

NLO and off-shell effects in top quark mass determinations

Theorist of the month seminar at DESY Hamburg.

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Work in collaboration with Gudrun Heinrich, Andreas Maier, Richard Nisius, Johannes Schlenk,
Markus Schulze and Ludovic Scyboz.



Mass of the top quark.

Importance of knowing this Standard Model parameter precisely:

- close link to Higgs physics because of largest Yukawa coupling
- Higgs potential and SM vacuum stability critically depend on this value
- processes involving top quarks used for SM precision tests and constraining PDFs
- many important backgrounds in new physics searches contain large portion of top quark physics

Measurements to extract this parameter are complicated:

- top quark production leads to complex hadronic & leptonic final states often plus additional jets
 - top hadrons do not exist: top quarks decay before they can hadronize
- challenging environment: many effects and subtle theoretical issues need be understood (top quark reconstruction, higher-orders, off-shell effects, interferences, colour reconnections, IR renormalon, Monte Carlo mass versus pole mass etc)



Why are we interested in $WbW\bar{b}$ predictions.

In the experiment we do not measure tops. We only have a handle on their decay products.
 $WbW\bar{b}$ is therefore the more realistic final state if you are interested in $t\bar{t}$ production.

~ quantum mechanical versus semi-classical treatment ~

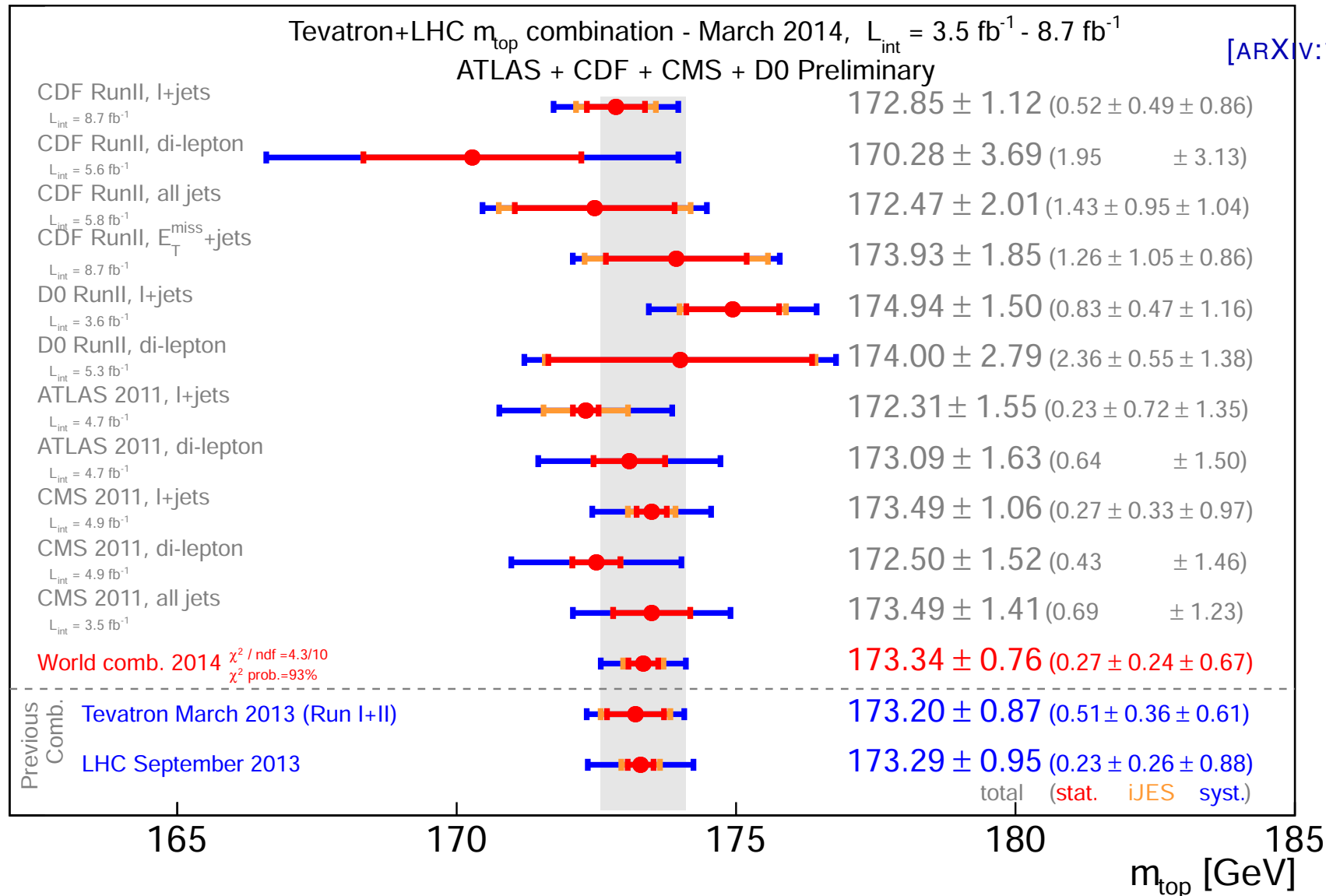
- important contributions to Wt and WW final states (tricky to disentangle at higher orders)
 - important background to BSM searches and SM measurements (e.g. population of N-jet bins in WW production)
 - at current precision, we start worrying about offshell effects, non-factorizable corrections, b-mass dependence etc.
- expect small ($\mathcal{O}(\Gamma_t/m_t)$) effects (wrt NWA) for inclusive $t\bar{t}$ observables
→ similar statements for more exclusive phase spaces?



Top quark mass measurements, March 2014

- first LHC+Tevatron result; total uncertainty on top quark mass < 1 GeV for combinations.

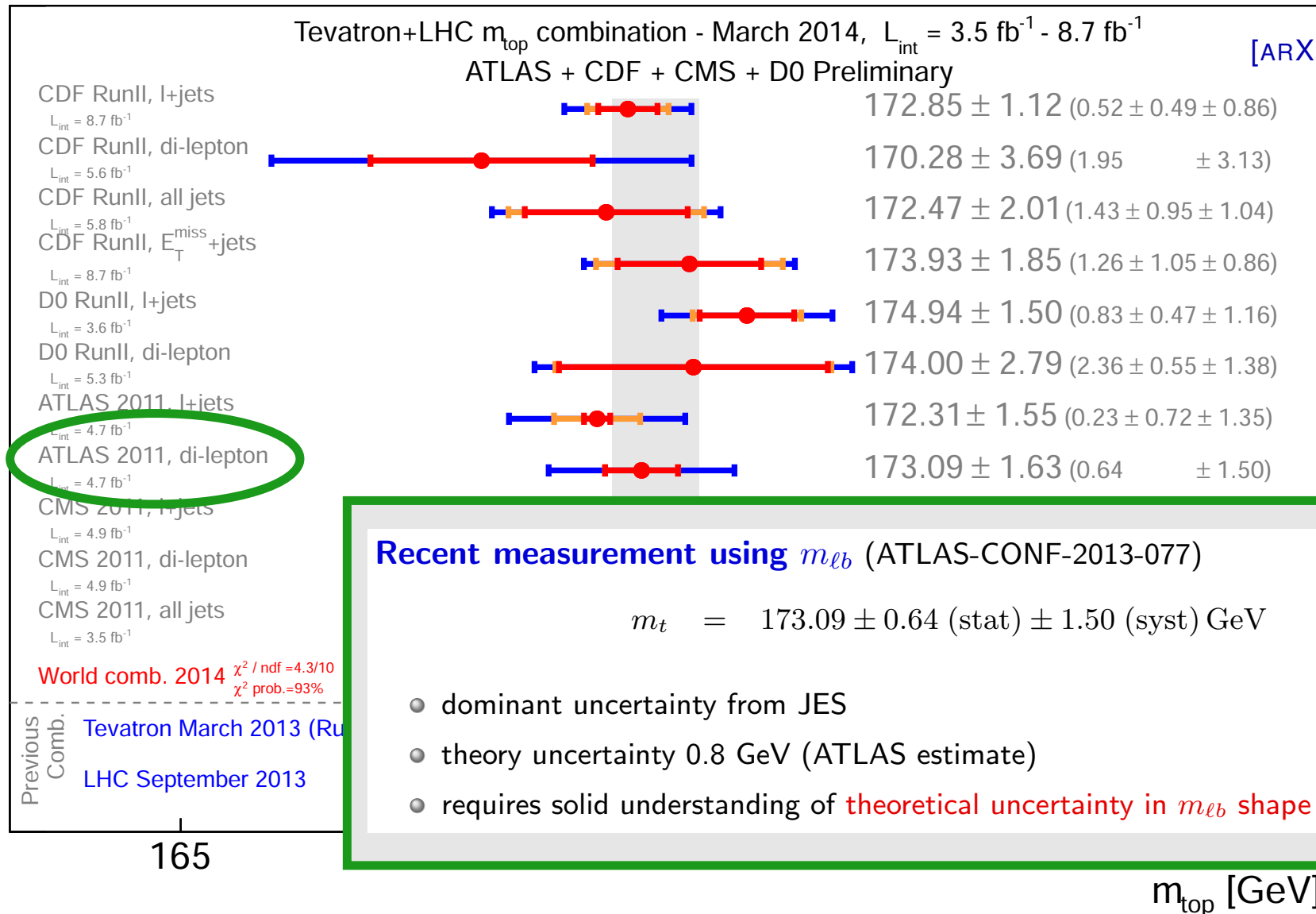
Rich and active experimental program (various complementary techniques).



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Outline.

Current experimental status

Template method for top quark mass sensitive observables &
Earlier results for LHC 7 TeV

New studies for LHC 13 TeV

$WbW\bar{b}$ calculations – quick recap

(Snippet of recent literature review)

(Recent developments in NLO+PS matched simulations for $WbW\bar{b}$)

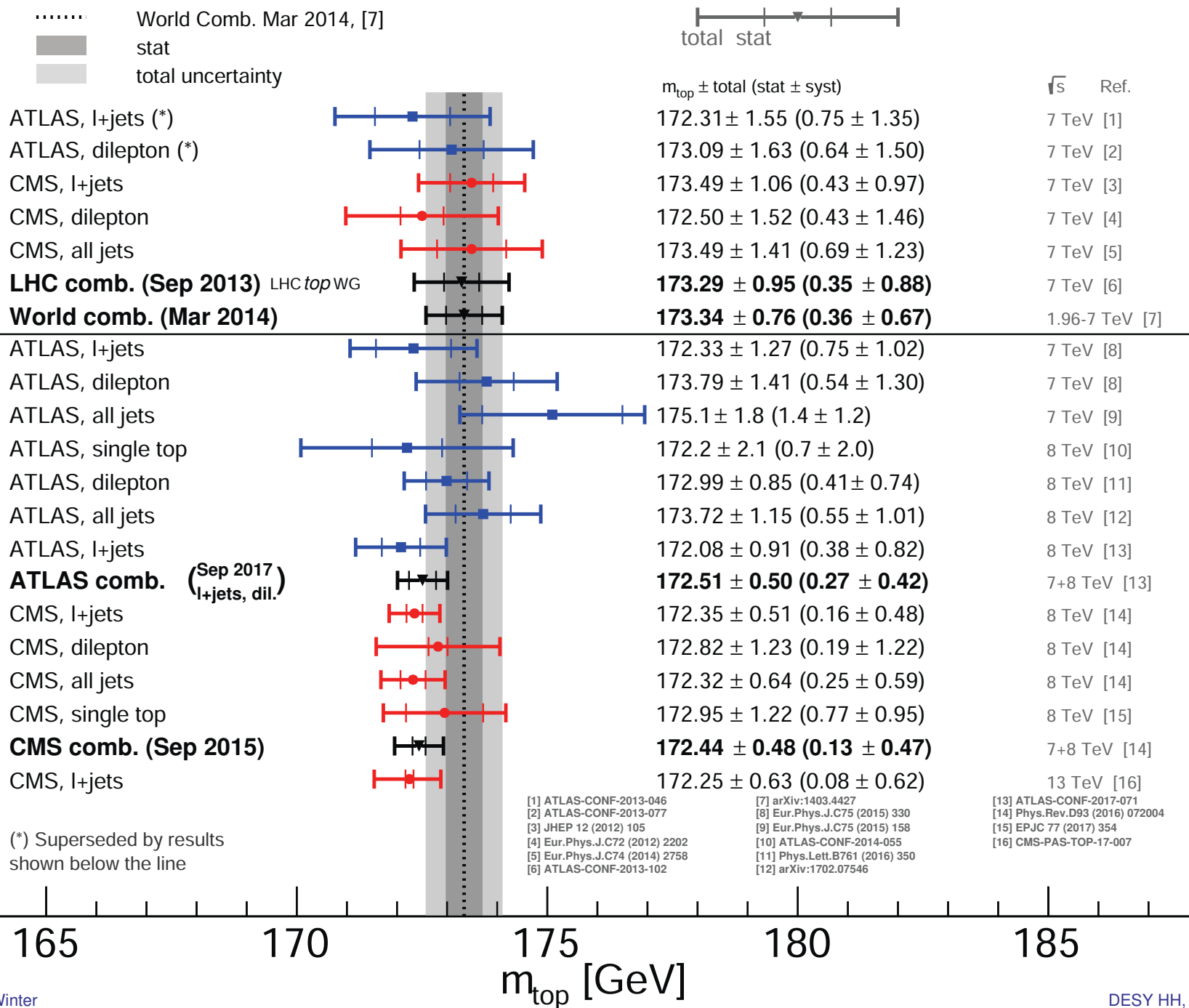


ATLAS+CMS Preliminary

LHC $t\bar{t}$ WG

$m_{t\bar{t}}$ summary, $\sqrt{s} = 7\text{-}13\text{ TeV}$

September 2017



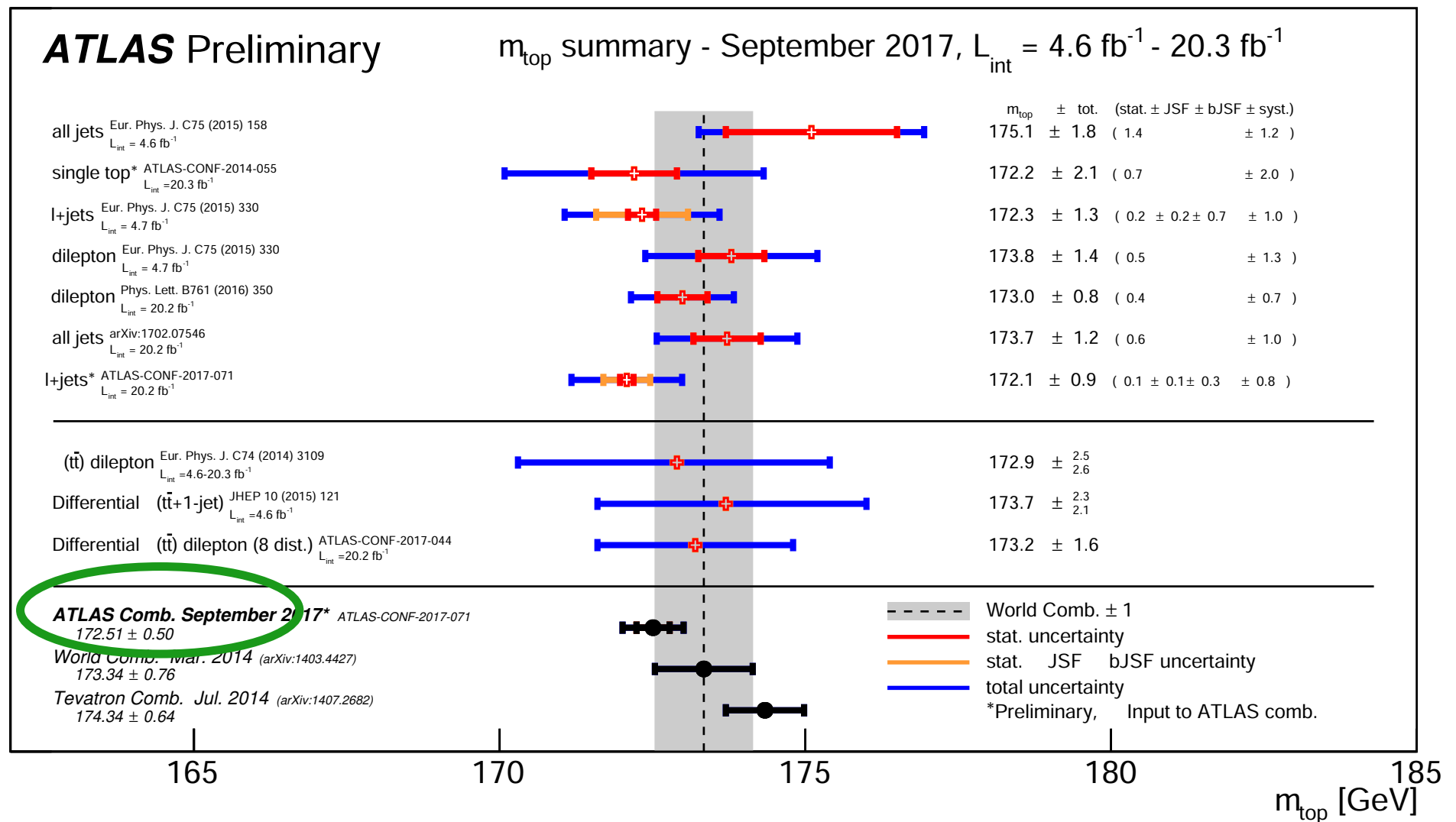
Current experimental status, September 2017

- first ATLAS combination result giving total uncertainty on top quark mass $\approx \frac{1}{2}$ GeV.

Dilepton channel results:

7 TeV: [EUR. PHYS. J. C **75** (2015) NO.7, 330; ARXIV:1503.05427]

8 TeV: [PHYS. LETT. B **761** (2016) 350; ARXIV:1606.02179]

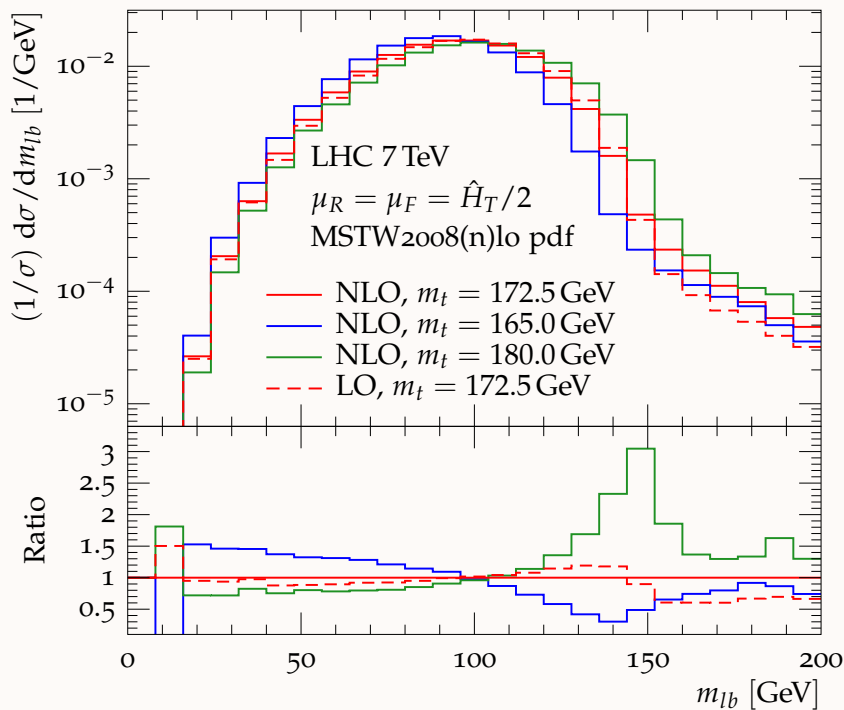


Top quark mass (parameter) determination from the shape of the invariant mass distribution of the charged lepton and b-jet in the dilepton channel at the LHC.

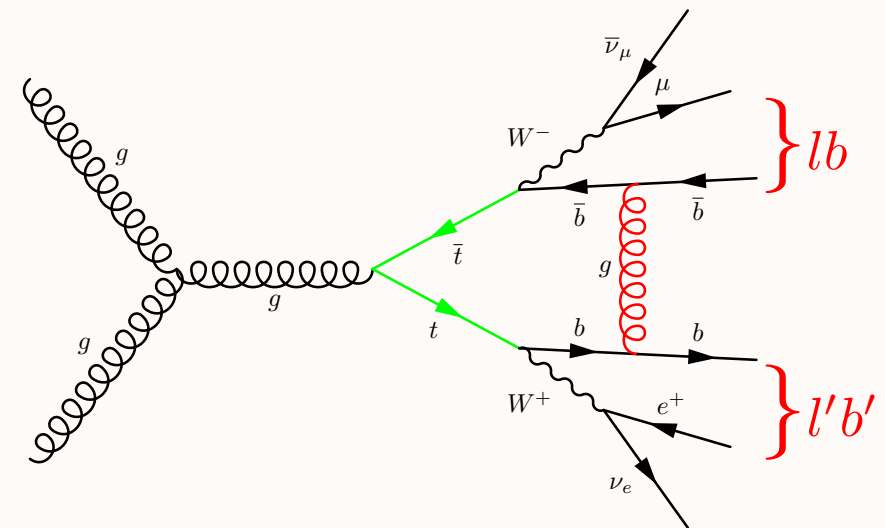


Invariant mass of b-jet and charged lepton

- Definition: $m_{lb}^2 = (p_{b\text{-jet}} + p_l)^2$ (and criteria to deal with combinatorial problem)
- The m_{lb} distribution is sensitive to varying the top quark mass $\rightarrow$ NLO calculations for different m_t .
- Analysis and kinematical requirements according ATLAS measurement as described in [\[ATLAS-CONF-2013-077\]](#); in particular: use lepton b-jet pairing minimizing sum of both m_{lb} and average
 $\rightarrow$ Suitable observable for template fit to precisely measure the (MC) top quark mass.



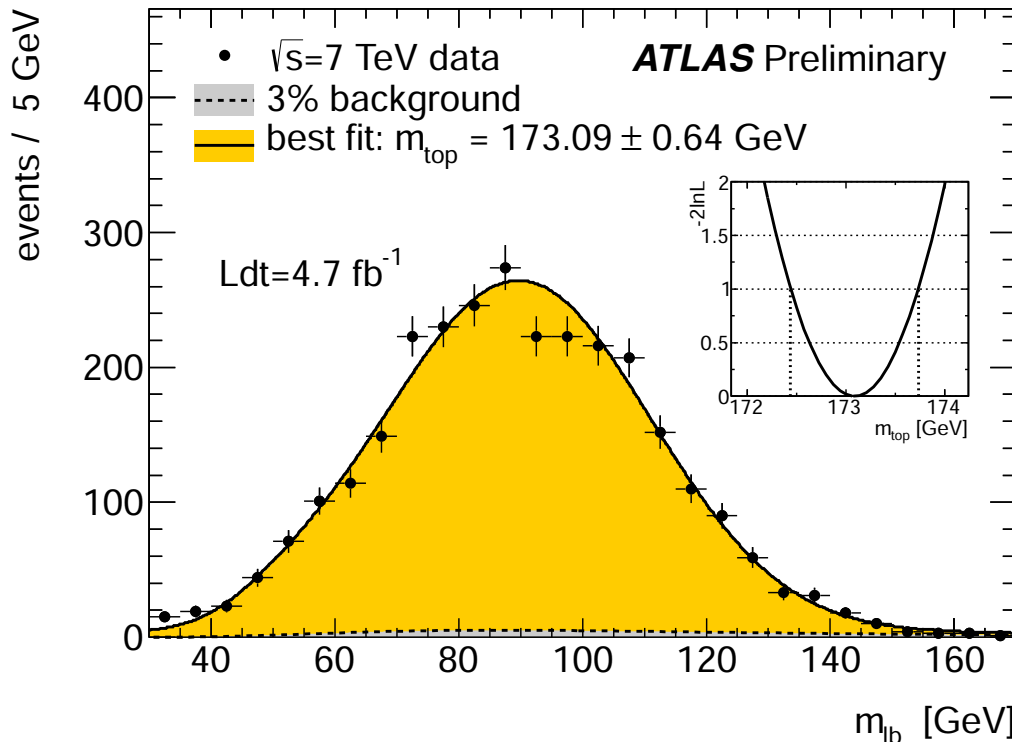
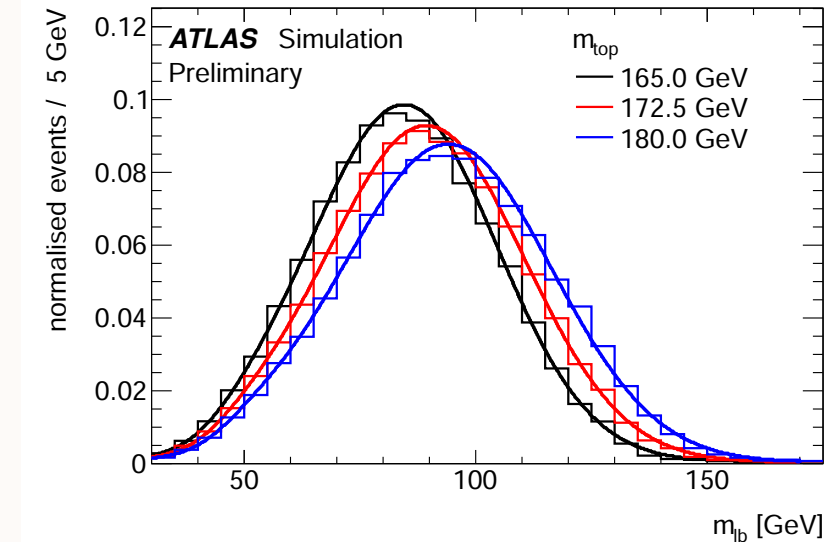
← shape predictions from full QCD NLO $WWb\bar{b}$
 from GoSam+Sherpa



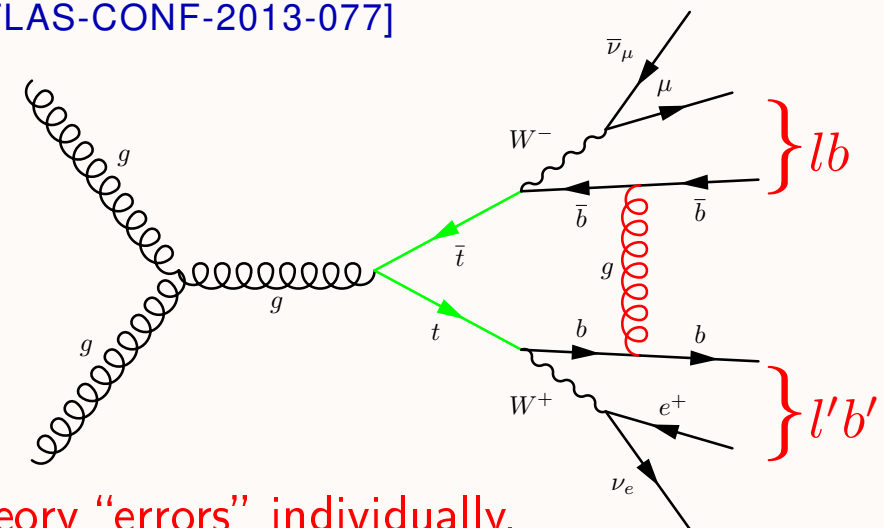
Top quark mass determination using the m_{lb} method

Parametrize “your” theory (provide m_{lb} predictions incl. as many physics effects as possible).

- Already at LO NWA or LO $W^+W^-b\bar{b}$ in dilepton channel: m_{lb} distribution is sensitive to top quark mass.
 - ATLAS uses one-dim. template method to determine m_t . Theory uncertainty had been estimated to 0.8 GeV at 7 TeV.
- Verify size of th. uncertainties using more advanced calc's!



[ATLAS-CONF-2013-077]



→ Use pseudo-data to study different types of theory “errors” individually.

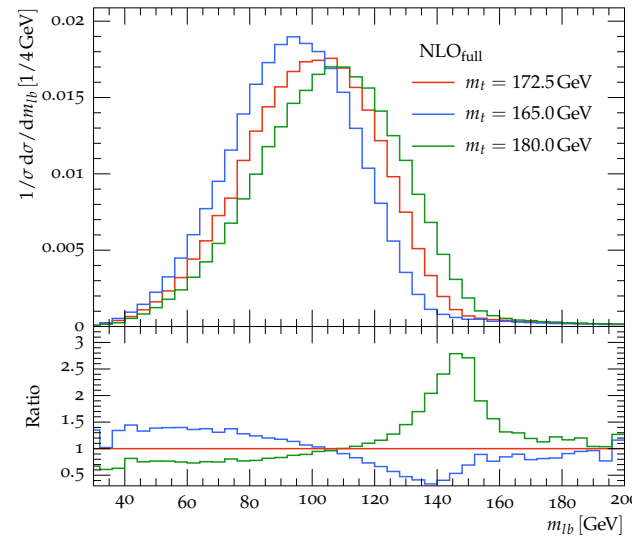


Calibration of template fit function

- 1. Choose distributions sensitive to the top-quark mass
- 2. Generate distributions for different input m_t^{in} :

$$m_t^{in} \in [165.0, 172.5, 180.0] \text{ GeV}$$

- 3. Calibrate the fit function
 - Choose appropriate fit functions
 - Verify, then impose linear dependence of function parameters on m_t
 - Fix the parameters by a simultaneous fit to all distributions
 - Top-quark mass m_t left as a floating parameter



- From a given underlying theory prediction, draw a number of events corresponding to a target luminosity.

⇒ “pseudo-data”

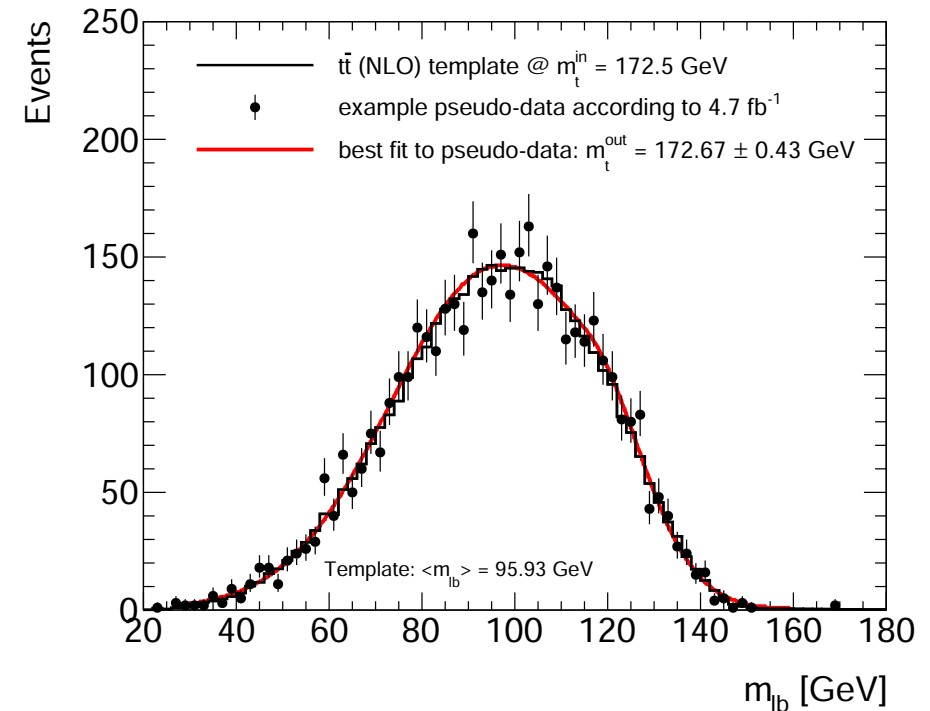
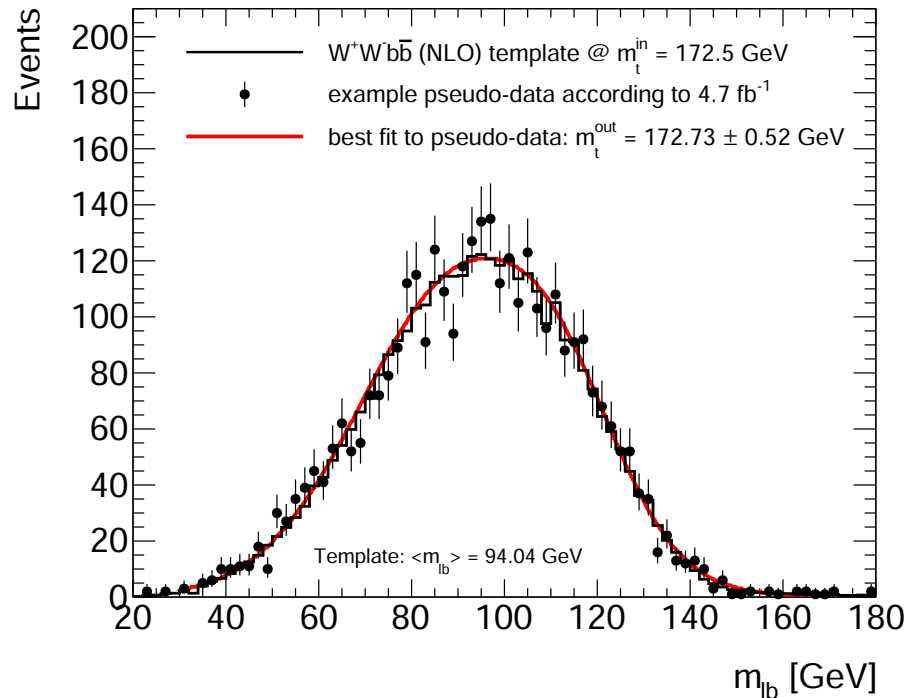
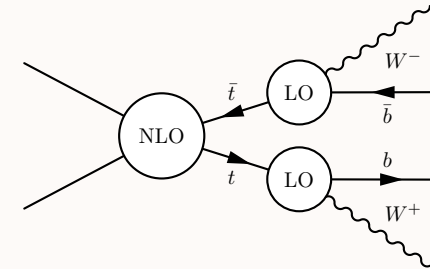
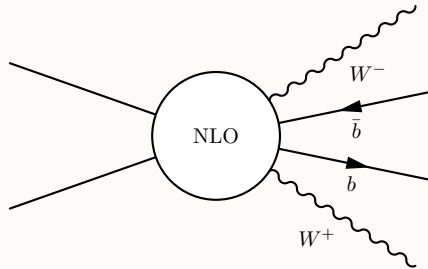
- Fit the pseudo-data with the calibration function → m_t^{out}

↪ “calibration”



NLO templates vs pseudo-data (same model)

➔ Representative examples for full (left) and factorized (right) NLO calculation.

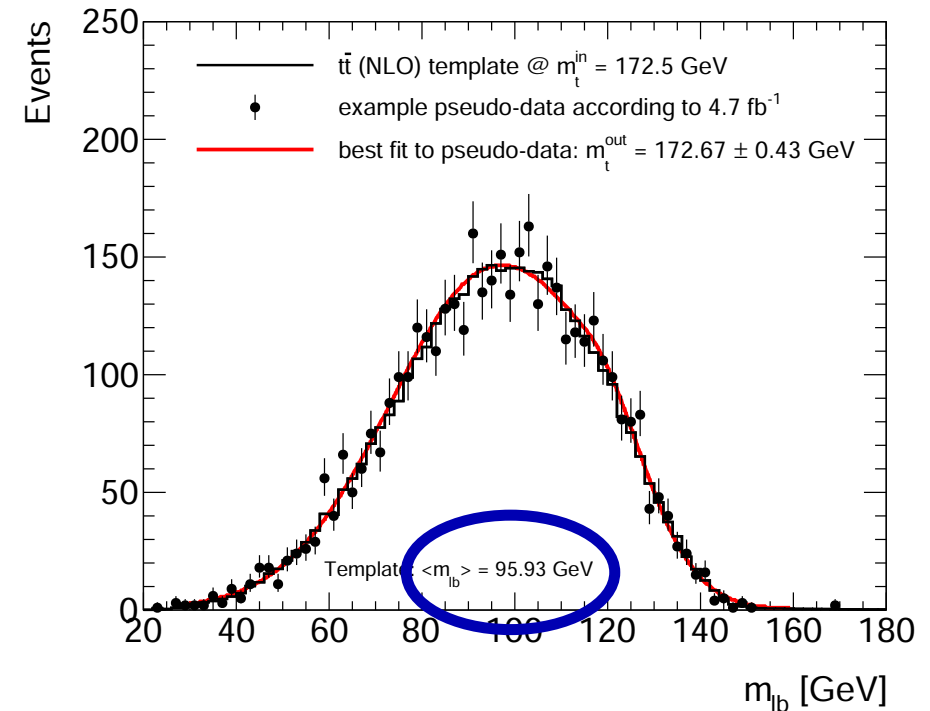
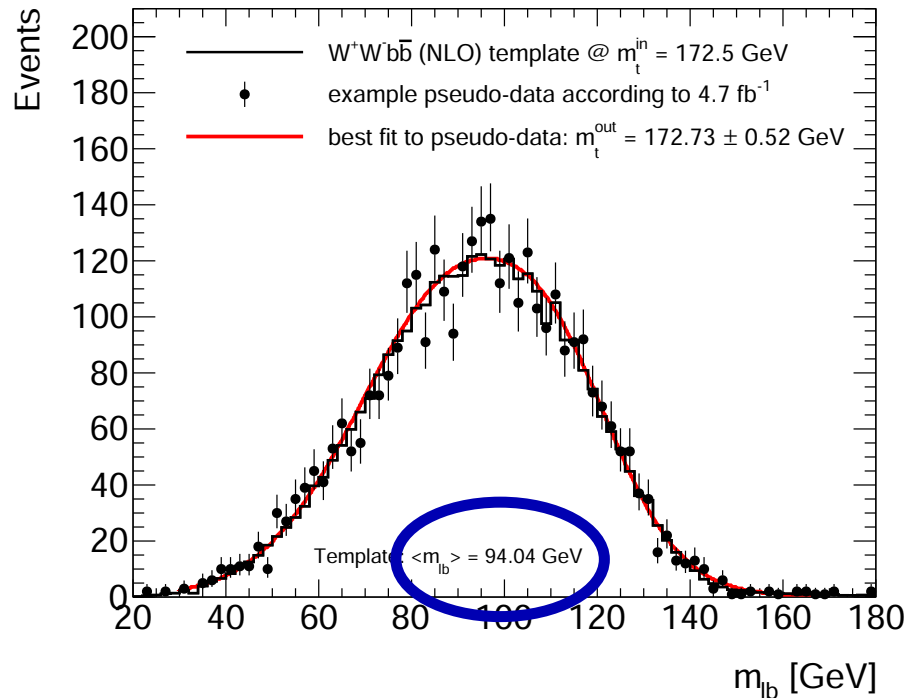
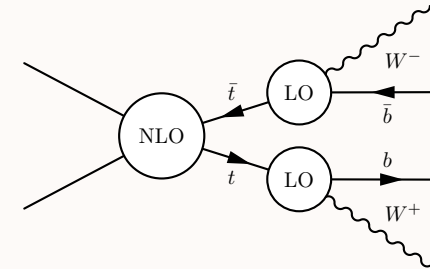
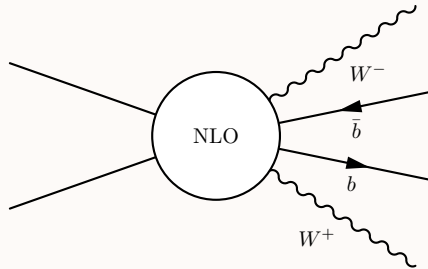


- pseudo-data (black points) are generated from the NLO distributions (black histograms) at m_t^{in}
- fit with NLO templates (parametrization) gives m_t^{out} → best fit to pseudo-data (red line)



NLO templates vs pseudo-data (same model)

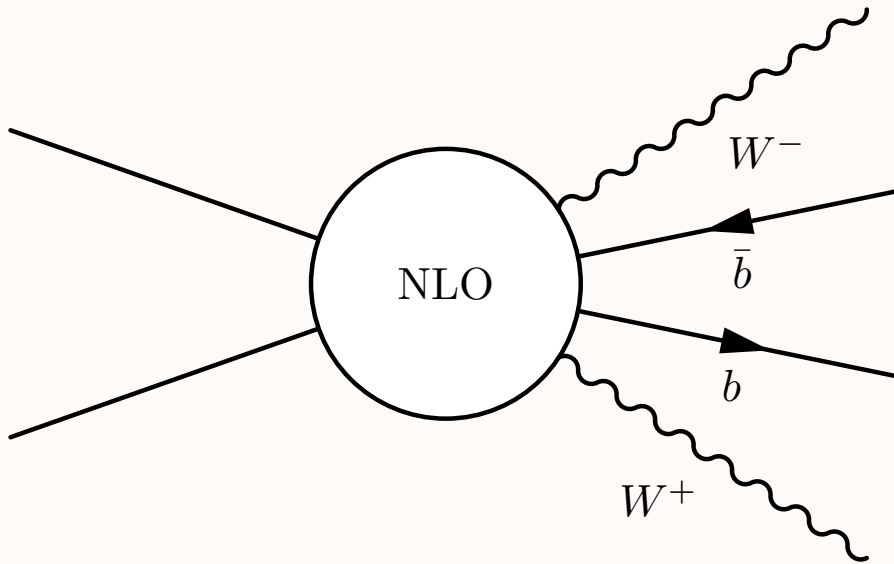
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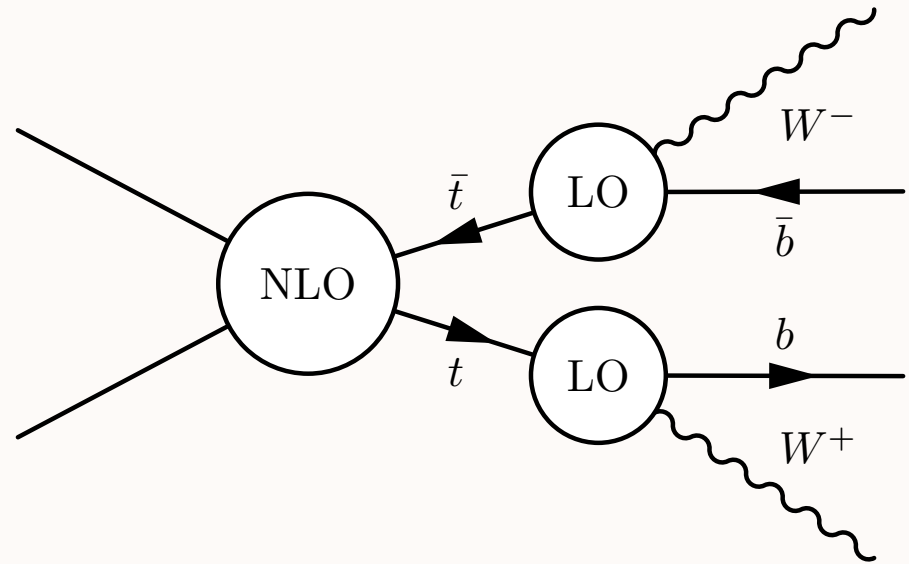


Full versus factorized approach



full ($WWb\bar{b}$)

- **full NLO** description of the $WWb\bar{b}$ final state ($2 \rightarrow 4$ processes)
- accounts for non-resonant/non-factorizing contributions, includes NLO effects in top quark decays



factorized ($t\bar{t}$)

- **NLO $t\bar{t}$** production ($2 \rightarrow 2$ processes) with **LO decays** attached and spin correlations preserved
- standard description for the NLO core in NLO+PS matching

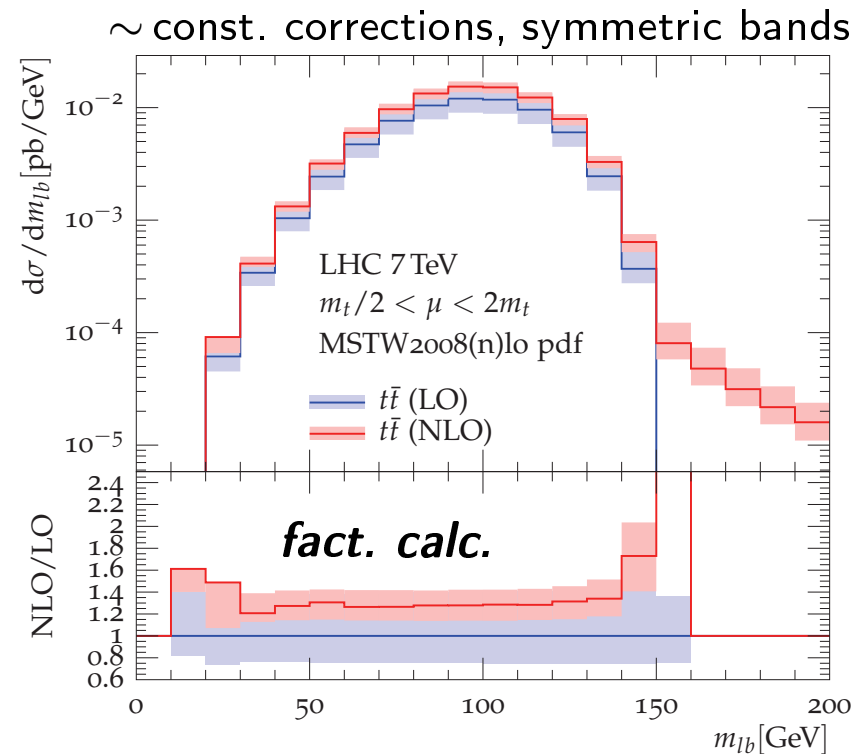
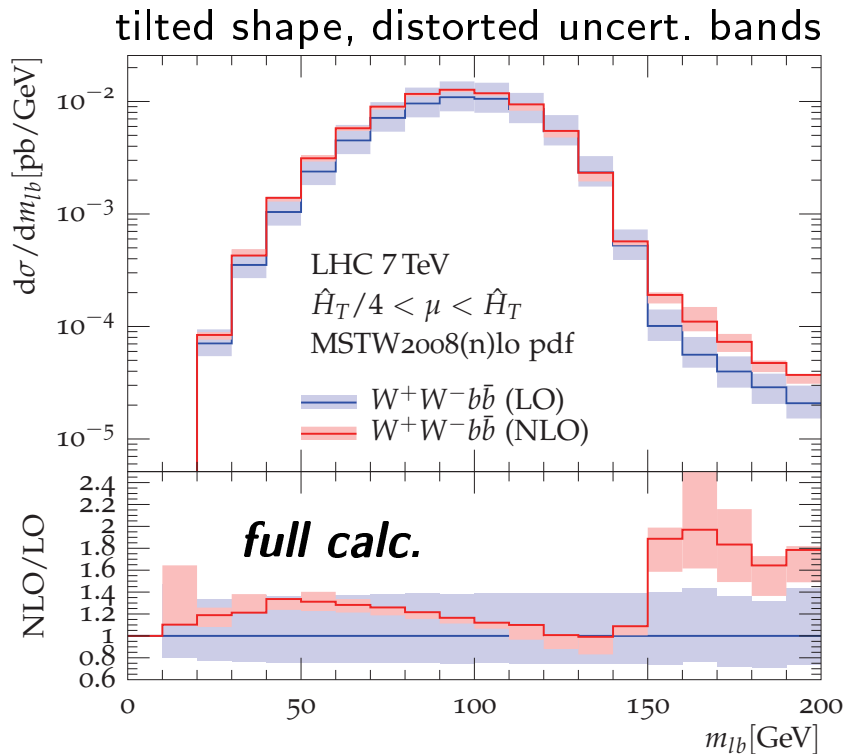
(Use these calculations for pure parton level analyses, i.e. m_t is no MC mass here, it is the pole mass.)



The m_{lb} distribution at NLO and scale variations

➔ *Parton-level NLO calculations for $W^+W^-b\bar{b}$ based on GoSam+Sherpa framework.*

(full & factorized calc., 5-flavour scheme, massless b-quarks, two resonant W decaying leptonically @ LO)



- Important NLO corrections to the shape of m_{lb}
- Values of m_{lb} larger than $\sqrt{m_t^2 - m_W^2}$ are kinematically forbidden in narrow width approximation at LO

• follow ATLAS strategy: use charged-lepton b-jet pairing minimizing sum of both m_{lb} and average.

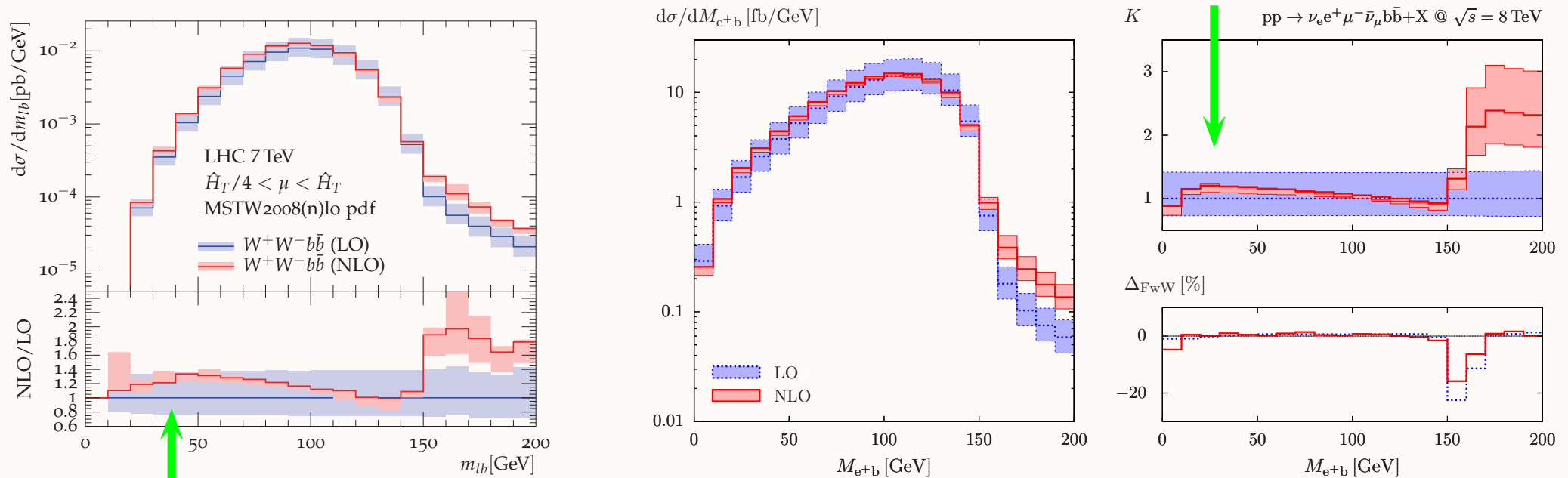


Qualitative comparison of m_{lb} predictions

[DENNER, DITTMAYER, KALLWEIT, POZZORINI, ARXIV:1207.5018]

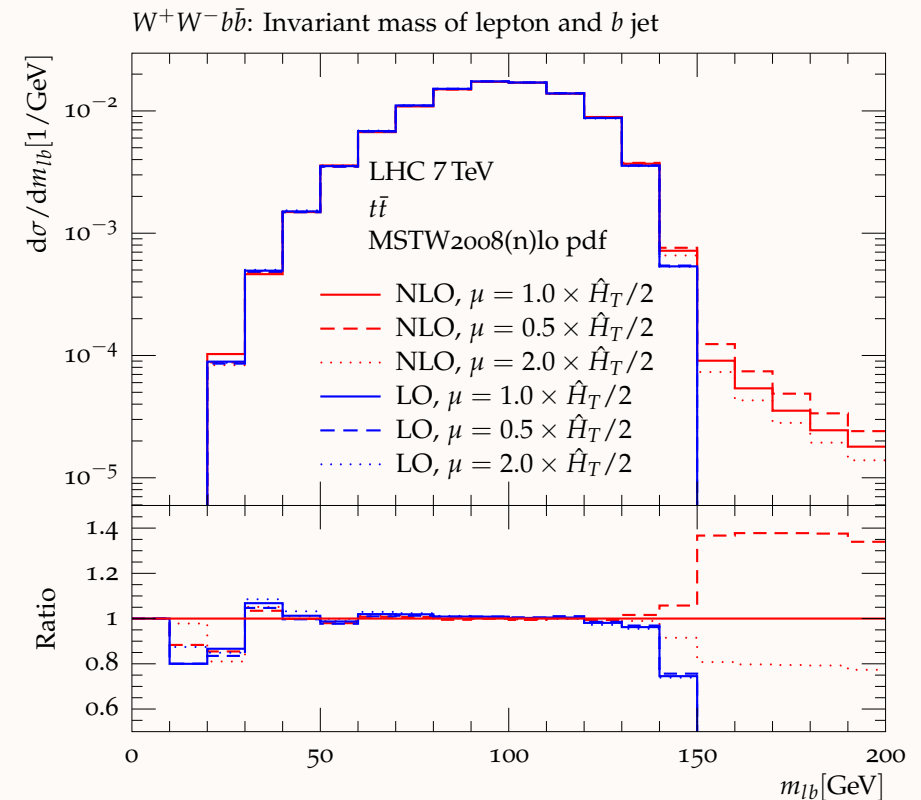
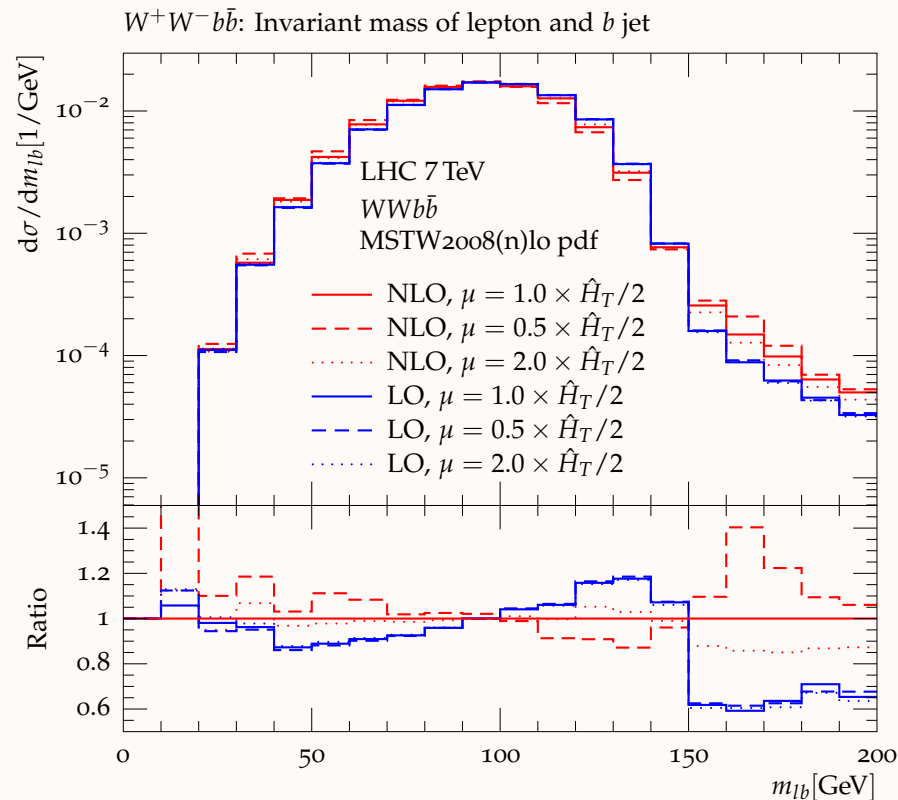
➔ $WWb\bar{b}$: NLO corrections strongly affect the shape of m_{lb}

- similar features → agreement on qualitative level only, noting the differences however:
- different LHC energies & kinematical constraints (cuts), slightly different observable (a truth m_{lb})
- different dynamical scale choice (transverse mass of tops)
- non-resonant and off-shell effects due to finite W boson width
- different treatment of b-quark initial states



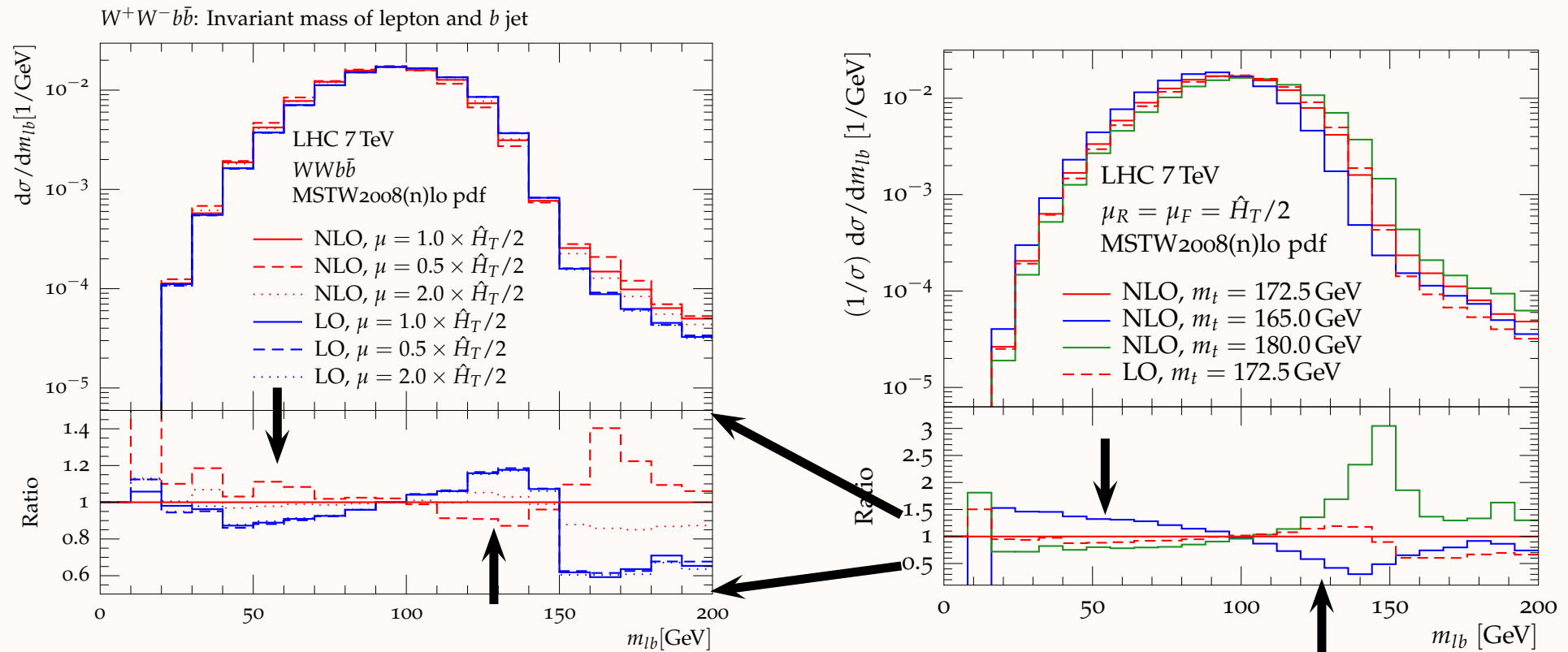
Normalized m_{lb} : scale versus m_t variation

- shape modifications resulting from variation of scales by factors of two
- left panel, for the full approach $\rightarrow$ visible • right panel, for the factorized approach $\rightarrow$ only in tails



Normalized m_{lb} : scale versus m_t variation

- shape modifications resulting from variation of scales by factors of two
- left panel, for the full approach → visible • right panel, shape changes due to m_t variation @ NLO
(m_t variation @ LO very similar!)



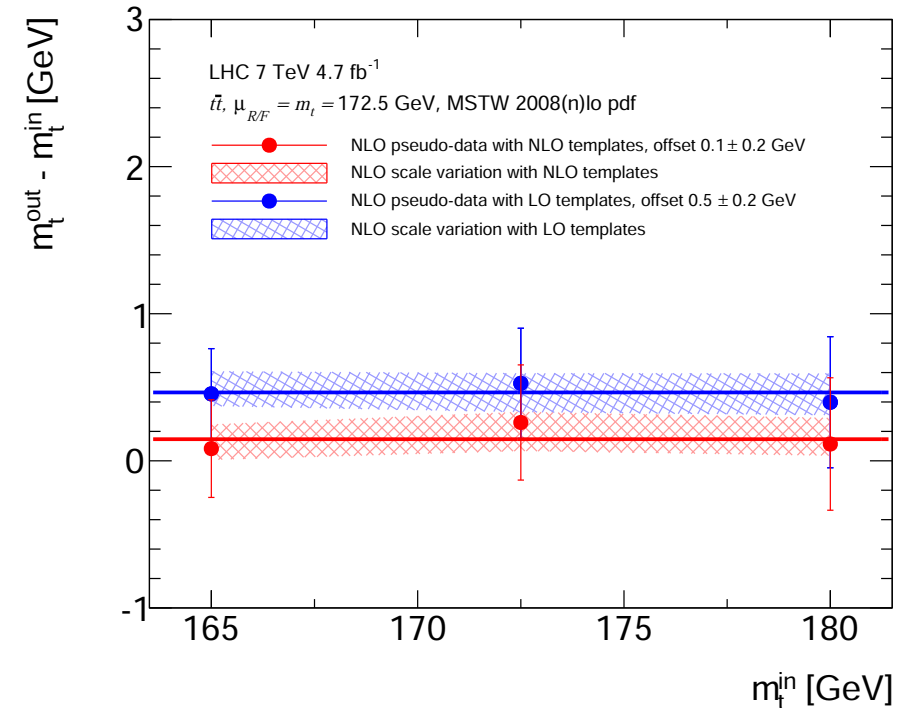
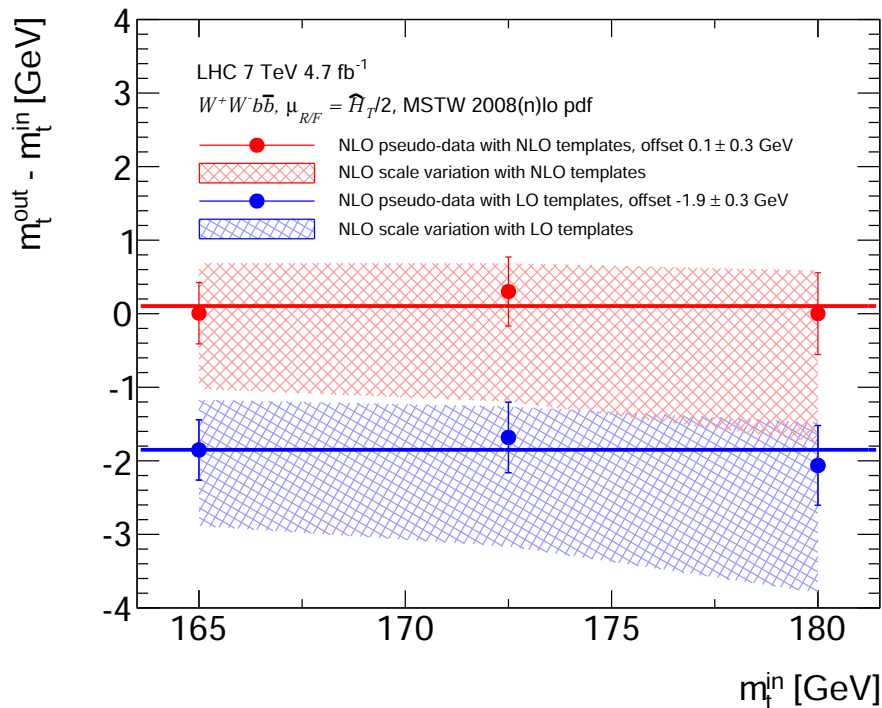
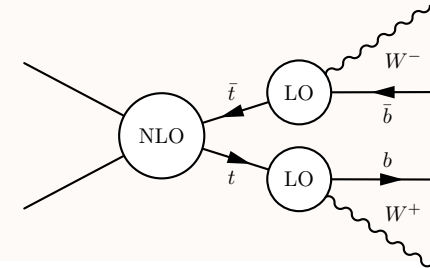
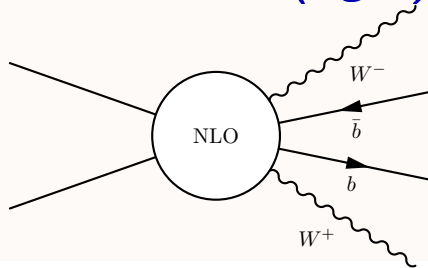
- scale factor variation mimics shape changes as induced by different m_t values → uncertainty
- @NLO: scale down corresponds to lower mass
- fit mass and scale simultaneously, but would resulting choice work for other distributions (eg. $m_{t\bar{t}}$)?



NLO templates vs pseudo-data (different models)

[HEINRICH, MAIER, NISIUS, SCHLENK, WINTER, ARXIV:1312.6659]

➔ Full (left) vs factorized (right) NLO calculation: results for mass shifts.



Use m_{lb} method in parton-level analysis assuming that data follows full / factorized theory [pseudo-data]. Test against the templates given by default scale choice **NLO** and **LO** theories [hypotheses].

- larger shift btwn NLO & LO description (~ 1.9 GeV) as compared to factorized approach (~ 0.5 GeV)
- significantly larger uncertainties from scale variations for full approach ($^{+0.6}_{-1.0}$ GeV vs ± 0.2 GeV)



Motivation for new studies.

(- mostly emerging from the open questions after the earlier studies.)

Larger scale uncertainties driven by radiative corrections in decays, or non-resonant/non-factorizing contributions?

Parton shower effects? To what extent are radiative decay corrections well modelled by shower in NLO+PS?

-> Include more calculations.

Observables with reduced scale dependence, and good top quark mass sensitivity?

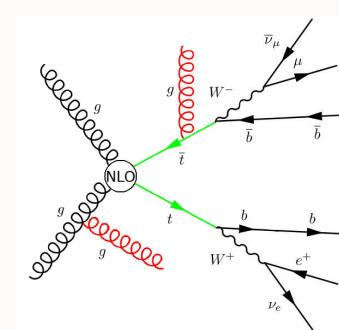
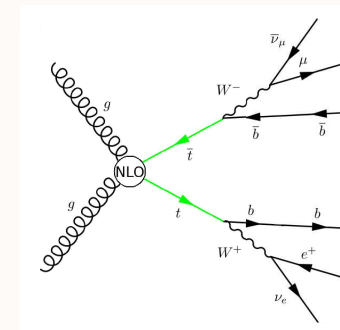
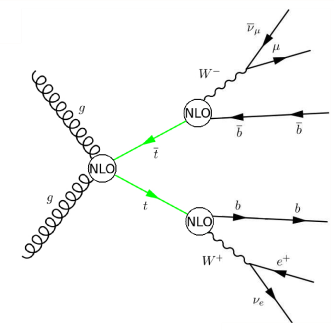
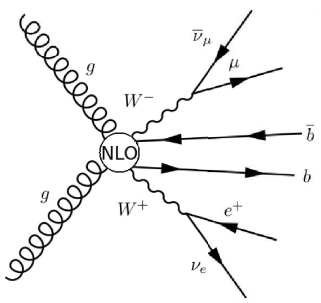
-> Study other observables/analysers than m_{lb} .

Different experimental setup:

LHC Run 2: higher center-of-mass energy, more luminosity, 'more gluons'

Active field of research (Nason et al, Bevilacqua et al, Hoang et al, Denner et al, ...)





Calculations used in the new studies.

NLO_{full} : full QCD NLO corrections to $WbW\bar{b}$ production with leptonic W decays

$\text{NLO}_{\text{NWA}}^{\text{NLOdec}}$: NLO $t\bar{t}$ production and NLO decays in narrow-width approximation

$\text{NLO}_{\text{NWA}}^{\text{LOdec}}$: NLO $t\bar{t}$ production and LO decays in narrow-width approximation

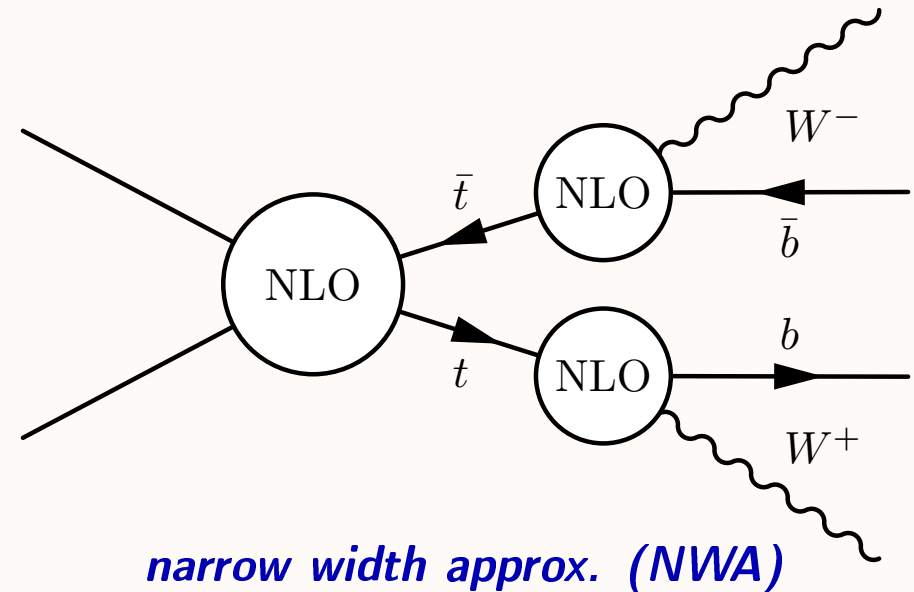
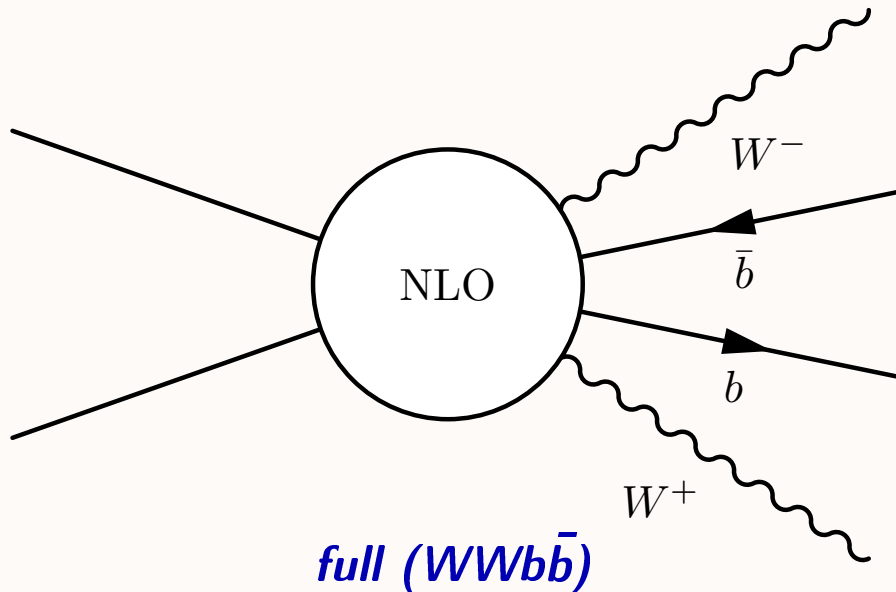
NLO_{PS} : NLO $t\bar{t}$ production plus parton shower and decays via parton showering

Also available here: associated LO calculations: LO_{full} , $\text{LO}_{\text{NWA}}^{\text{LOdec}}$, LO_{PS}

note NWA theory state-of-the-art: $\text{NNLO}_{\text{NWA}} + \text{NLO}_{\text{EW}}$, NNLO_{dec}

note shower state-of-the-art: $\text{NLO}_{\text{full}} + \text{PS} \rightarrow$ see T. Jezo's talk.

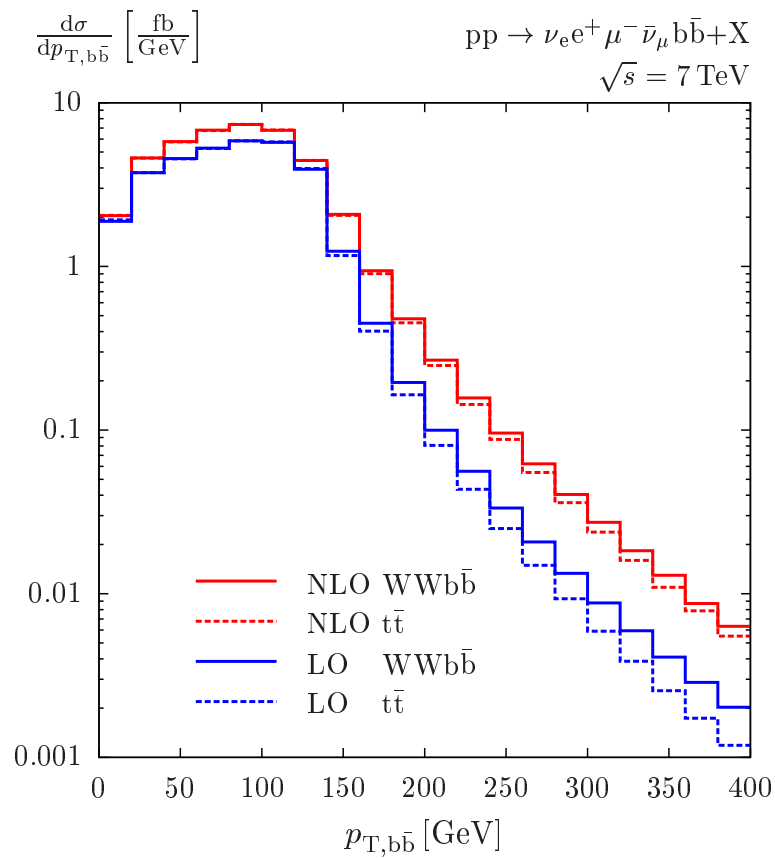
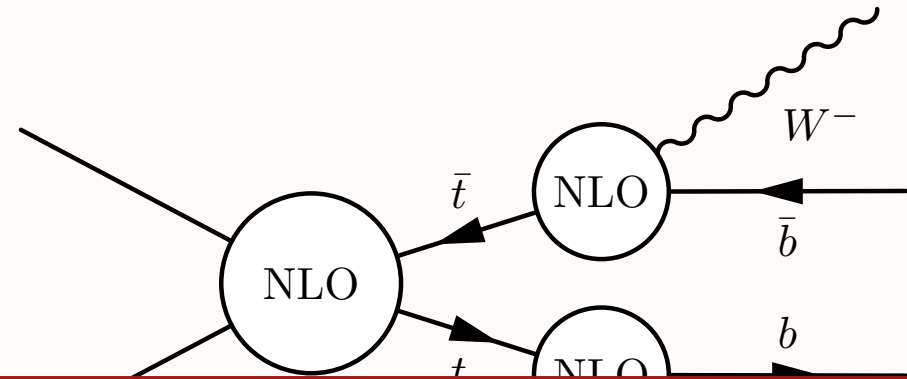
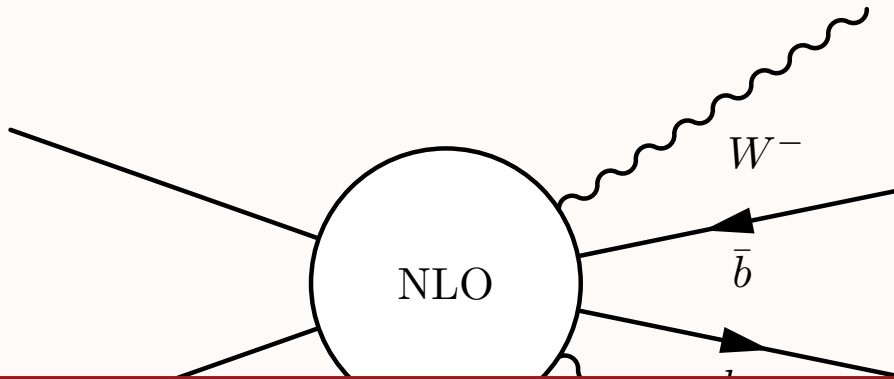
Full calculation versus NWA



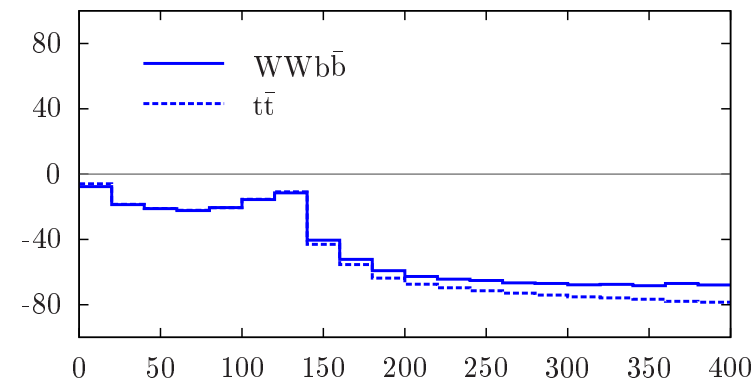
- full NLO description of the $WWb\bar{b}$ final state ($2 \rightarrow 4$ processes)
 - non-resonant/-factorizing contributions (quantum interferences)
 - NLO effects in top quark decays
- $\ell\ell$ channel comparison: [DENNER, DITTMAYER, KALLWEIT, POZZORINI, SCHULZE, LESHOUCHES2011, ARXIV:1203.6803]
 - no more than 1% deviations for inclusive cross sections (with experimental cuts), but effects can be (significantly) larger in differential distributions.
- Here, comparison will help disentangle effects from NLO decays and non-resonant/-factorizing contributions.



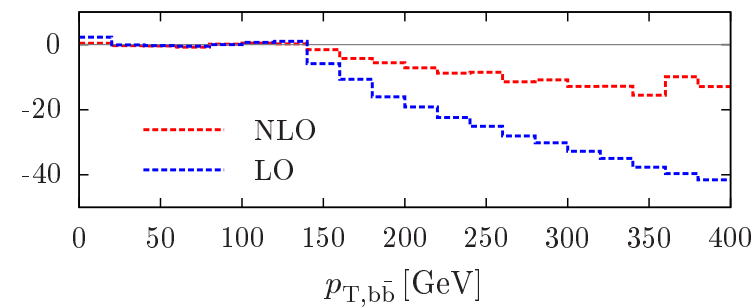
Full calculation versus NWA



LO/NLO - 1 [%]

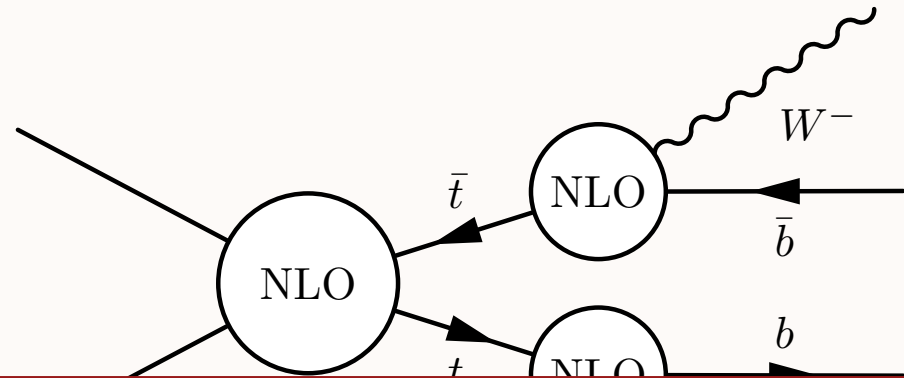
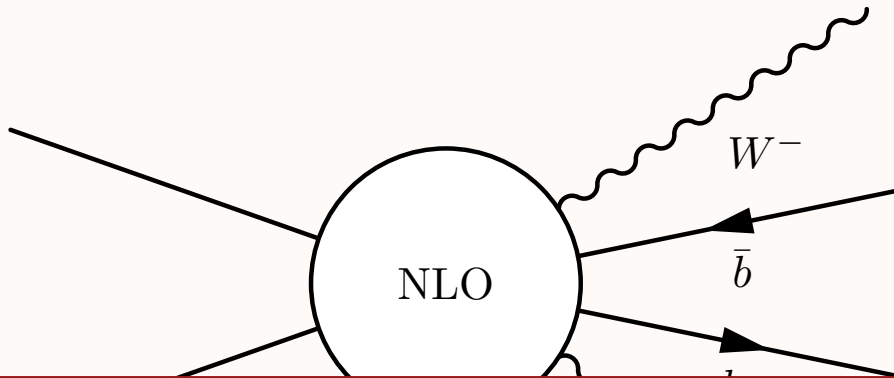


$t\bar{t}/WWb\bar{b} - 1$ [%]



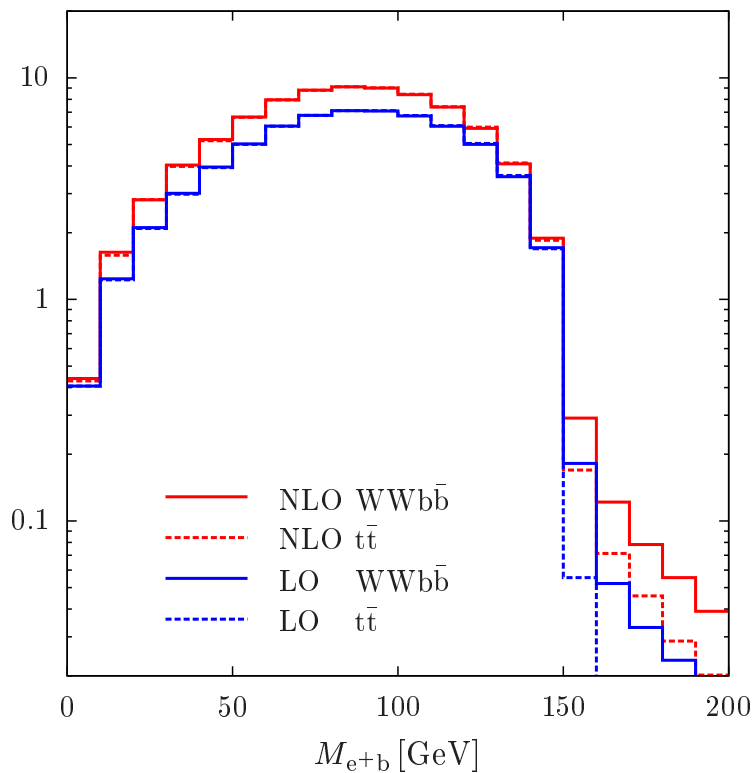
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Full calculation versus NWA

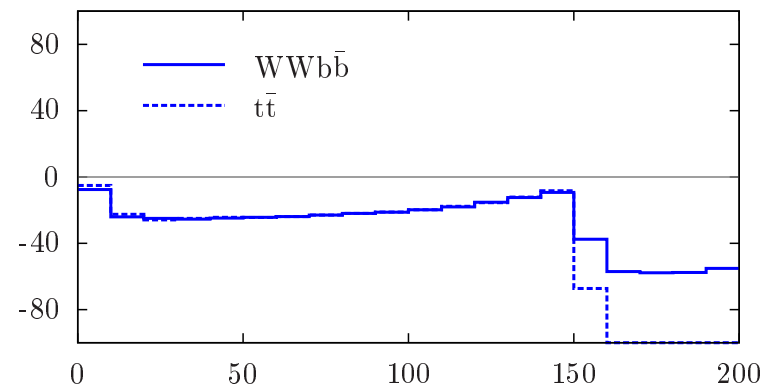


$$\frac{d\sigma}{dM_{e^+,b}} \left[\frac{\text{fb}}{\text{GeV}} \right] \quad pp \rightarrow \nu_e e^+ \mu^- \bar{\nu}_\mu b \bar{b} + X$$

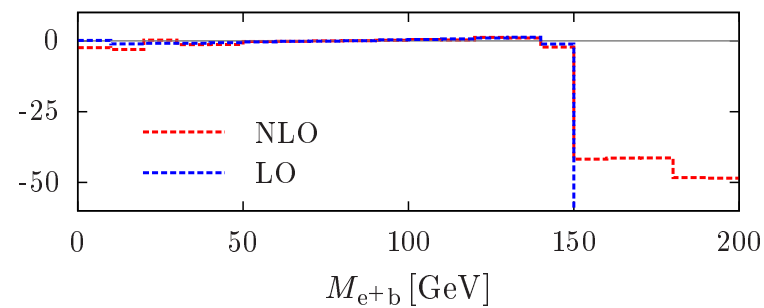
$$\sqrt{s} = 7 \text{ TeV}$$



$$\text{LO/NLO} - 1 [\%]$$



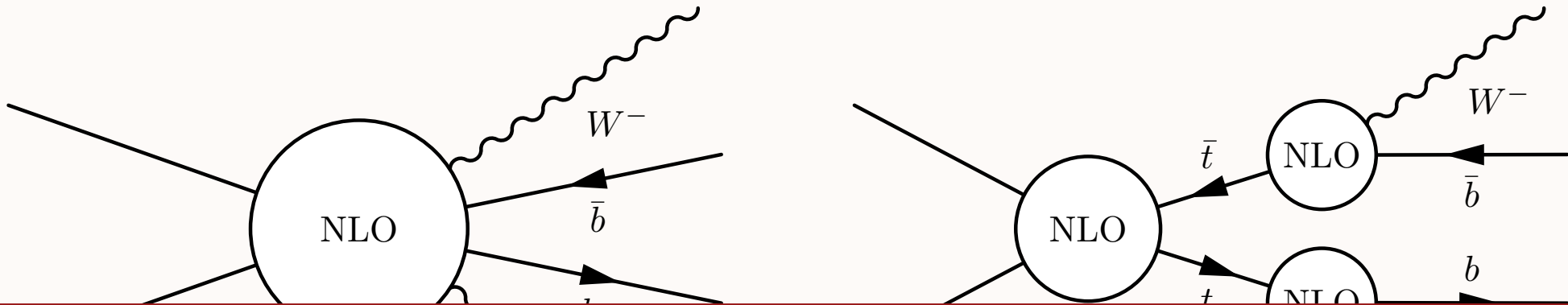
$$t \bar{t} / WWb \bar{b} - 1 [\%]$$



- Here, comparison will help disentangle effects from NLO decays and non-resonant/-factorizing contributions.



Full calculation versus NWA



➔ *Remark: state-of-the-art NNLO NWA.*

- On-shell $t\bar{t}$ NNLO QCD corrections to differential distributions, NLO EW corrections

[Czakon, Heymes, Mitov '15], [Czakon, Heymes, Mitov, Pagani, Tsinikos, Zaro '17], [Hollik, Pagani '11],
[Kühn, Scharf, Uwer '13]

- Top decay: NNLO+NNLL QCD corrections

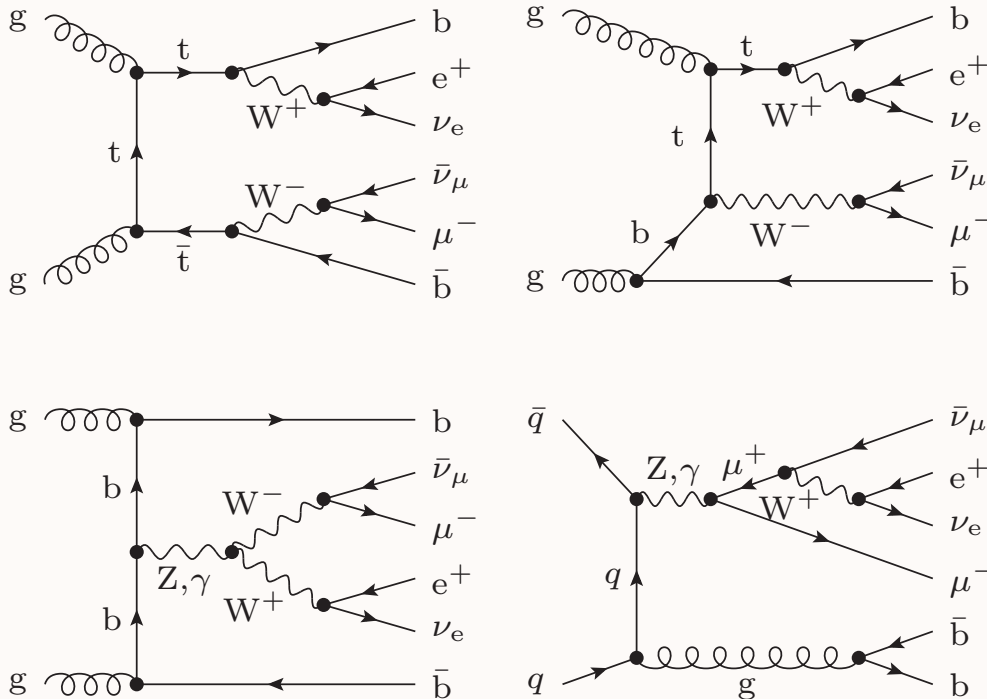
[Beneke, Falgari, Klein, Schwinn '11], [Cacciari, Czakon, Mangano, Mitov, Nason '11], [Pecjak, Scott, Wang, Yang '16], [Ferrogli, Marzani, Pecjak, Yang '13], [Broggio, Papanastasiou, Signer '14], [Kidonakis '15], [Gao, Papanastasiou '17]



● Here, comparison will help disentangle effects from NLO decays and non-resonant/-factorizing contributions.

$WWb\bar{b}$ production at NLO in QCD

➔ *Going beyond $t\bar{t}$ production to include offshell, non-factorizing and b-mass effects.*



- earlier done in NWA ($\Gamma_t \rightarrow 0$ limit) where production and decay factorize (neglected contributions are suppressed by powers of $\Gamma_t/m_t \lesssim 1\%$)

[BERNREUTHER, BRANDENBURG, SI, UWER, ARXIV:HEP-PH/0403035]

[MELNIKOV, SCHULZE, ARXIV:0907.3090]

[BISWAS, MELNIKOV, SCHULZE, ARXIV:1006.0910]

- full NLO treatment includes double-, single- and non-resonant contributions (DR – $t\bar{t}$ -like) (SR – Wt -like) (NR – VV -like)
- complex-mass scheme
- finite top-quark and W width effects (offshell DR, SR, NR and interferences)

- first done in massless b-quark approximation ($\rightarrow$ requires two hard b-jets) ($\rightarrow WWb\bar{b}$ in 5-flavour scheme)

[DENNER ET AL. ARXIV:1012.3975, ARXIV:1207.5018]

[BEVILACQUA ET AL. ARXIV:1012.4230]

[HEINRICH ET AL. ARXIV:1312.6659]

- also available in 4-flavour scheme

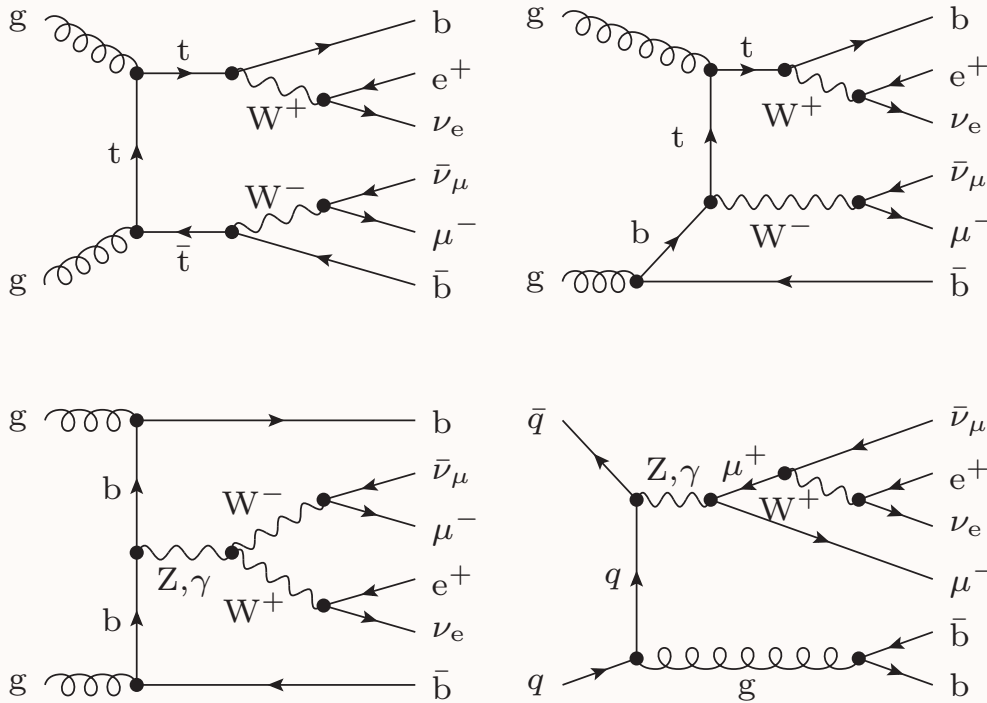
[FREDERIX, ARXIV:1311.4893] [CASCIOLI, KALLWEIT,

MAIERHÖFER, POZZORINI, ARXIV:1312.0546]

$$\lim_{\Gamma_t/m_t \rightarrow 0} \frac{1}{(p_t^2 - m_t^2)^2 + m_t^2 \Gamma_t^2} = \frac{\pi}{m_t \Gamma_t} \delta(p_t^2 - m_t^2)$$



$WWb\bar{b}$ production at NLO



- full NLO treatment includes double-, single- and non-resonant contributions
- complex-mass scheme
- finite top-quark and W width effects
- first done in massless b-quark approximation
[DENNER ET AL. ARXIV:1012.3975, ARXIV:1207.5018]
[BEVILACQUA ET AL. ARXIV:1012.4230]
- earlier done in NWA where production and decay factorize (neglected contributions are suppressed by powers of $\Gamma_t/m_t \lesssim 1\%$)
[BERNREUTHER ET AL. ARXIV:HEP-PH/0403035]
[MELNIKOV, SCHULZE, ARXIV:0907.3090]

→ Our parton level calculations ...

- use the **GoSam+Sherpa** combined generator package (current versions, GoSam 2.0 and Sherpa 2.2.3).
- **Sherpa** for calculating Born, real corrections and infrared subtractions [GLEISBERG ET AL, ARXIV:0811.4622]
- **GoSam** for calculating virtual corrections [CULLEN, VANDEURZEN, GREINER, HEINRICH ET AL, ARXIV:1404.7096]
- 5-flavour scheme, massless b-quarks, two resonant W decaying leptonically @ LO respecting spin correlations

⇒ [HEINRICH, MAIER, NISIUS, SCHLENK, WINTER, ARXIV:1312.6659]



Other observables

Top-quark mass $m_t^{(\text{MC})}$ extracted from fitting $\frac{d\sigma}{d\Omega}(m_t^{\text{MC}})$



- Differential distributions **particularly sensitive** to m_t :
 - $m_{\ell b}^2 = (p_\ell + p_b)^2$
 - $m_{T2}^2 = \min_{\mathbf{p}_T^{\nu 1} + \mathbf{p}_T^{\nu 2} = \mathbf{p}_T^{\text{miss}}} \left[\max \left\{ m_T^2 \left(\mathbf{p}_T^{(\ell b)1}, \mathbf{p}_T^{\nu 1} \right), m_T^2 \left(\mathbf{p}_T^{(\ell b)2}, \mathbf{p}_T^{\nu 2} \right) \right\} \right]$
 - $E_T^{\Delta R} = \frac{1}{2} \left(E_T^{\ell 1} \Delta R(\ell_1, b_1) + E_T^{\ell 2} \Delta R(\ell_2, b_2) \right)$
 - $m_{\ell\ell}^2 = (p_{\ell_1} + p_{\ell_2})^2$

where we choose the lepton- b -jet pairing (ℓ^+, b) , (ℓ^-, b') minimizing $m_{\ell+b} + m_{\ell-b'}$.

- Top quark mass sensitivity for $m_{\ell b}$ and m_{T2} is similar while it is lower for the other two variables.



- $\sqrt{s} = 13 \text{ TeV}$
- PDF4LHC15_nlo_30_pdfas
- Fastjet anti- k_T algorithm with $R = 0.4$
- EW parameters:

$$\begin{array}{ll} \Gamma_{\ell}^{\text{LO}} &= 1.4806 \text{ GeV} & \Gamma_{\ell}^{\text{NLO}} &= 1.3535 \text{ GeV}, \\ \Gamma_W^{\text{LO}} &= 2.0454 \text{ GeV} & \Gamma_W^{\text{NLO}} &= 2.1155 \text{ GeV}, \\ \Gamma_Z &= 2.4952 \text{ GeV} \end{array}$$

$$G_\mu = 1.16637 \cdot 10^{-5} \text{ GeV}^{-2}$$

$$M_W = 80.3850 \text{ GeV} \quad M_Z = 91.1876 \text{ GeV}$$

- 5FNS (massless b 's)

For the NLO_{full} , $\text{NLO}_{\text{NWA}}^{\text{LOdec}}$, NLO_{PS} calculations:

- Sherpa 2.2.3 interfaced to GoSam OLP

→ *Calculational/experimental setup.*

- Exactly two b -tagged jets with
 - $p_T^{\text{jet}} > 25 \text{ GeV}$ and $|\eta^{\text{jet}}| < 2.5$
- Exactly two oppositely charged leptons with
 - $p_T^\mu > 28 \text{ GeV}$, $|\eta^\mu| < 2.5$
 - $p_T^e > 28 \text{ GeV}$, $|\eta^e| < 2.47$ with the exclusion of $1.37 < |\eta^e| < 1.52$
 - $\Delta_R(\ell, \text{jets}) > 0.4$
- $p_T^{\ell b} > 120 \text{ GeV}$

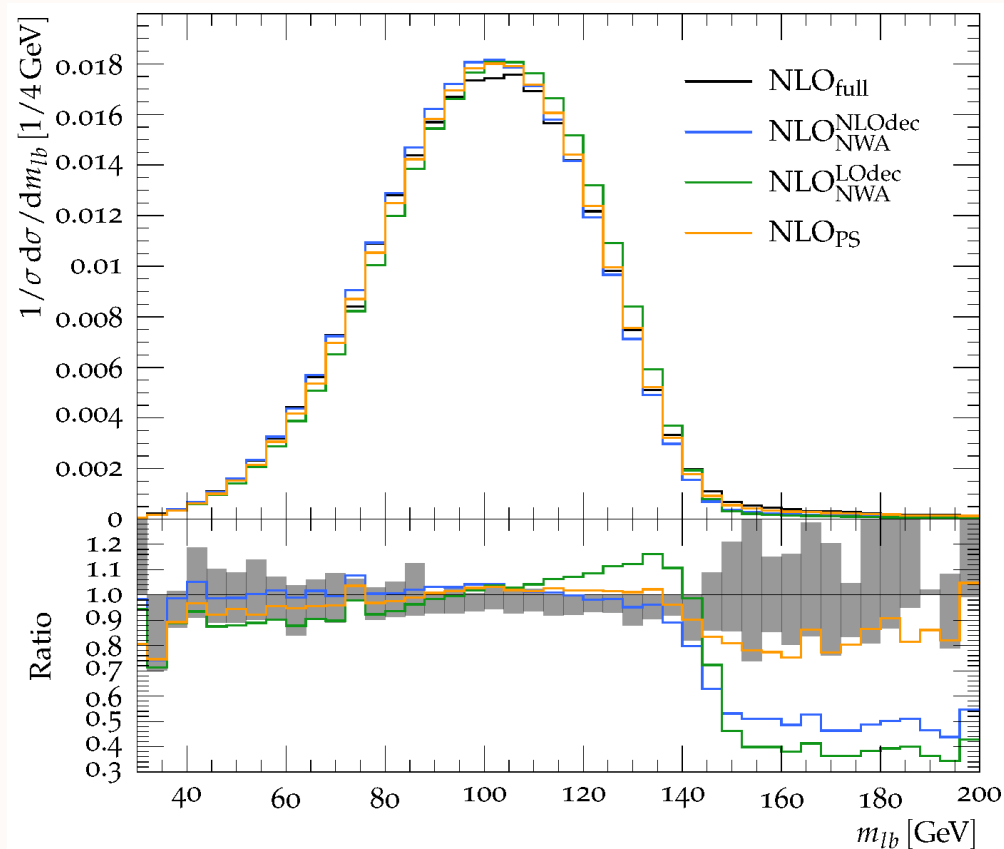


Dilepton final states at NLO

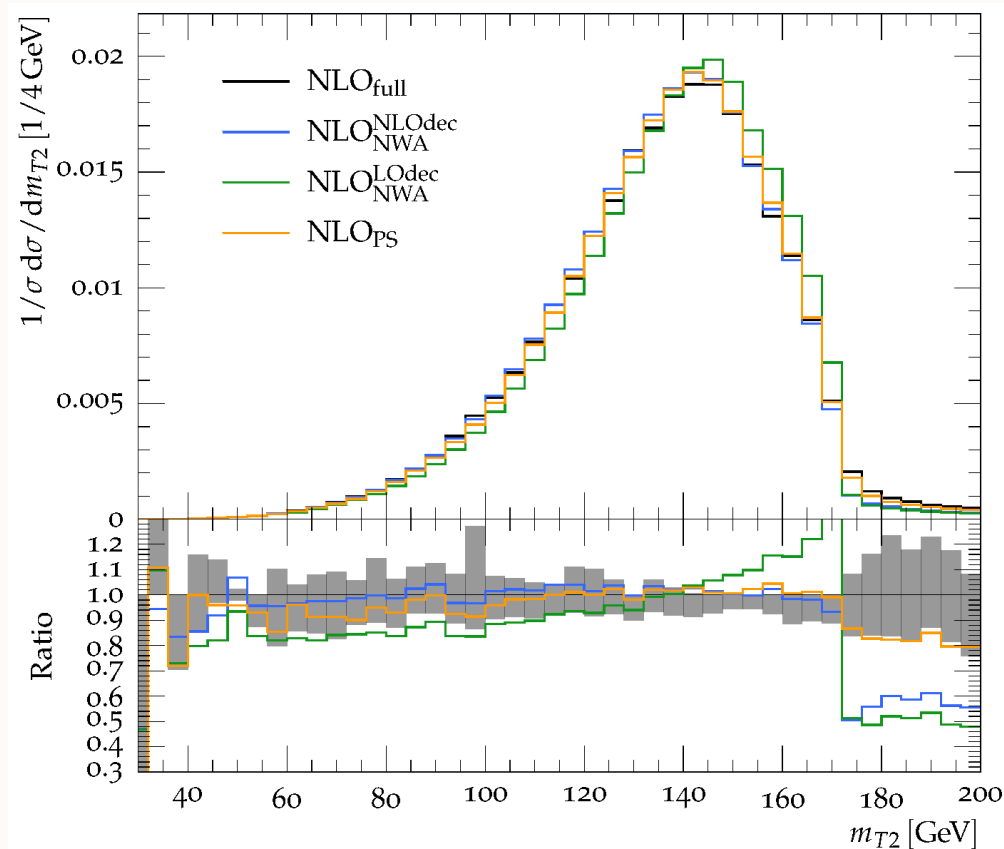
[HEINRICH ET AL, ARXIV:1709.08615]

➔ *Comparison of different theory predictions for analyser shape observables at LHC 13 TeV.*

● m_{lb} analyser



● m_{T2} analyser



- scale uncertainty band for full NLO calculation included in ratio plot
- top quark mass sensitive observables with ‘reasonable’ theory uncertainties

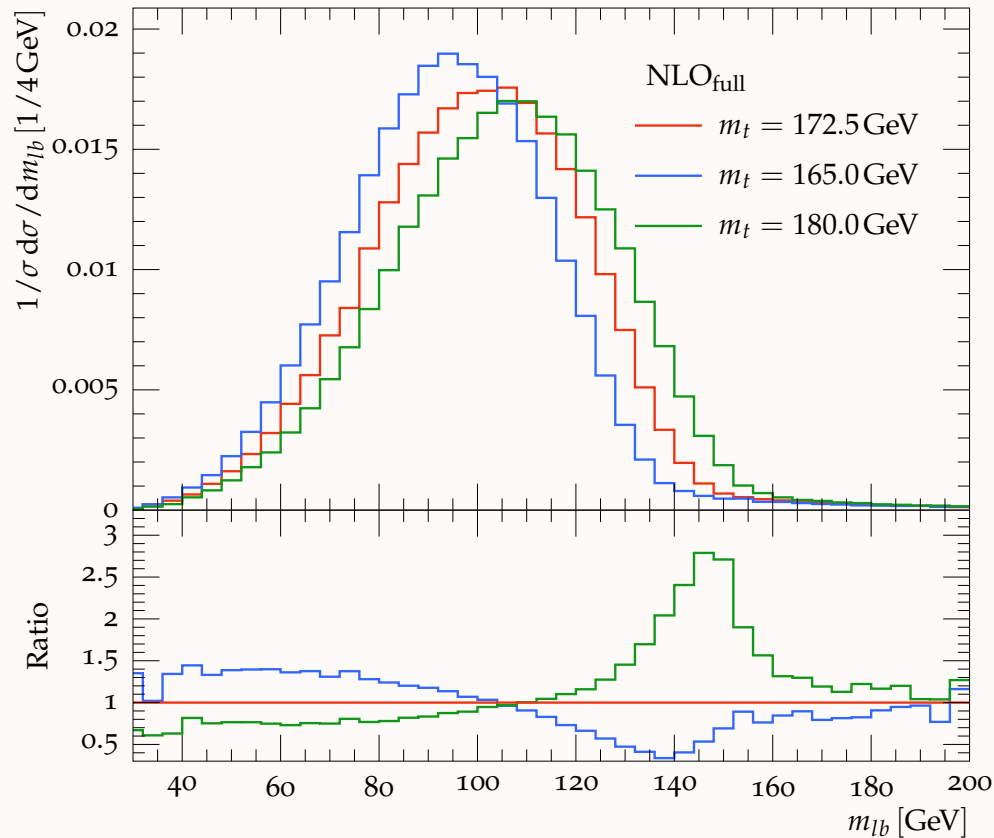


Top quark mass variation

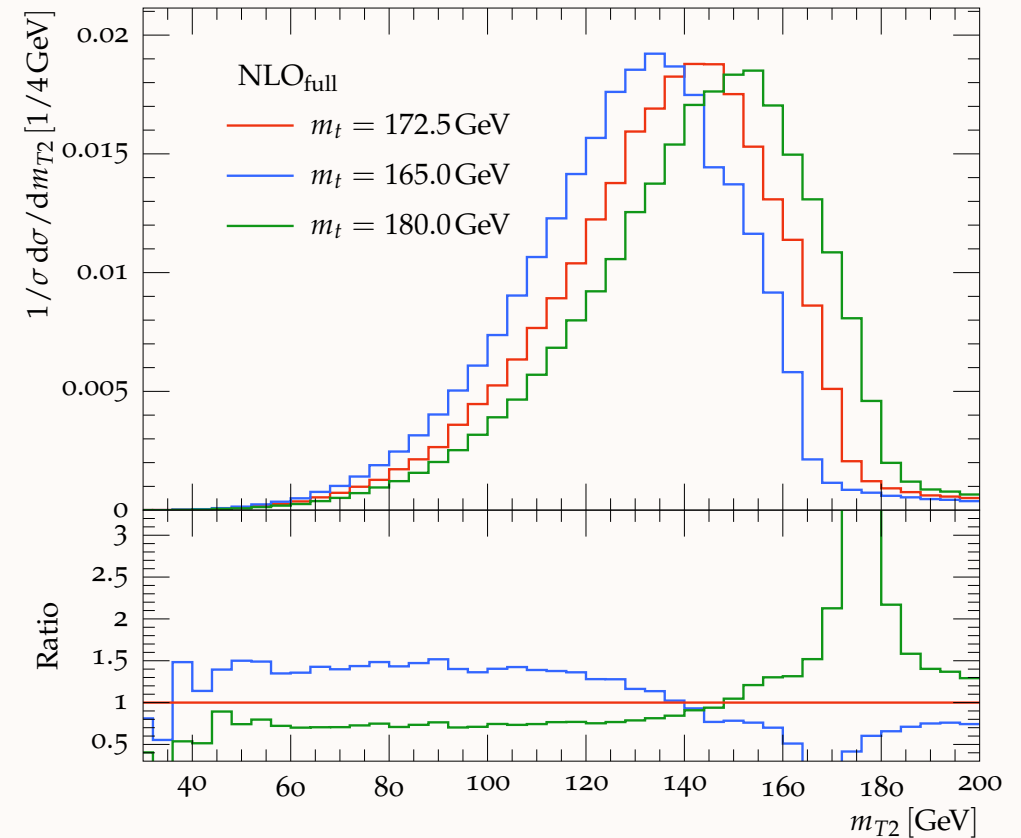
[HEINRICH ET AL, ARXIV:1709.08615]

➔ *Shape comparison of full NLO predictions for analyser shape observables at LHC 13 TeV.*

● m_{lb} analyser



● m_{T2} analyser

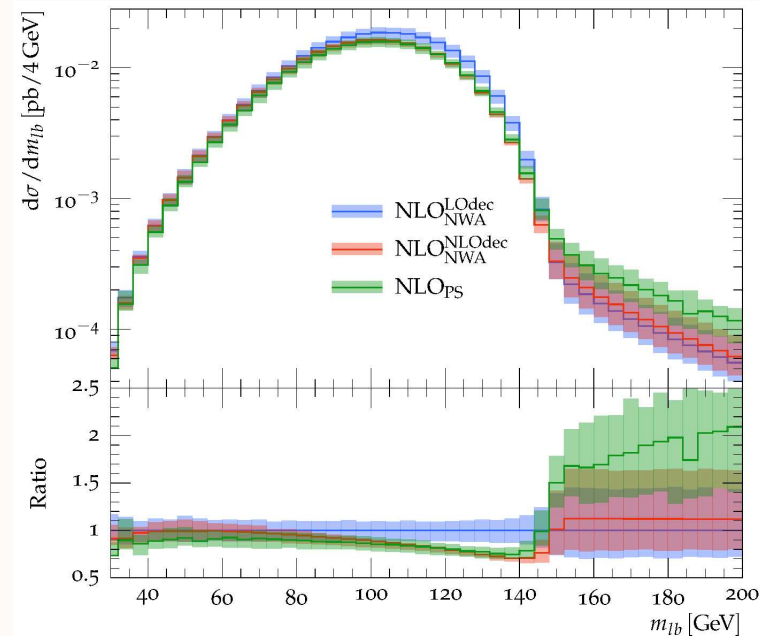
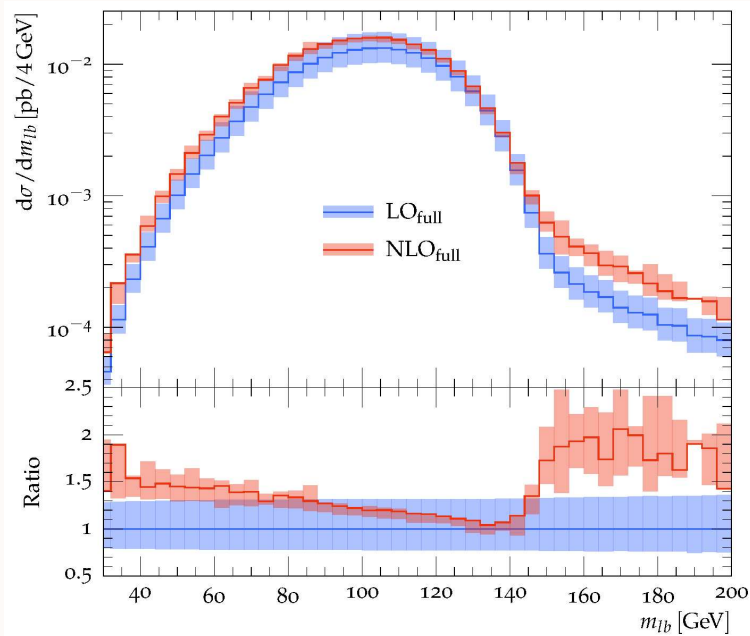


- harder shapes for increasing top quark masses
- top quark mass sensitive observables with ‘reasonable’ theory uncertainties

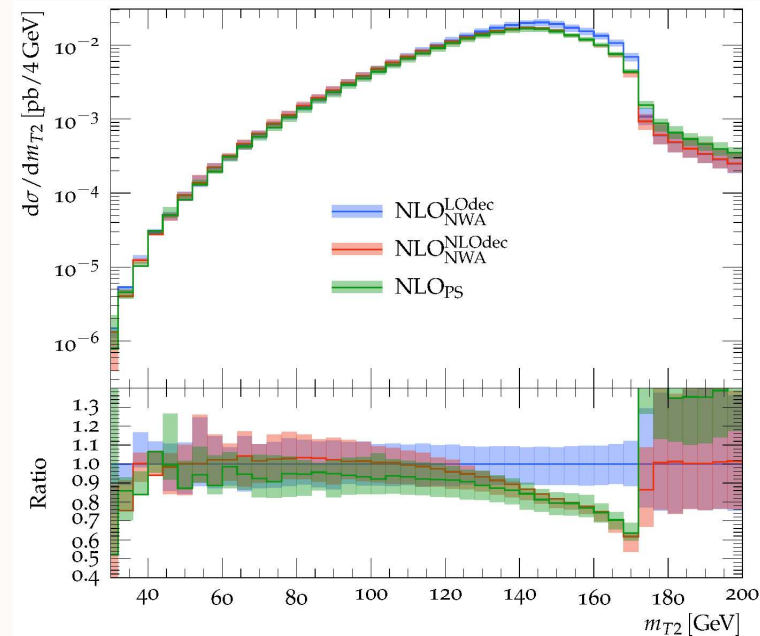
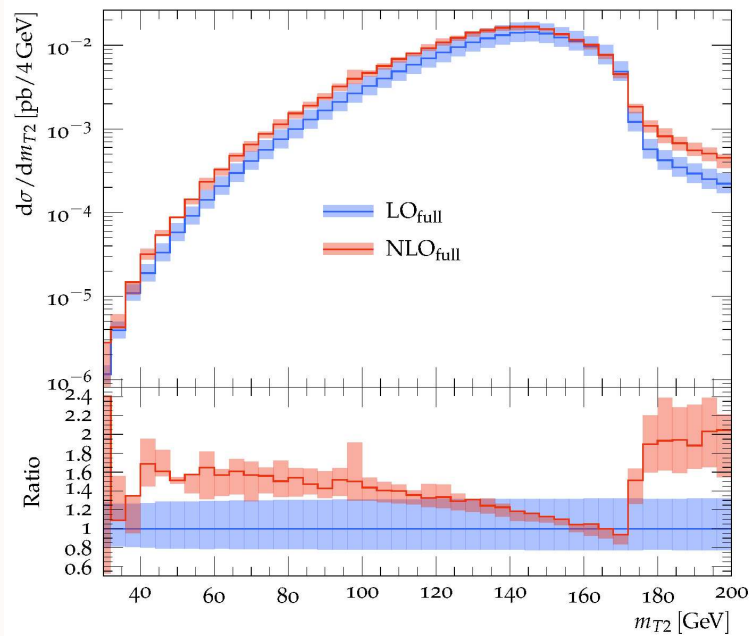


Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]



• m_{l_b} analyser

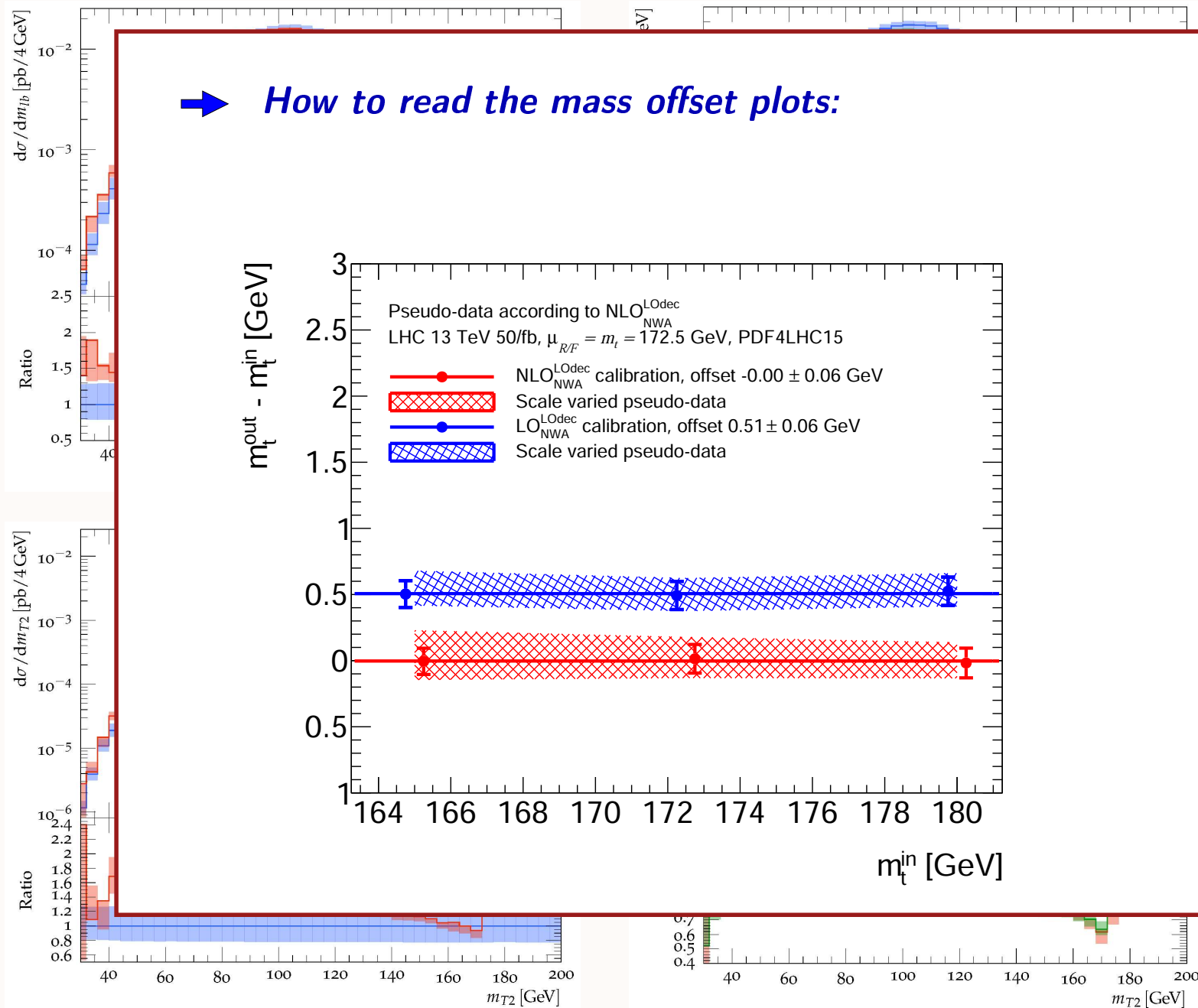


• m_{T2} analyser



Scale dependence

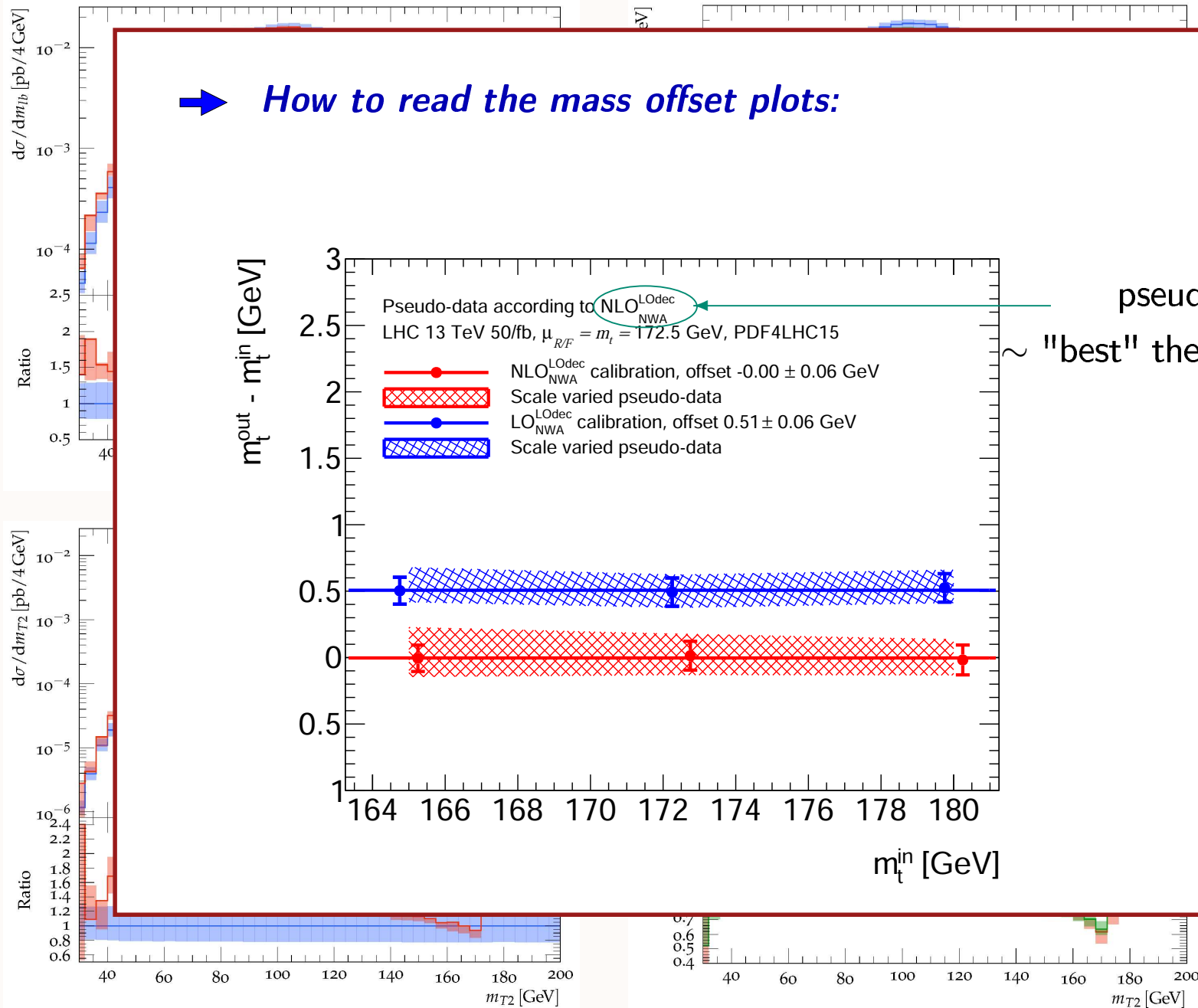
[HEINRICH ET AL, ARXIV:1709.08615]



Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]

➔ *How to read the mass offset plots:*



pseudo-data set
 ~ "best" theoretical model

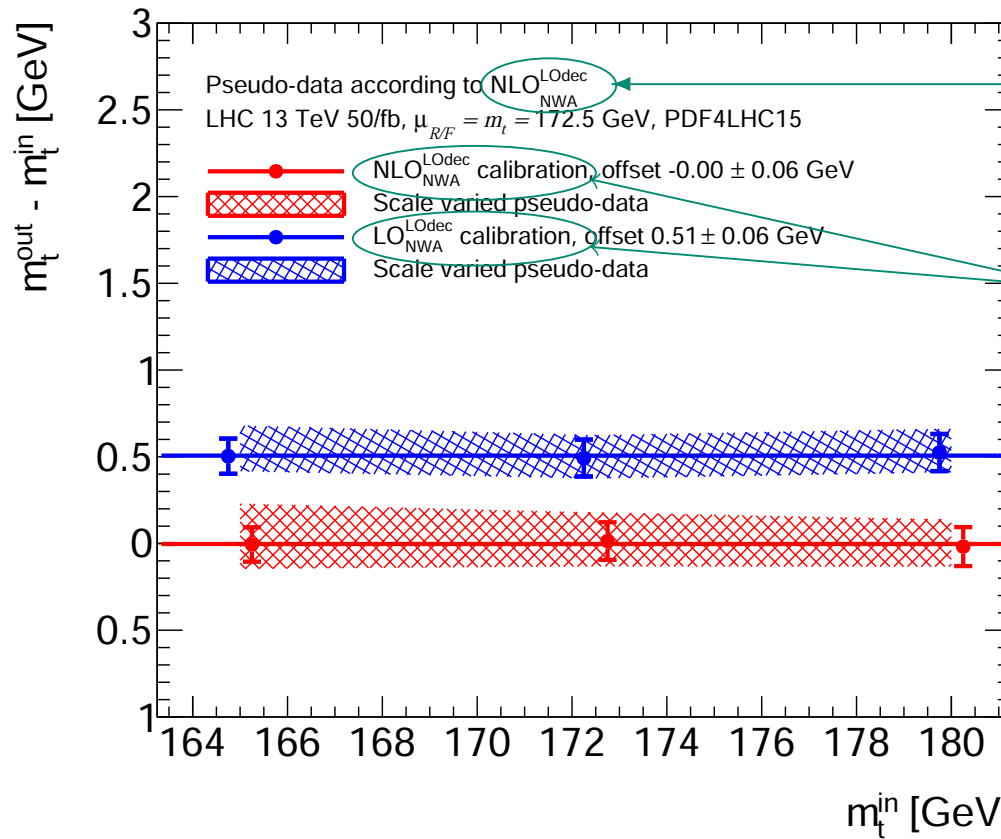


Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]

➔ *How to read the mass offset plots:*

Fit range: $40 \text{ GeV} \leq m_{lb} \leq 160 \text{ GeV}$



pseudo-data set

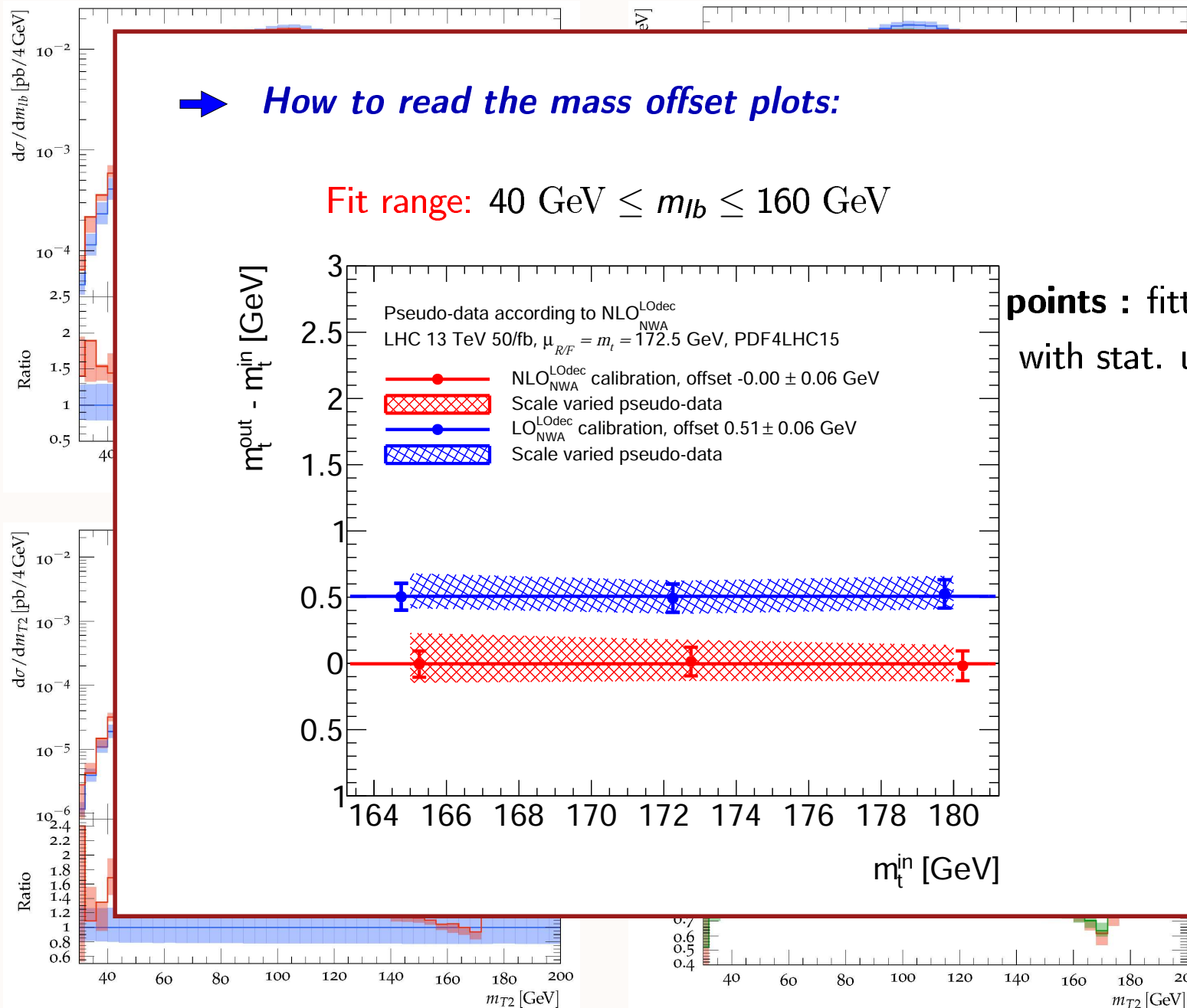
~ "best" theoretical model

theoretical model used for
the calibration fit function



Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]

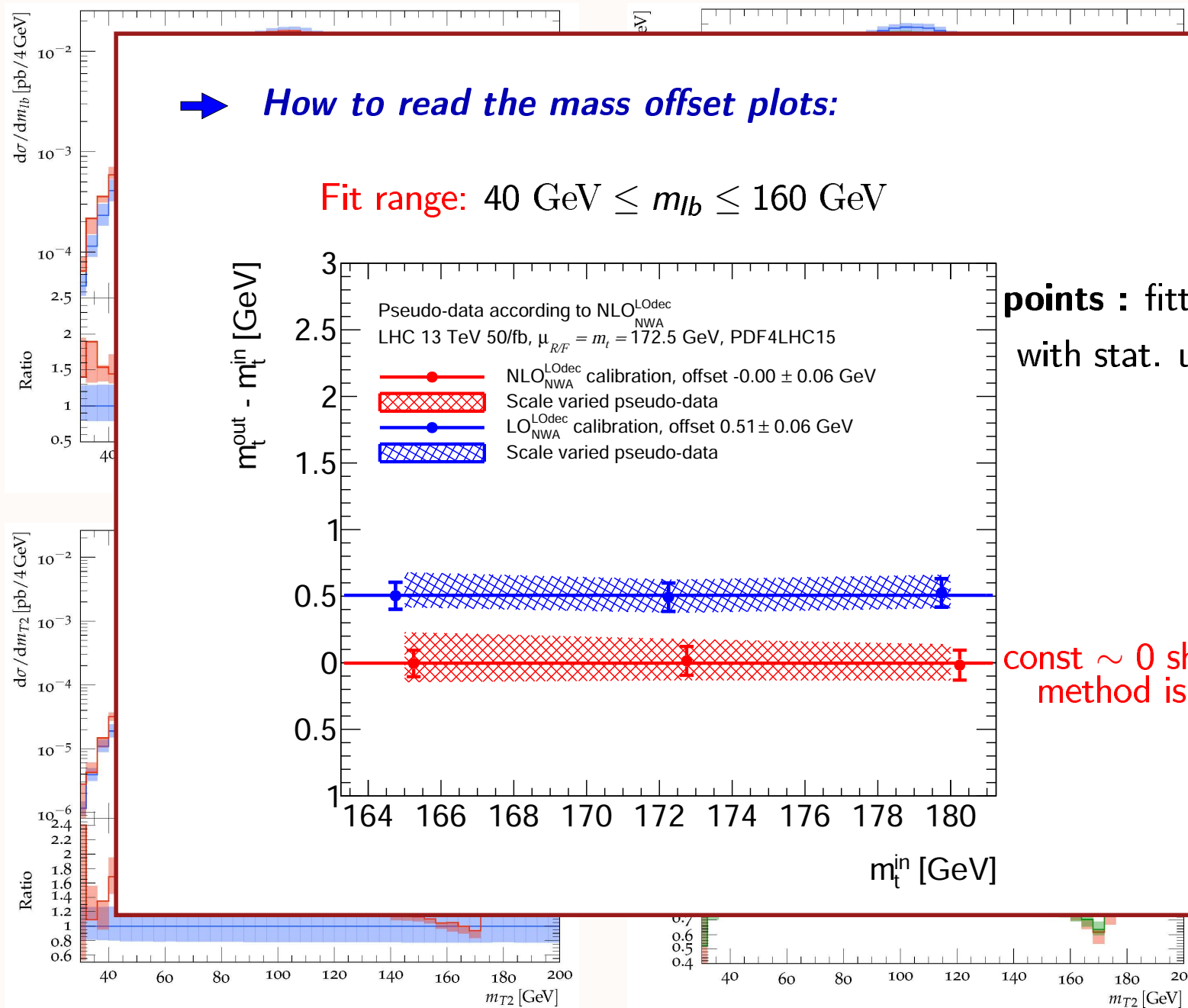


points : fitted difference
with stat. uncertainties



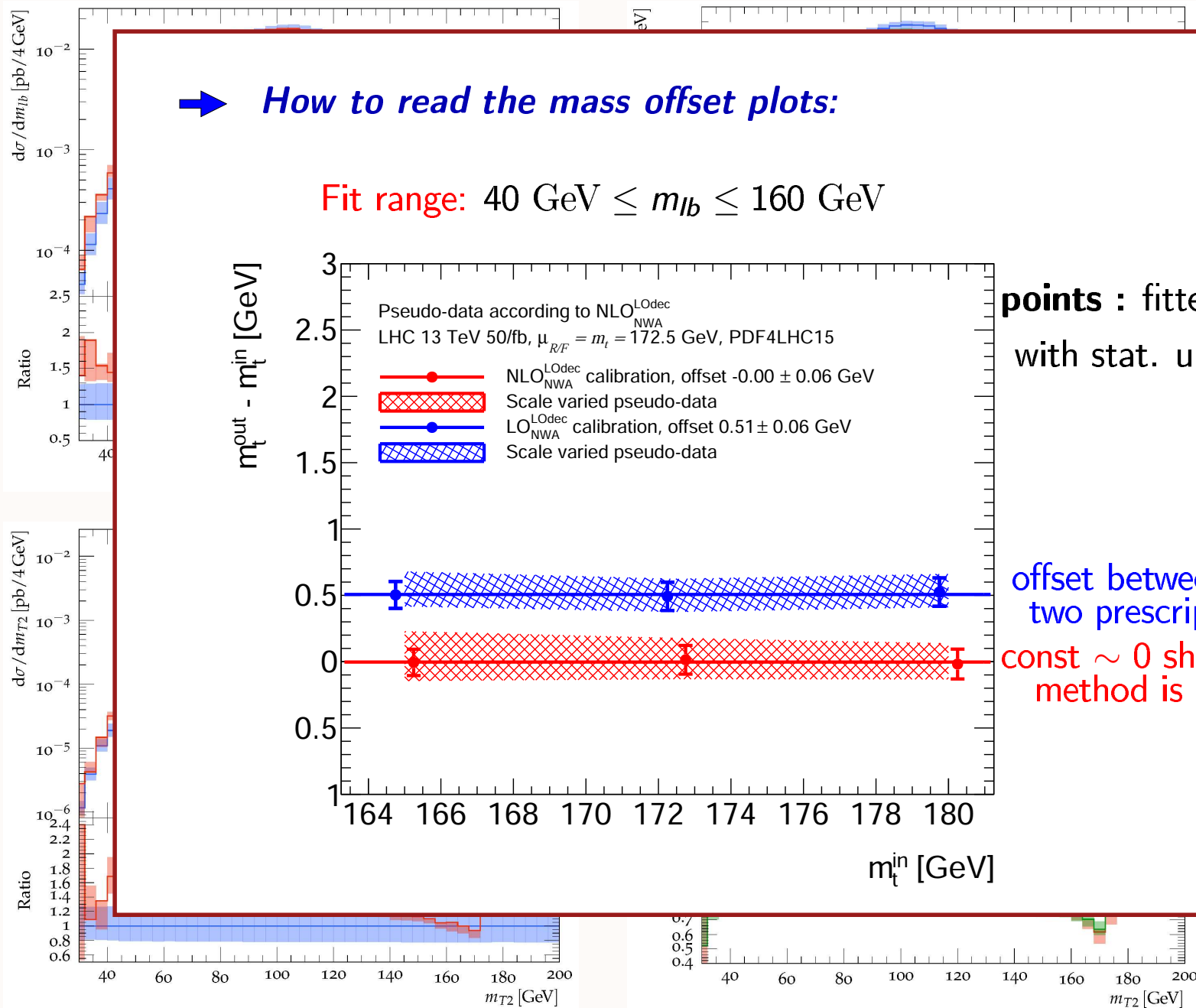
Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]



Scale dependence

[HEINRICH ET AL, ARXIV:1709.08615]

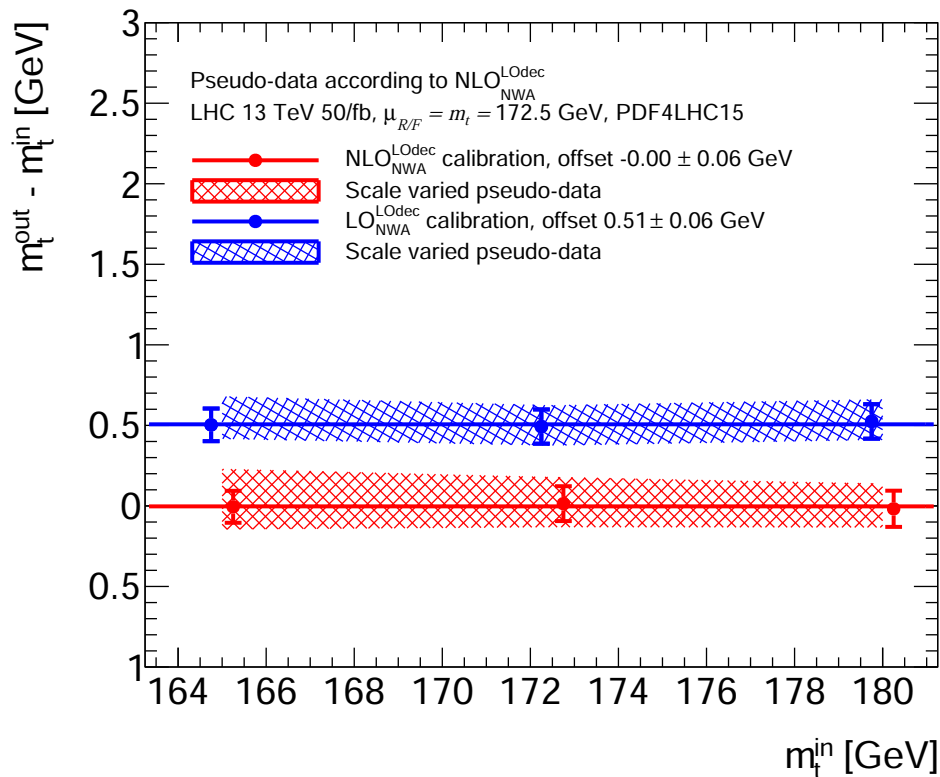


Fit results for m_{tb}

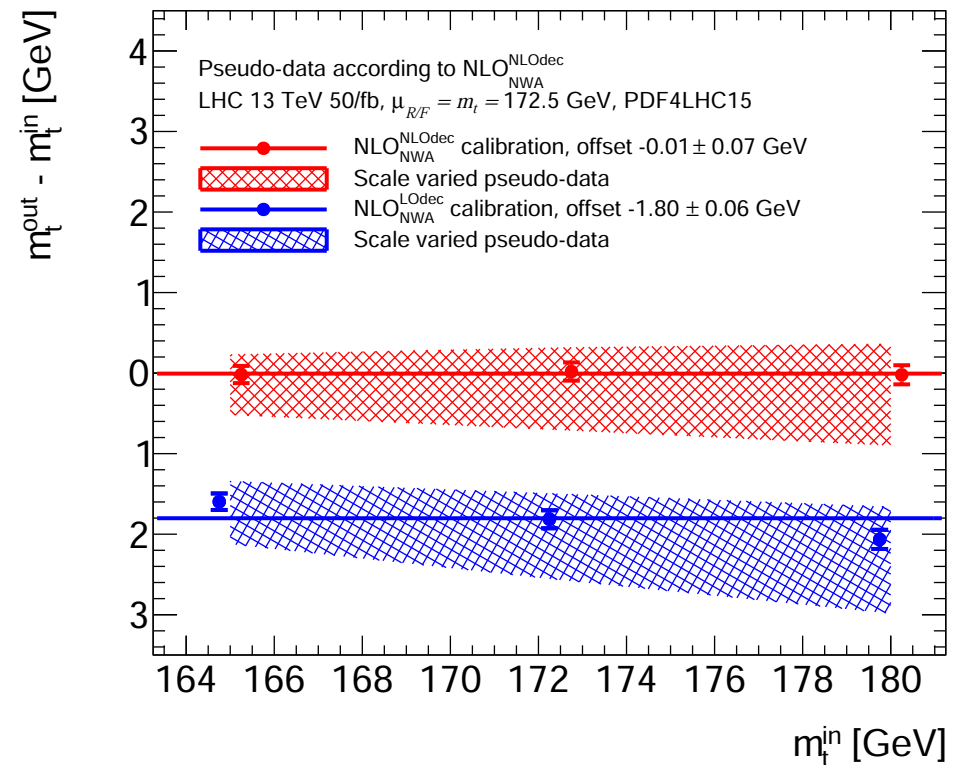
[HEINRICH ET AL, ARXIV:1709.08615]

➔ *Changing the order of perturbation theory in production or decay only.*

● NLO NWA vs LO NWA using LOdec



● NLO NWA NLOdec vs LOdec



- NLO description of the top quark decay makes a large difference
- much larger mass shifts and much larger uncertainties on mass determination due to scale variations

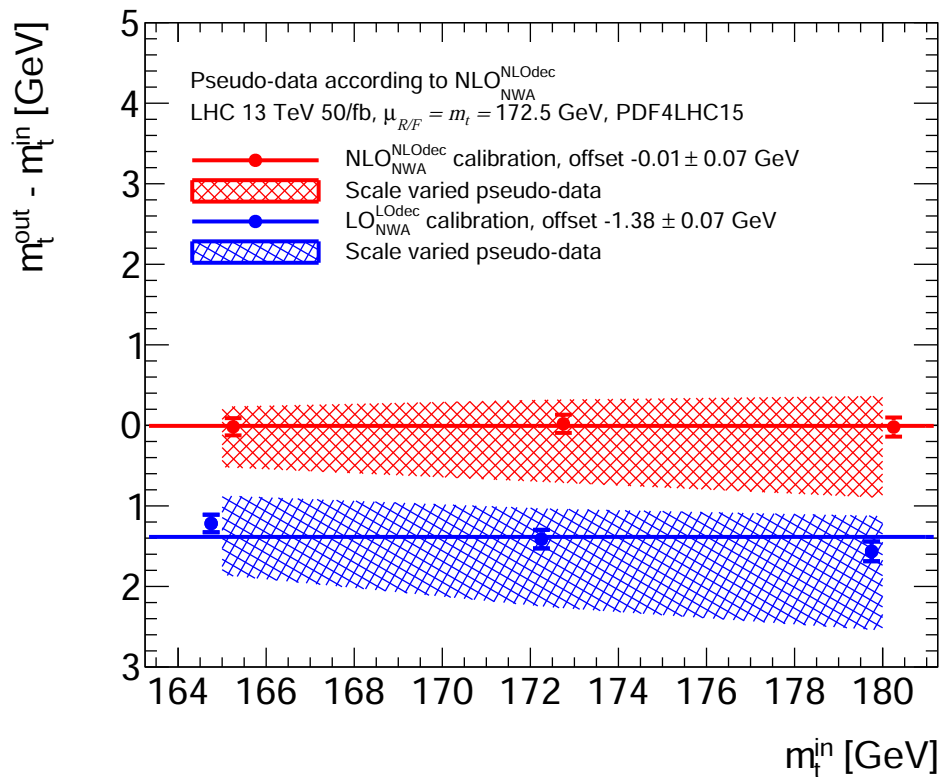


Fit results for m_{lb}

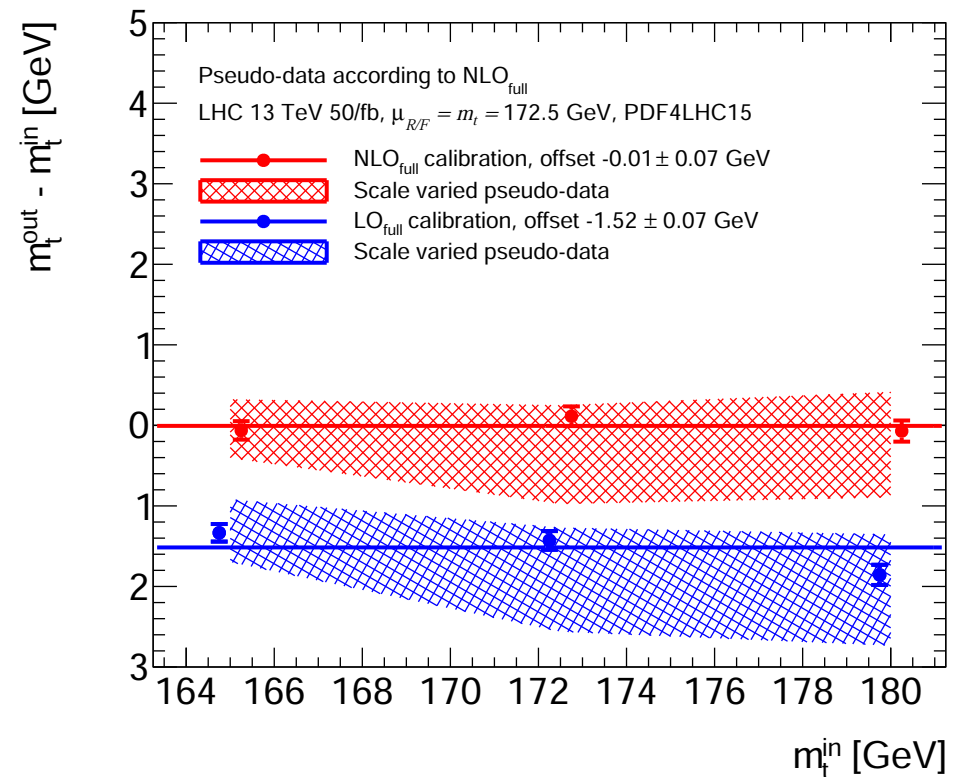
[HEINRICH ET AL, ARXIV:1709.08615]

➔ *Changing the order of perturbation theory in both production and decay.*

● NLO NWA NLOdec vs LO NWA LOdec



● NLO full vs LO full



- similar magnitude of mass shift
- similar size of scale uncertainties on mass determination

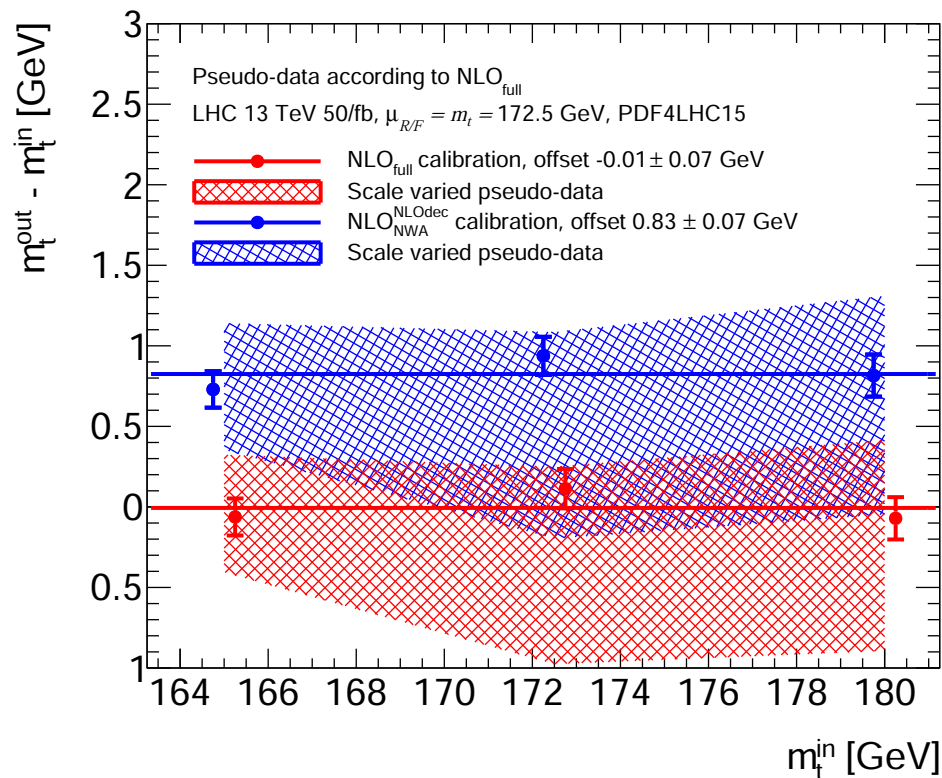


Fit results for m_{lb}

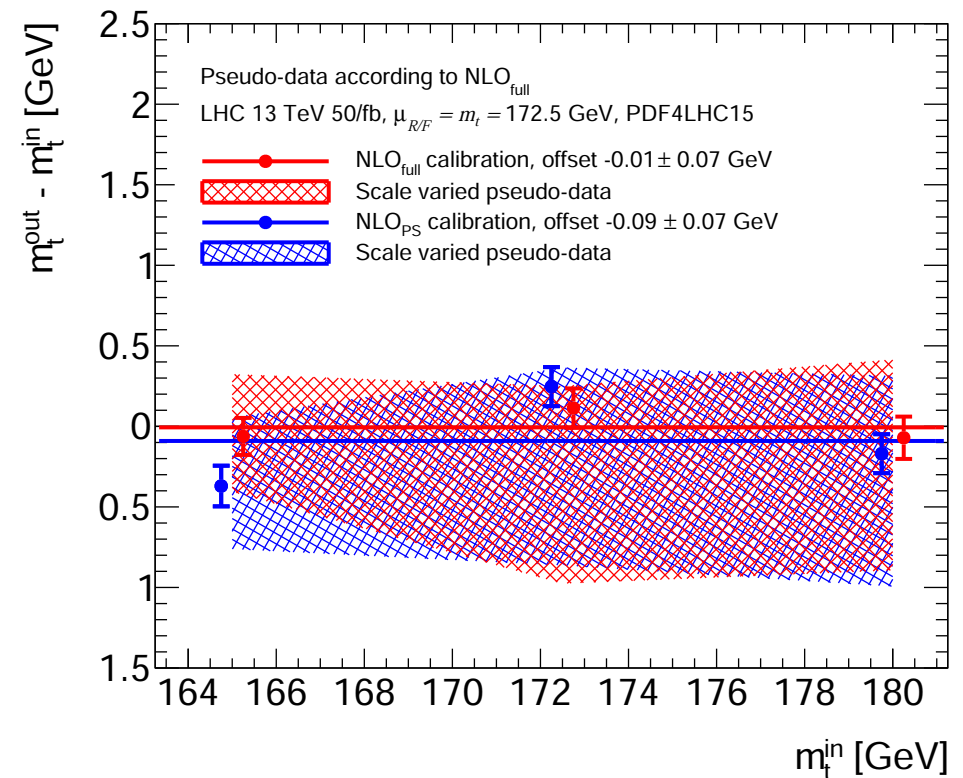
[HEINRICH ET AL, ARXIV:1709.08615]

➔ Comparing full with NWA calculations.

● NLO full vs NLO NWA NLOdec



● NLO full vs NLO+PS



- NLO full vs NLO NWA produces mass shift of the order of 0.8 GeV
- nearly zero mass shift between NLO full and NLO+PS is accidental (agreement despite different physics content)

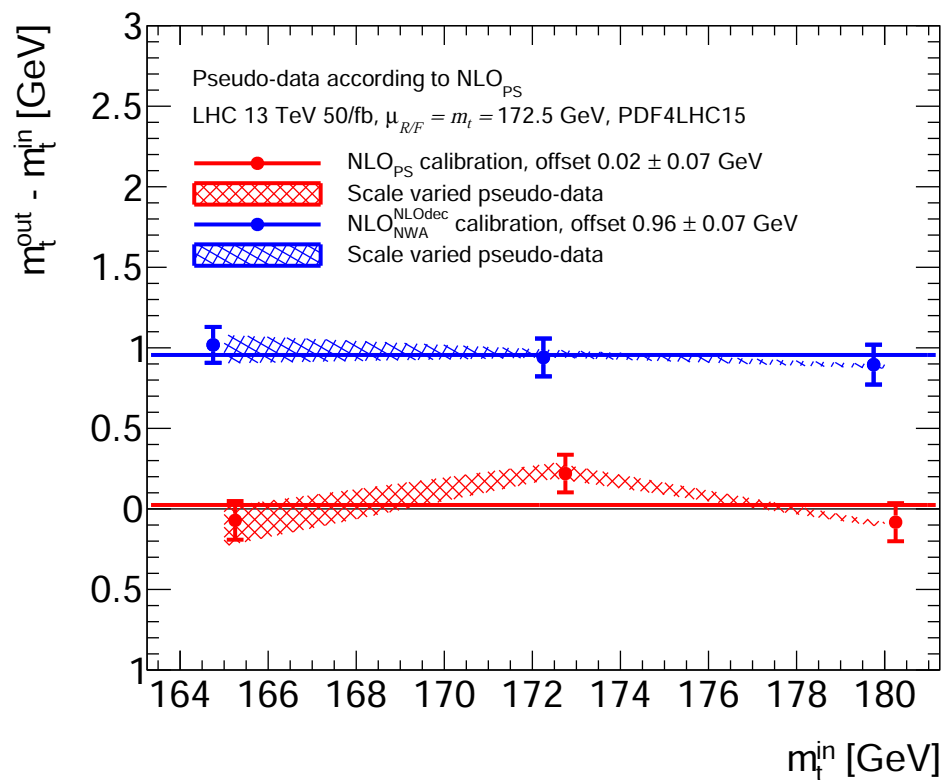


Fit results for m_{lb}

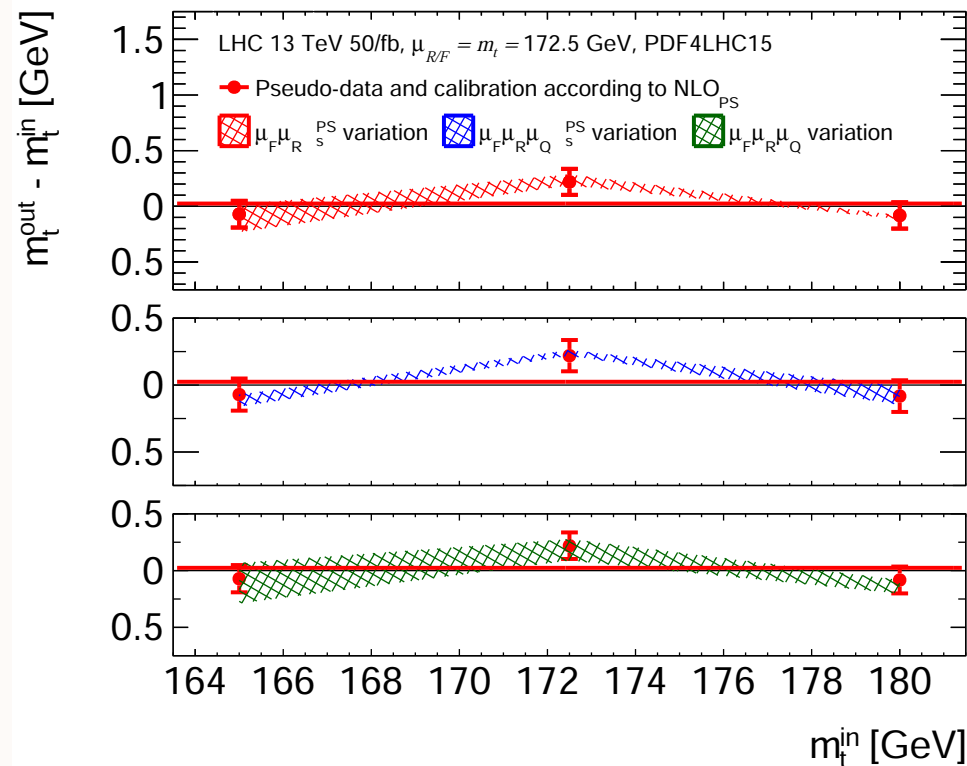
[HEINRICH ET AL, ARXIV:1709.08615]

➔ Comparing NLO+PS matching with the NLO NWA approach.

● NLO+PS vs NLO NWA NLOdec



● NLO+PS variation schemes



- further investigation needed to understand source of largish $\mathcal{O}(1 \text{ GeV})$ mass shift (resummation effect or ME corrections)
- rhs plot: different schemes of assessing the theory uncertainties of the NLO+PS shape predictions

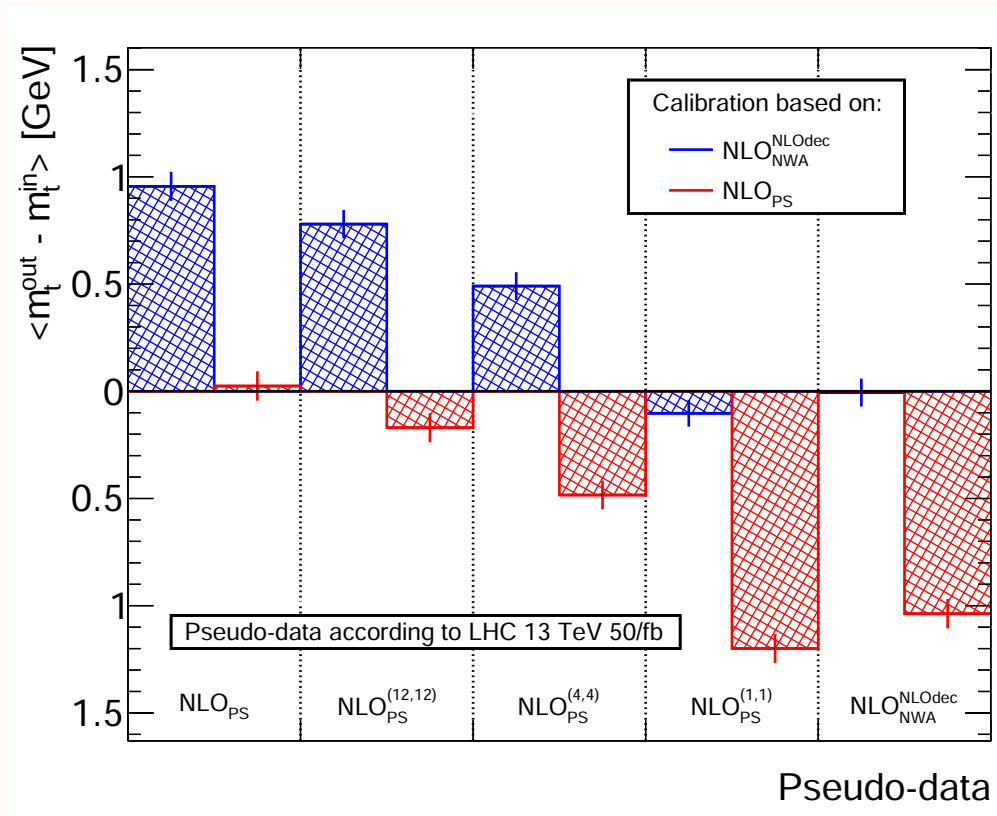


Fit results for m_{lb}

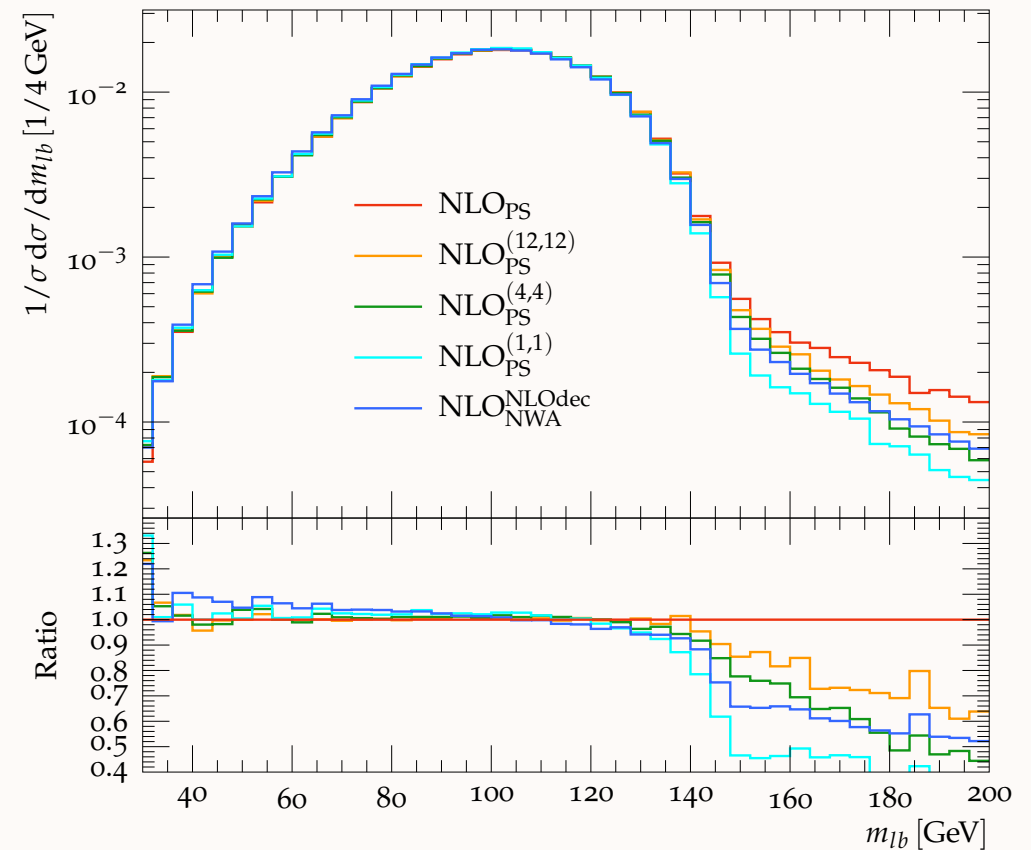
[HEINRICH ET AL, ARXIV:1709.08615]

➔ Estimating the impact of resummation effects.

- NLO+PS vs NLO NWA NLOdec
restricted showering



- associated m_{lb} distributions



- our studies suggest that most of the mass shift between NLO NWA and NLO+PS emerge due to resummation effects (fixed-order vs fully evolved radiation pattern)

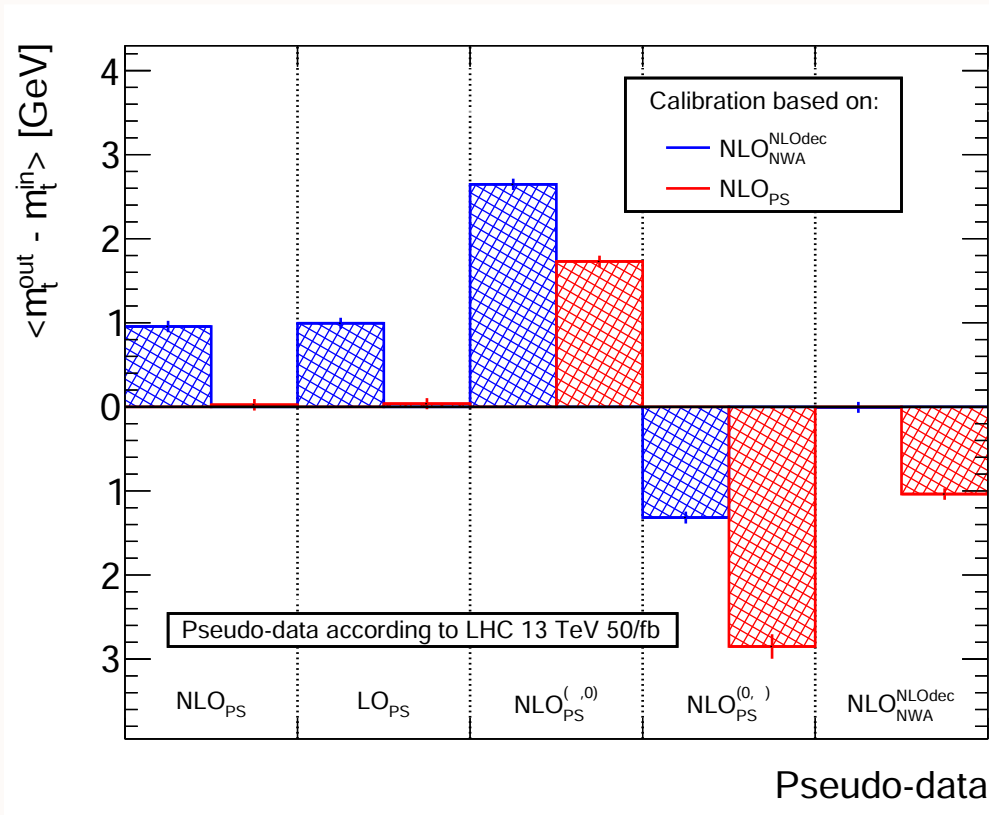


Fit results for m_{lb}

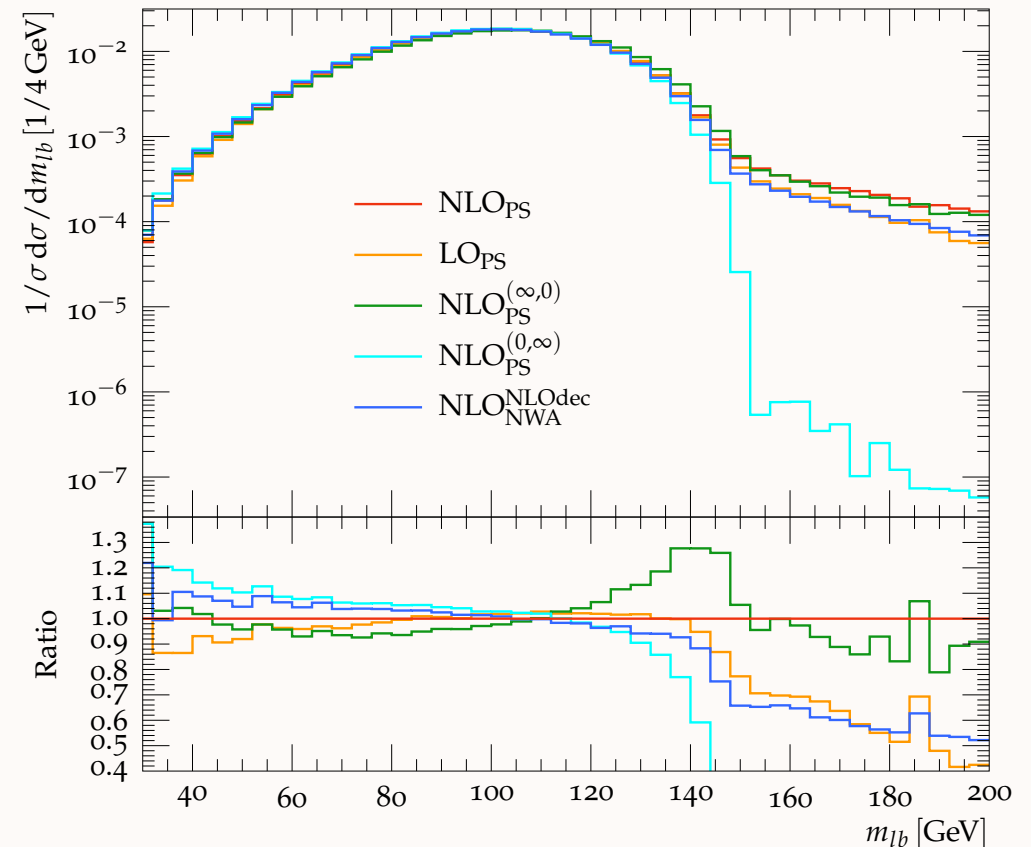
[HEINRICH ET AL, ARXIV:1709.08615]

➔ Effects of dissecting the radiation pattern in other ways.

- NLO+PS vs NLO NWA NLOdec
pure LO, production & decay showers



- associated m_{lb} distributions

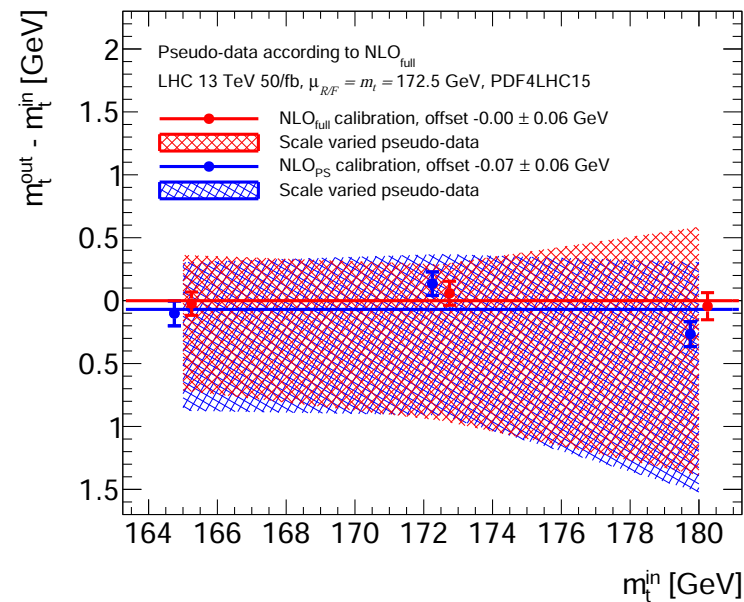
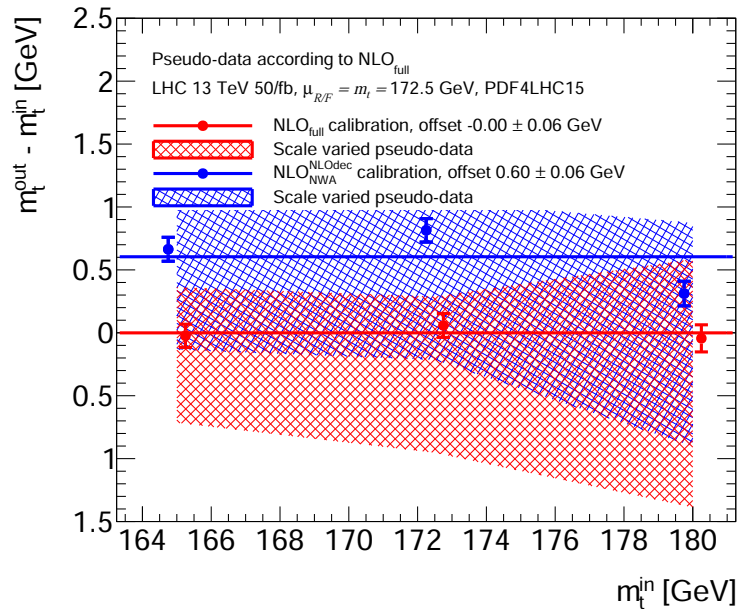
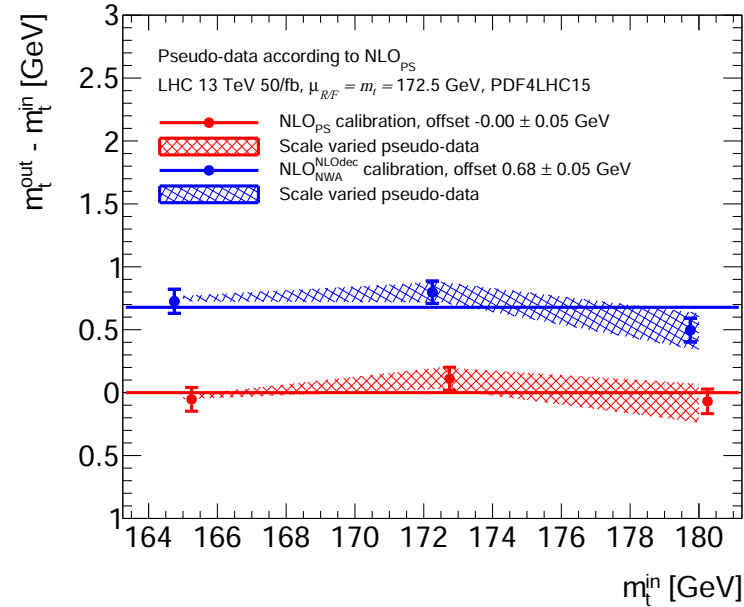
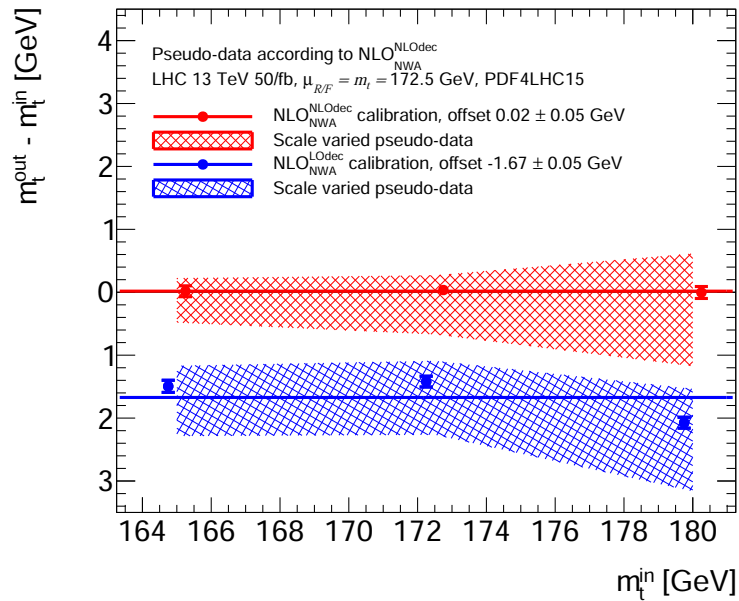


- similarity of NLO+PS & LO+PS because the same resummation corrections are applied in both showers
- large mass shifts for both pure prod. and decay showers (ISR for former / FSR for latter pull fits towards larger / smaller m_t); **blue offsets:** $\text{bin1} \approx \text{bin3} + \text{bin4} \Rightarrow$ generation-level factorization survives



Fit results for m_{T2}

[HEINRICH ET AL, ARXIV:1709.08615]



◀ radiation in decay

▶ resummation effects

◀ NLO full vs NWA

▶ NLO full vs NLO+PS



Snippet list of recent literature.

Nason et al. group

(1509.09071, 1607.04538, 1711.06281, 1801.03944)

resonance-aware subtraction and NLO+PS matching and the $b\bar{b}4l$ generator,
theoretical study comparing NLO+PS generator of increasing accuracy

Bevilacqua et al. group

(1509.09242, 1609.01659, 1710.07515)

comprehensive study of the $t\bar{t}j$ final state including off-shell, non-resonant effects and
its prospects for top quark mass extractions

Denner et al. group

(1607.05571, 1711.10359)

NLO EW corrections to dilepton final state and NLO QCD corrections to lepton+jets final state

Hoang et al. group

(1608.01318, 1704.01580)

discussing/quantifying the relation of the Monte Carlo generator mass with a short-distance mass, the MSR mass

Disclaimer: this is a personal selection which by no means is intended to be complete.



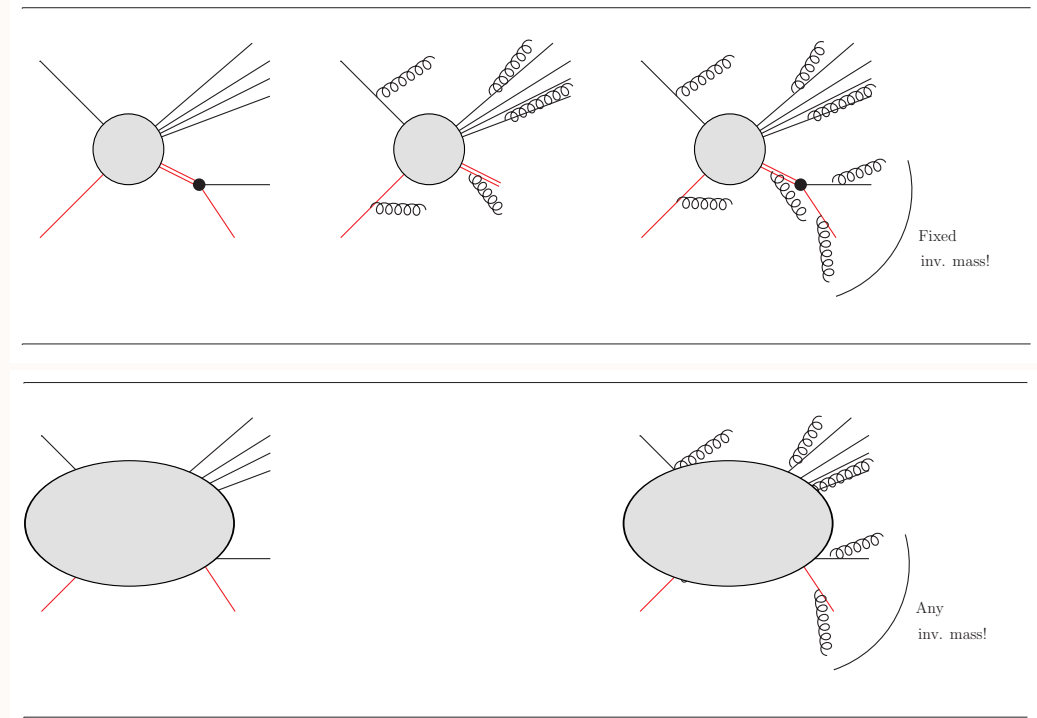
$WbW\bar{b}$ combined with parton showering.



Combination with parton showers.

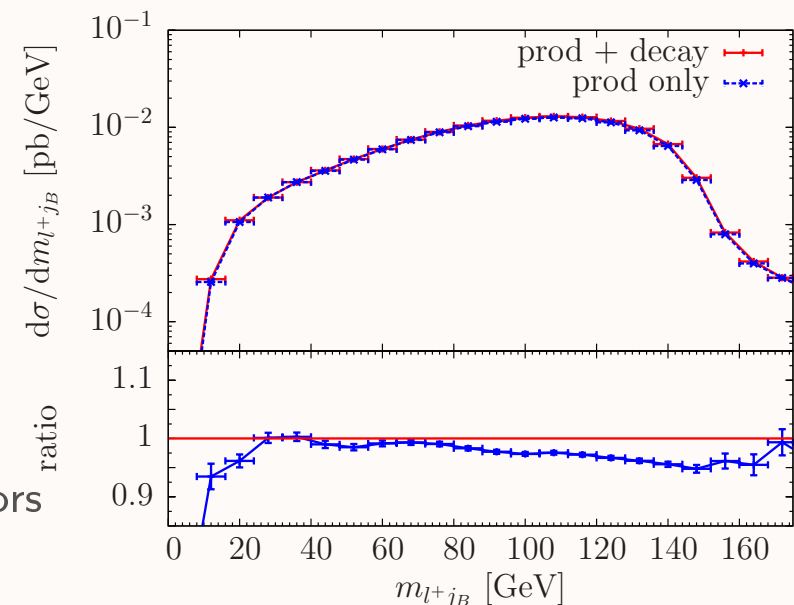
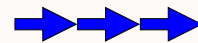
What about NLO+PS matching for $WWb\bar{b}$?

- ... to obtain more realistic, i.e. hadron level final states.
- first attempt and results using **PowHel**
[GARZELLI, KARDOS, TROCSANYI, ARXIV:1405.5859]
- however, the issue of intermediate resonances has not been addressed.
(Without a proper treatment of intermediate resonances, parton shower effects will distort the (NLO-accurate) Breit–Wigner shape.)



Recent developments towards a consistent treatment:

- consistent NLO+PS in the narrow-width limit
[CAMPBELL, ELLIS, NASON, RE, ARXIV:1412.1828]
 - resonance-aware subtraction and matching in Powheg
[JEZO, NASON, ARXIV:1509.09071]
- each xsec component (B, V, R) separated into their dominant resonance histories
subtr. procedure preserves offshellness of resonant $\hat{s}$ -propagators
resonance info on fs particles communicated to shower.



NLO+PS generator for $WWb\bar{b}$ production

[JEZO, LINDERT, NASON, OLEARI, POZZORINI, ARXIV:1607.04538], [JEZO, NASON, ARXIV:1509.09071]
 [20] [FRIXIONE, NASON, RIDOLFI, ARXIV:0707.3088], [35] [CAMPBELL, ELLIS, NASON, RE, ARXIV:1412.1828]

Significant theoretical
improvements:

NLO MEs for $pp \rightarrow \ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$ (i.e. up to $\mathcal{O}(\alpha_s^3 \alpha^4)$) in the 4FNS can be matched to parton shower using information about resonance histories and the resonance-aware subtraction and matching method of Powheg.

label	$t\bar{t}$	$t\bar{t} \otimes \text{decay}$	$b\bar{b}4\ell$
generator	hvg [20]	ttb_NLO_dec [35]	bb4l
framework	POWHEG-BOX	POWHEG-BOX-V2	POWHEG-BOX-RES
NLO matrix elements	$t\bar{t}$	$t(\rightarrow \ell^+ \nu_\ell b) \bar{t}(\rightarrow l^- \bar{\nu}_l \bar{b})$	$\ell^+ \nu_\ell l^- \bar{\nu}_l b \bar{b}$
decay accuracy	LO+PS	NLO+PS	NLO+PS
NLO radiation	single	multiple	multiple
spin correlations	approx.	exact	exact
off-shell $t\bar{t}$ effects	BW smearing	LO $b\bar{b}4\ell$ reweighting	exact
Wt & non-resonant effects	no	LO $b\bar{b}4\ell$ reweighting	exact
b -quark massive	yes	yes	yes

Physics features combined for the 1st time in one generator dubbed bb4l (POWHEG-BOX-RES framework with OpenLoops interface matched to Pythia8)

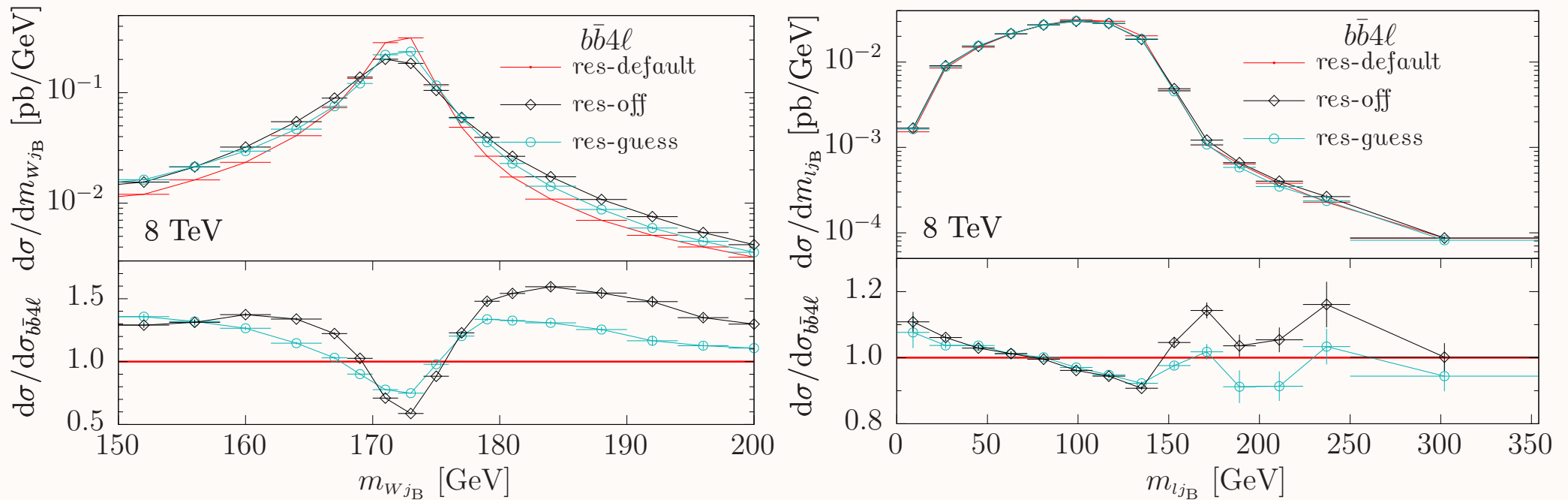
- consistent NLO+PS treatment of t resonances, including quantum corrections to t propagators and offshell t decay chains
- exact spin correlations at NLO, interference between NLO radiation from top production and decays, full NLO accuracy in $t\bar{t}$ production and decays
- unified treatment of $t\bar{t}$ and Wt production with interference at NLO
- improved modelling of b quark kinematics owing to b quark mass effects
- access to phase-space regions with unresolved b quarks and/or jet vetoes



NLO+PS generator for $WWb\bar{b}$ production

[JEZO, LINDERT, NASON, OLEARI, POZZORINI, ARXIV:1607.04538]

➔ comparison of resonance-aware with resonance-unaware, traditional ($res - off$), and resonance kinematic-guess based predictions ($res - guess$): the traditional approach gives wider mass peaks as it does not preserve the virtuality of top resonances, neither in Powheg nor in the shower.



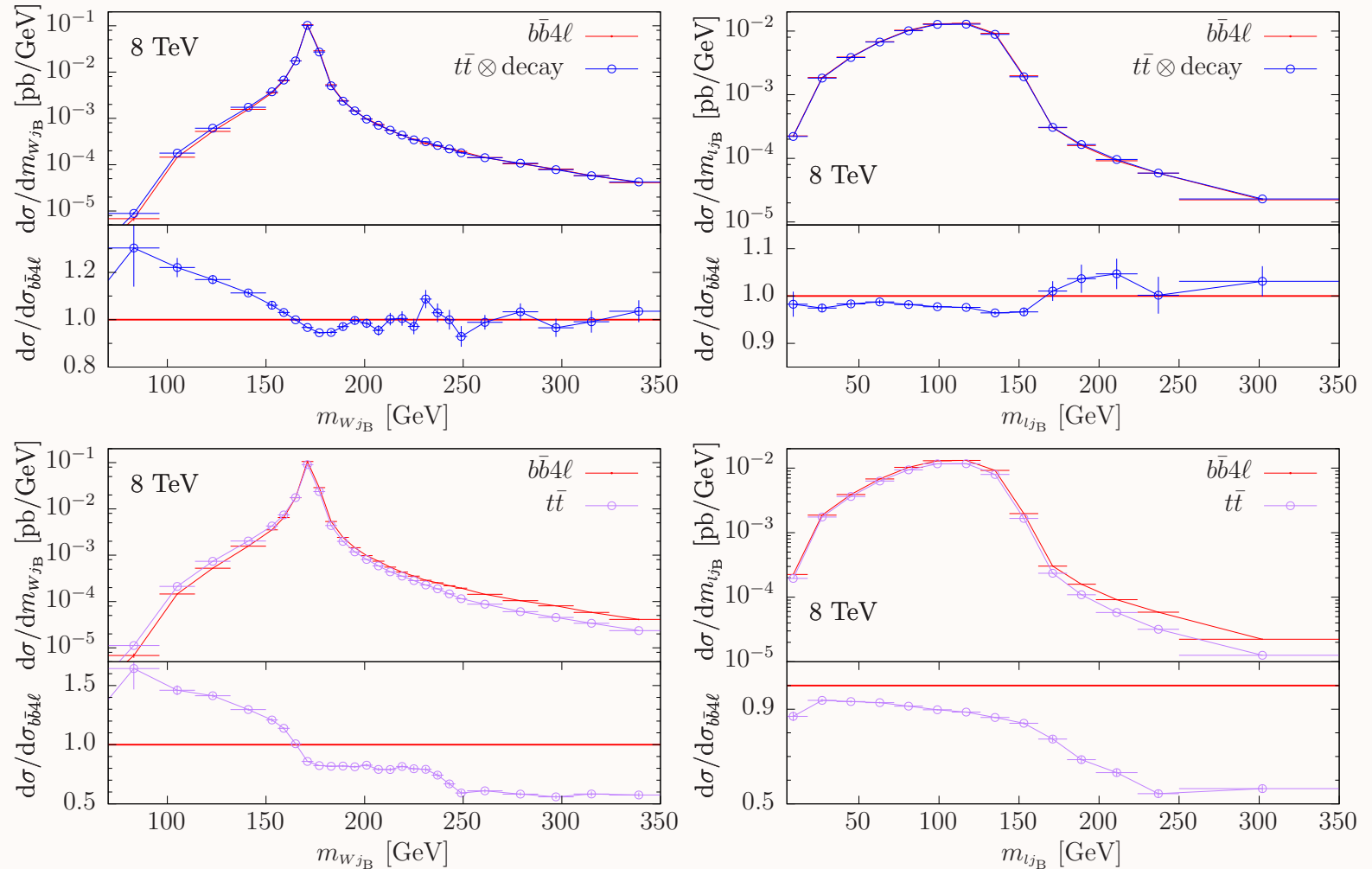
- j_B is b jet containing hardest B hadron, W reco. via corresponding offshell $\ell\nu$ pair in hard ME
- t reconstruction via charge and b flavour information at MC truth level



NLO+PS generator for $WWb\bar{b}$ production

[JEZO, LINDERT, NASON, OLEARI, POZZORINI, ARXIV:1607.04538]

➔ comparison of the predictions of the new generator, standard Powheg (lower row) and that obtained from operating in narrow-width limit (upper row) including offshellness and interference effects in an approximate way.



- j_B is b jet containing hardest B hadron, W reco. via corresponding offshell $\ell\nu$ pair in hard ME
- t reconstruction via charge and b flavour information at MC truth level



Summary / Outlook.

Studied variety of relevant theoretical descriptions and top quark mass sensitive observables (for a well-defined framework in the context of a dilepton analysis).

Prospects for top quark mass extraction are best for observables m_{lb} and m_{T2} . The details of the theoretical descriptions matter for the (quality of the) extraction, much more than finding the optimal choice for the fit range.

Confirmed: large effect of radiative corrections in the top quark decay increases the theoretical uncertainties of the extraction. Still: non-factorizing/non-resonant contributions lead to mass shift of the order of 1 GeV between full calculation and NWA at NLO.

NLO matched parton shower gives surprisingly small mass shift compared to full NLO calculation (agreement by chance $\rightarrow$ different physics effects, similar pheno impact) while the associated theoretical uncertainties very likely are underestimated. The offset compared to NLO NWA is larger than expected, again of the order of 1 GeV. Restricted-shower studies indicate that the resummation corrections included by the full showers lead to effects on the m_t determination that can be as large as 1 GeV. Fixed-order results can be expected to change accordingly once parton showering is included but their non-uniform scale dependence will likely survive.

Desired: comparison to NNLO production plus NNLO decay; comparison to real data; estimation of hadronization & detector effects in same framework; dedicated (maybe larger-scale) comparison between different parton shower (matched) approaches as useful exercise to assess related uncertainties on broader, more global scale.



The end. (:o) Thank You.

