

# preliminary request: Search for collective effects in electron-proton collisions with ZEUS

Jaap Onderwaater

Ilya Selyuzhenkov

Achim Geiser

Silvia Masciocchi

Stefan Floerchinger

07.05.2018

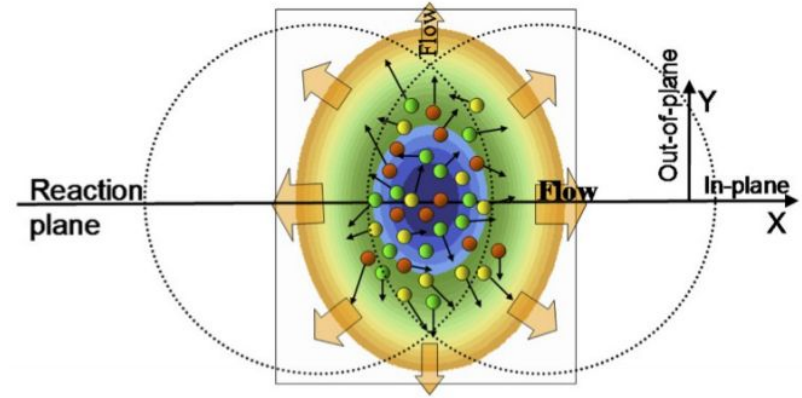
# Collectivity and related anisotropy in heavy-ion collisions

Response of matter produced in the heavy-ion collision to the geometry of the initial state.

Produced particles receive a stronger boost along the short axis of the geometry wrt to the long axis (see ellipse on the right)

The amplitude ( $v_n$ ) of the resulting anisotropy is quantified with a Fourier decomposition:

$$\frac{dN}{d(\varphi - \Psi_R)} = \frac{N_0}{2\pi} \left( 1 + 2 \sum_n v_n \cos[n(\varphi - \Psi_R)] \right)$$



# Analysis techniques

We report a measurement of 2-particle correlations:

$$c_n\{2\} = \langle\langle 2 \rangle\rangle \equiv \left\langle\left\langle e^{in(\phi_1 - \phi_2)} \right\rangle\right\rangle \qquad v_n\{2\} = \sqrt{c_n\{2\}}$$

The inner brackets denote the average in a single event, the outer brackets the average over all events.

The correlation are studied as a function of

- event multiplicity
- separation of particles in pseudorapidity
- particle's transverse momentum

# Different mechanisms resulting in 2-particle correlations

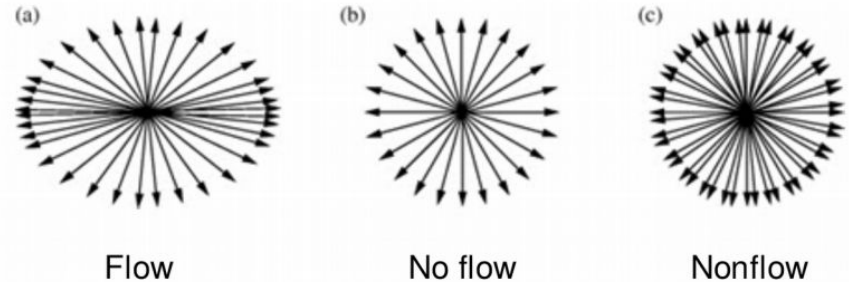
Multiple mechanisms contribute to (multi)particle correlations, from the initial state to response to the initial geometry.

Correlations contain flow, flow fluctuations and nonflow.

$$\langle\langle e^{in(\phi_a - \phi_b)} \rangle\rangle = \langle v_n^2 \rangle + \delta_n$$

Flow fluctuations:  $\sigma_{vn}^2 = \langle v_n^2 \rangle - \langle v_n \rangle^2$

Nonflow:  $\delta_n$ : resonances, jets, decays, momentum conservation



Suppression of  $\delta_n$  (suppression of few particle correlated clusters):

- High multiplicity  $\delta_2 \sim 1/M$
- Pseudo-rapidity gap (particles from jets and decays are mostly closeby in  $\eta$ )

# Analyzed data sets (common ntuples)

	Trigger events (x10 <sup>6</sup> )	
Period	All events in tree	DIS
03p	3.7	0.24
04p	47	4.6
05e	132	17
06e	44	7.0
06p	87	12
07p	41	5.4
All	355	45.8

DIS: Sinistra electron,  $Q^2 > 5 \text{ GeV}$ ,  $E_e > 10 \text{ GeV}$ ,  $47 < E-p_z < 69 \text{ GeV}$ ,  $\theta_e > 1$ ,  $e_p > 0.9$ , exclusion of some problematic detector areas

# DIS selection

Sinistra electron:

$$Q^2 > 5 \text{ GeV}$$

$$E_e > 10 \text{ GeV}$$

$$\theta_e > 1$$

Sinistra probability:

$$e_p > 0.9$$

$$47 < E - p_z < 69 \text{ GeV}$$

Exclusion of some problematic detector areas

Radius cut :

```
if(sqrt( TMath::Power((Sipos[0][0] + 1),2.0) + TMath::Power(Sipos[0][1],2.0) )  
< 15) isDIS = false;
```

Chimney cut (RCAL):

```
if ( (Sipos[0][0] > -10) && (Sipos[0][0] < 10 ) &&  
    (Sipos[0][1] > 110) && (Sipos[0][2] < -141.)){  
    isDIS = false;  
}
```

//HES fiducal/CAL crack cut (RCAL)

```
if ( ( ( (Sipos[0][0] > 5) && (Sipos[0][0] < 11) && (Sipos[0][1] > 0.) )  
    || ( (Sipos[0][0] > -15) && (Sipos[0][0] < -9) && (Sipos[0][1] < 0.) ) )  
    && (Sipos[0][2] < -141. ) ){  
    isDIS = false;  
    if ( (Sinrsl[0] < 3) ||  
        (Sitrkp[0] < 0.3 * Sicalene[0]) ||  
        (run < 45000)){  
        isDIS = false;  
    }  
}
```

# DIS trigger bits

## HPP20

```
if (((ibits(Tltw[3],16+0,1) > 0) && (ibits(Sltw[5],7-1,1) > 0))  
(ibits(Tltw[3],16+1,1) > 0)  
((ibits(Tltw[3],16+2,1) > 0) && (ibits(Sltw[5],7-1,1) > 0))  
((ibits(Tltw[3],16+3,1) > 0) && (ibits(Sltw[5],7-1,1) > 0))  
((ibits(Tltw[3],16+4,1) > 0) && (ibits(Sltw[5],7-1,1) > 0))  
(ibits(Tltw[3],16+5,1) > 0)  
((ibits(Tltw[3],16+10,1) > 0) && (ibits(Sltw[5],7-1,1) > 0))  
(ibits(Tltw[2],16+0,1) > 0)  
(ibits(Tltw[2],16+1,1) > 0)  
(ibits(Tltw[2],16+2,1) > 0)  
(ibits(Tltw[2],16+8,1) > 0)  
(ibits(Tltw[13],16+0,1) > 0)  
(ibits(Tltw[11],16+14,1) > 0)  
(ibits(Tltw[11],16+3,1) > 0))
```

TLT HFL2 (inclusive mesons in DIS) and HFL10 (e in DIS)

```
if ( (ibits(Tltw[9],16+1,1) > 0) || (ibits(Tltw[9],16+9,1) > 0) )
```

TLT HFL6 (dijets in DIS)

```
if (ibits(Tltw[9],16+5,1) > 0)
```

# Event selection

- DIS selection
- $-30 < \text{vertex } Z < 30 \text{ cm}$
- Fraction of tracks associated to event vertex  $> 0.1$
- $N_{\text{vtx}}$  tracks  $> 0$
- Event vertex from beam spot ( $R_{xy}$ )  $< 0.5$



# Track selection

- $0.1 < p_T < 5 \text{ GeV}/c$
- $-1.5 < \eta < 2.0$
- Tracks constrained to the vertex (`orange.Trk_prim_vtx = true`)
- Exclude scattered electron (`orange.Trk_id[itrack] != orange.Sitrknr[0]`)
- $\text{Trk\_Imppar} < 0.5 \text{ cm}$
- $\text{MVD hits} > 0$  ( $\text{MVD hits} = \text{orange.Trk\_nbr}[itrack] + \text{orange.Trk\_nbz}[itrack] + \text{orange.Trk\_nwu}[itrack] + \text{orange.Trk\_nwv}[itrack];$ )

# Simulation selection

## Event level

- $Q^2 > 5 \text{ GeV}^2$  , in code : `orange.Mc_q2 > 5`
- $47 < E - P_z < 69 \text{ GeV}$ , in code : `47 < (orange.Mc_esum - orange.Mc_ez) < 69`
- Final state lepton energy  $E > 10 \text{ GeV}$ , in code: `orange.Mc_pfsl[3] > 10`
- Final state lepton  $\theta > 1$

## Track level

- $0.1 < p_T < 5 \text{ GeV}/c$
- $-1.5 < \eta < 2.0$

For calculation of tracking efficiency, the same selection is applied on data and MC on the reconstruction level.

# Correcting for track reconstruction effects

Particles reconstruction efficiency as a function of  $p_T$ ,  $\eta$ ,  $\phi$ , *charge* and *event multiplicity* is considered.

Particle weights are extracted in two steps:

1.  $p_T$ - $\eta$ -*charge* efficiency is calculated by comparing generated and reconstructed yields in simulation
2.  $\phi$  weights are extracted from data, after filling  $\phi$ - $\eta$ -*charge*-*event multiplicity* maps with the weights from step 1

The product of 1. and 2. gives the track weight. Weights are calculated separately for each dataset.

The 2-particle correlation is modified to include weights:

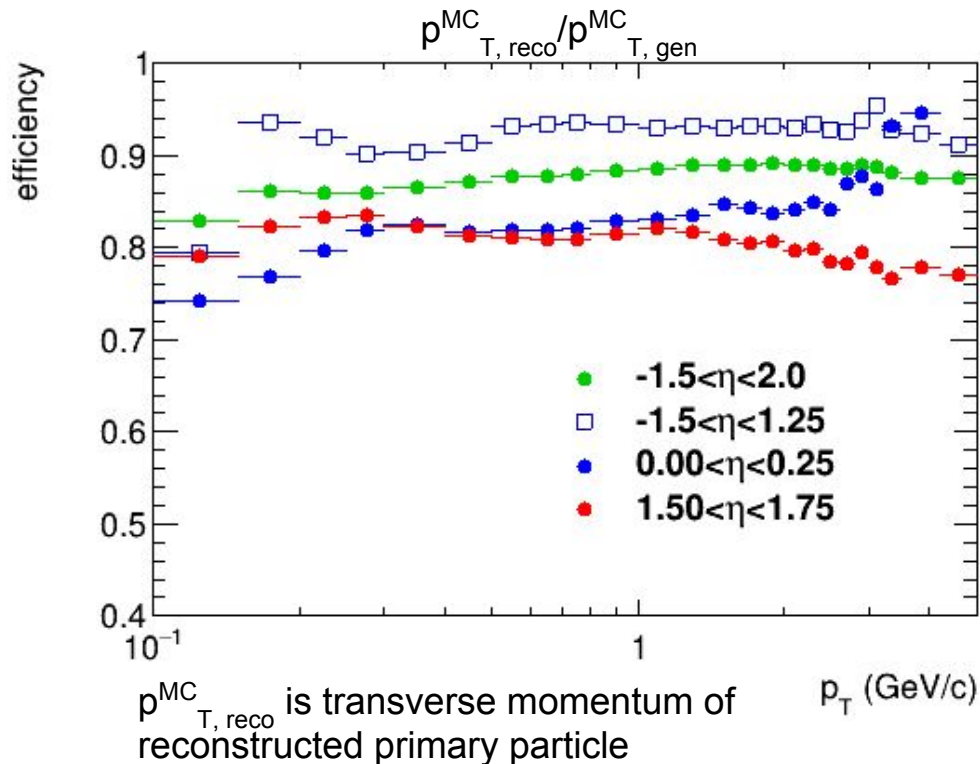
$$\langle c_n \rangle = \sum w_i w_j \cos(n\phi_i^a - n\phi_j^b) / \sum w_i w_j$$

# Determining $p_T$ - $\eta$ efficiency

## Charged primary particle:

- Charged particle with lifetime  $\tau > 1$  cm/c
- Particles with parents with lifetime  $\tau > 1$  cm/c in the decay chain are rejected

Specie	Width $\Gamma$ (GeV)	Mean proper lifetime $\tau$	
		(ps)	(cm/c)
$p^+$	0	$\infty$	$\infty$
$\gamma$	0	$\infty$	$\infty$
$K^0$	0	$\infty$	$\infty$
$e^-$	0	$\infty$	$\infty$
$n$	$7.478 \times 10^{-28}$	$8.861 \times 10^{+14}$	$2.656 \times 10^{+13}$
$\mu^-$	$2.996 \times 10^{-19}$	$2.212 \times 10^{+06}$	$6.63 \times 10^{+04}$
$K_L^0$	$1.287 \times 10^{-17}$	$5.148 \times 10^{+04}$	1543
$\pi^+$	$2.528 \times 10^{-17}$	$2.621 \times 10^{+04}$	785.7
$K^+$	$5.317 \times 10^{-17}$	$1.246 \times 10^{+04}$	373.6
$\Xi^0$	$2.27 \times 10^{-15}$	291.9	8.751
$\Lambda$	$2.501 \times 10^{-15}$	264.9	7.943
$\Xi^-$	$4.02 \times 10^{-15}$	164.8	4.941
$\Sigma^-$	$4.45 \times 10^{-15}$	148.9	4.464
$K_S^0$	$7.351 \times 10^{-15}$	90.14	2.702
$\Omega^-$	$8.071 \times 10^{-15}$	82.1	2.461
$\Sigma^+$	$8.209 \times 10^{-15}$	80.72	2.42



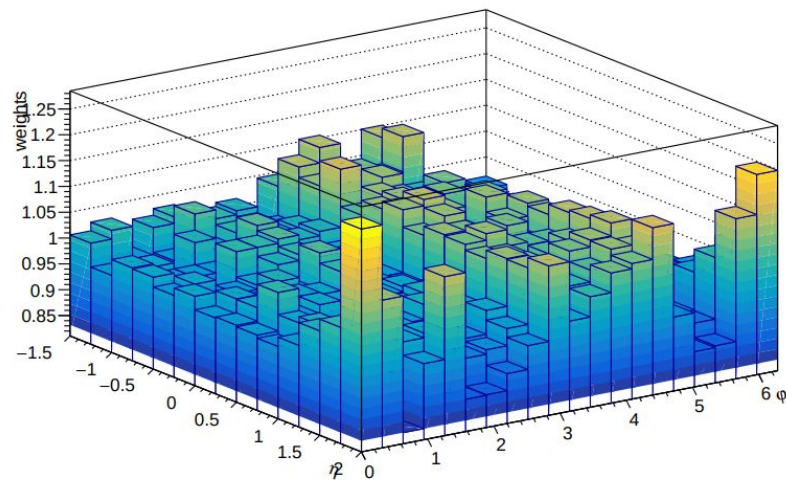
The efficiency is calculated for DIS events, same event selection for generated and reconstructed particles

# Determining $\phi$ -weights from data

Particle yields are measured in  $\eta$ - $\phi$ -charge- $M$  bins, after weighting with acquired  $p_T$ - $\eta$ -charge weights in the previous slides.

In each  $\eta$ -charge- $M$  slice, weights are calculated to make  $\phi$  uniform while maintaining the integral in the slice.

05e, M=10, positive



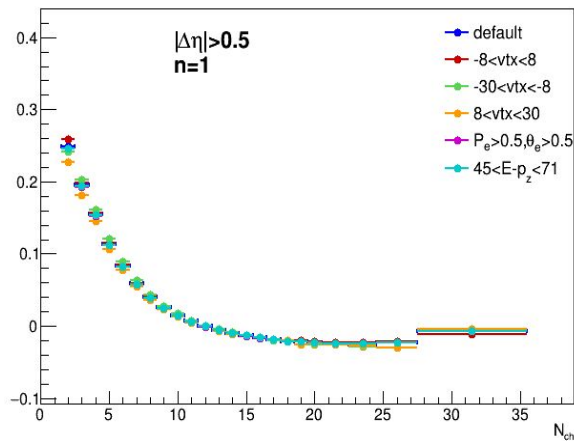
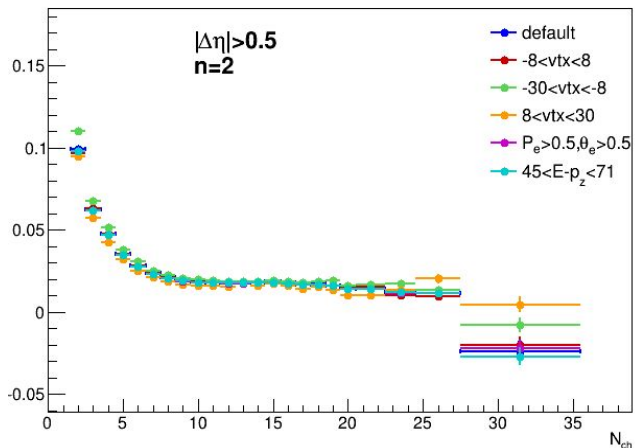
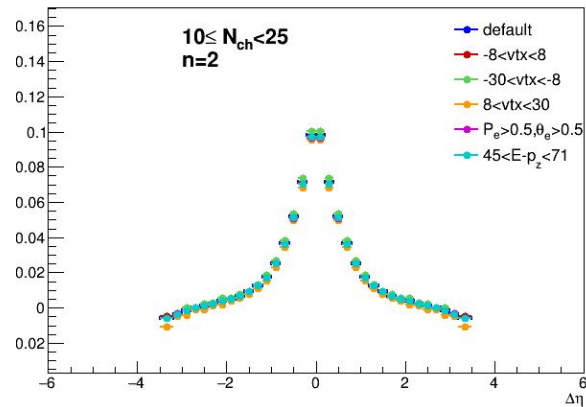
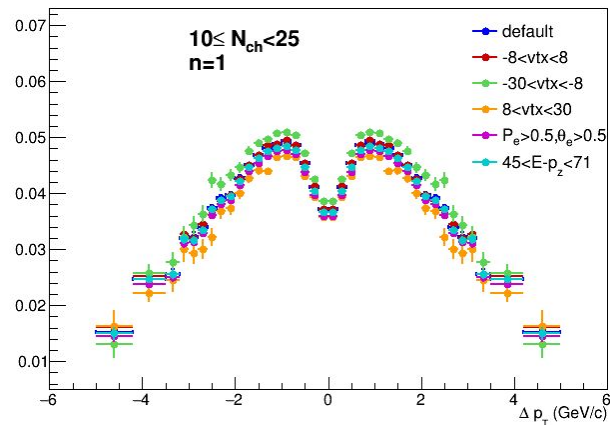
# Systematic uncertainties

# Study of systematics

Class	Default	Variation
DIS event selection	$47 < E-p_z < 69$	$45 < E-p_z < 71$
	$\theta_e > 1.0$ $P_e > 0.9$	$\theta_e > 0.5$ $P_e > 0.8$
	Chimney cut, radius cut, CAL crack cut	
Event quality selection	$-30 < Z_{\text{vtx}} < 30 \text{ cm}$	$-30 < Z_{\text{vtx}} < -8 \text{ cm}$ $-8 < Z_{\text{vtx}} < 8 \text{ cm}$ $8 < Z_{\text{vtx}} < 30 \text{ cm}$
MC closure		Check generated vs reconstructed correlations
Trigger efficiency		Check generated correlations for MC events vs generated for reconstructed events
Consistency of periods	Sum of all periods	Periods individually

# Event selection

The variation from the event selection is relatively minor. Variations are added to the systematic uncertainty (z vertex variations are taken as a group with largest deviation).



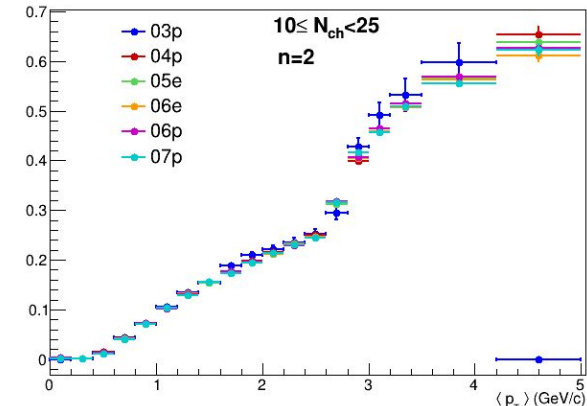
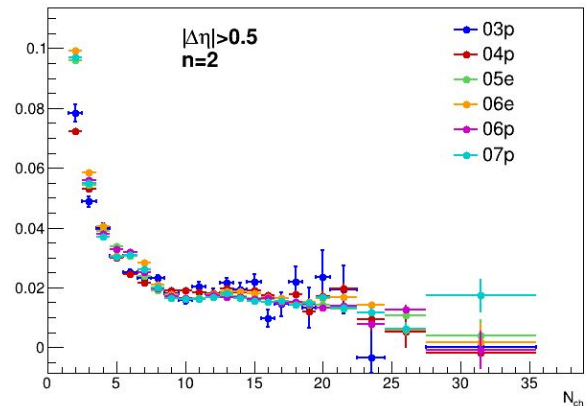
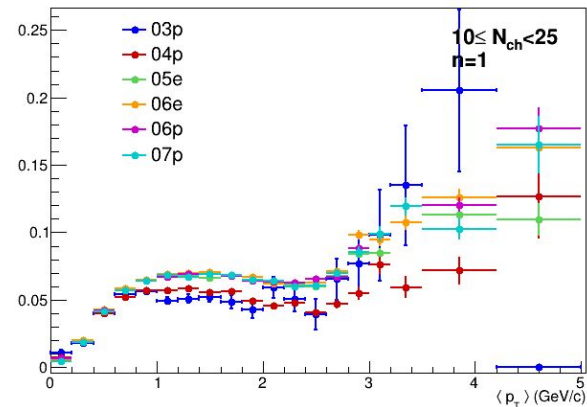
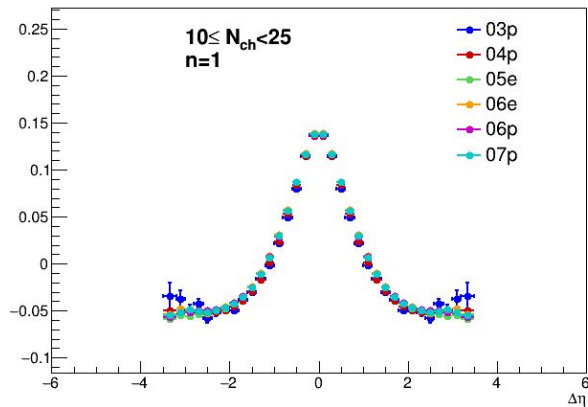


# Period consistency

The results from different periods should be consistent.

Some deviation for 03p/04p is observed, as for very low multiplicity correlations.

Deviations are added to the systematic uncertainty, excluding 03p due to low significance.



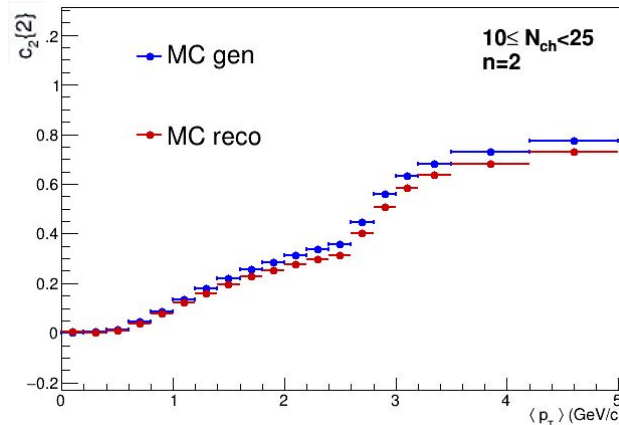
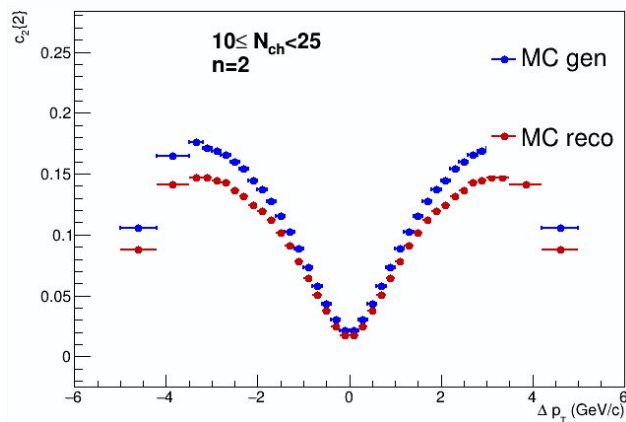
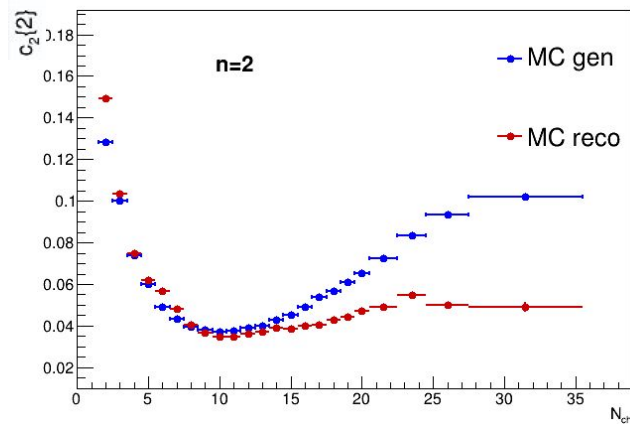
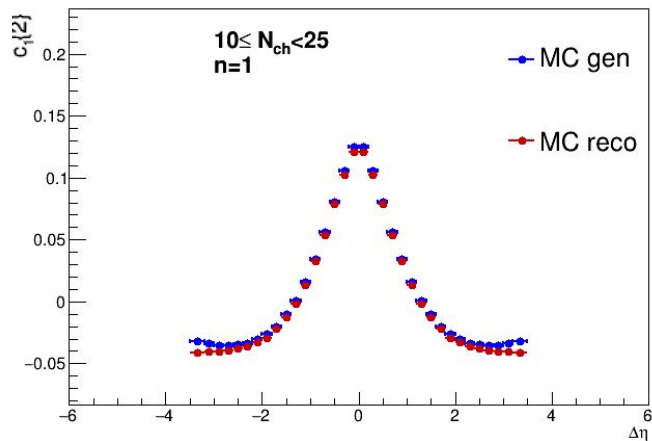
# MC closure

The correlations on true level are compared to the reconstructed correlations.

If the trigger is efficient, contamination is low and corrections for detector acceptance are effective, the correlations should match.

In several places significant deviations are observed, as is visible on the right, while others are smaller. The discrepancy is added as a systematic uncertainty on the data.

For final results this has to be more carefully studied.



# DIS/secondaries

MCDIS:

True level cuts on the electrons

DIS:

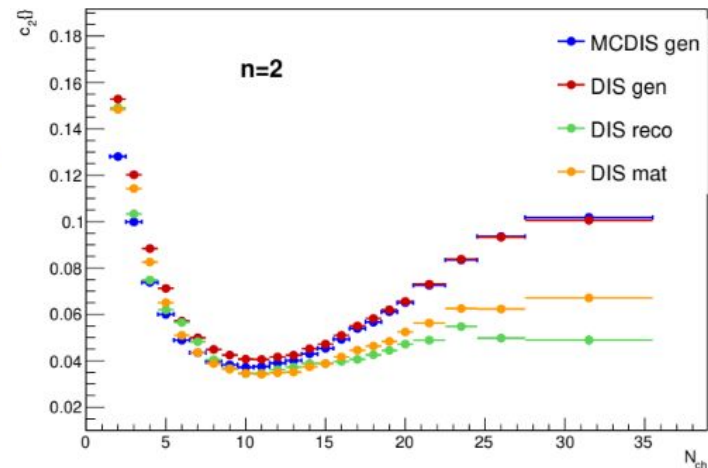
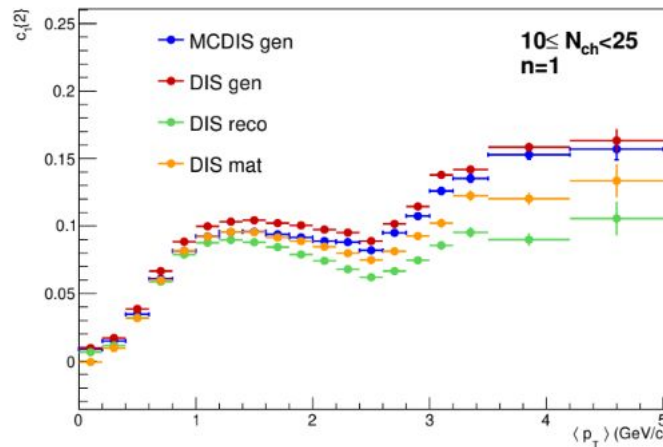
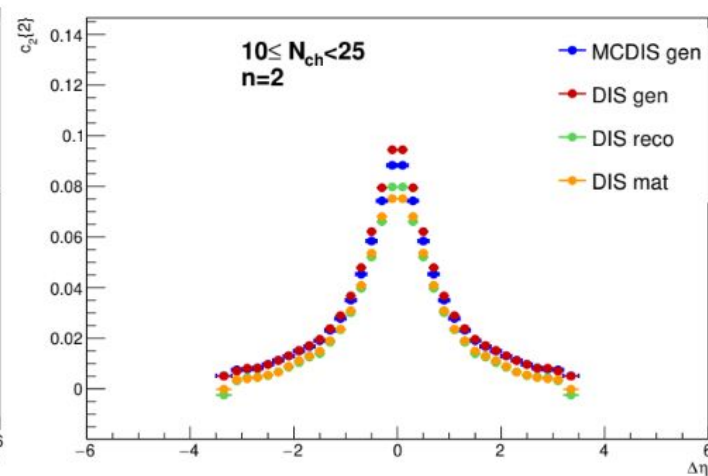
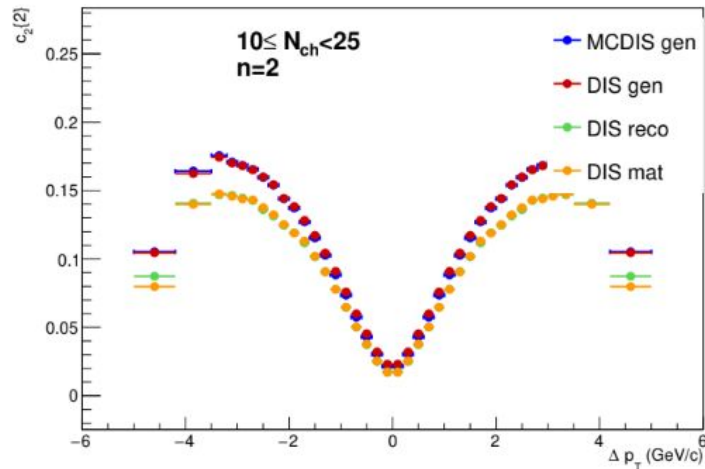
In addition DIS selection on the reconstruction parameters (and event quality selection)

DIS reco:

Measured as in data

DIS mat:

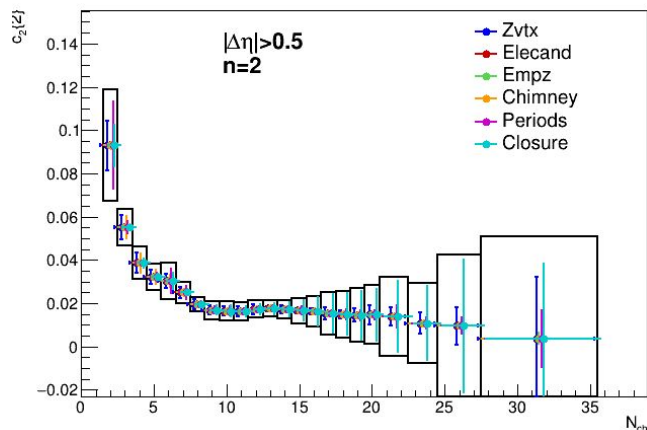
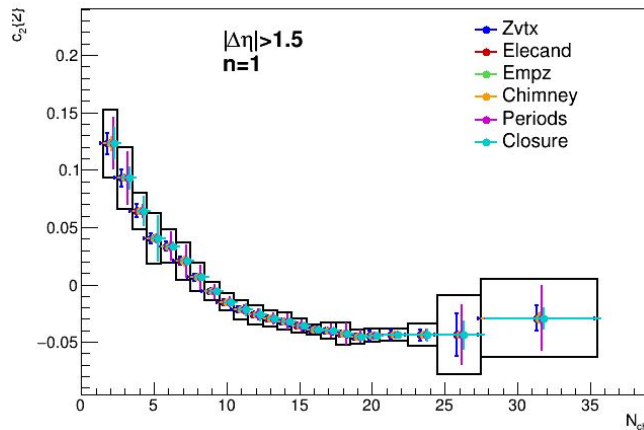
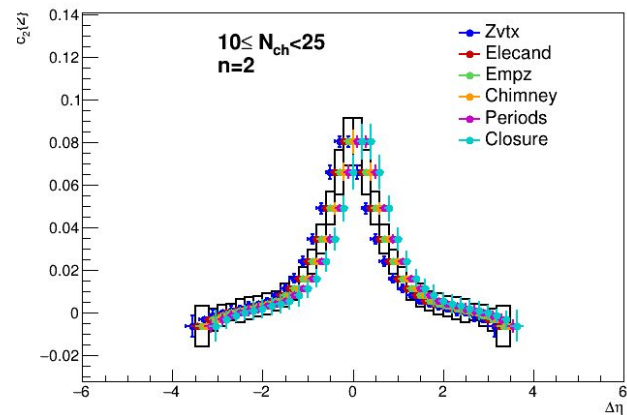
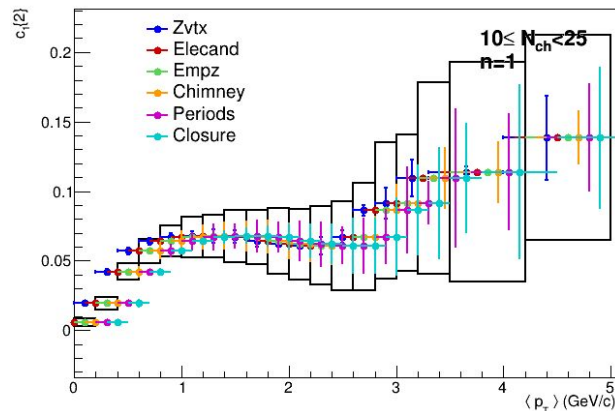
Measured as in data but only with tracks matched to primary tracks



# Overview

A compilation with the contributions to the systematic uncertainty. The black boxes show the total uncertainty. Points are shifted for clarity.

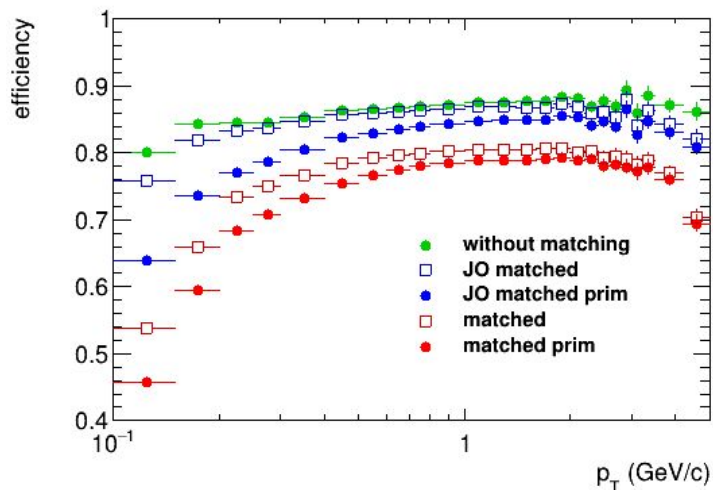
It is clear that the MC closure and DIS event selection effects are largest, although not always in the same places.



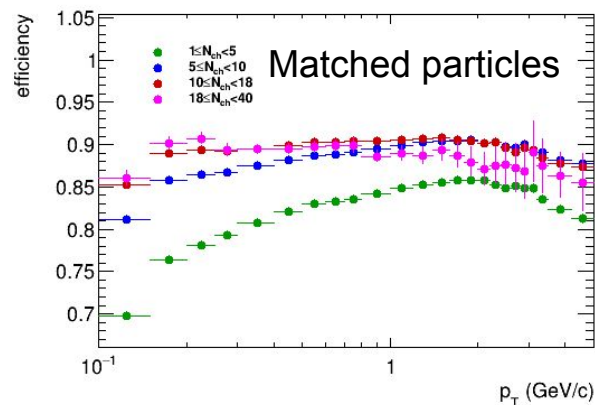
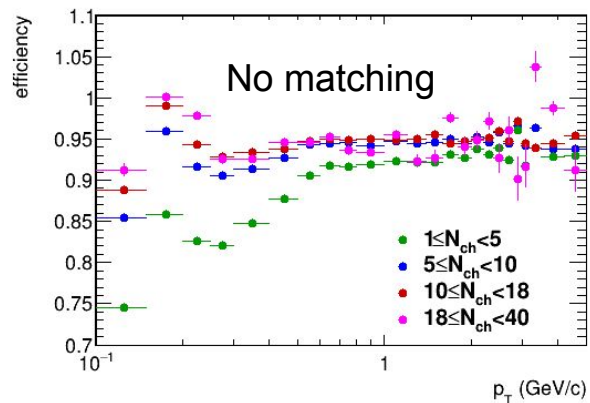
# Variation of tracking efficiency

I intended to run with this efficiency as a check but unfortunately something went wrong

Difference between open squares and solid circles is the rejection of secondary particles



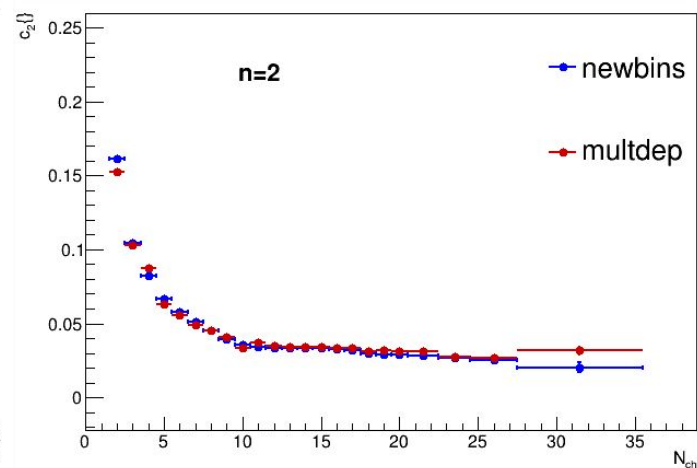
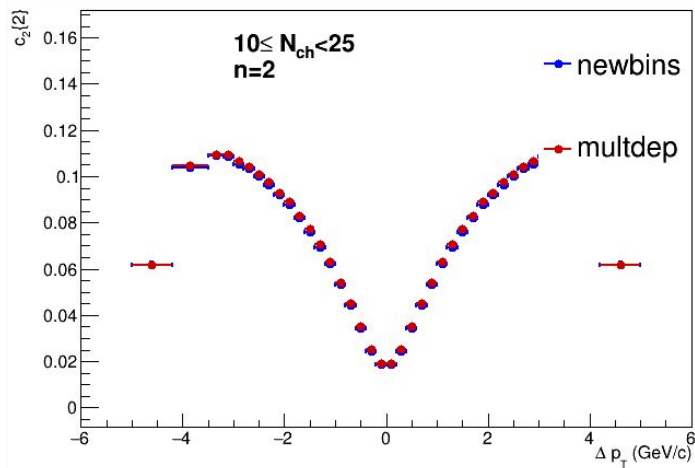
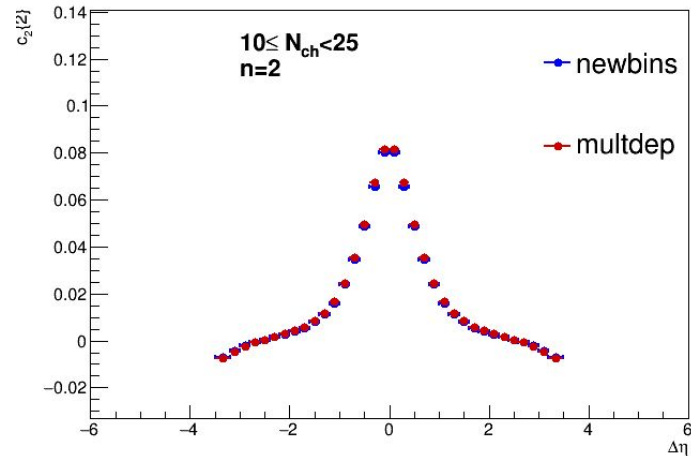
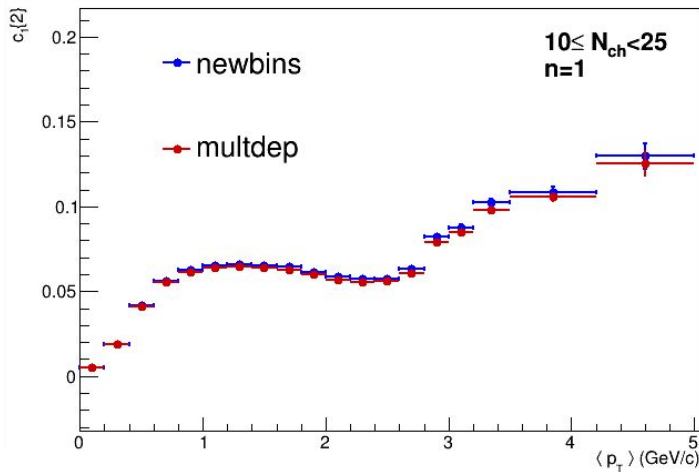
# Multiplicity dependent $p_T$ -correction



Reran corrections with multiplicity dependence

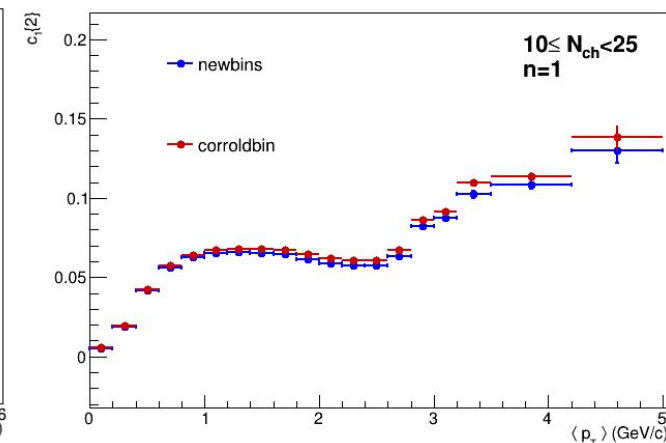
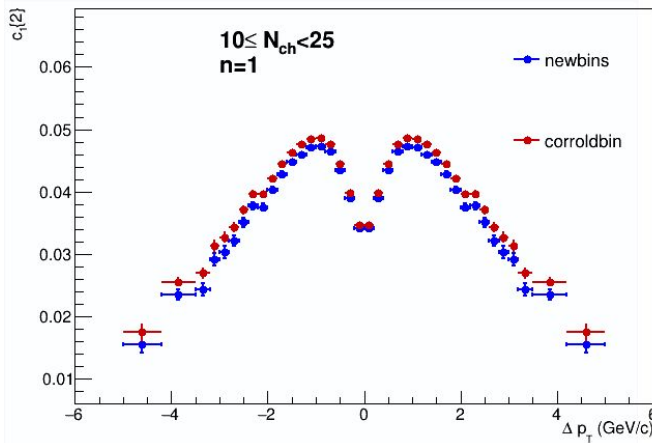
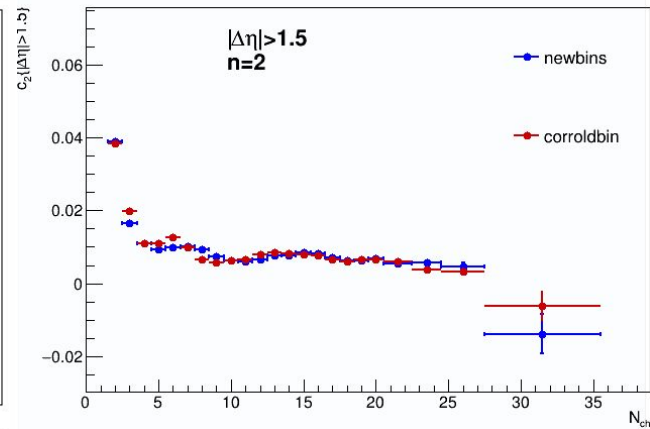
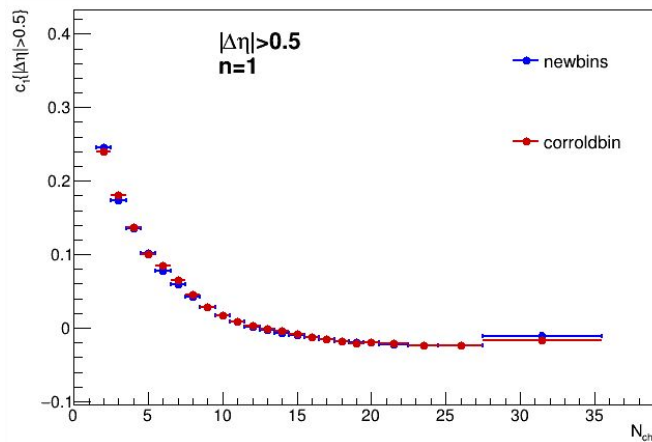
# Mult-dependence

The effect of corrections using multiplicity dependent weights is small



# Correction binning

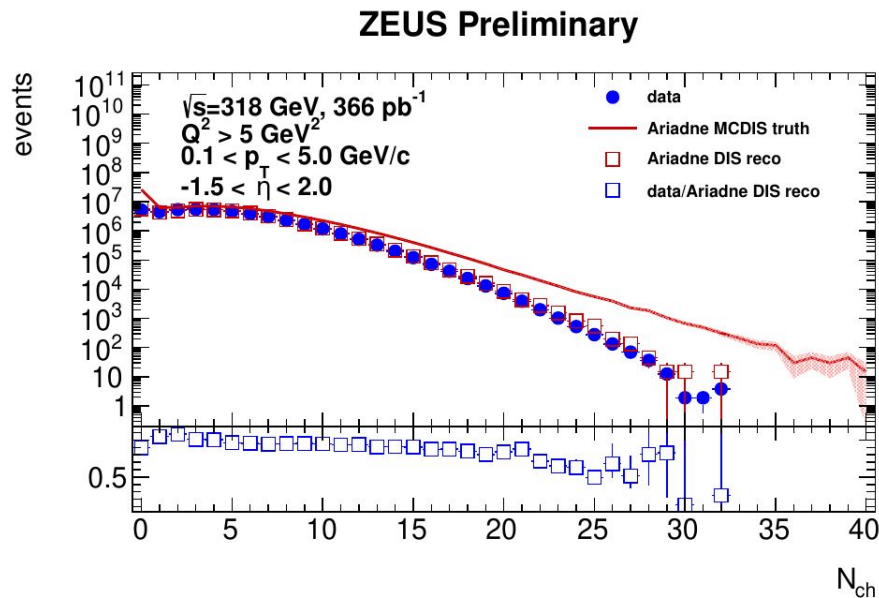
Using finer binning at low  $p_T$  (50 MeV/c instead of 100 MeV/c for  $p_T < 300$  MeV/c) and coarser  $\eta$  bins (0.25 instead of 0.2) gives largely consistent results.





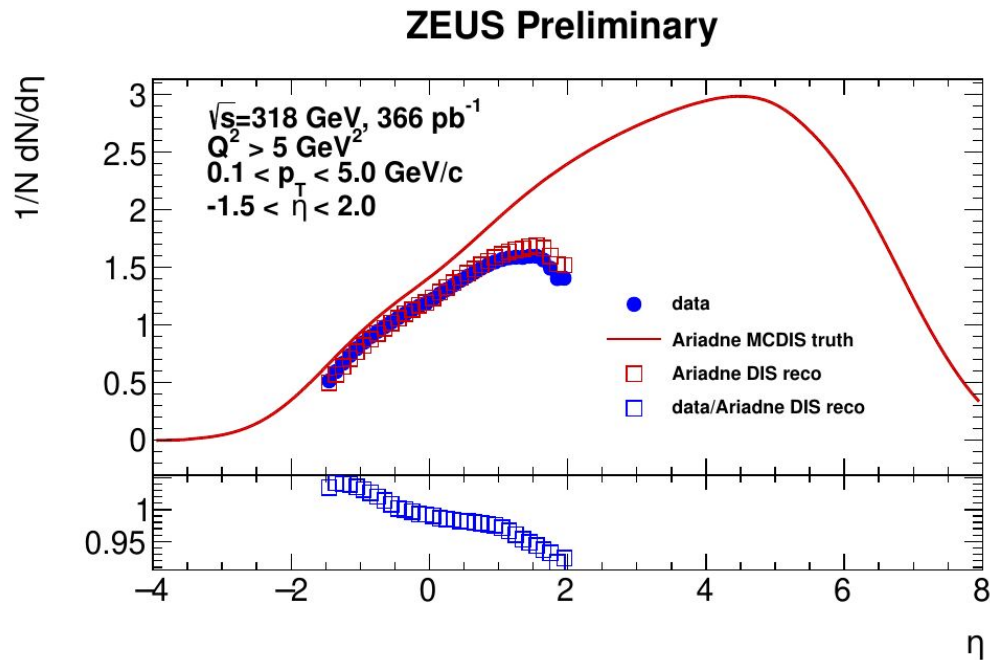
# Control figures

# Multiplicity distribution



Open squares show the agreement with data (it is normalized to data for  $N_{ch} > 2$ )  
Red line is the true  $N_{ch}$ . It is normalized relative to Ariadne DIS reco.

# Eta distribution

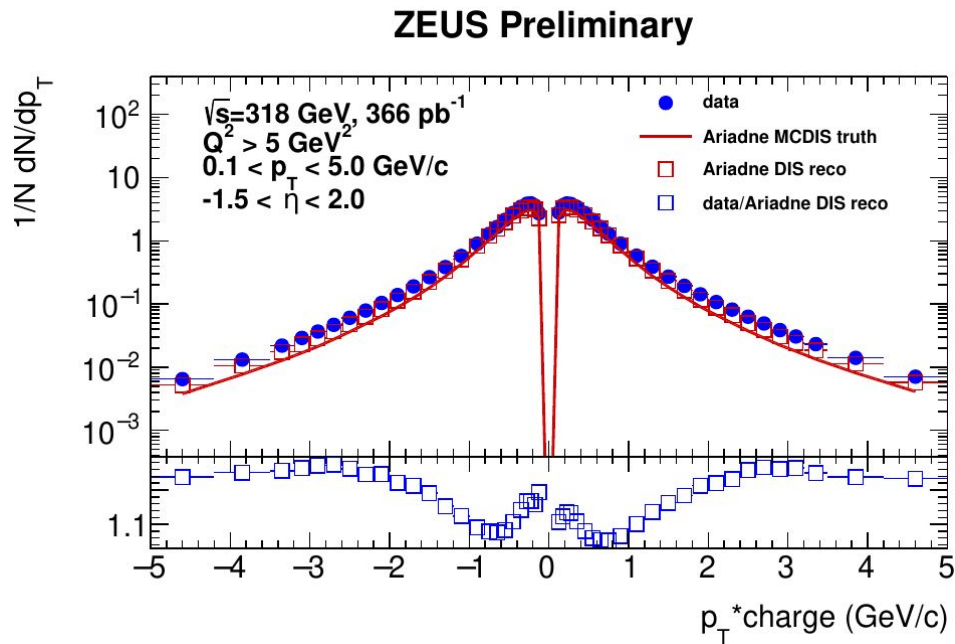


Open squares show the agreement with data.

Red line is the true  $N_{ch}$ .

All histograms are scaled by the number of events in which the tracks were measured.

# $p_T$ distribution



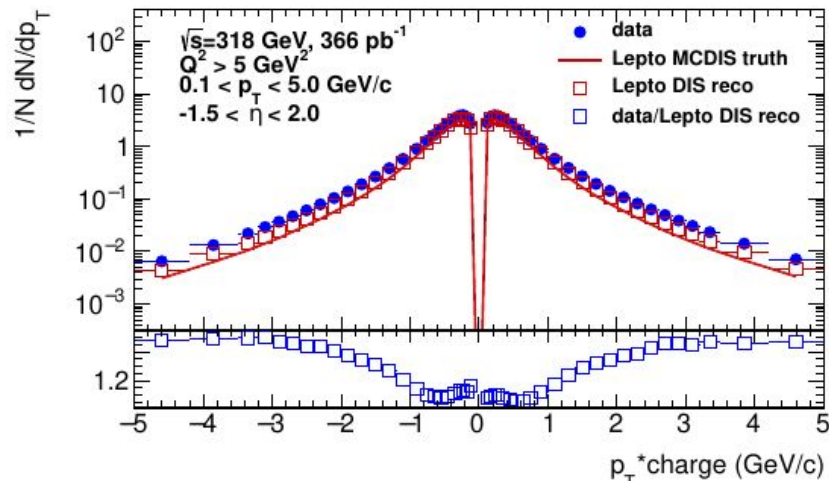
Open squares show the agreement with data.

Red line is the true  $N_{ch}$ .

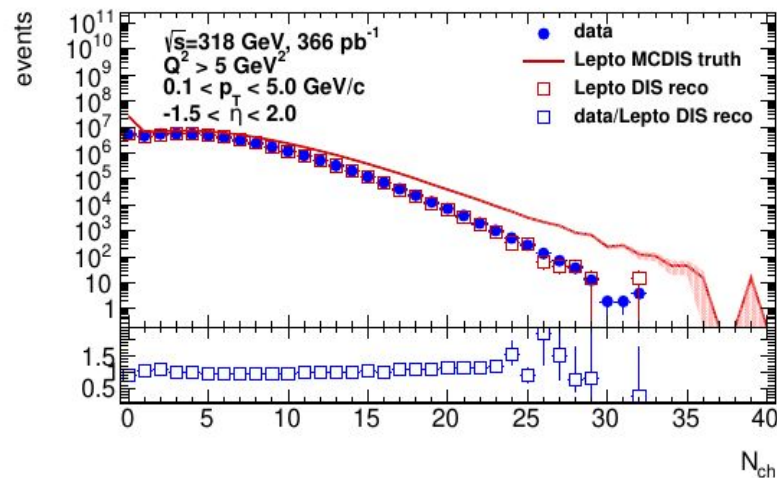
All are histograms are scaled by the number of events in which the tracks were measured.  
(remove 0-100 MeV/c for theory)

# Same for Lepto

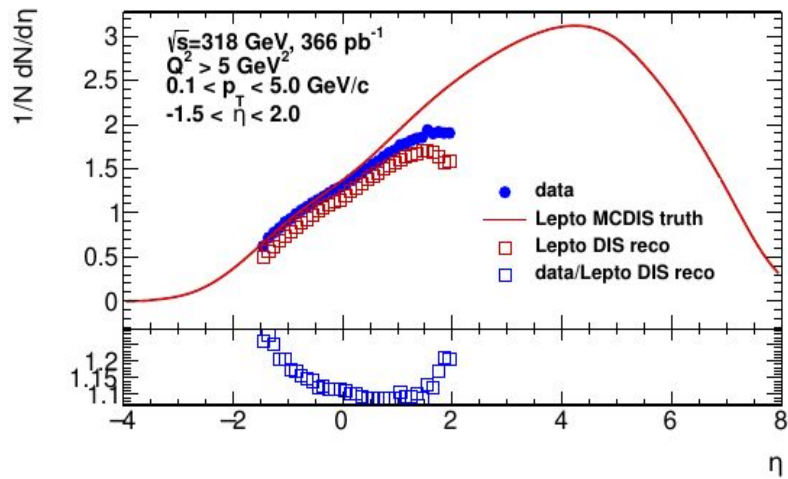
ZEUS Preliminary



ZEUS Preliminary

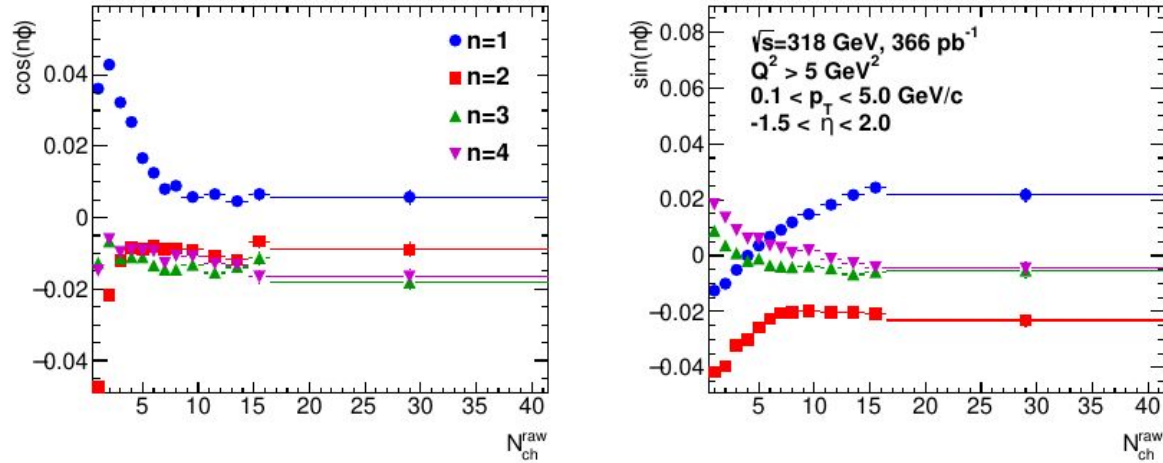


ZEUS Preliminary



# Magnitude of non-uniform acceptance in the correlation

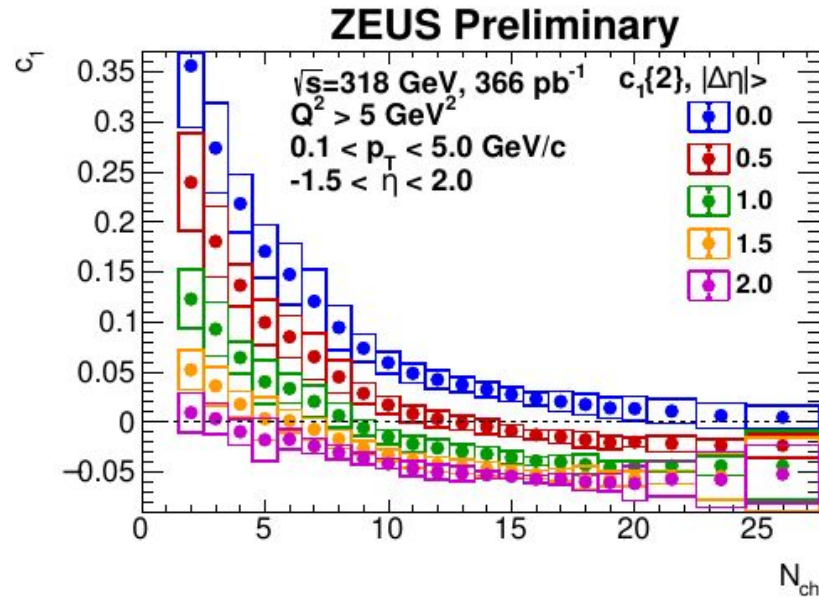
ZEUS Preliminary



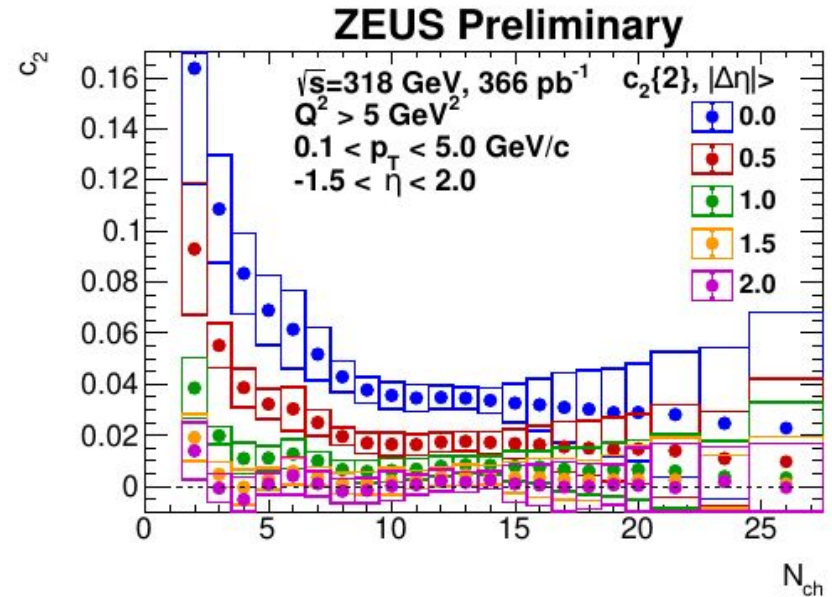
Average of the sine/cosine of the angles show the magnitude of the acceptance correction that is applied, which is of the order  $\langle \cos(n\phi)^2 \rangle + \langle \sin(n\phi)^2 \rangle \approx 10^{-4}$

# Results

# Multiplicity dependent correlations with pseudorapidity gaps (1st and 2nd harmonic)



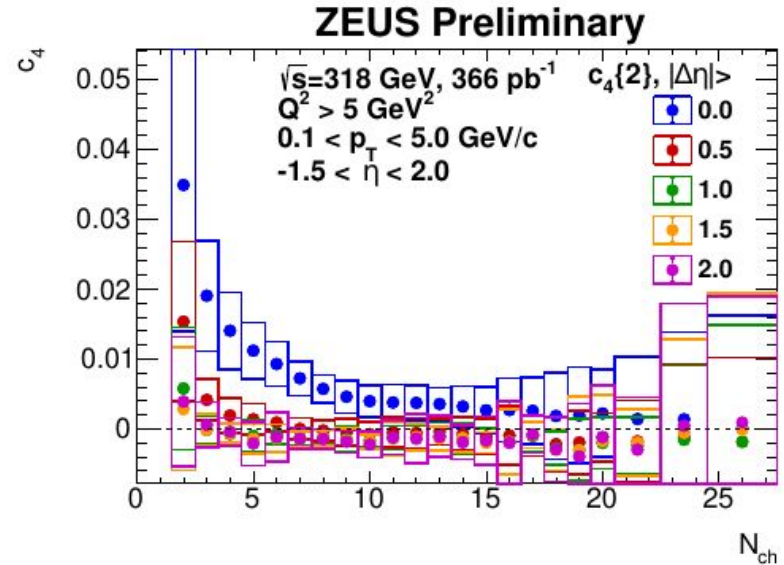
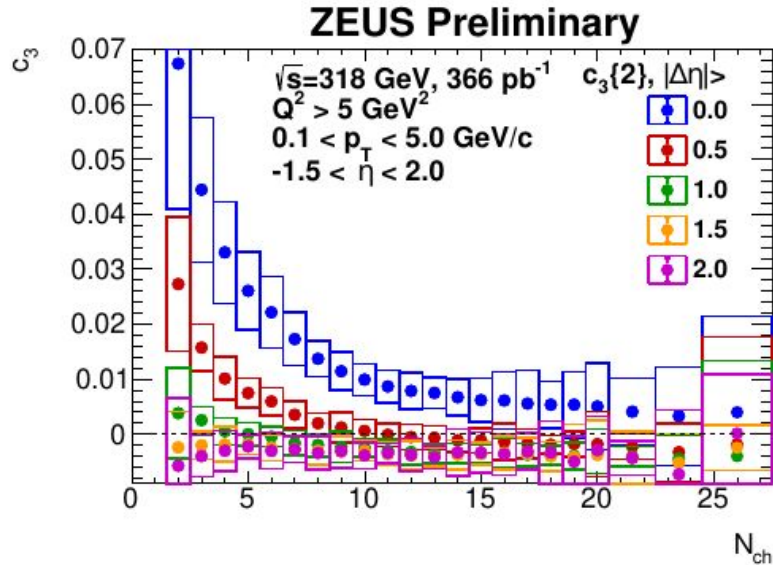
Zoomed in a bit, cleaned up legend



Increasing pseudo rapidity separation suppresses correlations. Consistent with 0 for  $|\Delta\eta| > 1.0$

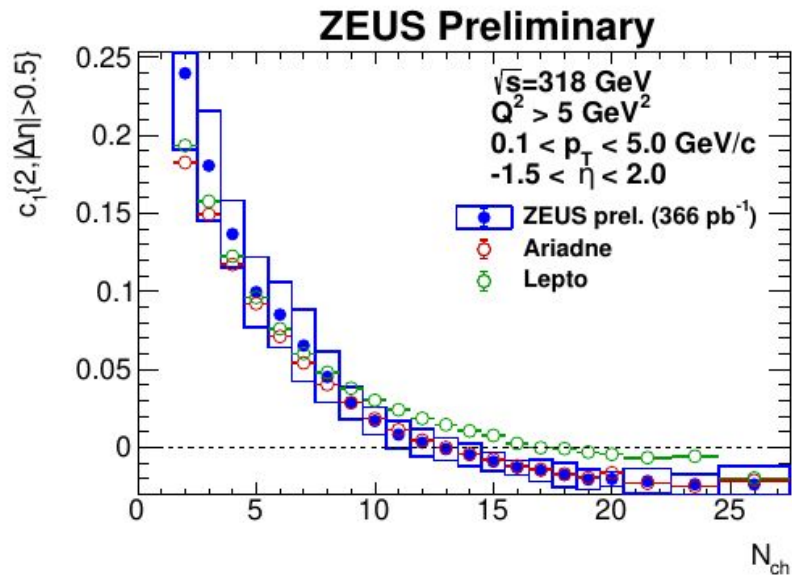


# Multiplicity dependent correlations with pseudorapidity gaps (3st and 4nd harmonic)

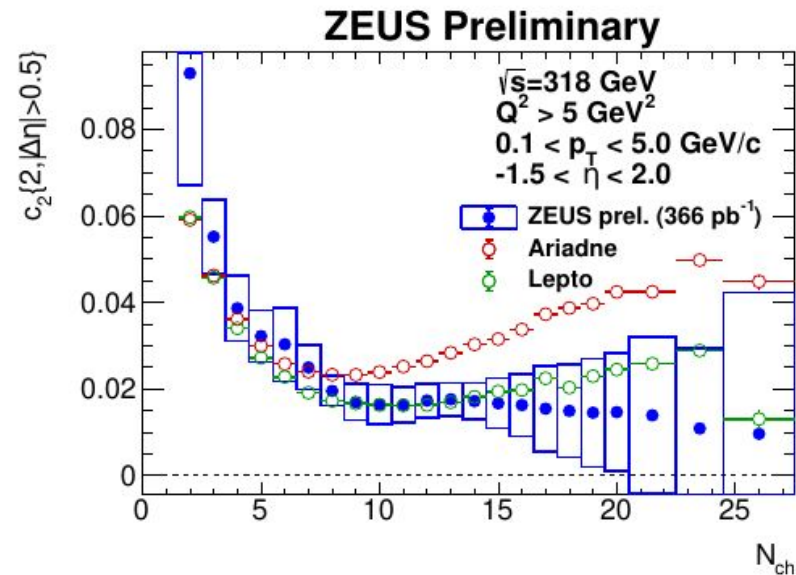


Increasing pseudo rapidity separation suppresses correlations. Consistent with 0 for  $|\Delta\eta| > 0.5$

# Multiplicity dependent correlations with simulations



First harmonic is well described by Ariadne



Second harmonic favors Lepto.

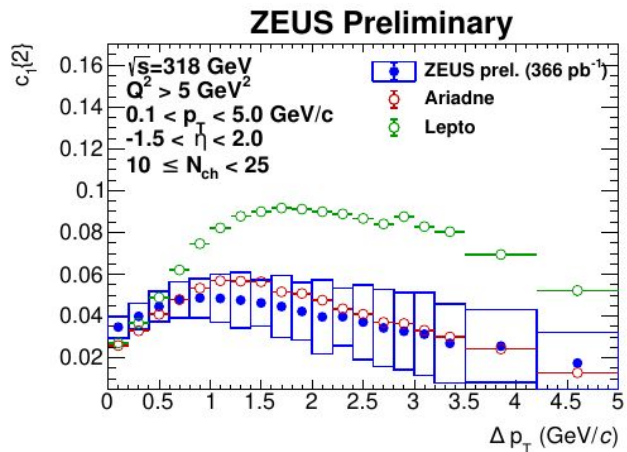
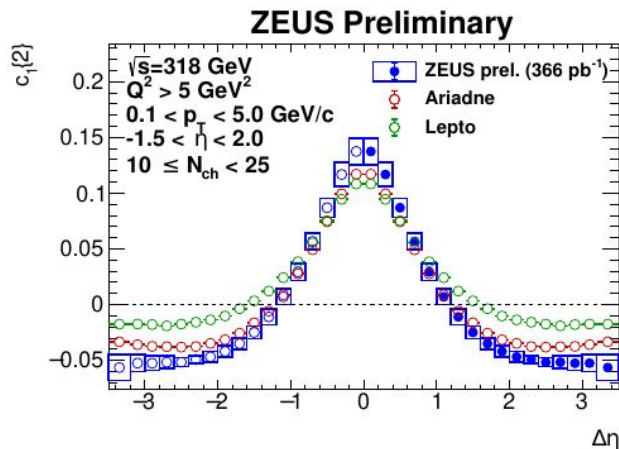
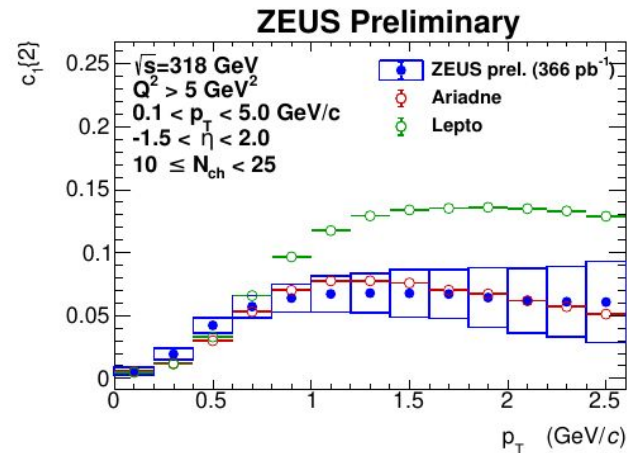
# Differential $c_1\{2\}$ comparisons with MC

First harmonic has good agreement with Ariadne simulation.

Restricted mean  $p_T \rightarrow$  no artefact from upper  $p_T$  cut, better visibility at low  $p_T$

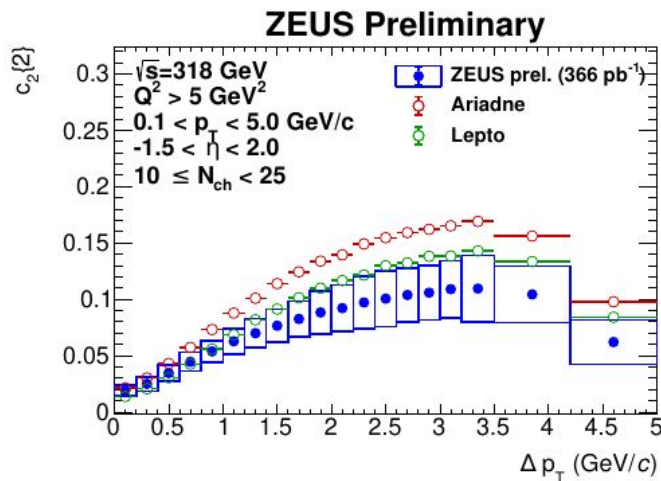
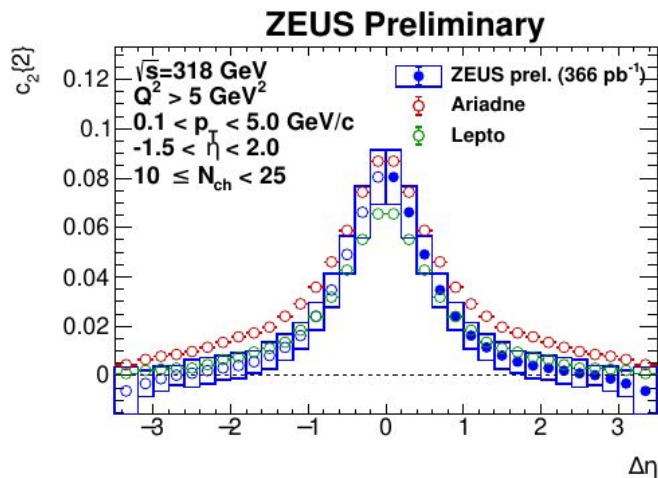
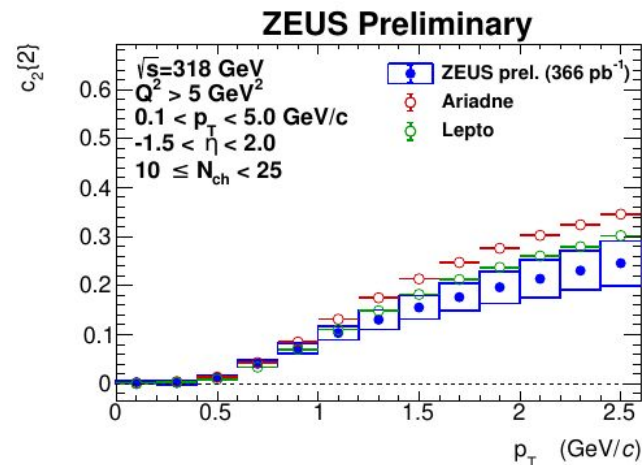
For  $\Delta\eta$  : open circles below  $\Delta\eta < 0$

Why: reflection of positive  $\Delta\eta$ , no independent datapoints (style used more often in heavy ion field)



# Differential $c_2\{2\}$ comparisons with MC

Second harmonic has better agreement with Lepto, especially for for larger pseudorapidity separation



# Physics messages

- Reported measurement for the correlations for different harmonics, and as a function of multiplicity, pair pseudorapidity, pair transverse momentum, pair  $\Delta p_T$
- Correlations as a function of rapidity separation approach 0 for large  $\Delta\eta$  except for  $c_1\{2\}$ , where it changes sign towards negative, which is a signature of momentum conservation.
- New data from ZEUS adds new information to the ongoing efforts at LHC and RHIC for the search for collective effects in high multiplicity events for small collision systems.
- Comparisons to different Monte Carlo generators tuned to Hera data are able to reproduce overall features of the multiplicity dependence of the correlations.

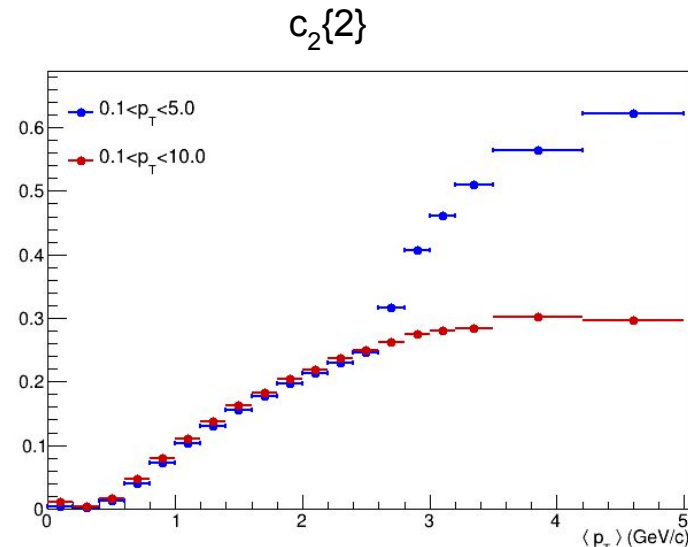
# Backup

# Kink in $c_n$ vs $\langle p_T \rangle$

It was asked why there is a kink observed in the correlation vs pair mean transverse momentum.

The momentum range of the selected particles has the effect that for  $\langle p_T \rangle > 2.5$  GeV/c, both particles in the pair need high momentum.

If the upper momentum is extended to 10 GeV/c, the kink disappears.



# MC track matching



# Track matching

Matching in the orange tree with  $\text{Mcmatquality}==1$  is fairly inefficient, ~10% of reconstructed tracks can't find a match to a generated particle.

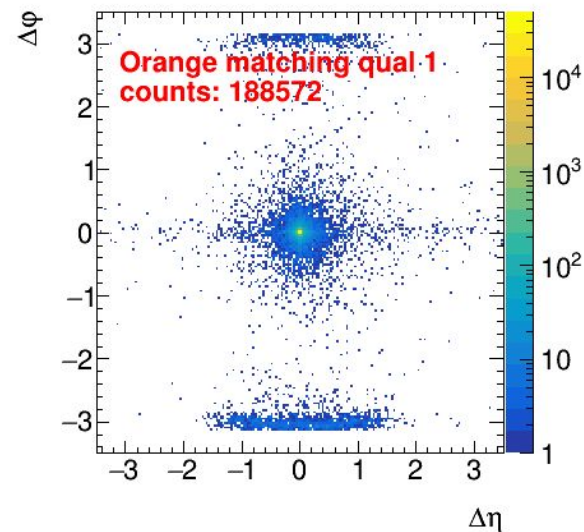
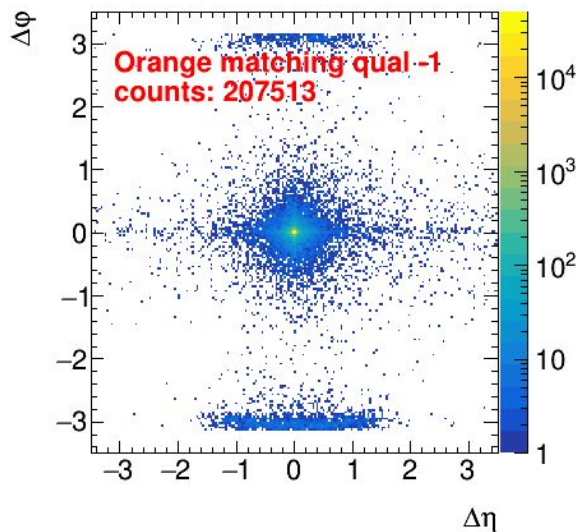
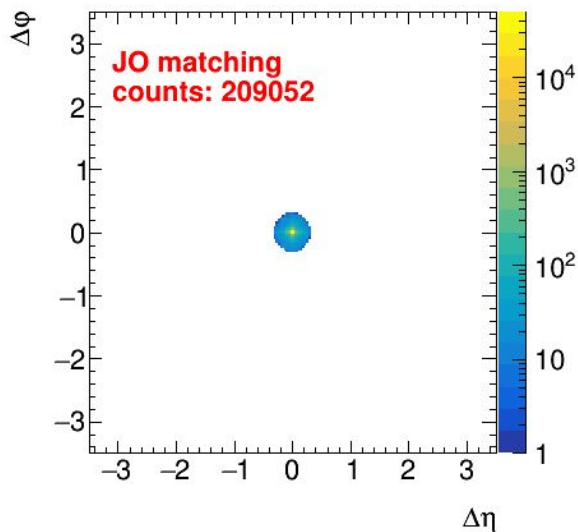
Attempt at a new matching algorithm:

Look for the closest match in  $(\eta^T - \eta^R)^2 + (\phi^T - \phi^R)^2$ ,

with requirement  $(p_T^T - p_T^R)/p_T^T < 0.3$

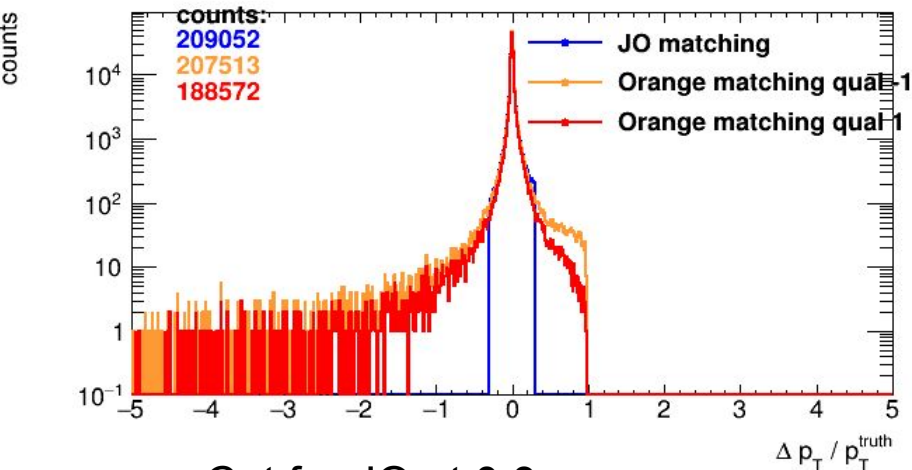
# MC track matching $\Delta\eta\Delta\phi$

For generated primary particles  
with  $p_T > 0.1$  GeV/c and  $-1.5 < \eta < 2$

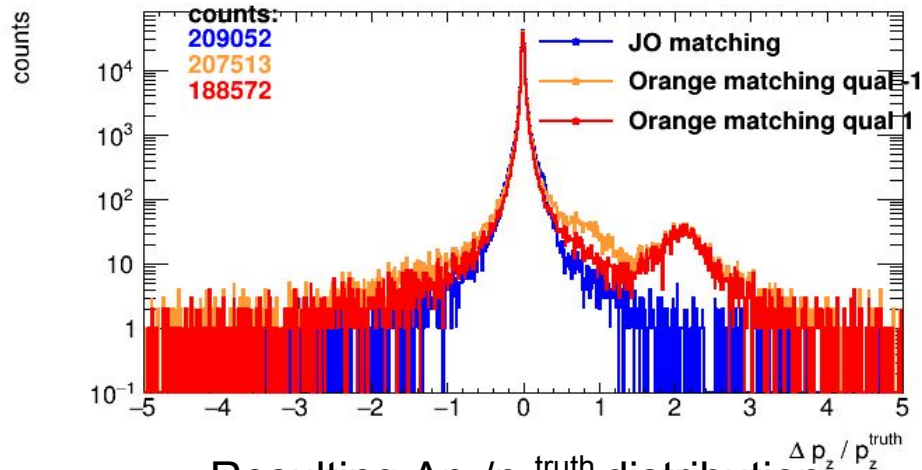


# MC track matching $\Delta p_T/p_T^{\text{truth}}$ , $\Delta p_z/p_z^{\text{truth}}$

For generated primary particles  
with  $p_T > 0.1$  GeV/c and  $-1.5 < \eta < 2$



Cut for JO at 0.3



Resulting  $\Delta p_z/p_z^{\text{truth}}$  distribution  
for JO narrower

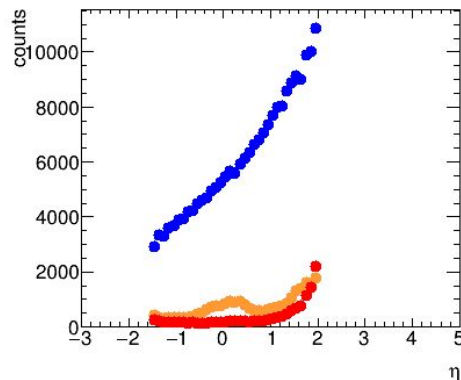
# Unmatched tracks

For reconstructed particles with  $p_T > 0.1$  GeV/c and  $-1.5 < \eta < 2$

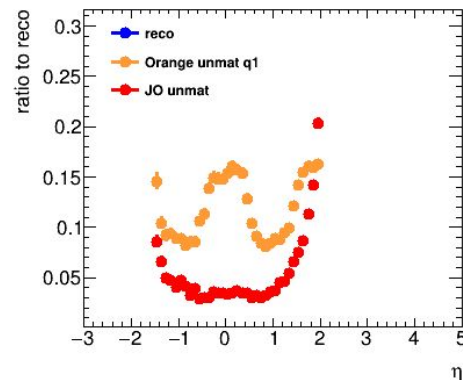
Matching efficiency for selected tracks at ~87% for orange quality = 1.

Matching efficiency for selected tracks at ~95% for JO.

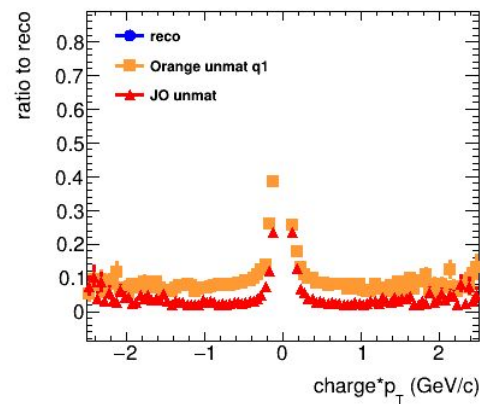
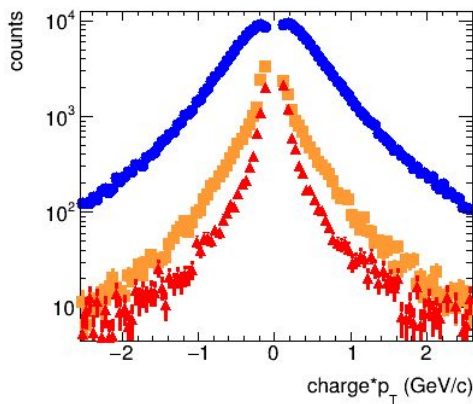
$\eta$



ratio



$p_T^* \text{charge}$



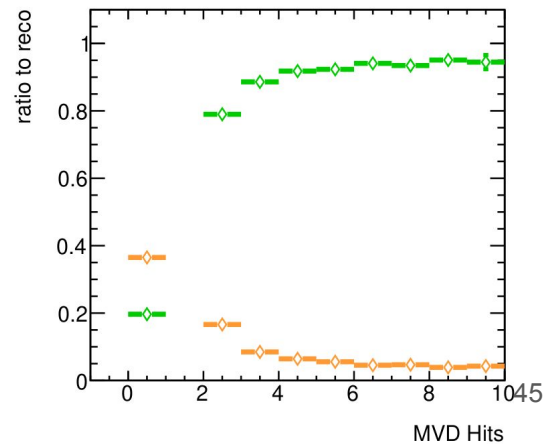
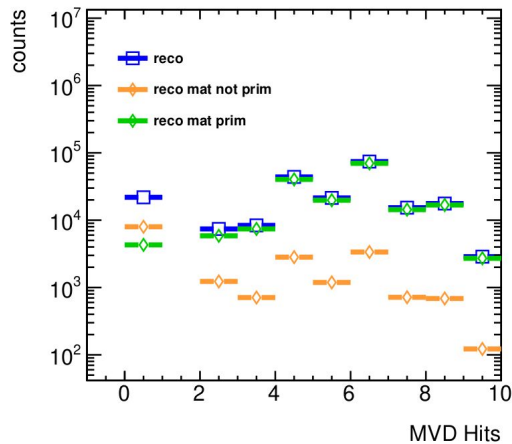
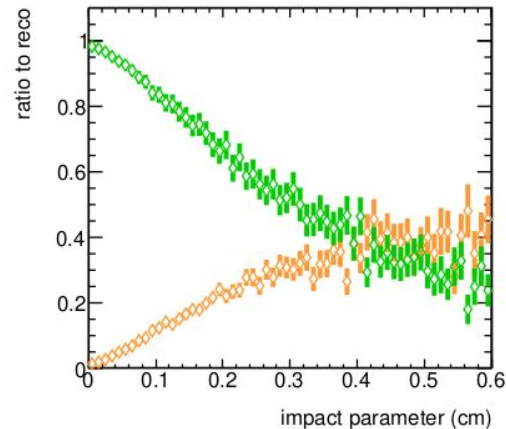
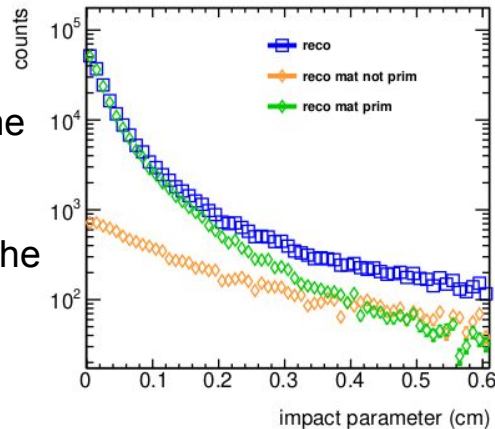
# Contamination of secondaries in reconstructed tracks

Check whether the **selected reconstructed particles** fulfill the definition of **primary particle**, if not, the particle is a **secondary particle**.

At an impact parameter of  $\sim 0.5$  cm, the fraction of primary drops below secondary

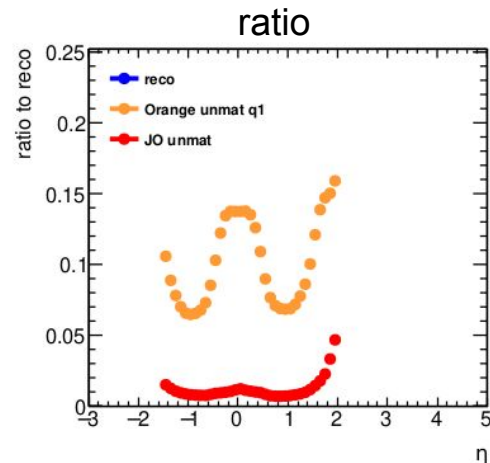
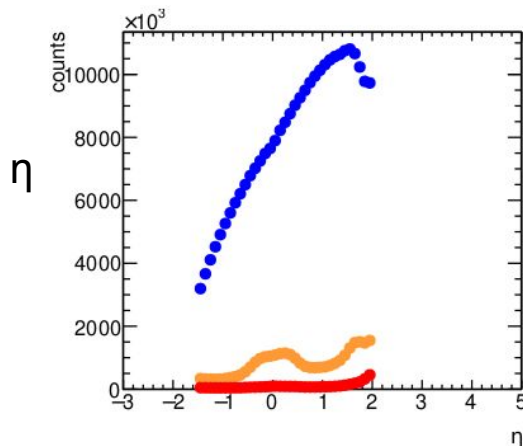
→ cut at 0.5 (previously 1.0) cm.

For 0 MVD hits, there are more secondary than primary particles  
→ require MVD hits  $> 0$   
(previously no cut)



# Unmatched tracks

For reconstructed particles with  $p_T > 0.1$  GeV/c and  $-1.5 < \eta < 2$ , impact parameter  $< 0.5$  cm and MVD hits  $> 0$



Matching efficiency for selected tracks at  $\sim 90\%$  for orange quality = 1.

Matching efficiency for selected tracks at  $\sim 98\%$  for JO.

$p_T^* \text{charge}$

