

NLO effects on stop searches in the compressed region.

by

[Marco Hufnagel](#)

in collaboration with

Jan Heisig, Benjamin Nachman

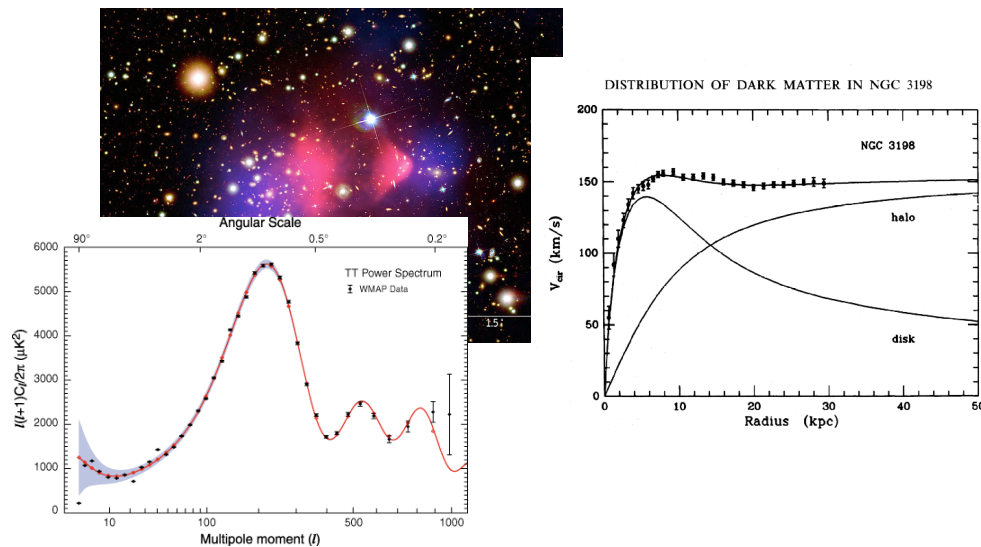


Monday, September 24, 2018

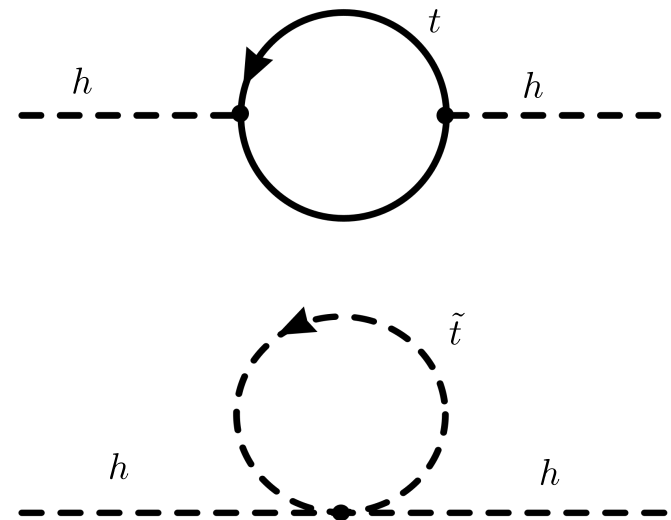
Introduction

Shortcomings of the Standard Model:

Dark Matter:



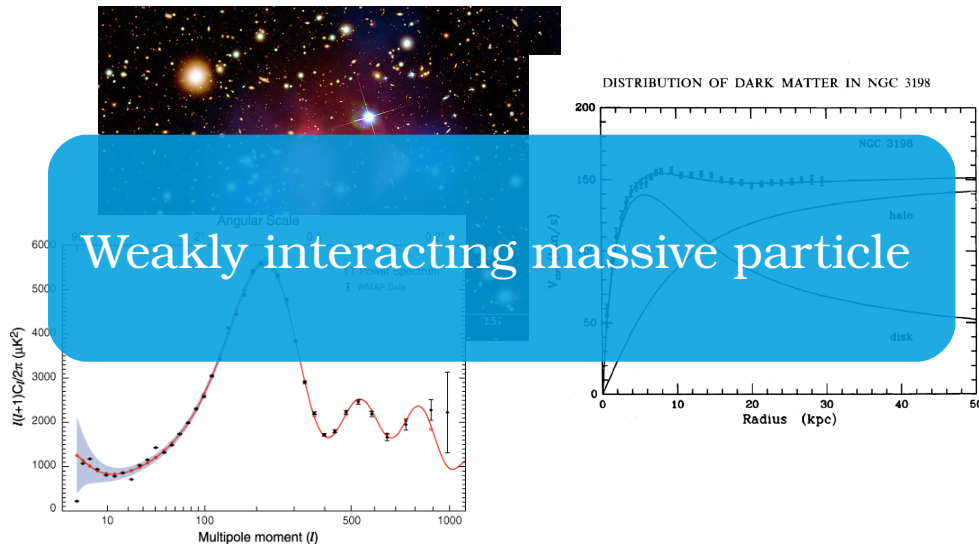
Hierarchy Inconvenience:



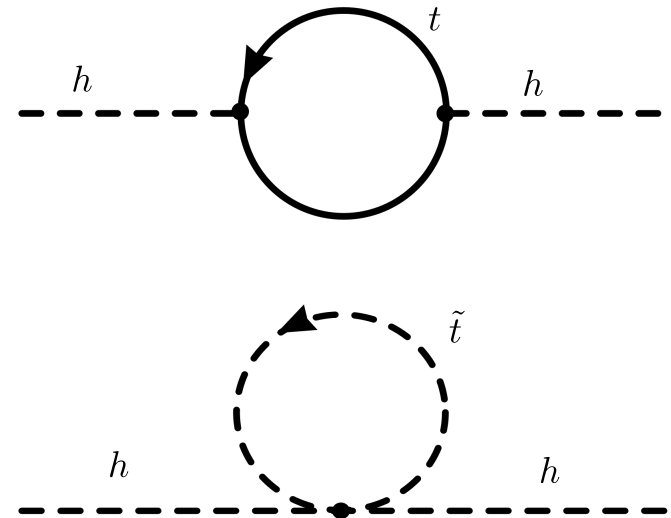
Introduction

Shortcomings of the Standard Model:

Dark Matter:



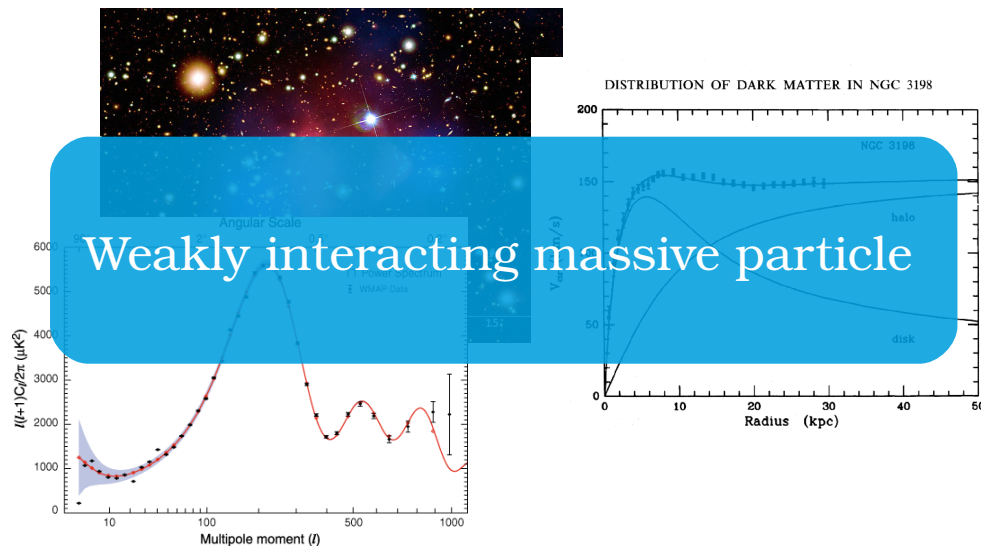
Hierarchy Inconvenience:



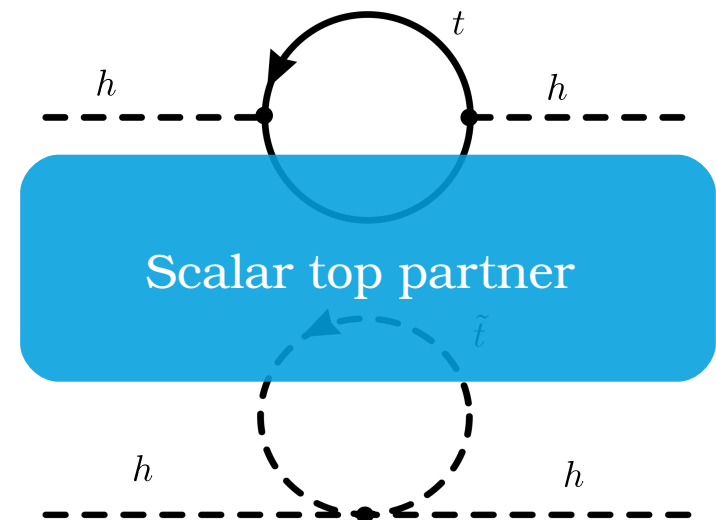
Introduction

Shortcomings of the Standard Model:

Dark Matter:



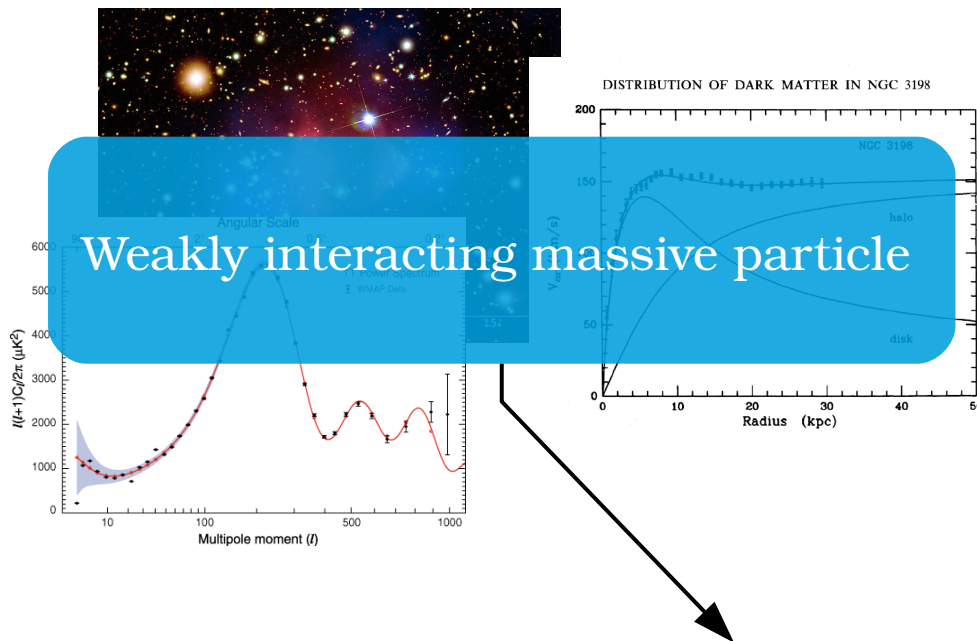
Hierarchy Inconvenience:



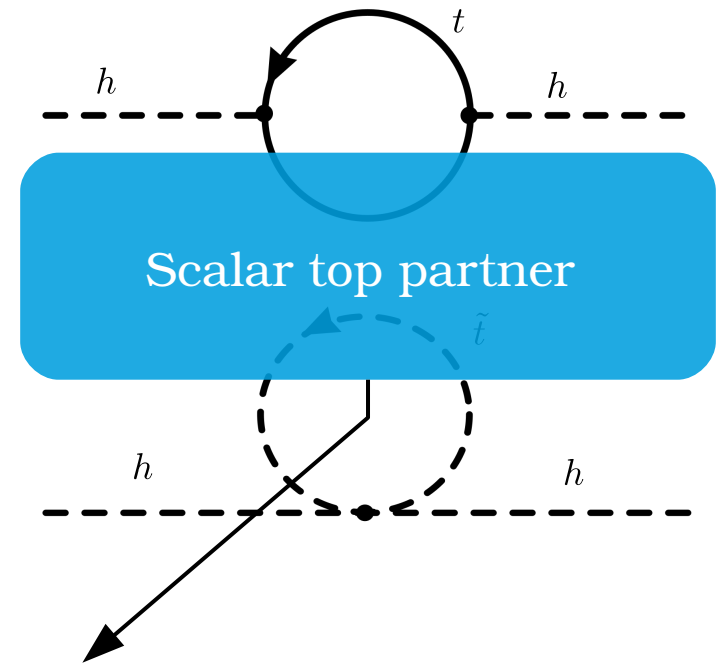
Introduction

Shortcomings of the Standard Model:

Dark Matter:



Hierarchy Inconvenience:

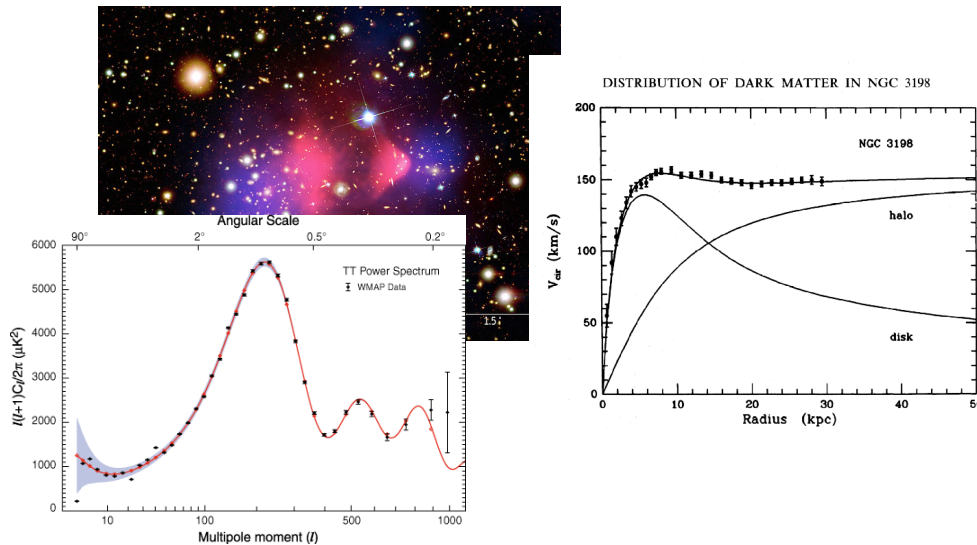


Possible Solution: Lightest neutralino & top partner from the MSSM

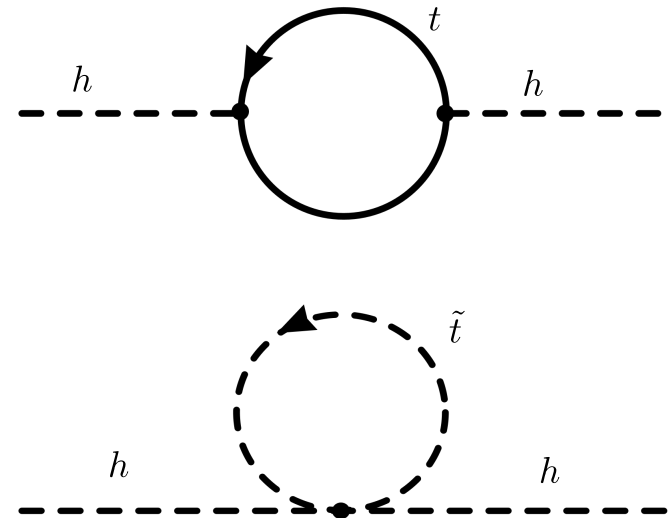
Introduction

Shortcomings of the Standard Model:

Dark Matter:



Hierarchy Inconvenience:



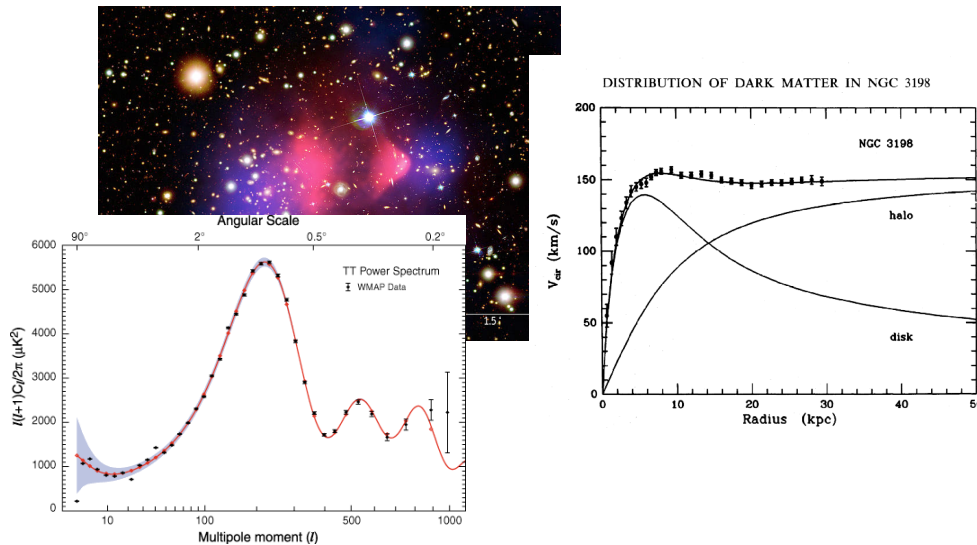
Possible Solution: Lightest neutralino & top partner from the MSSM

$$\mathcal{L} - \mathcal{L}_{\text{SM}} = |D_\mu \tilde{t}|^2 - m_{\tilde{t}}^2 |\tilde{t}|^2 + 1/2 \bar{\chi} (i \not{\partial} - m_\chi) \chi + [y_\chi \bar{t}_R \tilde{t} \chi_L + \text{h.c.}]$$

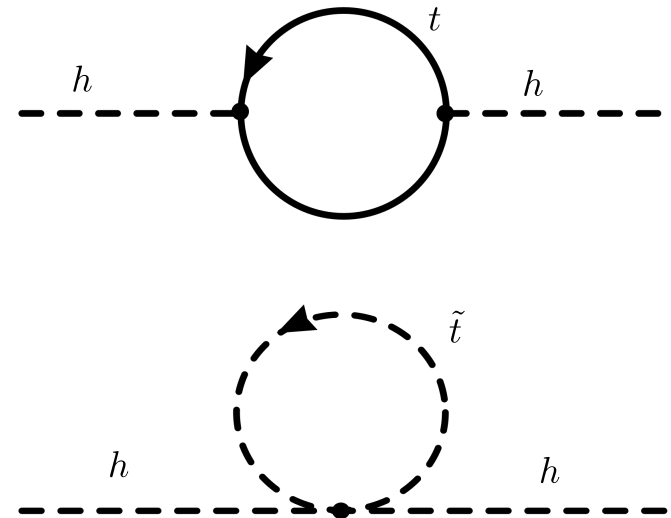
Introduction

Shortcomings of the Standard Model:

Dark Matter:



Hierarchy Inconvenience:



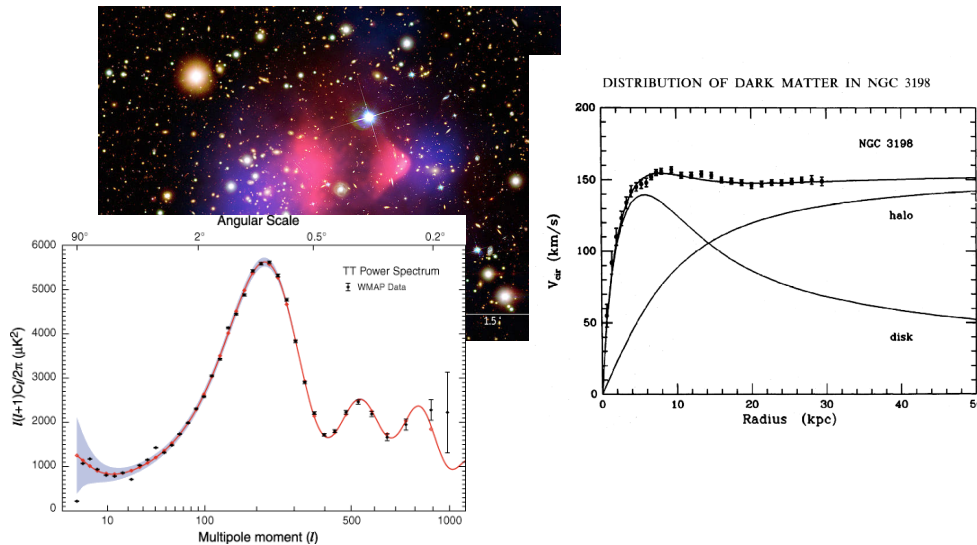
Possible Solution: Lightest neutralino & top partner from the MSSM

$$\mathcal{L} - \mathcal{L}_{\text{SM}} = \boxed{|D_\mu \tilde{t}|^2} - m_{\tilde{t}}^2 |\tilde{t}|^2 + 1/2 \bar{\chi} (i \not{\partial} - m_\chi) \chi + [y_\chi \bar{t}_R \tilde{t} \chi_L + \text{h.c.}]$$

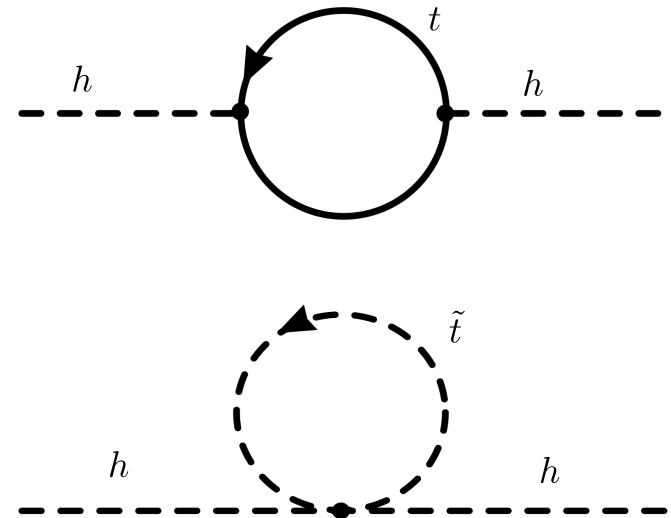
Introduction

Shortcomings of the Standard Model:

Dark Matter:



Hierarchy Inconvenience:



Possible Solution: Lightest neutralino & top partner from the MSSM

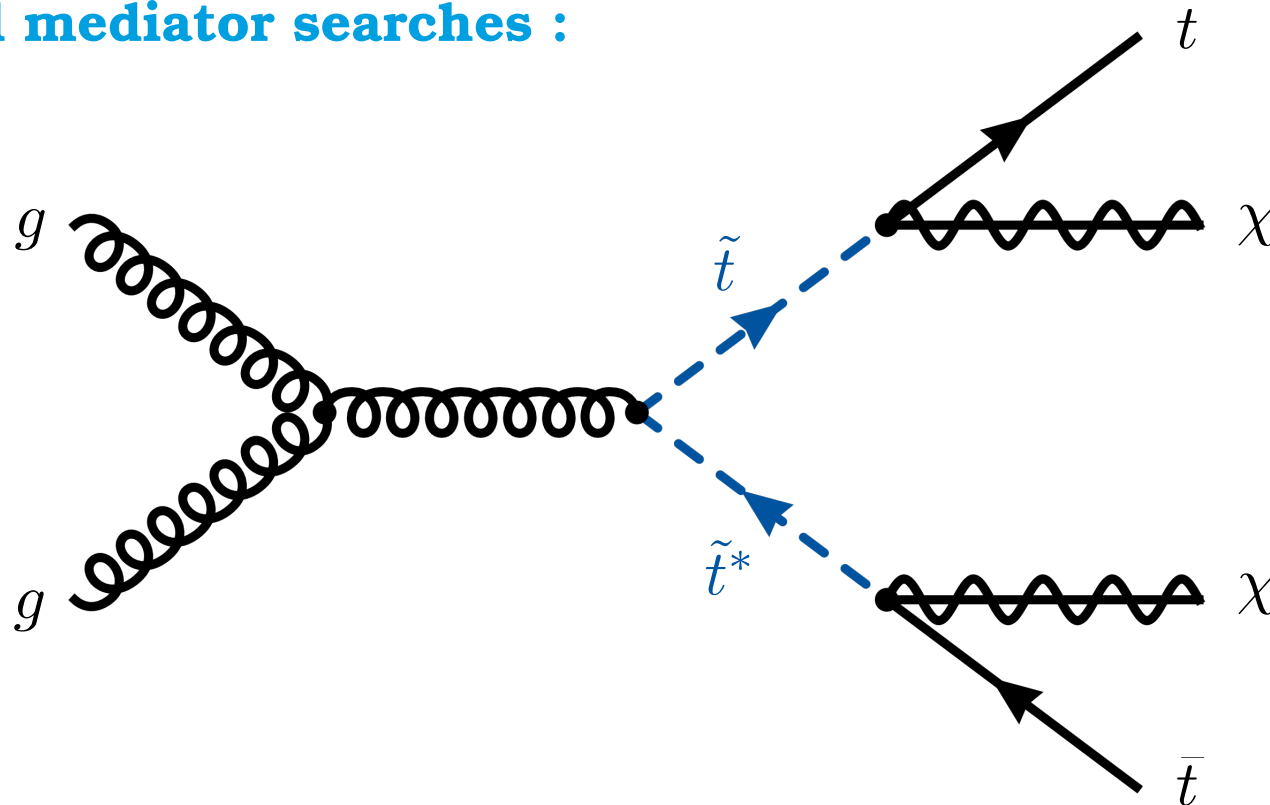
$$\mathcal{L} - \mathcal{L}_{\text{SM}} = |D_\mu \tilde{t}|^2 - m_{\tilde{t}}^2 |\tilde{t}|^2 + 1/2 \bar{\chi} (i \not{\partial} - m_\chi) \chi + [y_\chi \bar{t}_R \tilde{t} \chi_L + \text{h.c.}]$$

Existing searches at the LHC

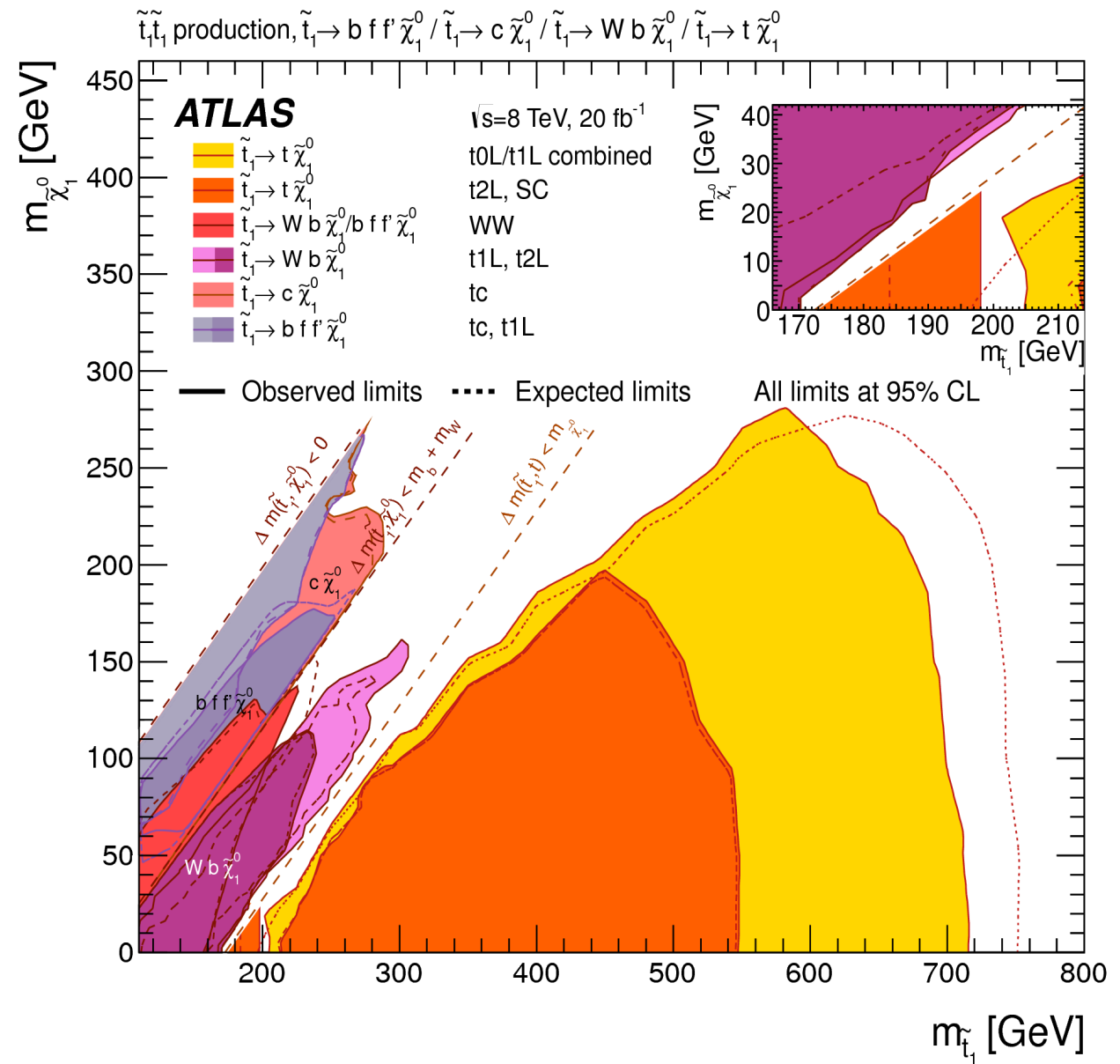
Existing searches at the LHC utilize the production of stop pairs

ATLAS @ [arXiv:1506.08616](#), CMS @ [arXiv:1603.00765](#)

Standard mediator searches :



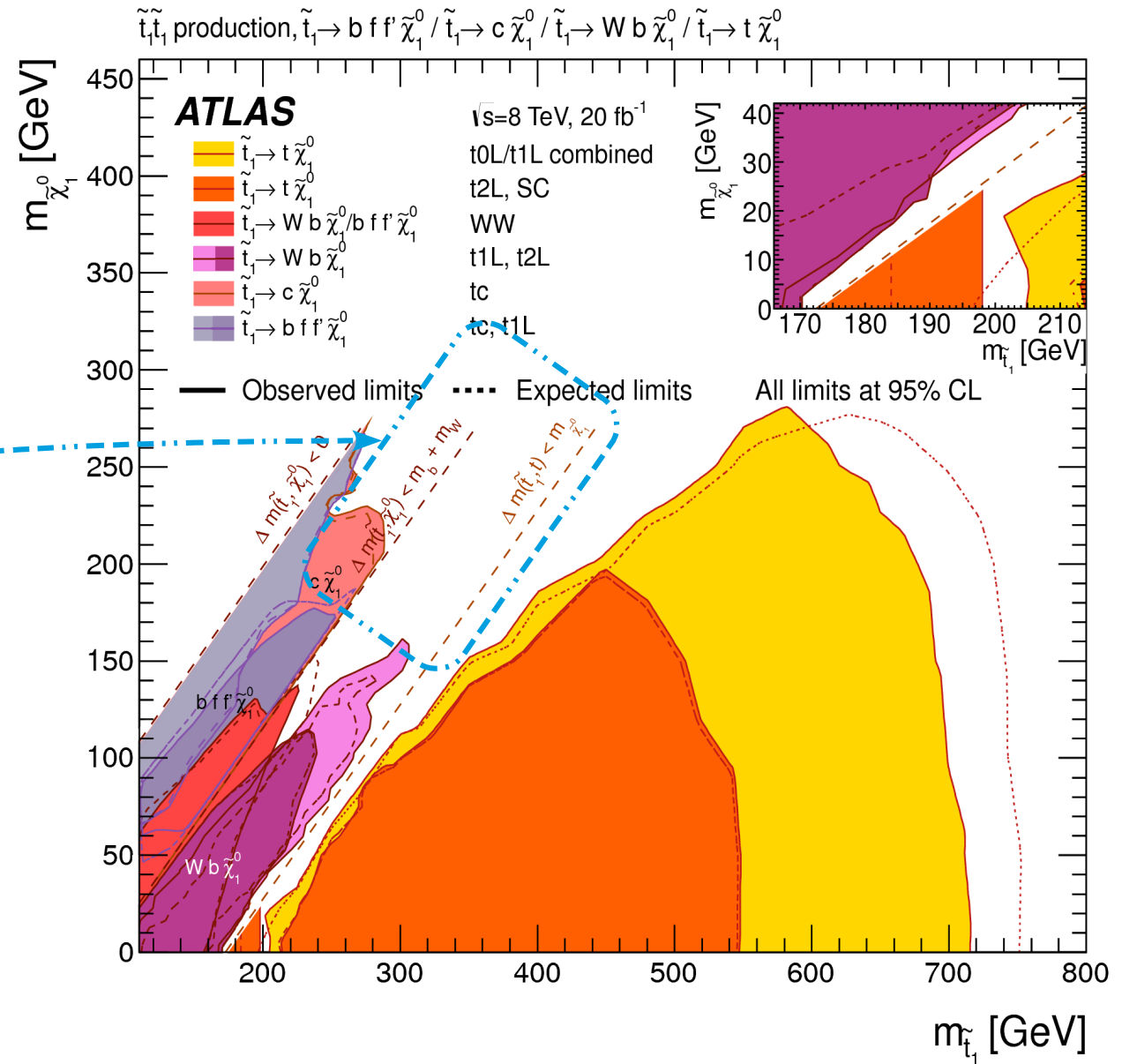
2015



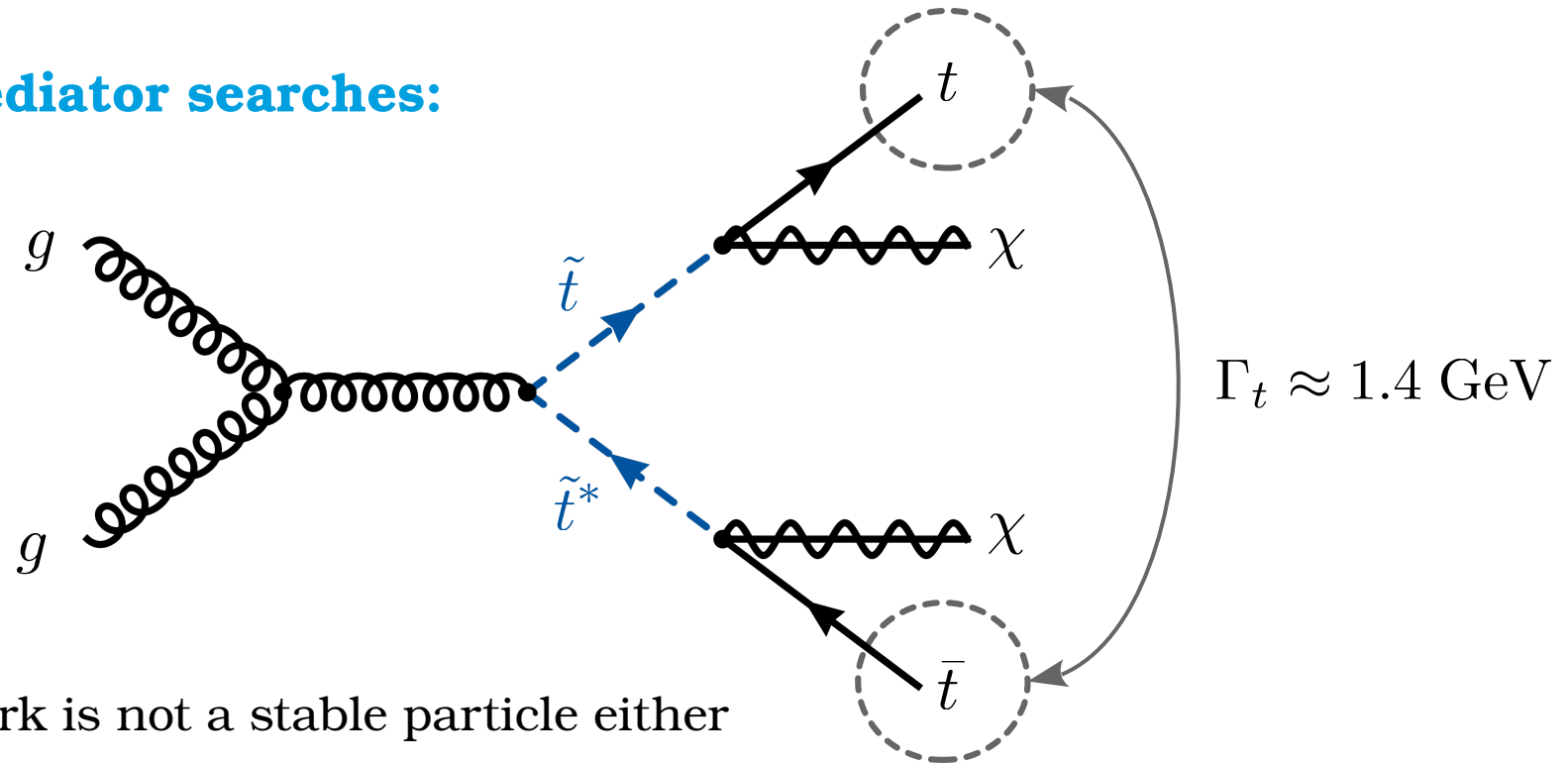
2015

$$\Delta m(\tilde{t}, \chi) \simeq m_W + m_b$$

$$\Delta m(\tilde{t}, \chi) \simeq m_t$$

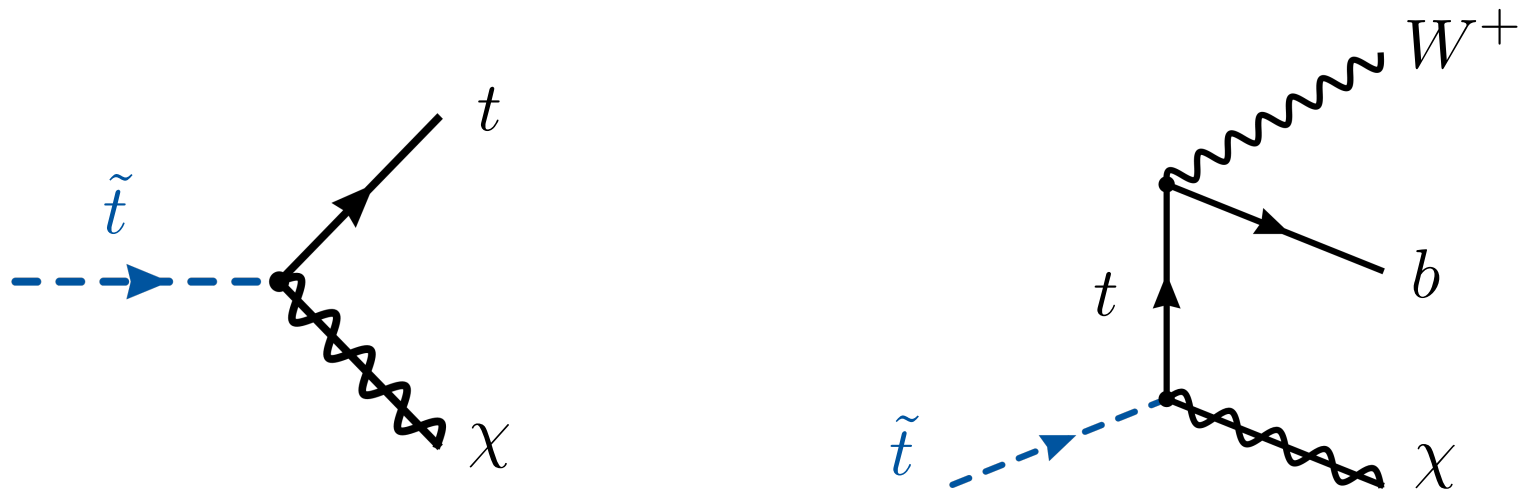


Standard mediator searches:

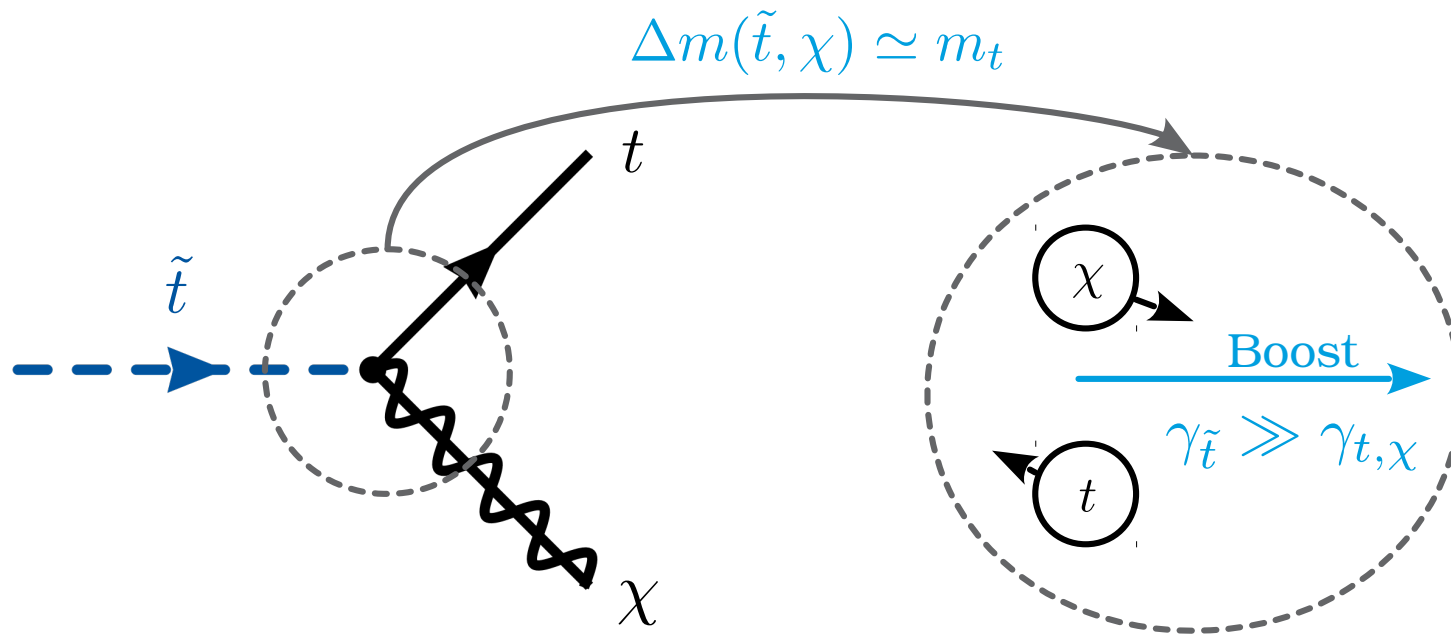


- The top quark is not a stable particle either

→ Might be on- / off-shell depending on the **mass configuration**



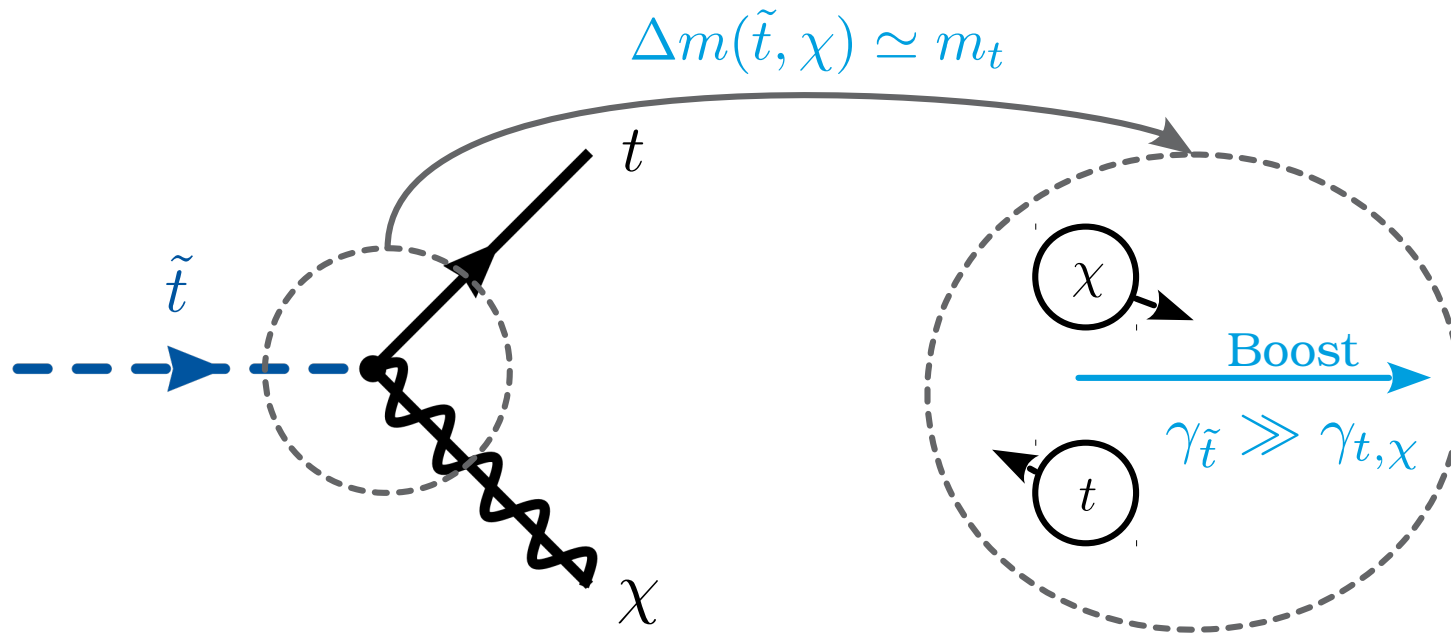
Kinematics in the compressed region (on-shell top)



$$\left. \begin{aligned}
 \underline{p}(\chi)_{\text{RF}} &= \begin{pmatrix} m_\chi \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_\chi \gamma_{\tilde{t}} \\ m_\chi \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix} \\
 \underline{p}(t)_{\text{RF}} &= \begin{pmatrix} m_t \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_t \gamma_{\tilde{t}} \\ m_t \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix}
 \end{aligned} \right\} \Rightarrow \boxed{\frac{\vec{p}(t)_{\text{LF}}}{m_t} = \frac{\vec{p}(\chi)_{\text{LF}}}{m_\chi}}$$

Parallel alignment!

Kinematics in the compressed region (on-shell top)

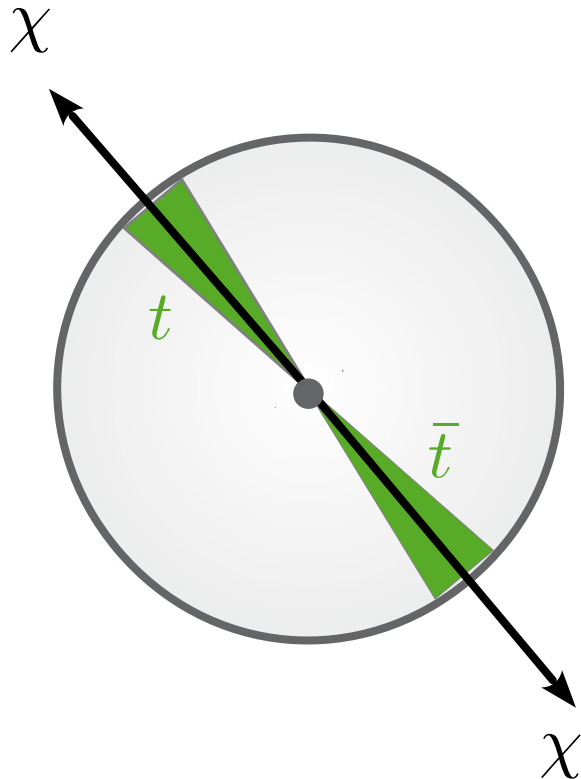


$$\left. \begin{aligned} \underline{p}(\chi)_{\text{RF}} &= \begin{pmatrix} m_\chi \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_\chi \gamma_{\tilde{t}} \\ m_\chi \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix} \\ \underline{p}(t)_{\text{RF}} &= \begin{pmatrix} m_t \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_t \gamma_{\tilde{t}} \\ m_t \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix} \end{aligned} \right\} \Rightarrow \boxed{\frac{\vec{p}(t)_{\text{LF}}}{m_t} = \frac{\vec{p}(\chi)_{\text{LF}}}{m_\chi}}$$

➡ Parallel alignment!

Kinematics at the detector level : $pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow t\bar{t}\chi\chi$

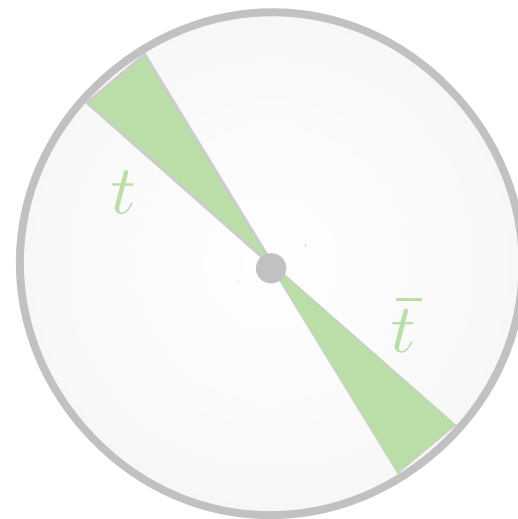
$$\frac{\vec{p}(t)_{\text{LF}}}{m_t} = \frac{\vec{p}(\chi)_{\text{LF}}}{m_\chi}$$



Looks just like

$$pp \rightarrow t\bar{t}$$

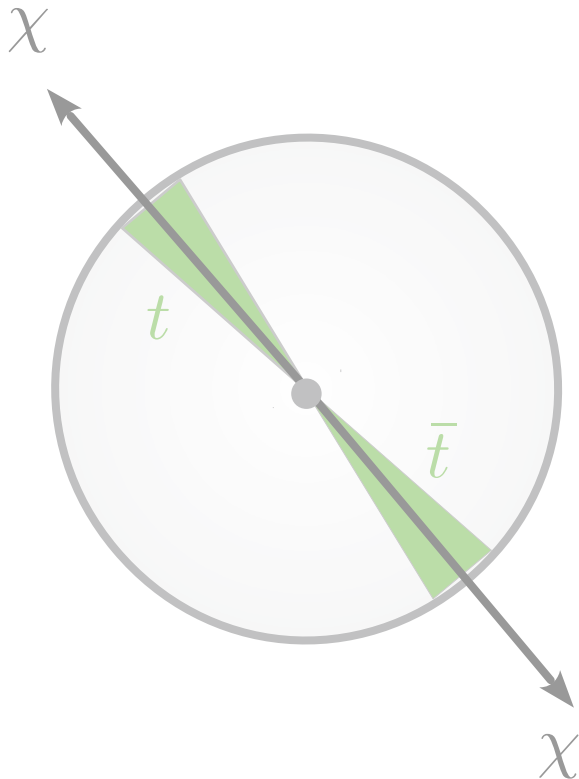
$$\text{MET} \simeq 0$$



➡ No experimental difference to top pair production for
 $\Delta m(\tilde{t}, \chi) \simeq m_t$

Kinematics at the detector level : $pp \rightarrow \tilde{t}\tilde{t}^* \rightarrow t\bar{t}\chi\chi$

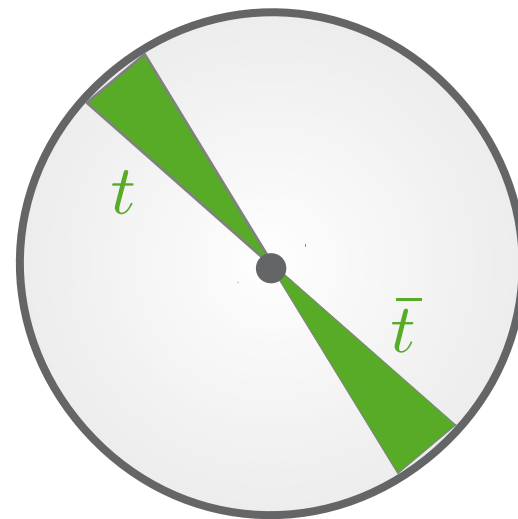
$$\frac{\vec{p}(t)_{\text{LF}}}{m_t} = \frac{\vec{p}(\chi)_{\text{LF}}}{m_\chi}$$



Looks just like

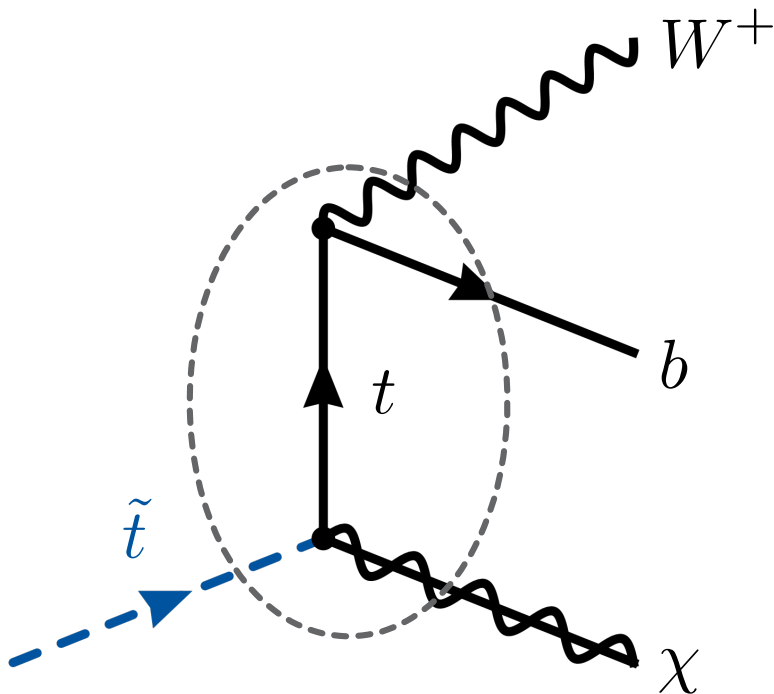
$$pp \rightarrow t\bar{t}$$

$$\text{MET} \simeq 0$$



➡ No experimental difference to **top pair production** for
 $\Delta m(\tilde{t}, \chi) \simeq m_t$

Kinematics in the compressed region (off-shell top) :



Similar result :

$$\frac{\vec{p}(W)_{\text{LF}}}{m_W} = \frac{\vec{p}(b)_{\text{LF}}}{m_b} = \frac{\vec{p}(\chi)_{\text{LF}}}{m_\chi}$$

➡ Parallel alignment!

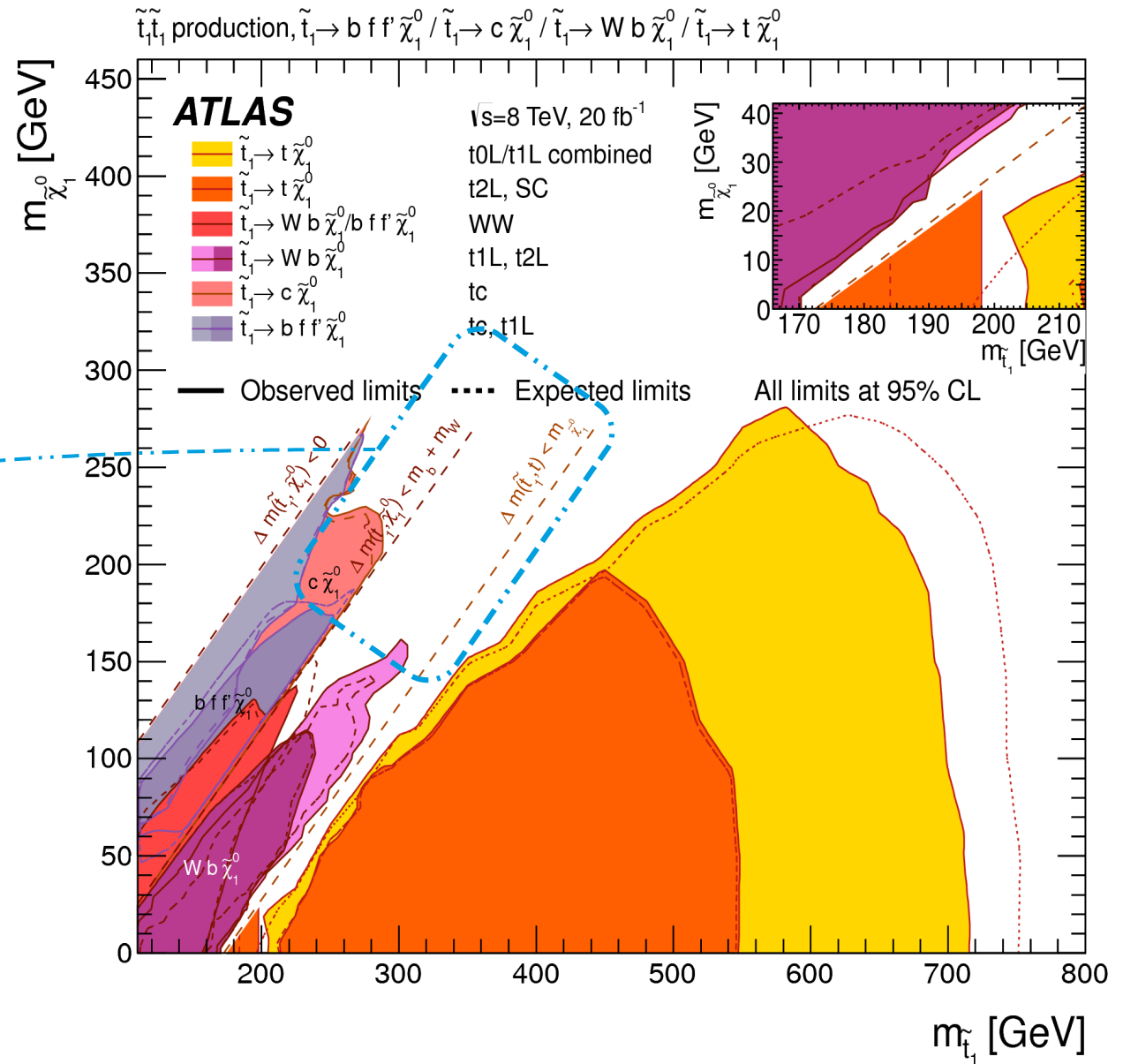
➡ No experimental difference to **top pair production** for

$$\Delta m(\tilde{t}, \chi) \simeq m_W + m_b$$

2015

$$\Delta m(\tilde{t}, \chi) \simeq m_W + m_b$$

$$\Delta m(\tilde{t}, \chi) \simeq m_t$$

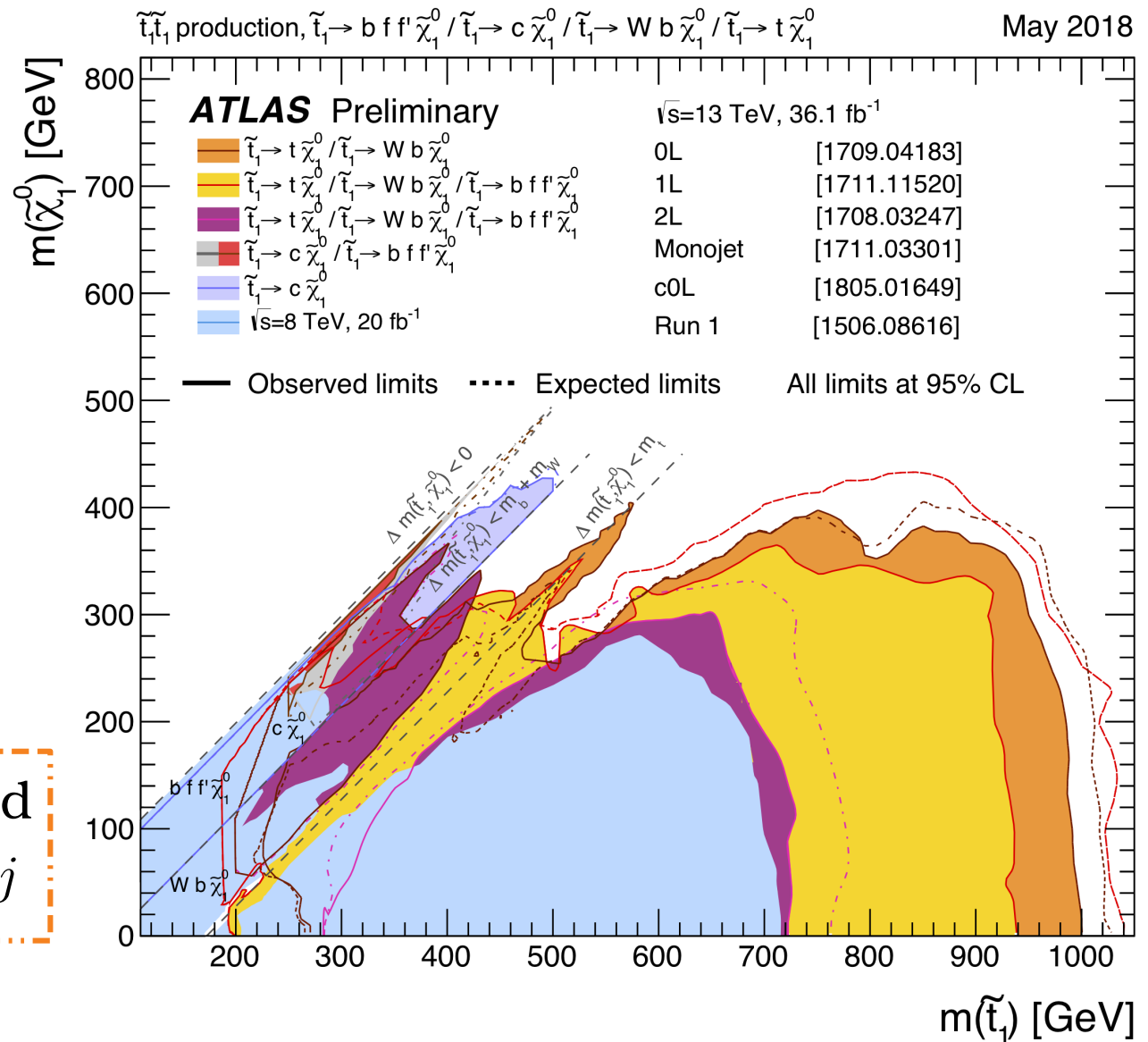


Exclusion limits for light stops (ATLAS Preliminary)

2018

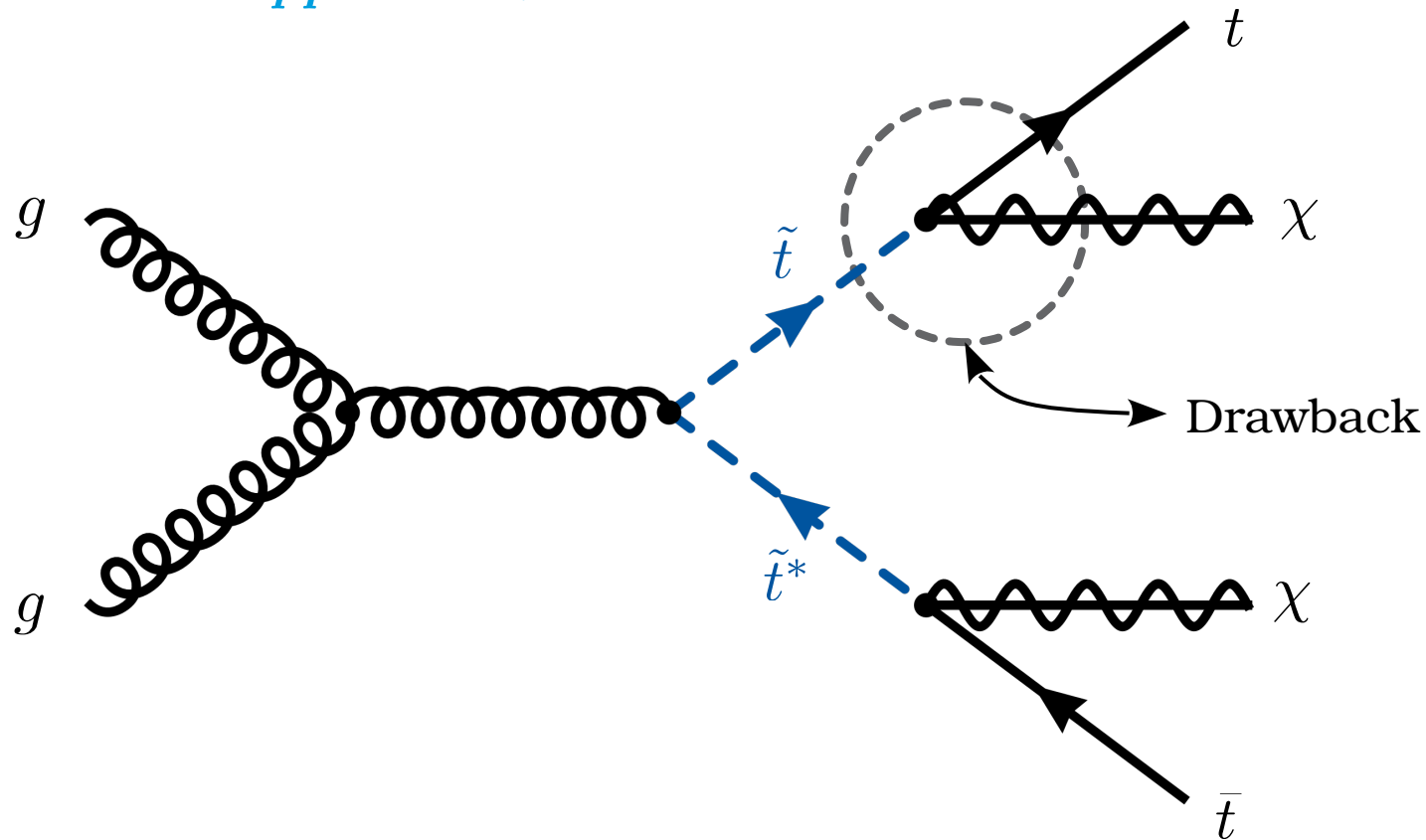
New method:

Probe the compressed region via $pp \rightarrow t\bar{t}^* + j$



Probing the compressed region in $pp \rightarrow \tilde{t}\tilde{t}^* j$

Search channel $pp \rightarrow \tilde{t}\tilde{t}^* j$:

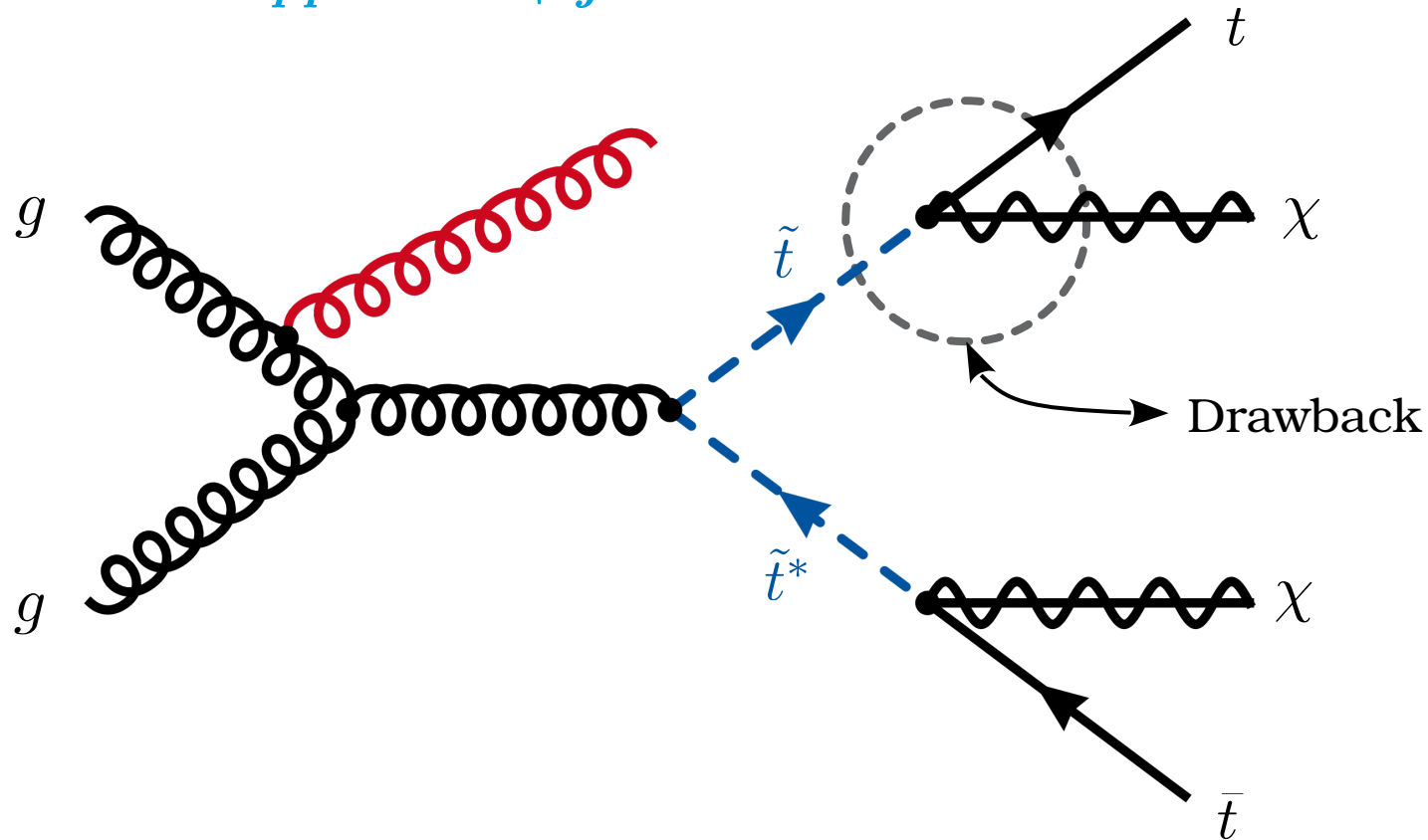


Loophole: Enrich the kinematics by an additional (hard) jet

→ Take advantage of the disadvantage (arXiv:1506.00653)

Probing the compressed region in $pp \rightarrow \tilde{t}\tilde{t}^* j$

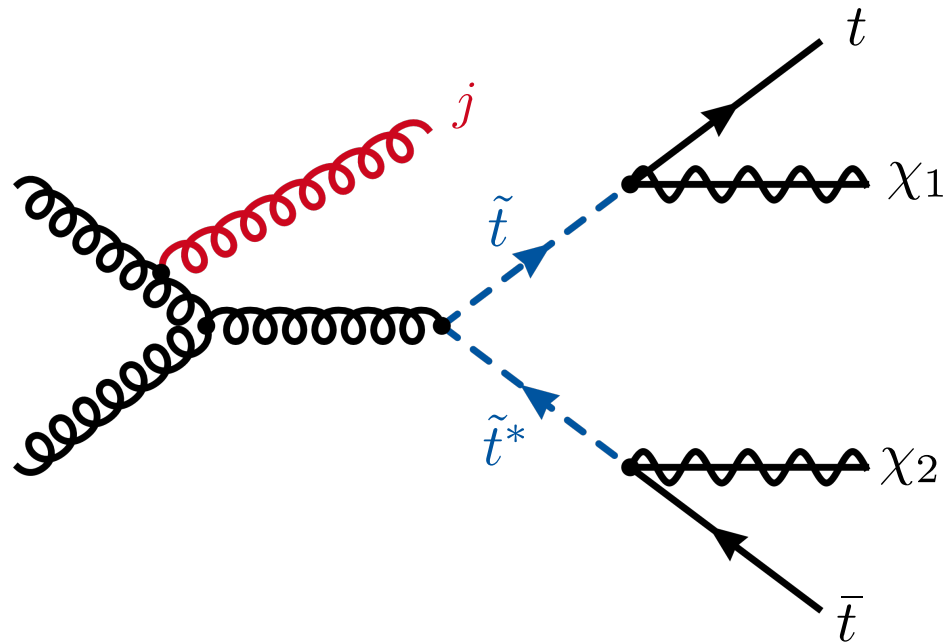
Search channel $pp \rightarrow \tilde{t}\tilde{t}^* + j$:



Loophole: Enrich the kinematics by an additional (hard) jet

→ Take advantage of the disadvantage (arXiv:1506.00653)

Influence of the additional jet



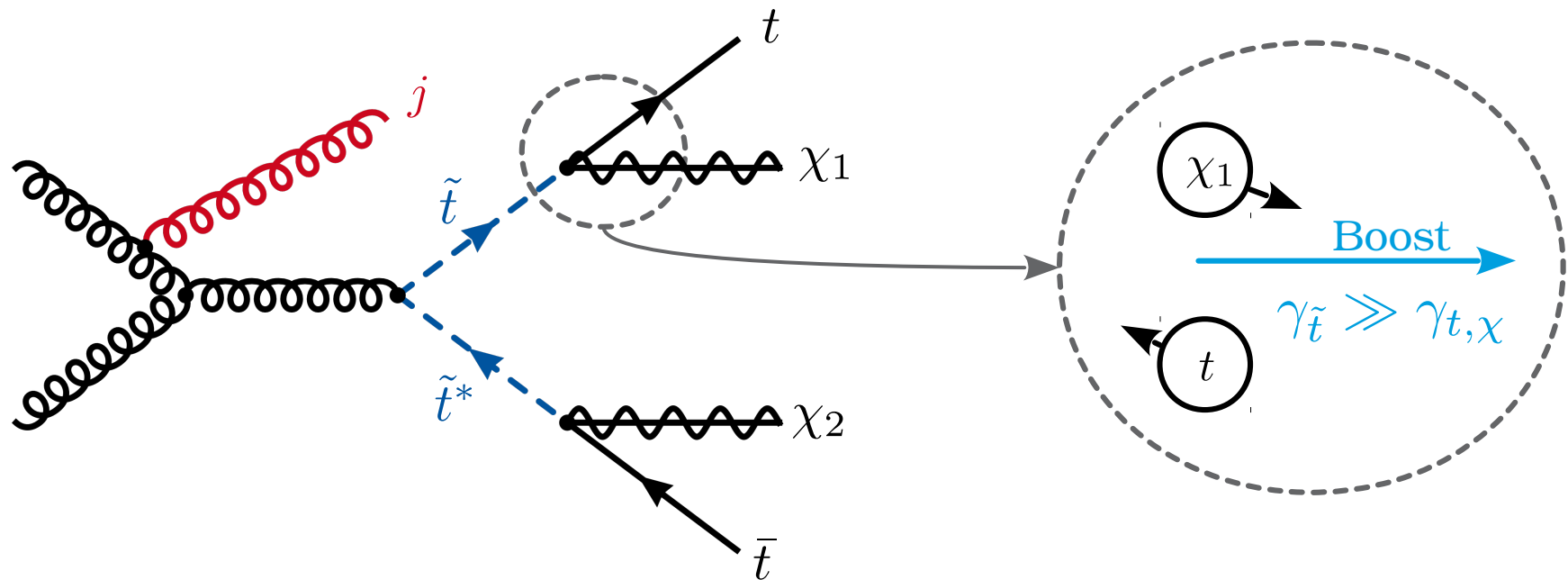
$$-\vec{p}_T(j) = \underbrace{\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)}_{\text{momentum conservation}}$$

$$\begin{aligned} \vec{p}_{T,\text{miss}} &= \vec{p}_T(\chi_1) + \vec{p}_T(\chi_2) \\ &= \frac{m_\chi}{m_{\tilde{t}}} (\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)) \end{aligned}$$

$\Rightarrow$

$$R_M \equiv \frac{|\vec{p}_{T,\text{miss}}|}{|\vec{p}_T(j)|} \simeq \frac{m_\chi}{m_{\tilde{t}}}$$

Influence of the additional jet



$$-\vec{p}_T(j) = \underbrace{\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)}_{\text{momentum conservation}}$$

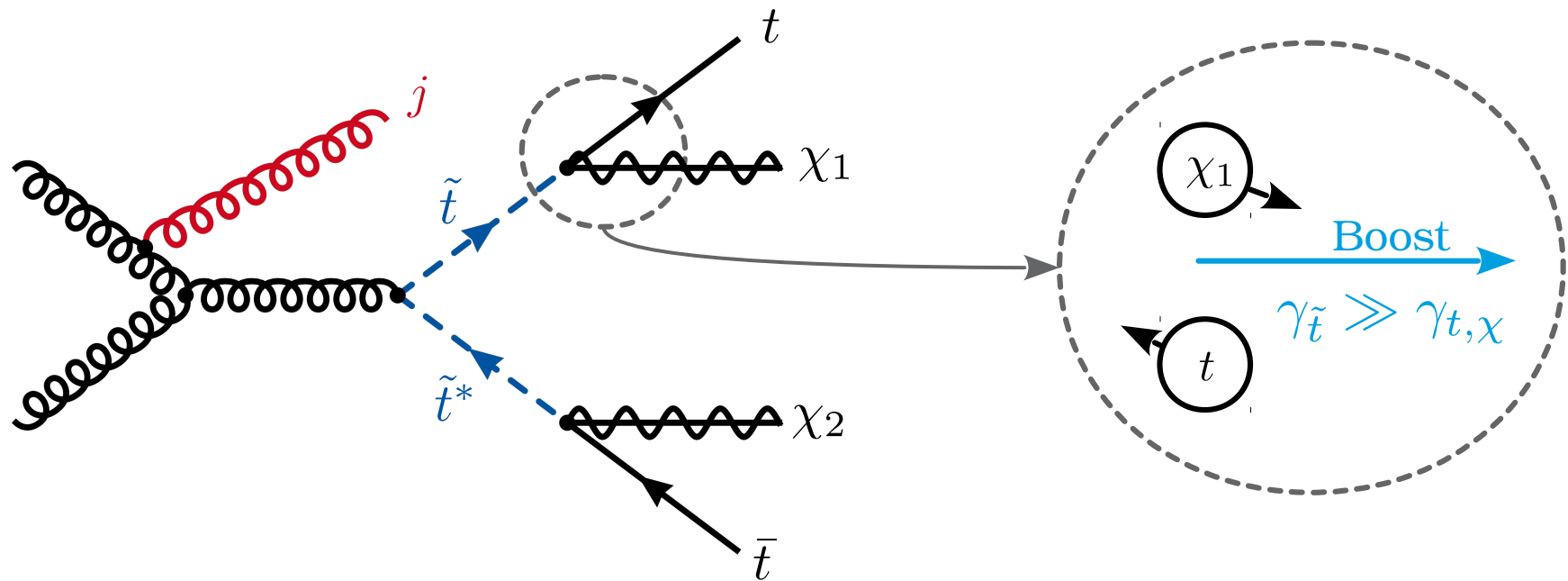
$$\begin{aligned} \underline{p}(\chi_1)_{\text{RF}} &= \begin{pmatrix} m_\chi \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_\chi \gamma_{\tilde{t}} \\ m_\chi \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix} = \underline{p}(\chi_1)_{\text{LF}} \\ &= \frac{m_\chi}{m_{\tilde{t}}} \begin{pmatrix} E_{\tilde{t}} \\ \vec{p}_{\tilde{t}} \end{pmatrix} = \frac{m_\chi}{m_{\tilde{t}}} \underline{p}(\tilde{t})_{\text{LF}} \end{aligned}$$

$$\begin{aligned} \vec{p}_{T,\text{miss}} &= \vec{p}_T(\chi_1) + \vec{p}_T(\chi_2) \\ &= \frac{m_\chi}{m_{\tilde{t}}} (\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)) \end{aligned}$$

$\Rightarrow$

$$R_M \equiv \frac{|\vec{p}_{T,\text{miss}}|}{|\vec{p}_T(j)|} \simeq \frac{m_\chi}{m_{\tilde{t}}}$$

Influence of the additional jet



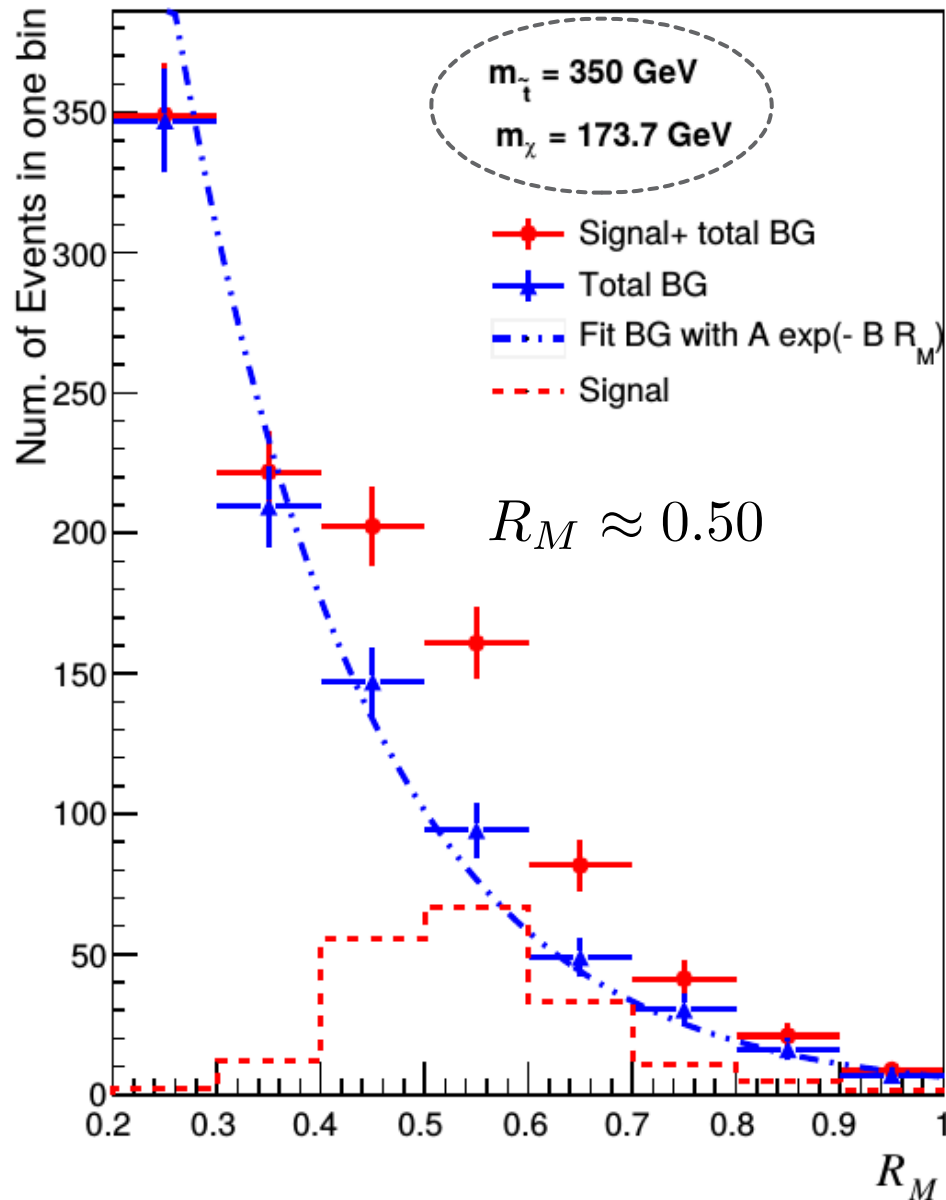
$$-\vec{p}_T(j) = \underbrace{\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)}_{\text{momentum conservation}}$$

$$\begin{aligned} \underline{p}(\chi_1)_{\text{RF}} &= \begin{pmatrix} m_\chi \\ \vec{0} \end{pmatrix} \xrightarrow{\text{Boost}} \begin{pmatrix} m_\chi \gamma_{\tilde{t}} \\ m_\chi \gamma_{\tilde{t}} \vec{\beta}_{\tilde{t}} \end{pmatrix} = \underline{p}(\chi_1)_{\text{LF}} \\ &= \frac{m_\chi}{m_{\tilde{t}}} \begin{pmatrix} E_{\tilde{t}} \\ \vec{p}_{\tilde{t}} \end{pmatrix} = \frac{m_\chi}{m_{\tilde{t}}} \underline{p}(\tilde{t})_{\text{LF}} \end{aligned}$$

$$\begin{aligned} \vec{p}_{T,\text{miss}} &= \vec{p}_T(\chi_1) + \vec{p}_T(\chi_2) \\ &= \frac{m_\chi}{m_{\tilde{t}}} (\vec{p}_T(\tilde{t}) + \vec{p}_T(\tilde{t}^*)) \end{aligned}$$

$\Rightarrow$

$$R_M \equiv \frac{|\vec{p}_{T,\text{miss}}|}{|\vec{p}_T(j)|} \simeq \frac{m_\chi}{m_{\tilde{t}}}$$



Background :

$t\bar{t}$ with hard ISR jet

QCD multijet

$W + j$, $Z + j$, single top

Detector-level cuts :

$$N_{\text{jets}} \geq 4$$

$$p_T(j_1) > 700 \text{ GeV}$$

$$N_e, N_{\mu} = 0$$

$$p_T(j_{2..4}) > 60 \text{ GeV}$$

$$|\phi(j_1) - \phi_{\text{MET}} - \pi| < 0.15$$

$$|\phi_{\text{MET}} - \phi(j_{2..4})| > 0.2$$

$$N_{\text{b-tags}} \neq 0$$

$$j_1:\text{b-tag} = \text{False}$$

➡ The peak-like structure of the signal results in a **clear excess**

Shortcomings of the analysis

Exclusion/Discovery limits have already been calculated

However: Only at leading order with inclusive K-factors

Aim:

Investigate the effects of **next-to-leading order** corrections to the signal distributions

→ Especially analyze the effect of the cuts

Shortcomings of the analysis

Exclusion/Discovery limits have already been calculated

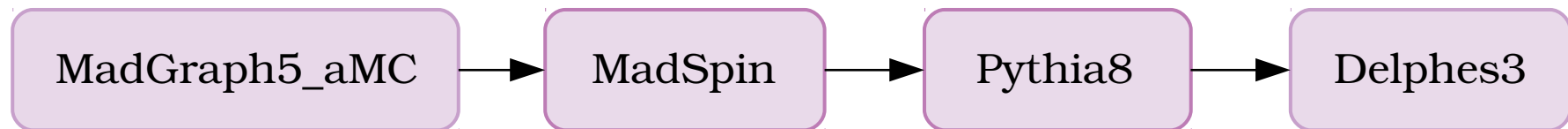
However: Only at leading order with inclusive K-factors

Aim:

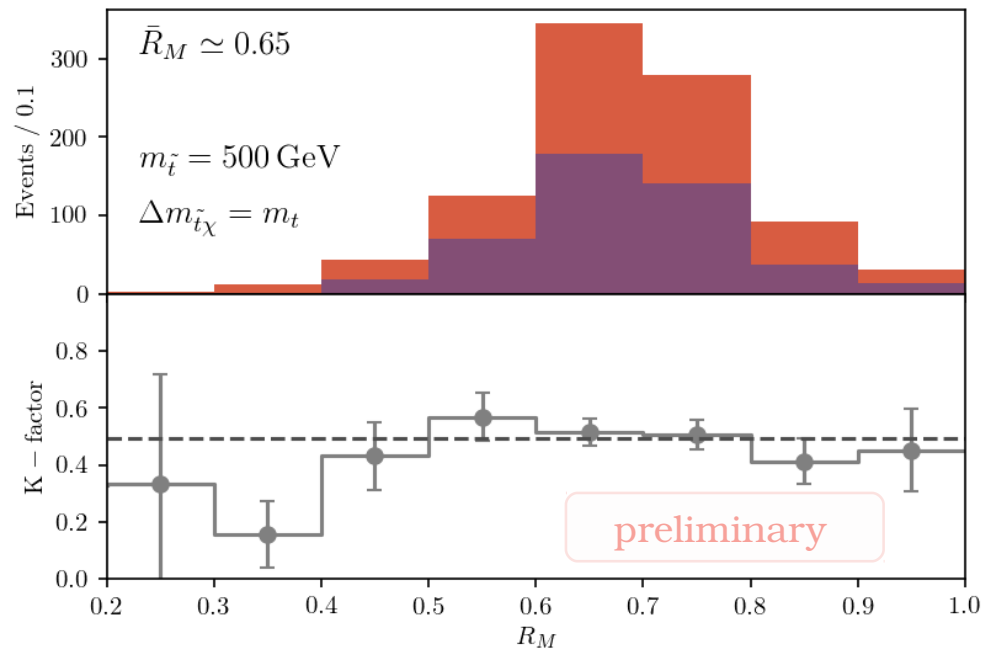
Investigate the effects of **next-to-leading order** corrections to the signal distributions

→ Especially analyze the effect of the cuts

Event generation (Monte Carlo samples):

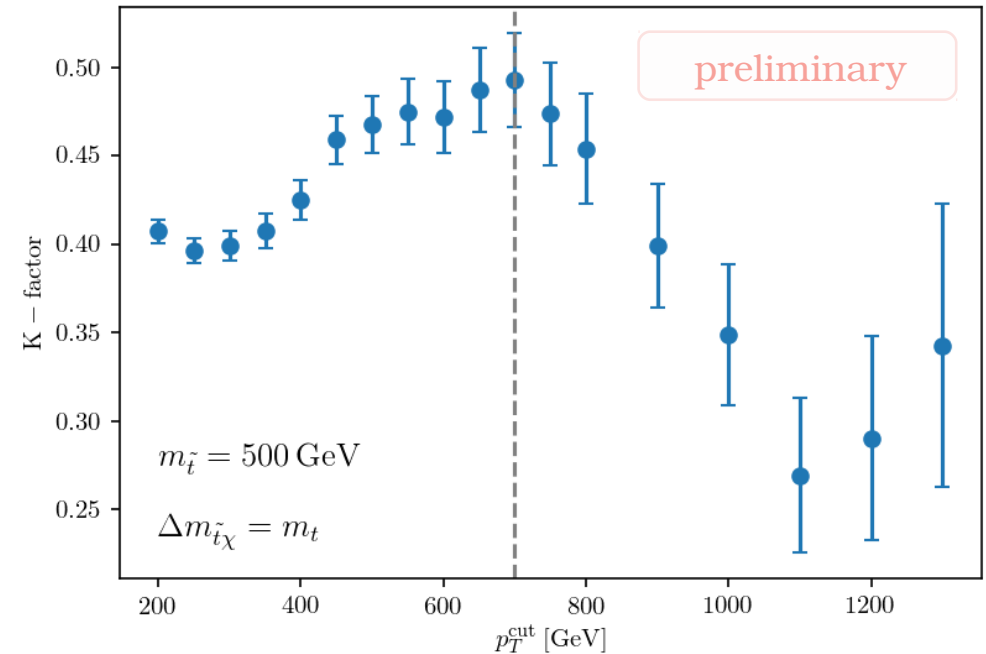
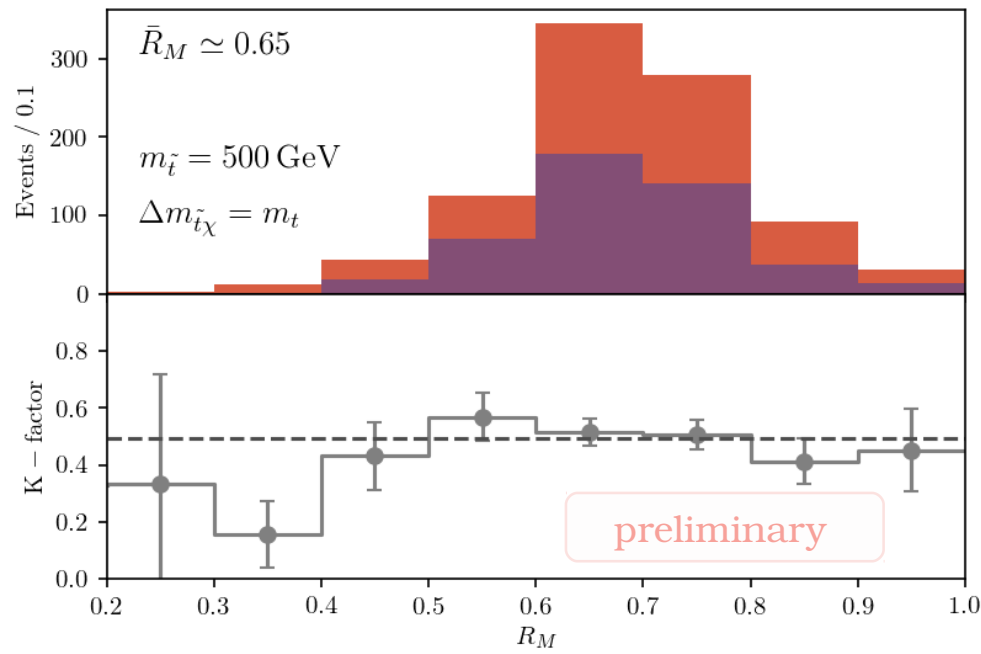


Effects of NLO corrections



- NLO distribution does not seem to exhibit any new physical features
→ Cross-section seems to be scalable by a global K-factor

Effects of NLO corrections

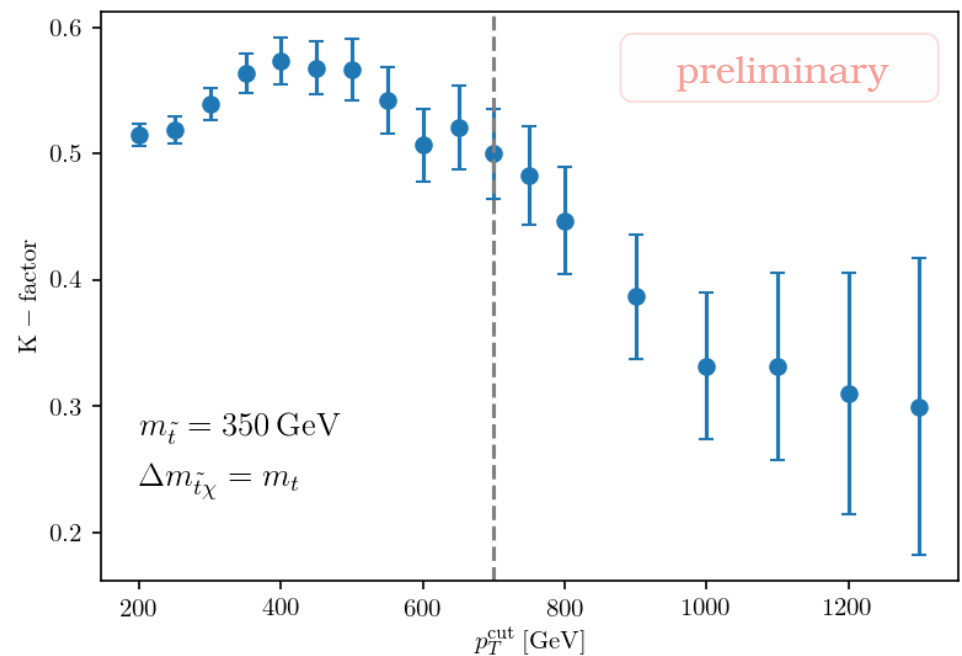
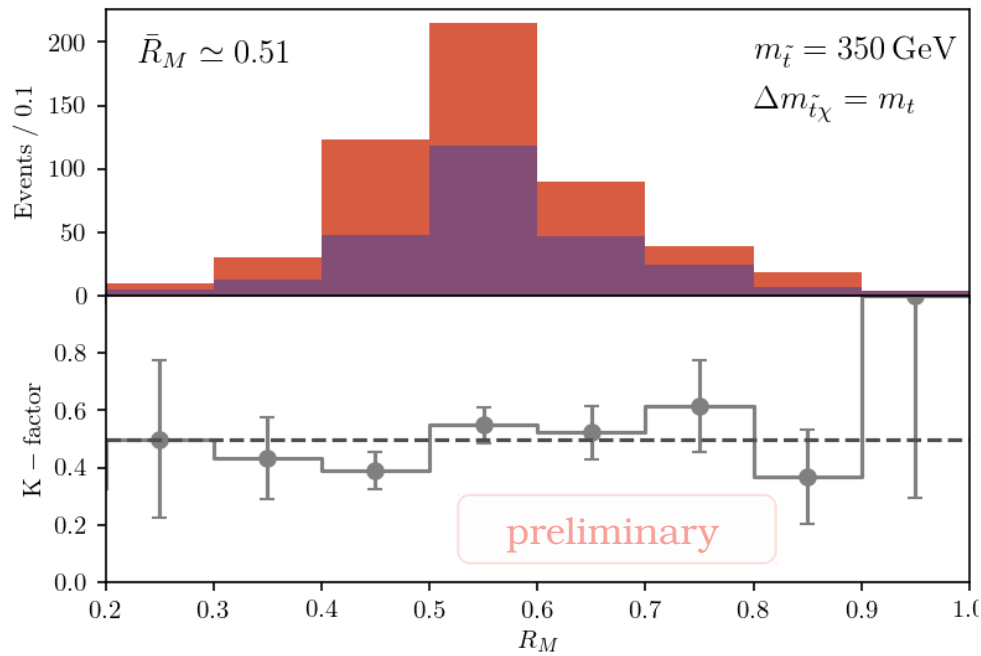


- NLO distribution does not seem to exhibit any new physical features
→ Cross-section seems to be scalable by a global K-factor



- However: The K-factor depends on the exact nature of the cut, p_T^{cut}

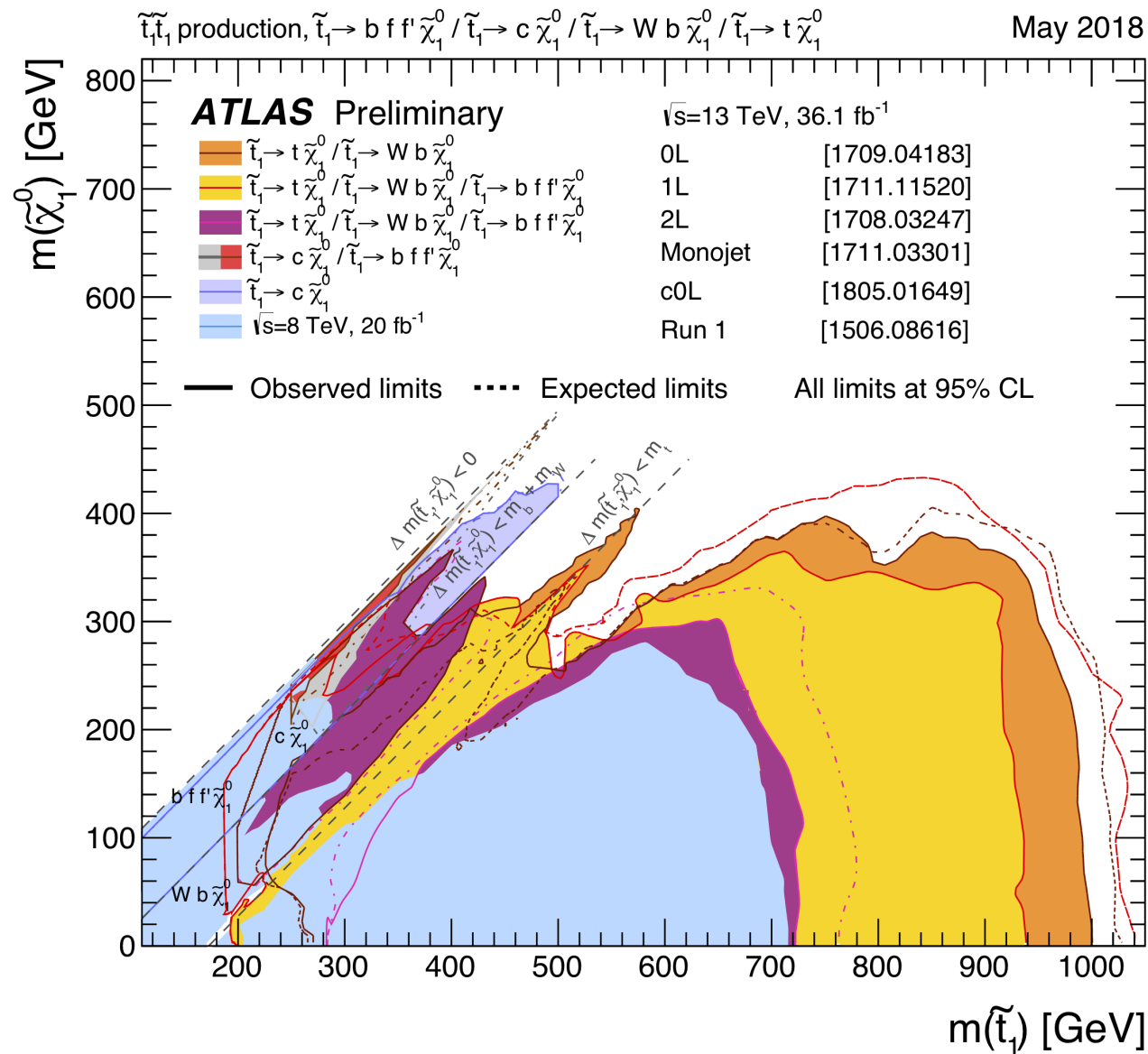
Effects of NLO corrections



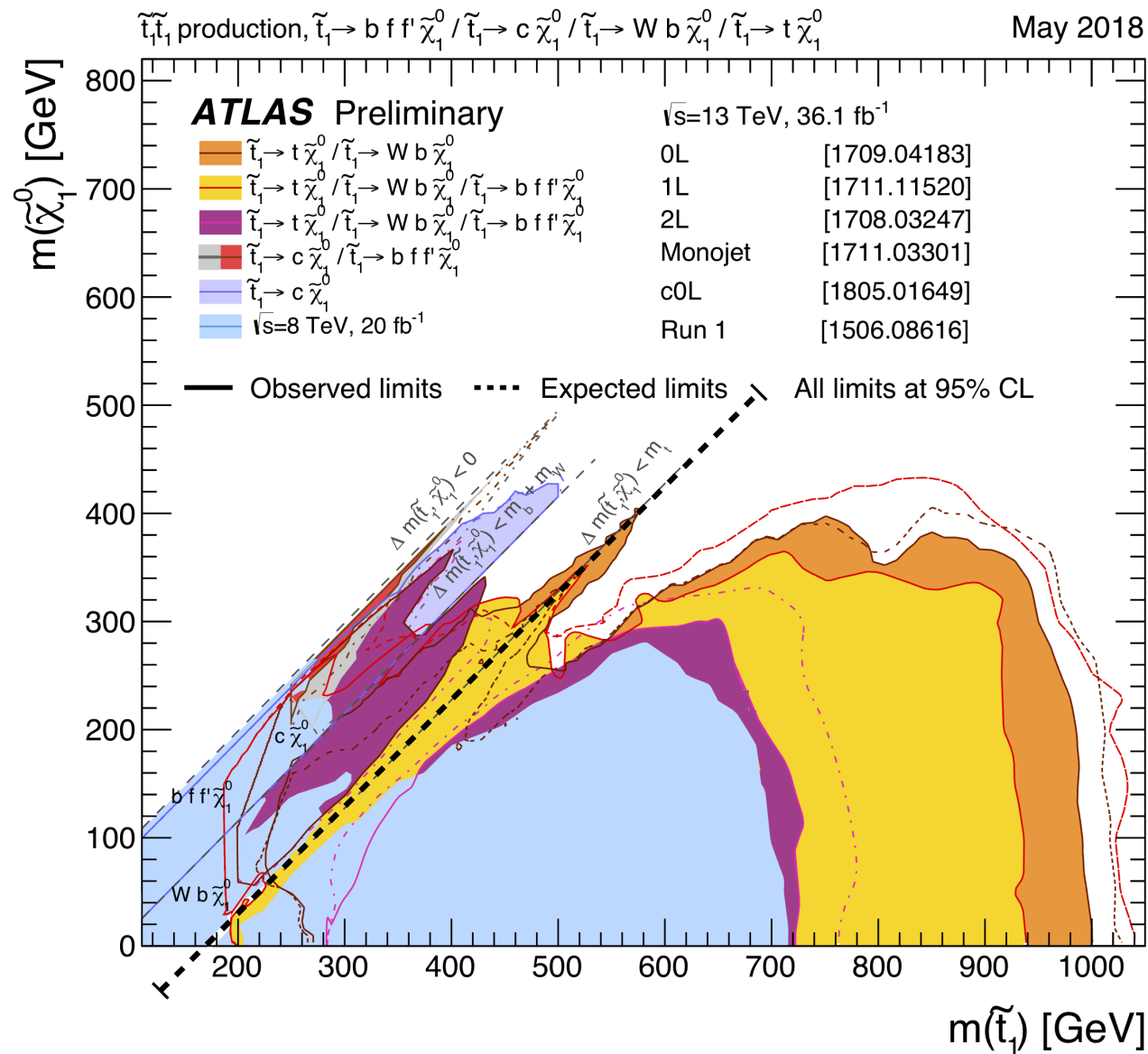
- NLO distribution does not seem to exhibit any new physical features
→ Cross-section seems to be scalable by a global K-factor
- However: The K-factor depends on the exact nature of the cut, p_T^{cut}



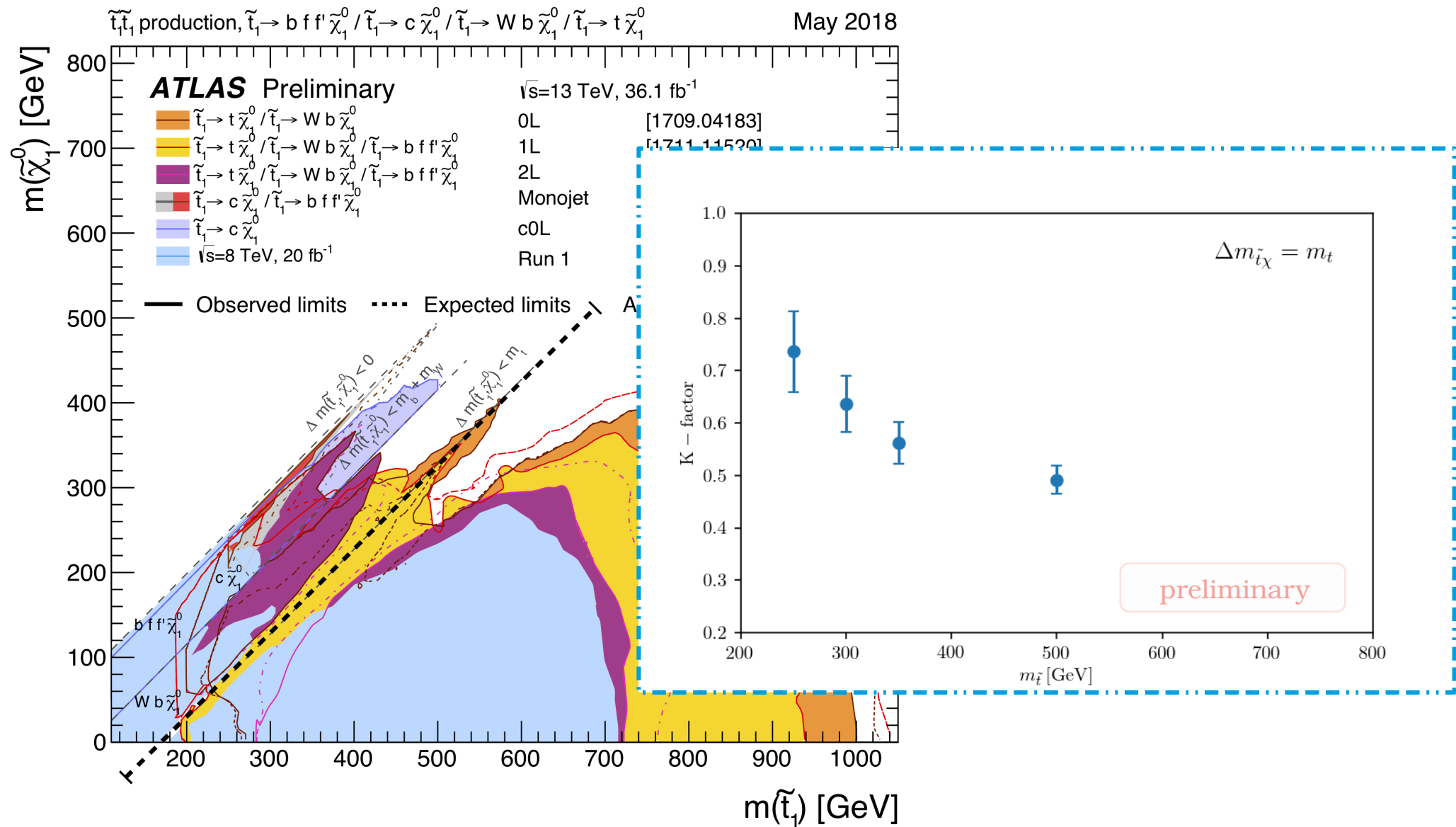
Effects of NLO corrections



Effects of NLO corrections



Effects of NLO corrections



Conclusions

- Using new techniques it has been possible to completely close the compressed region of stop searches
- However: The current limits are sensitive to the exact nature of the cuts → More detailed analysis needed

Conclusions

- Using new techniques it has been possible to completely close the compressed region of stop searches
- However: The current limits are sensitive to the exact nature of the cuts → More detailed analysis needed

Thank you for
your attention!