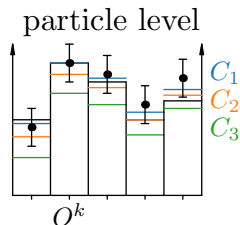
- Compute EFT predictions to the particle level

$$O^k = B_I^k + \frac{C_i}{\Lambda^2} S_i^k + \frac{C_i C_j}{\Lambda^4} S_{ij}^k + \dots$$

SM bkg composition      linear dim-6 contributions (EFT-SM interf.)      quadratic dim-6 contributions      higher powers, and higher-dim. operators

- Obtain and release likelihoods in the full  $\{C_i\}$  space

≡ **global** constraints  
to combine with other measurements



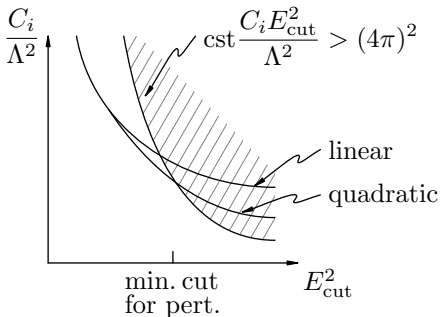
- also quote **individual** constraints
  - information about sensitivity and the magnitude of approximate degeneracies

- quote both the **linear** and **quadratic** dim-6 approx.
  - information about the importance of higher powers of dim-6 coeff.

- quote limits as functions of  $E_{\text{cut}}$  on a characteristic energy scale

- valid interpretation for models with lower scales, with  $[\text{dim}>6] > [\text{dim}-6]$  without  $E_{\text{cut}}$
- perturbativity possibly ensured by minimal  $E_{\text{cut}}$

[Contino et al '16]



$$\sum_n \left[ \text{diagram of } n \text{ loops} \right] \sim \sum_n \left[ \text{cst} \frac{C_i E_{\text{cut}}^2}{(4\pi)^2 \Lambda^2} \right]^n$$

# Summary

# Left aside

Higher-order corrections

especially in QCD, based on existing results

Theory uncertainties

especially the intrinsic EFT ones

EFT treatment of unstable tops

Process and observable sensitivity studies

- largest strength in specific directions
- complementarity to cover the whole EFT param. space

...

# Covered

## Delimit an initial scope

- address only processes involving top quarks
- decide which contributions are relevant, in principle
- prioritize the study of flavour structures

## Fix notation

- define d.o.f. natural for top physics at the LHC
- fix notation, normalization, and indicative allowed ranges
- provide simulation tools as TH/EXP interface

## Discuss analysis strategies (one example)

- address the challenges of a global EFT
- highlight useful experimental outputs

A wide agreement was reached among theorists.  
A summary note will be released in days.