

Prediction for charm in CC in xFitter

Oleksandr Zenaiev (DESY)

ZAF, 15 May 2018

(slides from xFitter developers meeting on 09 May 2018, with contributions from Valerio Bertone and Fred Olness)

Introduction

- In the old interface we had:
 - ▶ NC DIS inclusive
 - ▶ NC DIS charm
 - ▶ NC DIS beauty
 - ▶ CC DIS inclusive
- In the new interface it is straightforward to implement all possible variants:
 - ▶ reaction: NC DIS or CC DIS
 - ▶ final state: inclusive, charm or beauty (beauty is not of interest in CC)
 - ▶ cross section type: F_2 , F_L , σ_{red} , $\frac{d^2\sigma}{dx dQ^2}$

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- However, the definition of charm CC cross sections seems to be not straightforward from theory point of view

Implementation

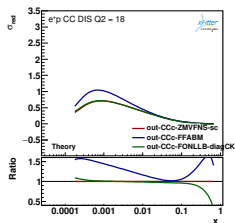
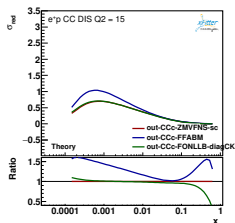
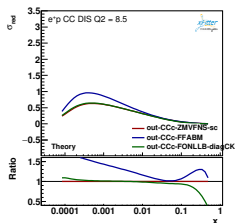
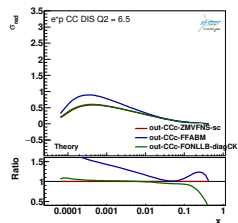
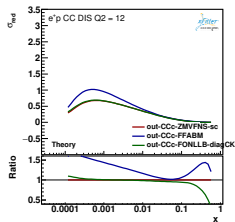
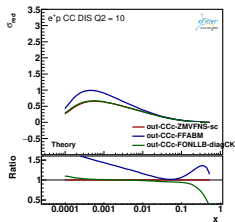
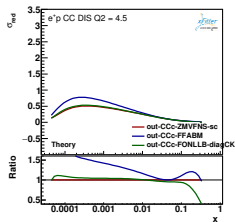
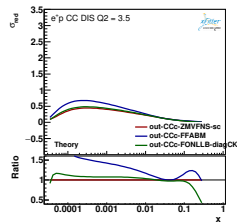
- Task is of practical interest for ongoing analysis in ZEUS
- Implementation will be done in the new interface only
- Branch `charm-CC-XFITTER-70` (JIRA task # 70)
- Implementation can be done for:
 - ▶ ZMVFNS (QCDNUM)
 - ▶ FFABM (OPENQCDRAD)
 - ▶ FONLL (APFEL)
 - ▶ RTOPT implementation in xFitter uses ZMVFNS for CC
 - ▶ ACOT? not yet in new interface
- Inclusive CC prediction verified vs old interface for 3 schemes: exact agreement
- Charm CC prediction:
 - ▶ depends on CKM matrix
 - ▶ for massless treatment, can be defined only in terms of couplings, not partonic processes [1001.2312, section 5]
 - ⇒ both $s \rightarrow c$, $c \rightarrow s$ need to be included in FONLL and (?) ZMVFNS
 - ⇒ this does not apply to FFABM? only $s \rightarrow c$ included in OPENQCDRAD
 - ▶ paper [1001.2312] discusses the impact of different theoretical and experimental definition of charm contribution, but only for NC
 - ★ concluded that it is at (sub)percent level for HERA
 - ★ seems to be not the case for CC: differences appear at LO

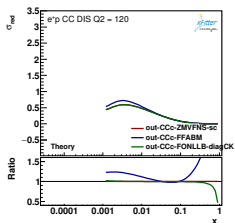
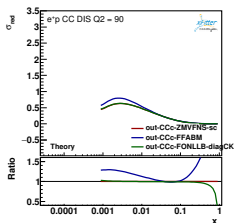
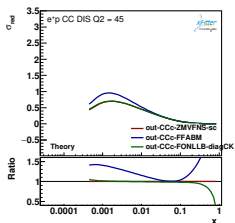
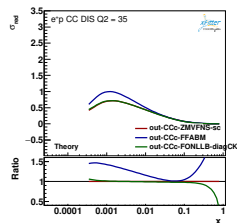
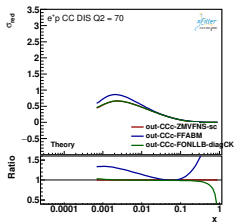
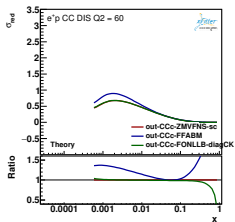
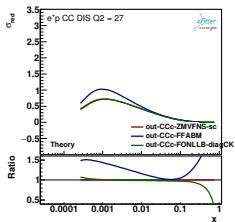
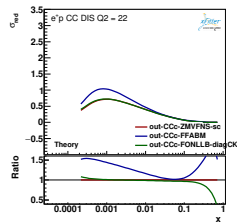
Implementation

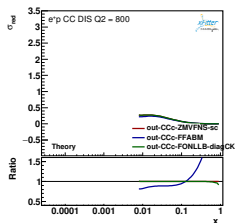
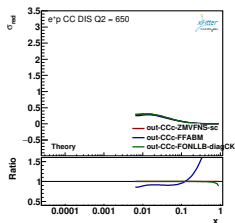
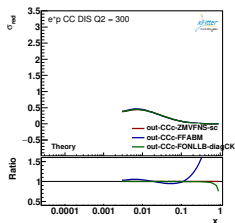
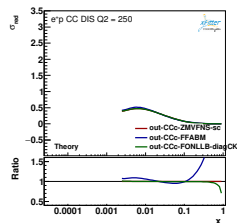
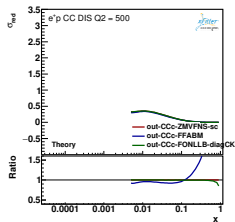
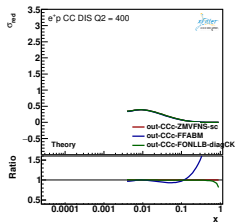
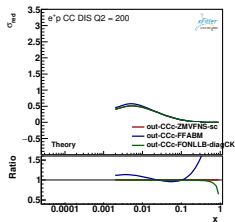
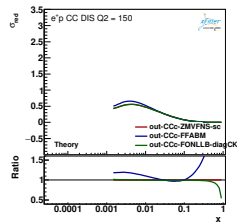
- Current status of charm CC implementation:
 - ▶ ZMVFNS (QCDNUM): implemented, but results are weird: to be checked with Michiel Botje
 - ▶ FFABM (OPENQCDRAD): straightforward to take from OPENQCDRAD (to be checked with Sergey Alekhin)
 - ▶ FONLL (APFEL): straightforward to take from APFEL, checked with Valerio
- Dependence on CKM matrix:
 - ▶ ZMVFNS so far uses diagonal CKM matrix
 - ▶ FFABM so far uses its own non-diagonal CKM matrix (setup inside OPENQCDRAD)
 - ▶ FONLL presently is the only reaction which properly takes CKM matrix from parameters.yaml
 - ⇒ dependence on CKM checked by using diag. vs non-diag. CKM with FONLL
 - ⇒ setting of CKM from parameters.yaml to be implemented for ZMVFNS and FFABM
- Contribution from different processes:
 - ▶ checked in ZMVFNS by switching them on/off

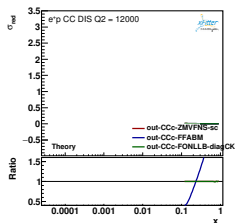
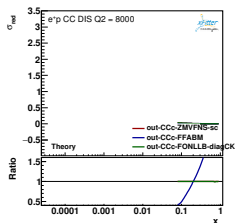
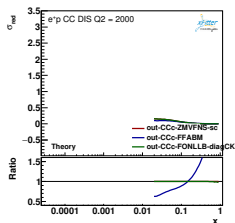
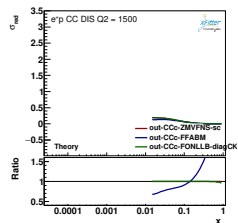
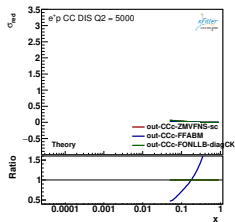
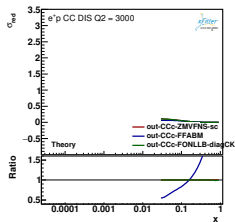
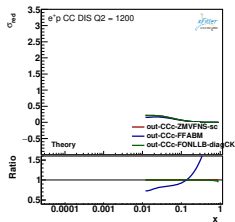
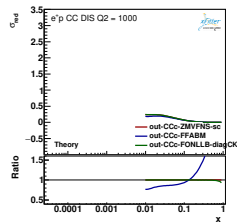
- Comparison of charm CC prediction done for ZMVFNS, FFABM and FONLL:
 - ▶ calculation at NLO
 - ▶ same PDFs at starting scale 1.9 GeV^2 (but different evolution, plots for PDFs in backup)
 - ▶ same charm mass (pole mass option in FFABM and FONLL)
 - ▶ FONLL-B with diagonal CKM (non-diag. CKM only in FFABM)
 - ▶ scales set to Q^2
 - ▶ pseudodata file used to provide predictions for $3.5 < Q^2 < 30000 \text{ GeV}^2$ and full accesible x_{Bj} range
- Additionally, FONLL with diag. and non-diag. CKM compared
- Additionally, ZMVFNS with three sets of processes compared:
 - ▶ $c \rightarrow s$ only ('ZMFVNS-1')
 - ▶ $s \rightarrow c$ only ('ZMFVNS-2')
 - ▶ $s \rightarrow c$ and $c \rightarrow s$ ('ZMFVNS-3')

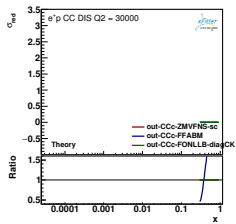
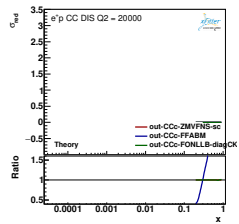
CC charm in three schemes (CC inclusive in backup)



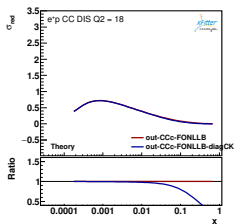
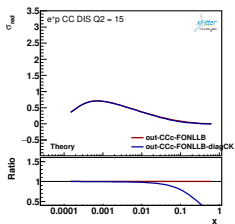
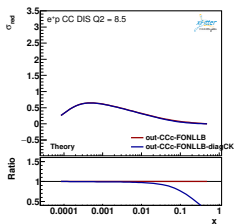
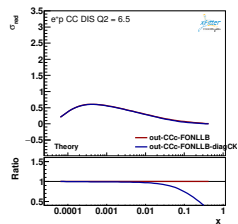
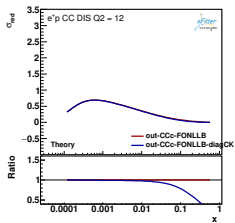
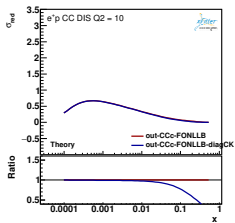
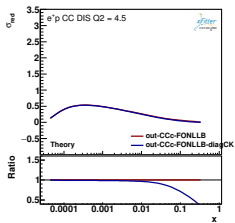
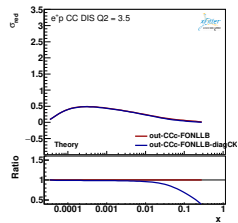


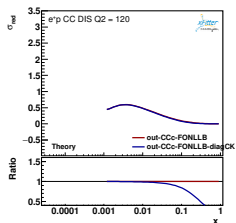
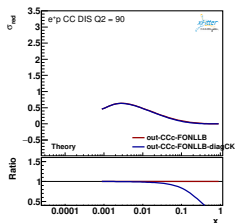
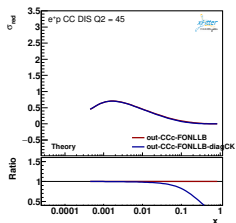
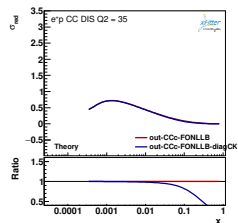
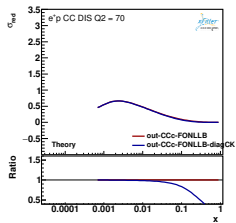
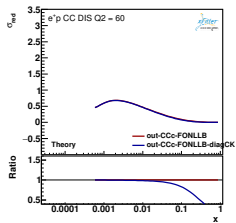
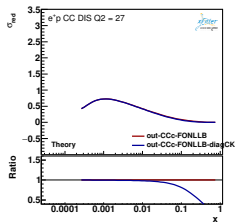
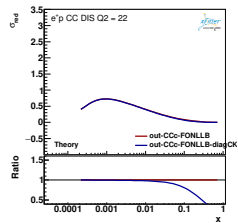


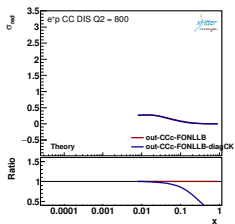
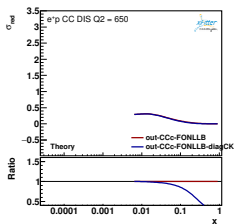
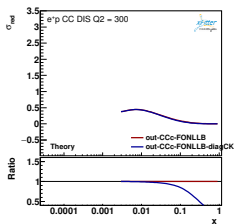
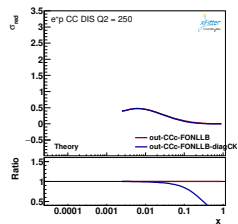
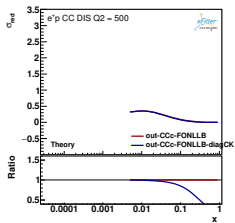
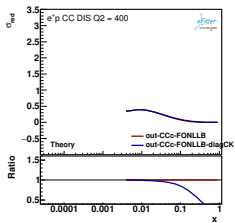
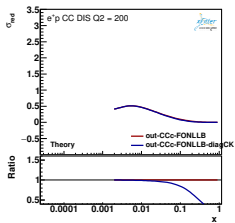
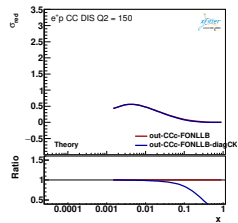


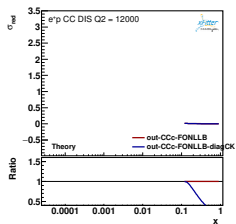
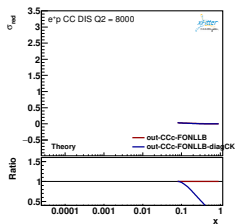
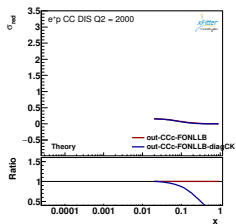
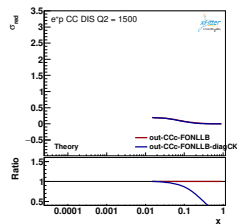
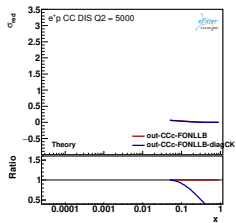
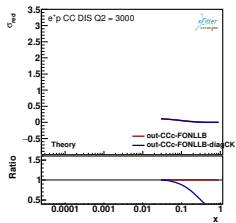
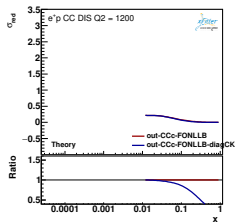
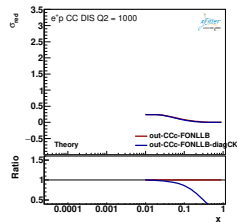


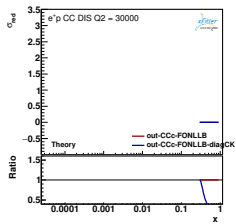
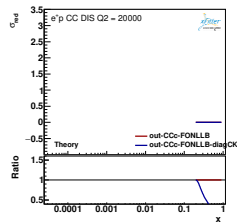
CC charm in FONLL, diag. vs non-diag. CKM





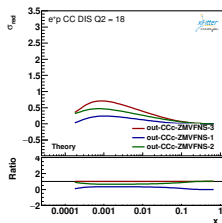
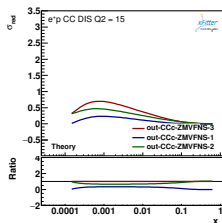
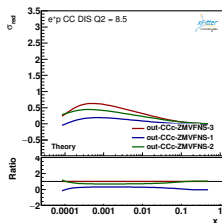
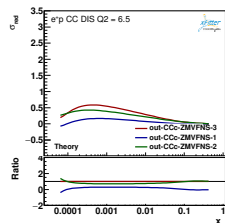
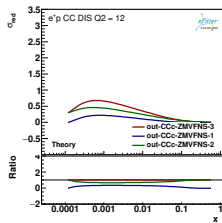
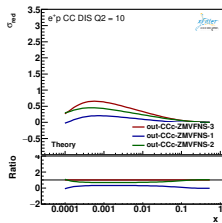
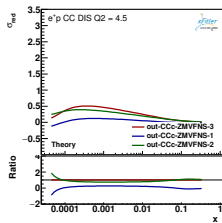
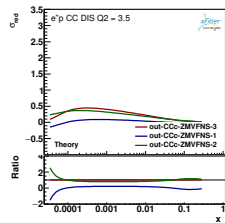


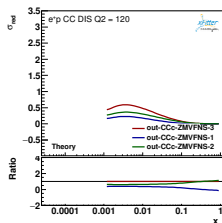
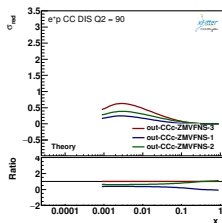
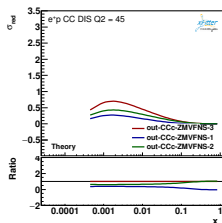
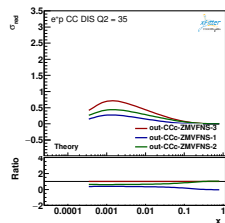
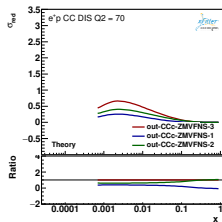
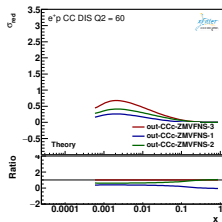
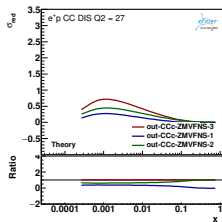
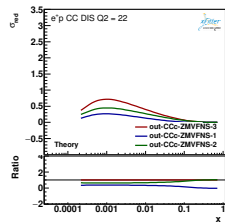


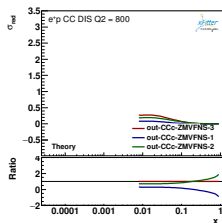
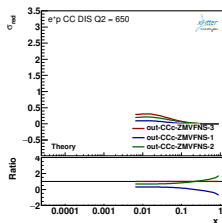
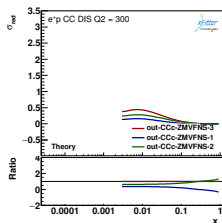
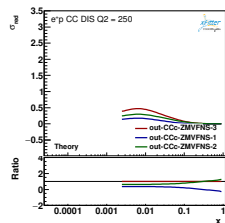
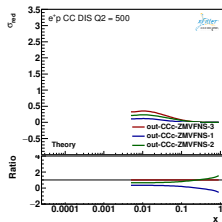
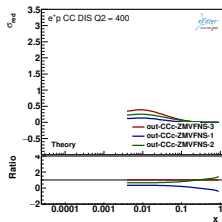
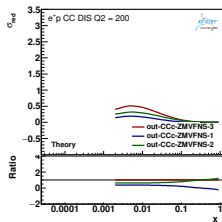
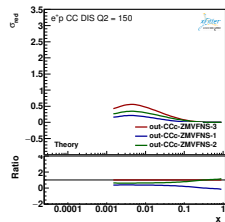


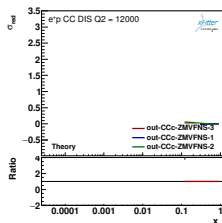
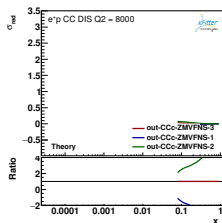
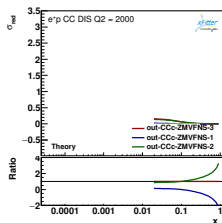
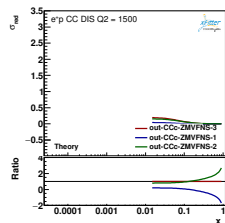
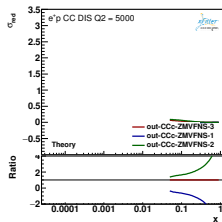
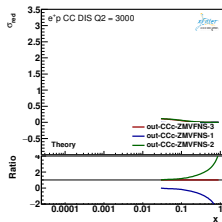
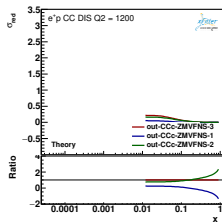
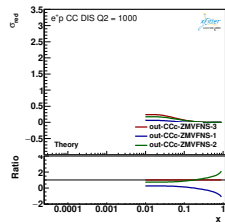
CC charm in ZMVFNS, different partonic processes:

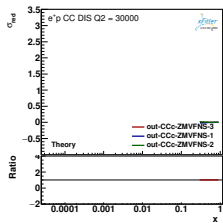
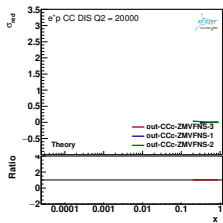
- $c \rightarrow s$ only ('ZMFVNS-1')
- $s \rightarrow c$ only ('ZMFVNS-2')
- $s \rightarrow c$ and $c \rightarrow s$ ('ZMFVNS-3')











Summary

- Significant differences ($\sim 50\%$ and larger at high Q^2) between schemes:
 - ▶ good agreement between ZMVFNS and FONLL
 - ▶ difference vs FFABM due to:
 - ★ massive vs massless treatment, number of flavours in PDF evolution
 - ★ different definition ($s \rightarrow c$ only vs $c \rightarrow s$, $s \rightarrow c$): perhaps dominant?
 - Dependence on CKM becomes relevant for $x \gtrsim 0.05$
 - If only $s \rightarrow c$ and $c \rightarrow s$ makes sense in ZMVFNS/VFNS, perhaps the difference from experimental definition ($s \rightarrow c$ only) **can not be ignored**
 - ▶ $\Rightarrow$ is ZMVFNS/VFNS prediction for charm CC usable?
 - any advises/recommendations?
-

- Current plan:
 - ▶ we can have FFABM and FONLL inclusive predictions (function of Q^2 , x_{Bj}) checked with the authors but defined differently
 - ▶ estimate contributions $s \rightarrow c$ and $c \rightarrow s$ at LO in ZMVFNS using QCDNUM (cannot be done in FFNS, VFNS or higher orders?)
 - $\rightarrow$ might be useful in light of differences for s in global PDF analysis
 - $\rightarrow$ so far weird results (negative $c \rightarrow s$), in contact with Michiel Botje

BACKUP

Implementation in ZMVFNS (QCDNUM)

- Current status of charm CC implementation:
 - ▶ ZMVFNS (QCDNUM): seems to be straightforward once decided which processes to include

```
// Helpers for QCDNUM (CC):

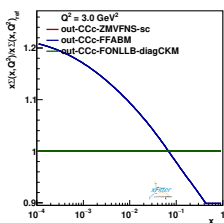
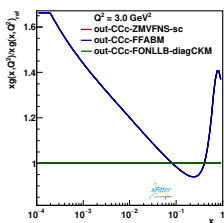
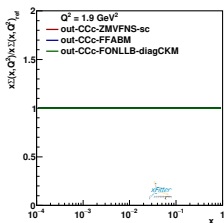
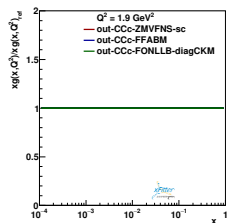
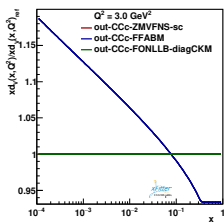
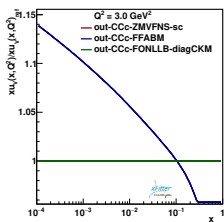
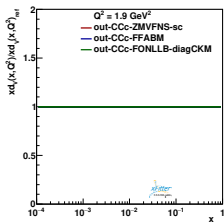
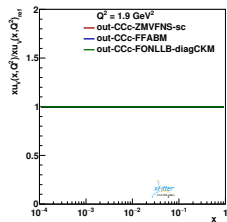
//! full
const double CCEP2F[] = {0.,0.,1.,0.,1.,0., 0. ,1.,0.,1.,0.,0.,0.} ;
const double CCEM2F[] = {0.,0.,0.,1.,0.,1., 0. ,0.,1.,0.,1.,0.,0.} ;

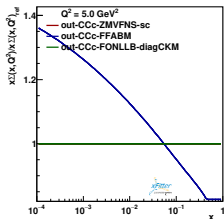
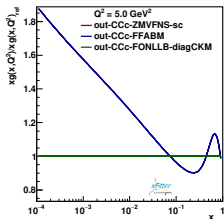
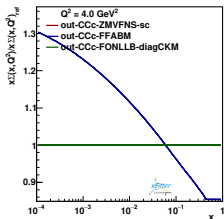
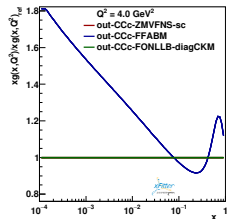
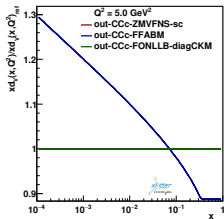
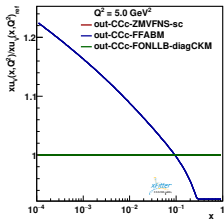
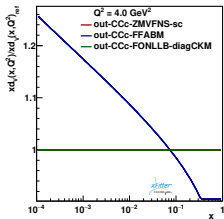
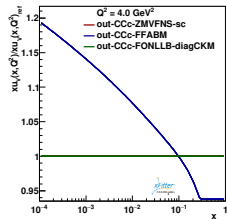
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const double CCEM3F[] = {0.,0. ,0.,-1.,0.,-1.,0.,0.,1.,0.,1.,0.,0.};

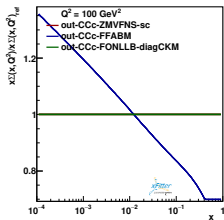
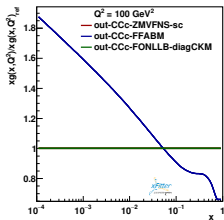
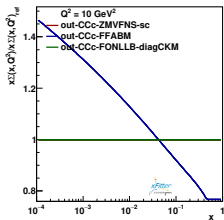
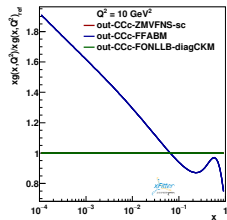
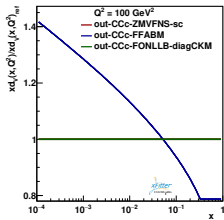
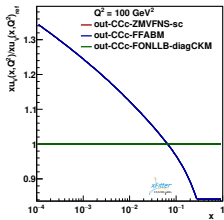
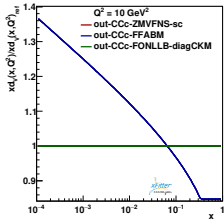
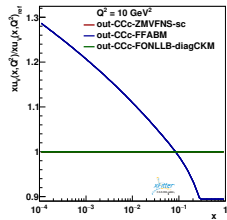
//! c
// work in progress: according to 1001.2312 section 5,
// in ZM only the sum of contributions s + c makes sense
// three different options are below for checks, uncommented one is for s + c
//
// only c
//const double CCEP2Fc[] = {0.,0.,1.,0.,0.,0.,0.,0.,0.,0.,0.,0.,0.} ;
//const double CCEM2Fc[] = {0.,0.,0.,0.,0.,0.,0.,0.,0.,0.,0.,1.,0.,0.} ;
// only s
//const double CCEP2Fc[] = {0.,0.,0.,0.,0.,0.,0.,0.,0.,0.,1.,0.,0.,0.} ;
//const double CCEM2Fc[] = {0.,0.,0.,1.,0.,0.,0.,0.,0.,0.,0.,0.,0.,0.} ;
// only s,c
const double CCEP2Fc[] = {0.,0.,1.,0.,0.,0.,0.,0.,0.,0.,1.,0.,0.,0.} ;
const double CCEM2Fc[] = {0.,0.,0.,1.,0.,0.,0.,0.,0.,0.,0.,0.,1.,0.,0.} ;
```

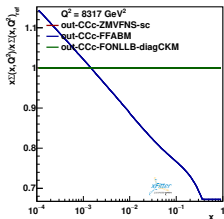
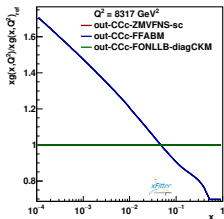
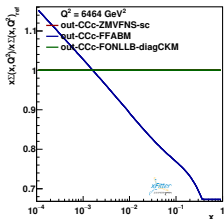
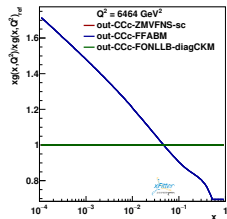
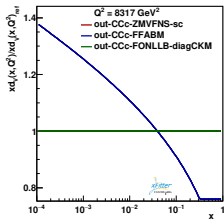
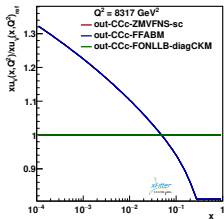
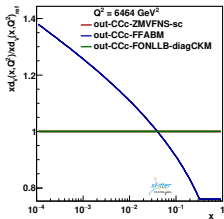
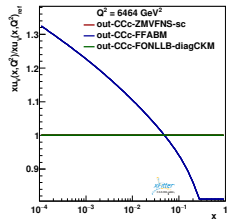
(same for CCEM3F,CCEP3F)

PDFs in three schemes

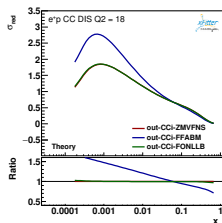
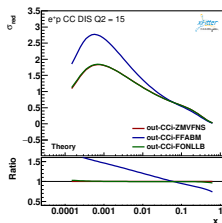
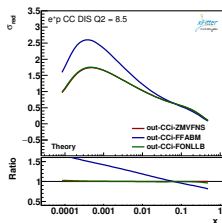
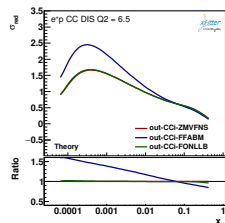
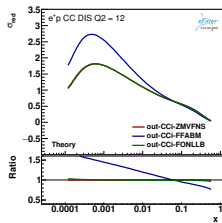
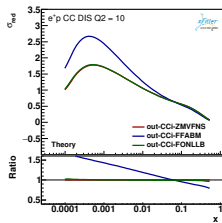
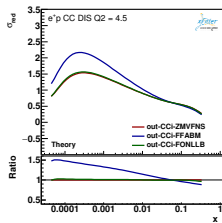
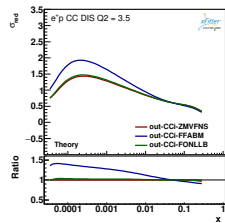


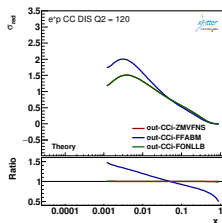
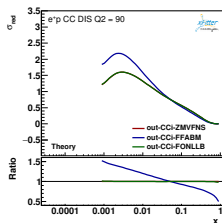
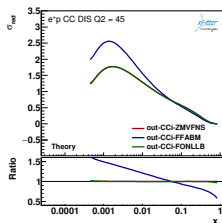
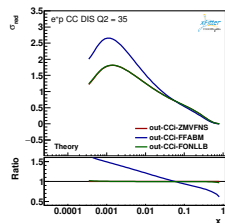
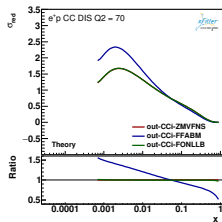
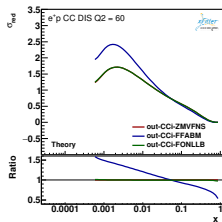
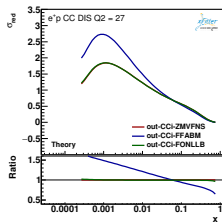
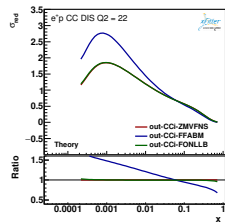


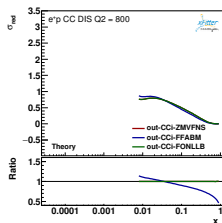
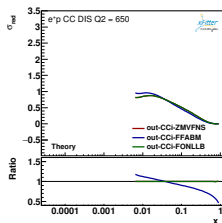
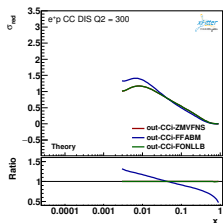
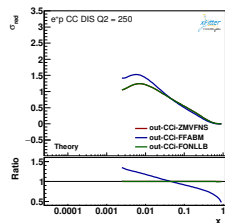
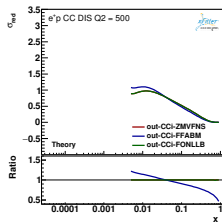
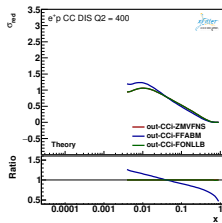
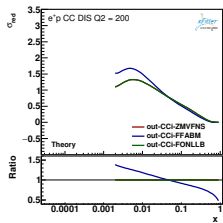
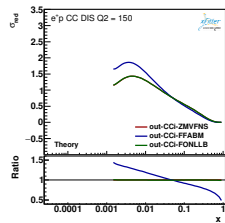


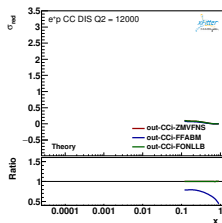
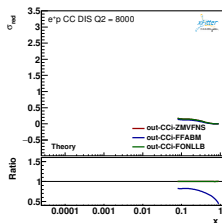
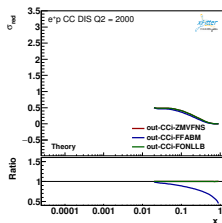
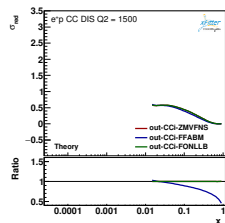
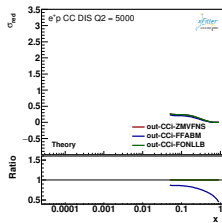
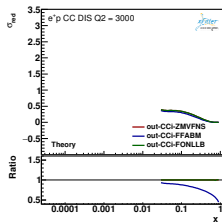
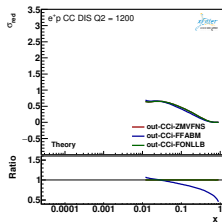
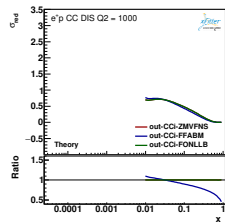


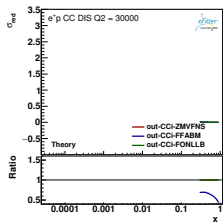
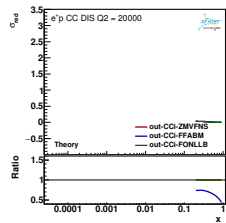
CC inclusive in three schemes











Notes from Fred Olness

A. Very Rough Notes on charged current with Variable Flavor Number Scheme.

I've attached some diagrams to illustrate the problem. Suppose we want to compute the inclusive F_2^{charm} with Charged Current (CC) charm production at NLO. The obvious LO diagram is ($sW^+ \rightarrow c$).

What is not so obvious is we also need ($cW^+ \rightarrow s$). This is because the $\bar{c}$ comes from a gluon splitting to $c\bar{c}$, and the c goes down the beam pipe. Experimentally, this is undesirable, but if you want to do a consistent truly "inclusive" F_2^{charm} , this is essential. If you ignore it, the problems will appear with the NLO terms. [More precisely, a truly "inclusive" F_2^{charm} is ill-defined; Collins' proof of factorization for heavy quarks specifically focused on a totally inclusive F_2 for this reason.]

Note, if you make a "non-inclusive" F_2^{charm} (or "exclusive") measurement, this can be well defined; however, you must add a resolution scale (or regulator) to define how close to the beam fragments your charm can be.

To keep things simple, we'll focus on just the NLO gluon-initiated diagrams. The arguments follow through in a similar manner for the quark-initiated graphs.

1. t-channel at NLO

The t-channel at NLO is straightforward. We add the ($gW^+ \rightarrow cs$) diagram. This exchanges an s quark in the t-channel, and thus will have a $\ln(m_s^2/Q^2)$ divergence for large Q . This is resolved by the subtraction term $f_g \otimes \bar{P}_{g \rightarrow s} \otimes \sigma_{W^+ \rightarrow cs}$ where the $\bar{P}_{g \rightarrow s}$ represents a perturbative splitting of $g \rightarrow s$ and will cancel the double counting between the LO and NLO graphs in the limit where the exchanged s quark becomes collinear. At NLO, the SUB term is proportional to $\bar{P}_{g \rightarrow s} \sim \frac{\alpha_s}{2\pi} P_{g \rightarrow s}^{(1)} \ln(m_s^2/Q^2)$. The logarithmic divergence will cancel between the NLO and SUB terms as $Q^2 \rightarrow \infty$, resulting in a finite result for the NLO t-channel contribution.

2. u-channel at NLO

The u-channel at NLO is more subtle. We add the ($gW^+ \rightarrow cs$) diagram with a $\bar{c}$ quark in the u-channel, and thus will have a $\ln(m_s^2/Q^2)$ divergence for large Q . This is resolved by the subtraction term $f_g \otimes \bar{P}_{g \rightarrow s} \otimes \sigma_{W^+ \rightarrow cs}$ where the $\bar{P}_{g \rightarrow s}$ represents a perturbative splitting of $g \rightarrow s$ and will cancel the double counting between the LO and NLO graphs in the limit where the

exchanged $\bar{c}$ quark becomes collinear. At NLO, the SUB term is proportional to $\bar{P}_{g \rightarrow s} \sim \frac{\alpha_s}{2\pi} P_{g \rightarrow s}^{(1)} \ln(m_s^2/Q^2)$. The logarithmic divergence will cancel between the NLO and SUB terms as $Q^2 \rightarrow \infty$, resulting in a finite result for the NLO u-channel contribution.

The above description parallels the t-channel analysis.

Why do we need the LO ($cW^+ \rightarrow s$)?

The SUB term: It is essential we include the subtraction term $f_g \otimes \bar{P}_{g \rightarrow s} \otimes \sigma_{W^+ \rightarrow cs}$ so that we get a finite answer at large energies.

But, what if we did NOT include the LO ($cW^+ \rightarrow s$)?

At energy scales $Q \sim m_c$, the LO and SUB terms remove the double counting between the LO and NLO processes. This is most apparent when you plot the individual terms versus the Q scale (or more properly, it is the μ scale). [See Figure]

In the region of $Q \sim m_c$, the charm PDF f_c (and hence, the LO contribution) rises very quickly as it is driven by the very large gluon, and coupled with a large $\alpha_S(m_c)$.

The SUB subtraction also rises quickly as this is driven by the logarithmic term $\ln(m_s^2/Q^2)$.

The difference LO-SUB is the physical contribution to the total (TOT-LO-NLO-SUB), and it is this combination which is smooth across the "turn on" of the charm PDF.

We now see that if we neglect the LO ($cW^+ \rightarrow s$) we lose the cancellation between LO and SUB in the $Q \sim m_c$ and our structure function (or cross section) would have an anomalous rise at the location where we arbitrarily turn on the charm PDF.

B. The bottom line:

A truly "inclusive" F_2^{charm} is ill-defined. Instead, we necessarily must an "experimentally" defined F_2^{charm} where we specify conditions so that the final state charm is isolated from the hadron remnants.

We can talk about a fully inclusive F_2 where we include all flavors; this was the subject of Collins' proof.

If we compute a "pseudo-inclusive" F_2^{charm} in the Variable Flavor Number Scheme, we do need to include the LO ($cW^+ \rightarrow s$) and the associated SUB ($gW^+ \rightarrow cs$).

We can compute in the Fixed Flavor Number Scheme, but in the large energy limit, we encounter $\ln(m_s^2/Q^2)$ divergences. In practice, our Q scales are not large enough to generate infinities, but they are large enough where we see the resummed logs included in the VFS charm PDF become important.

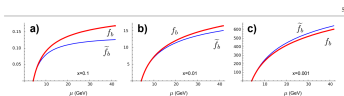


Fig. 2 The comparison of the DGLAP evolved PDF $f_2(c, \mu)$ and the perturbatively calculated $L_2(c, \mu)$ as a function of μ for selected s values. For $\mu \rightarrow m_c$ we find the functions match precisely: $f_2(c, \mu) \rightarrow L_2(c, \mu)$. We have used NNPDF31_m,118,nn1.9 as the base PDF set.

Figure 1. Figure taken from our ePrint paper "Impact of the heavy quark matching scales in PDF fits." The LO contributions correspond to f_s , and the SUB to $\tilde{f}_s$. The cancellation (LO-SUB) is quite precise. If we were to remove LO or SUB, our TOT result would have anomalous contributions in the region $Q \sim m_c$.

