

Inversion formula for Defects I: Defect Channel

Pedro Liendo



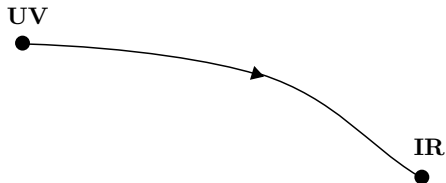
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DESY Theory Workshop

1712.08185 with M. Lemos, M. Meineri, S. Sarkar

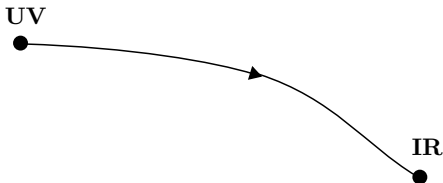
Motivation

Renormalization group flow



Renormalization group (RG) flow.

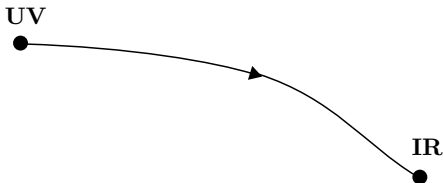
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Conformal symmetry is extremely powerful.

It includes **translations**, **rotations**, and **scale transformations** (+ more)

$$x \rightarrow x + a, \quad x \rightarrow R \cdot x, \quad x \rightarrow h x.$$

CFT

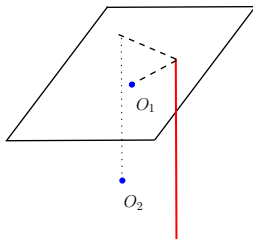
Defect CFT

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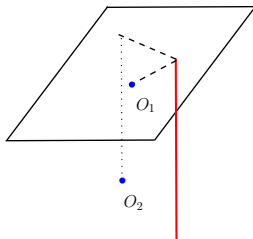
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Local operators in the presence of a defect.

We have $SO(1, d + 1) \rightarrow SO(1, p + 1) \times SO(q)$ where $q + p = d$.

Defect CFT correlators

The $SO(1, p+1) \times SO(q)$ symmetry preserved by the defect implies that one-point functions are non-zero:

$$\langle \mathcal{O}(x) \rangle = \frac{a_{\mathcal{O}}}{(x^i)^{\Delta}}.$$

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Two-point functions depend on two conformal invariants

$$\langle \phi(x_1) \phi(x_2) \rangle = \frac{1}{(z\bar{z})^{\Delta_{\phi/2}}} g(z, \bar{z}),$$

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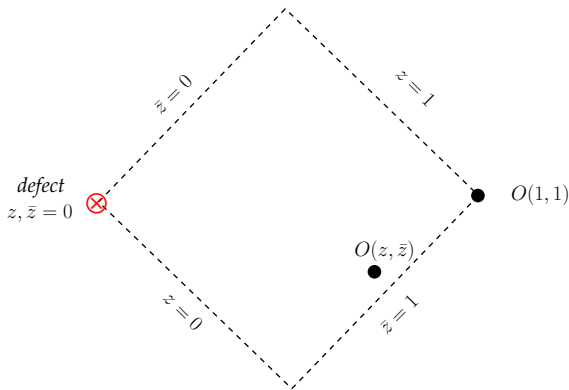
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Remark. Compare with the four-point function in standard CFT.

Two-point function configuration



Configuration of the system in the plane orthogonal to the defect.

Bulk OPE

Bulk channel:

We had

$$\phi(x)\phi(0) \sim \sum_{\mathcal{O}} c_{\phi\phi\mathcal{O}} d(x, \partial) \mathcal{O}(0).$$

recall that in the presence of a defect a scalar can have a non-zero one-point function

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The expansion for the two-point function is

$$\langle \phi(x_1)\phi(x_2) \rangle = \left(\frac{(1-z)(1-\bar{z})}{(z\bar{z})^{1/2}} \right)^{-\Delta_\phi} \sum_{\Delta, J} C_{\phi\phi\mathcal{O}} a_{\mathcal{O}} f_{\Delta, J}(z, \bar{z})$$

where the **sum** goes over the **bulk spectrum**.

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Defect channel:

We can also write a bulk operator as a sum of defect operators

$$\phi(x) = \sum_{\hat{O}} b_{\phi\hat{O}} D(x^i, \partial_{\vec{x}}) \hat{O}(\vec{x})$$

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Plugging this expansion into the two-point function,

$$\langle \phi(x_1) \phi(x_2) \rangle = \sum_{\hat{\Delta}, s} b_{\phi\hat{O}}^2 \hat{f}_{\hat{\Delta}, s}(z, \bar{z}).$$

where the **sum** goes over the **boundary spectrum**.

Crossing symmetry

Equality of both expansions implies

$$\left(\frac{(1-z)(1-\bar{z})}{(z\bar{z})^{1/2}} \right)^{-\Delta_\phi} \sum_{\Delta, J} C_{\phi\phi O} a_O f_{\Delta, J}(z, \bar{z}) = \sum_{\widehat{\Delta}, s} b_{\phi\widehat{O}}^2 \widehat{f}_{\widehat{\Delta}, s}(z, \bar{z})$$

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Pictorially

$$\sum_{\Delta, J} C_{\phi\phi O} a_O \text{ (diagram of two external legs meeting at a vertex } O \text{ with a horizontal line below)} = \sum_{\hat{\Delta}, s} b_{\phi\hat{O}}^2 \text{ (diagram of two external legs meeting at a vertex } \hat{O} \text{ with a horizontal line below)}$$

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The **defect blocks** are known in closed-form.

[Billo, Goncalvez, Lauria, Meineri (2016)]

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The **bulk blocks** are Calogero-Sutherland wave-functions.

[Isachenkov, PL, Linke, Schomerus (2018)]

Inversion Formula

From Euclidean to Lorentzian

The idea...

$$z = rw, \quad \bar{z} = \frac{r}{w}$$

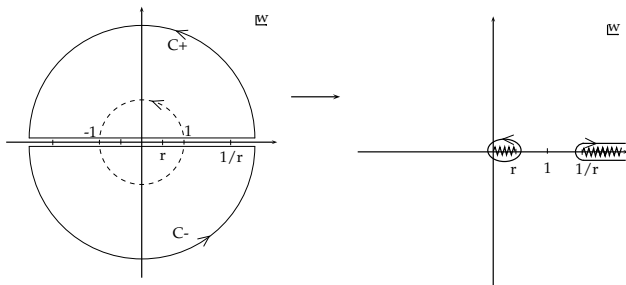
$$g(r, w) = \int b(\hat{\Delta}, s) h(r, w) \quad \rightarrow \quad b(\hat{\Delta}, s) \sim \int g(r, w) \bar{h}(r, w)$$

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Contour deformation from Euclidean to Lorentzian configuration.

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Long story short: we need an *infinite number of defect operators*.

[Lemos, PL, Meineri, Sarkar (2018)]

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$$\hat{\Delta} = \Delta_{\phi} + s + 2m + \mathcal{O}(s^{-\alpha}) , \quad s \rightarrow \infty ,$$

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$$b_{s,m}^2 = s^{\Delta_\phi-1} \left(\frac{1}{\Gamma(\Delta_\phi)} \binom{m - \frac{d}{2} + \Delta_\phi}{m} + \mathcal{O}(s^{-\beta}) \right) , \quad s \rightarrow \infty ,$$

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This result is **universal**!

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- We concentrated on the defect channel and prove universality at large transverse spin s .
- In the next talk Yannick will tell you about the bulk channel!

Thank you!