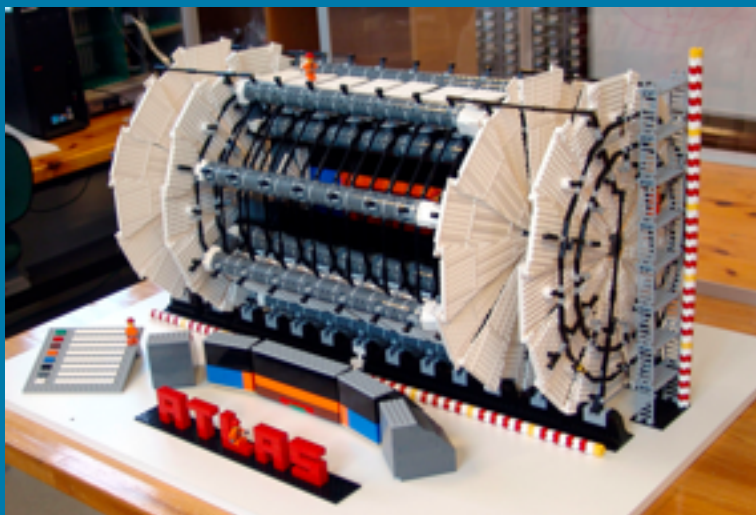




DETECTORS FOR HIGH ENERGY PHYSICS

Ingrid-Maria Gregor, DESY



DESY Summer Student Program 2018

Hamburg

July 26th/ Aug 1st

OVERVIEW

I. Detectors for Particle Physics

II. Interaction with Matter

III. Calorimeters

IV. Tracking Detectors

● Gas detectors

● Semiconductor trackers

V. Examples from the real life

}

Thursday

}

Wednesday

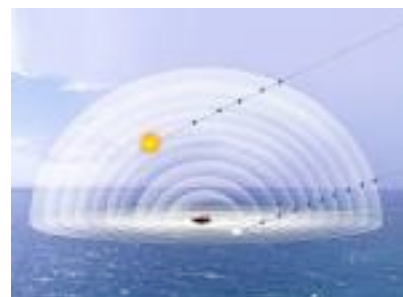
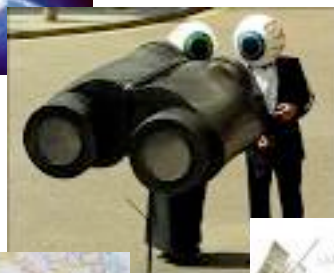
IV. TRACKING DETECTORS

TRACKING

🔴 “tracking” in google image search:

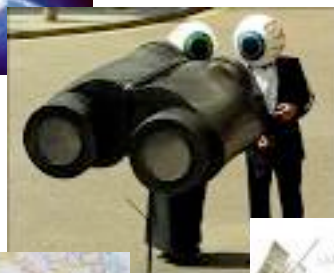
TRACKING

📍 “tracking” in google image search:



TRACKING


📍 “tracking” in google image search:



Tracking



TRACKING DETECTOR

 “tracking detector” in google image search



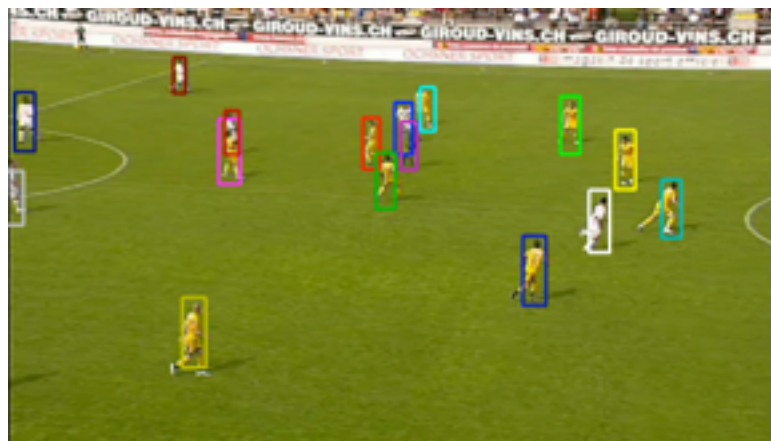
GPS Tracking Detector

TRACKING DETECTOR

🔴 “tracking detector” in google image search



GPS Tracking Detector



Online Multi-Person Tracking-by-Detection
from a Single, Uncalibrated Camera

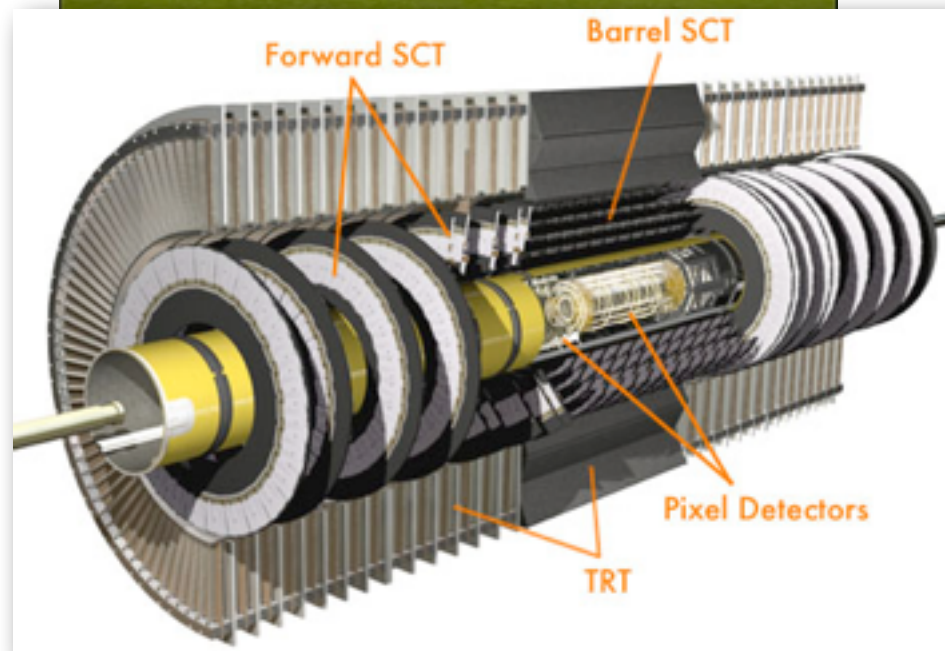
TRACKING DETECTOR

🔴 “tracking detector” in google image search



GPS Tracking Detector

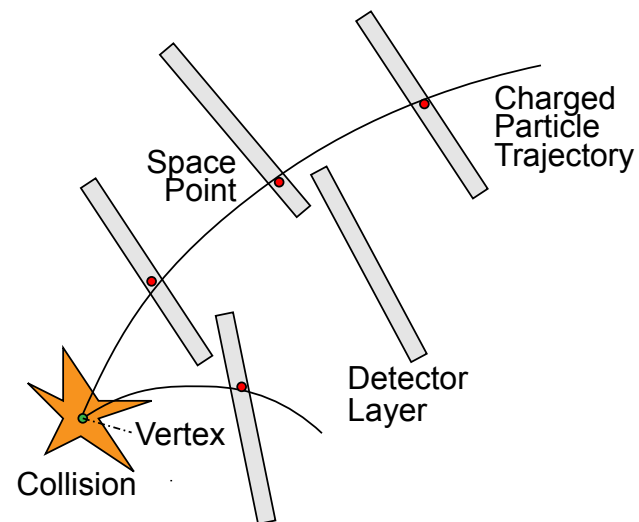
But the 1st image on list is:



Pic: ATLAS Collaboration

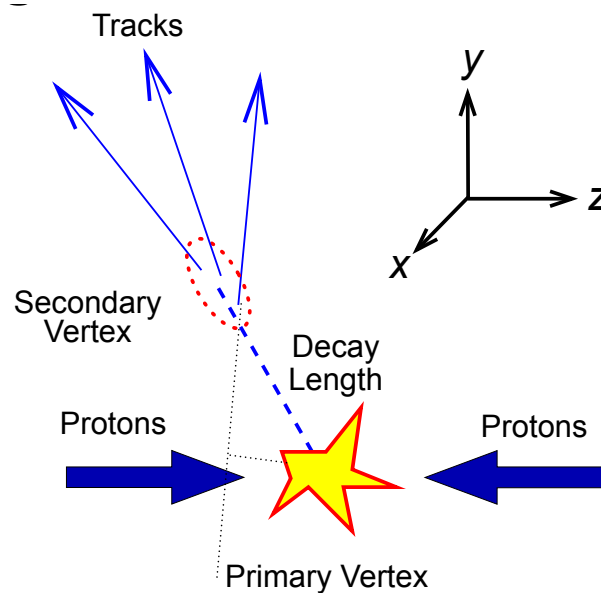
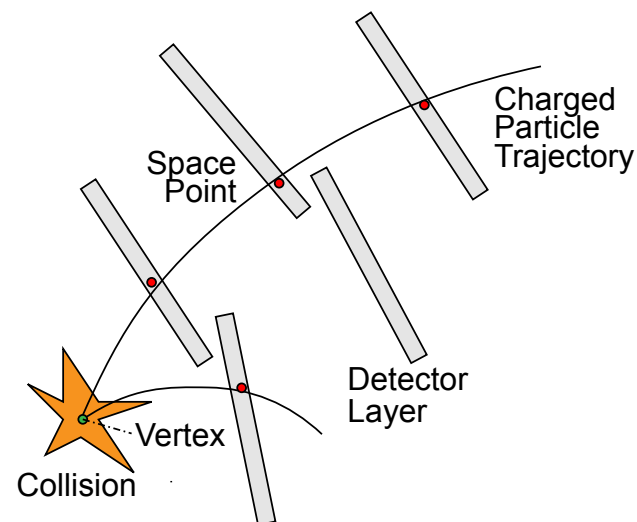
TRACKING DETECTORS

- Precise measurement of track and momentum of charged particles due to magnetic field.
- The trajectory should be disturbed minimally by this process (reduced material)



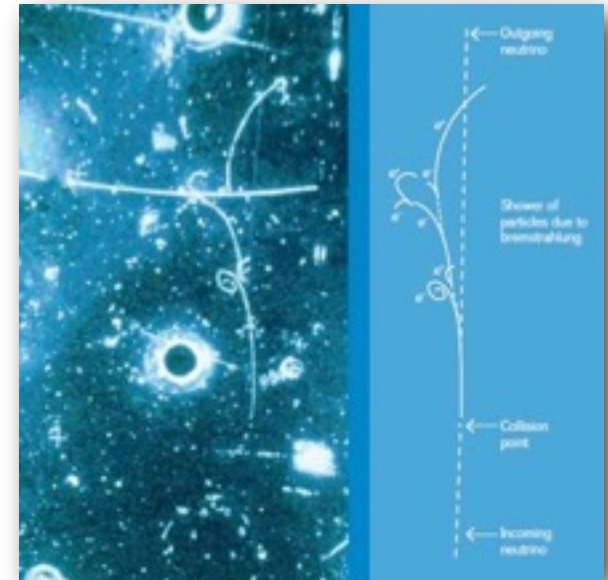
TRACKING DETECTORS

- Precise measurement of track and momentum of charged particles due to magnetic field.
- The trajectory should be disturbed minimally by this process (reduced material)
- Charged particles ionize matter along their path.
- Tracking is based upon detecting ionisation trails.
- An “image” of the charged particles in the event

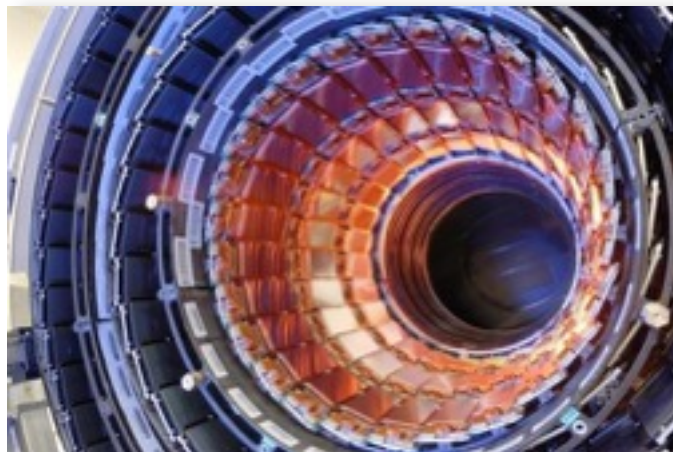


TRACKING DETECTORS - TECHNOLOGIES

- **“Classic”**: Emulsions, cloud, and bubble chambers
 - Continuous media
 - Typically very detailed information but slow to respond and awkward to read out
- **“Modern”**: Electronic detectors, wire chambers, scintillators, solid state detectors
 - Segmented
 - Fast, can be read out digitally, information content is now approaching the “classic” technology
 - Mostly used solid state detector -> Silicon (pixels and strips)



Discovery of neutral currents
Gargamelle, 1972

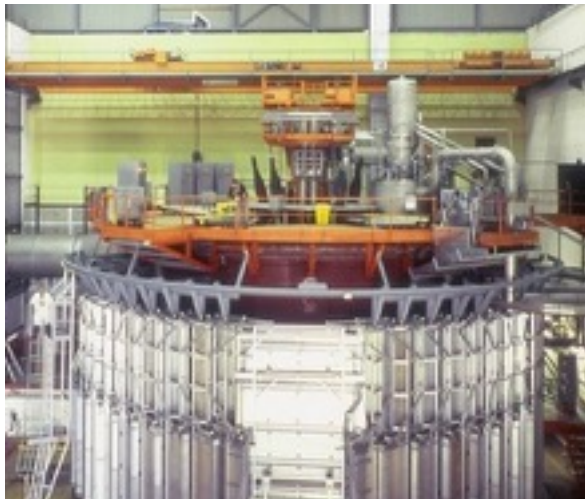


CMS Inner barrel Si Tracker:
Single-Sided Si-Strip



Pictures: CERN

VERY “CLASSIC”: BUBBLE CHAMBER



- The biggest: Big European Bubble Chamber
- 3.7 m diameter
- Until 1984 used at CERN for the investigation of neutron hadron interactions

Early report on bubble chamber analysis:



Second United Nations
International Conference
on the Peaceful Uses of
Atomic Energy

A/CONF.15/P/730
U.S.A.
June 1958

ORIGINAL: ENGLISH

ON THE ANALYSIS OF BUBBLE CHAMBER TRACKS

©
Hugh Bradner and Frank Solmitz

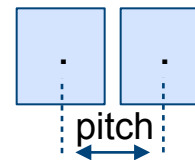


“... the large number of possible reactions, the variability of appearance of interaction, and the importance of being alert to possible new phenomena make it very important for a **trained physicist** to look at the bubble chamber pictures....”

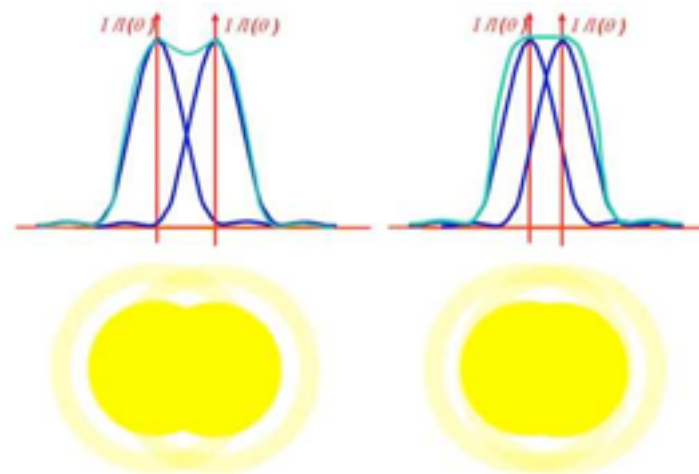
TRACKER: IMPORTANT PARAMETER



- An important figure of merit is the **spatial resolution** of a tracking detector
- Depending on detector geometry and charge collection
 - Pitch (distance between channels)
 - Charge sharing between channels



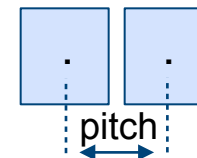
can be tubes,
strips, wires,
pixels



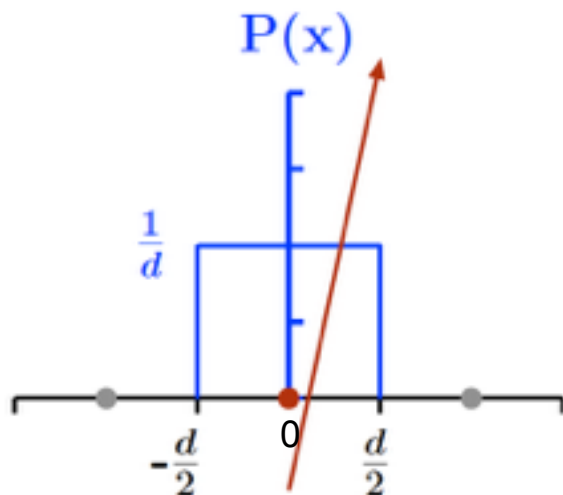
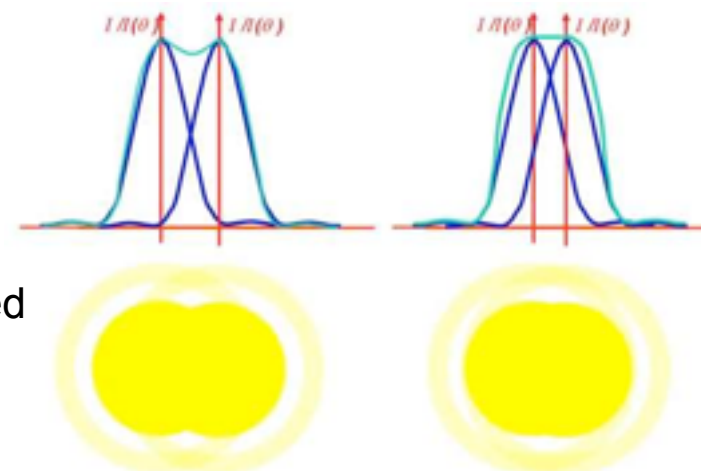
TRACKER: IMPORTANT PARAMETER



- An important figure of merit is the **spatial resolution** of a tracking detector
- Depending on detector geometry and charge collection
 - Pitch (distance between channels)
 - Charge sharing between channels
- Simple case: all charge is collected by one strip
- Traversing particle creates signal in hit strip (binary)
- Flat distribution along strip pitch; no area is pronounced
- ➔ Probability distribution for particle passage:



can be tubes,
strips, wires,
pixels



$$P(x) = \frac{1}{d} \quad \Rightarrow \quad \int_{-d/2}^{d/2} P(x) dx = 1$$

The reconstructed point is always the middle of the strip:

$$\langle x \rangle = \int_{-d/2}^{d/2} x P(x) dx = 0$$

TRACKER: IMPORTANT PARAMETER

- Calculating the resolution orthogonal to the strip:

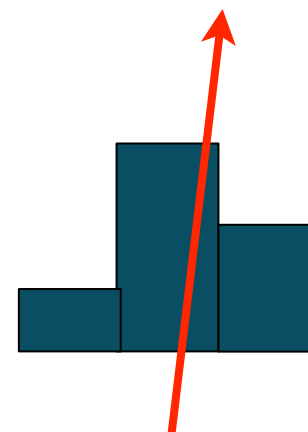
$$\sigma_x^2 = \langle (x - \langle x \rangle)^2 \rangle = \int_{-d/2}^{d/2} x^2 P(x) dx = \frac{d^2}{12}$$

- Resulting in a general term (valid for tracking detectors with a pitch d):

$$\sigma = \frac{d}{\sqrt{12}} \quad \leftarrow \text{very important !}$$

- For a silicon strip detector with a strip pitch of $80 \mu\text{m}$ this results in a minimal resolution of $\sim 23 \mu\text{m}$
- In case of charge sharing between the strip (signal size decreasing with distance to hit position) and information about signal size
 - resolution improved by additional information of adjacent channels

$$\sigma \propto \frac{d}{(S/N)}$$



TRACKING: DETERMINATION OF THE MOMENTUM IN MAGNETIC FIELD

- A tracking detector is typically placed within a B-field to enable momentum measurements
- Charged particles are deflected in a magnetic field:
 - takes only effect on the component perpendicular to the field

Radius of the circular path is proportional to the transversal momentum

$$F = qvB$$

$$ma = qvB$$

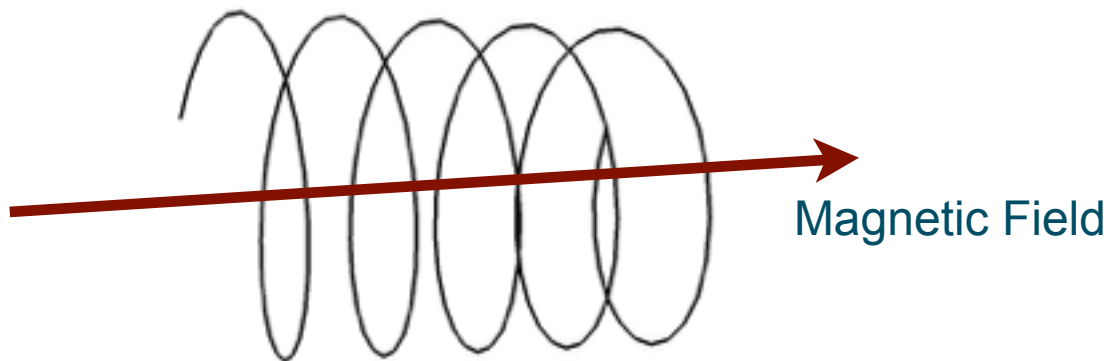
$$m\left(\frac{v^2}{r}\right) = qvB$$

$$p = 0.3Br$$

- parallel to the field is no deflection:

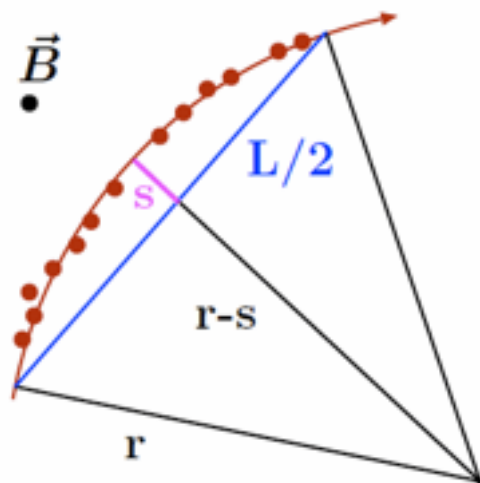
when converting in HEP units and assuming that all particles have the [electron charge]

⇒ particle is moving on a helix, the radius is determined by the field and p_T





DETERMINATION OF THE MOMENTUM IN MAGNETIC FIELD II



● In real applications usually only slightly bent track segments are measured

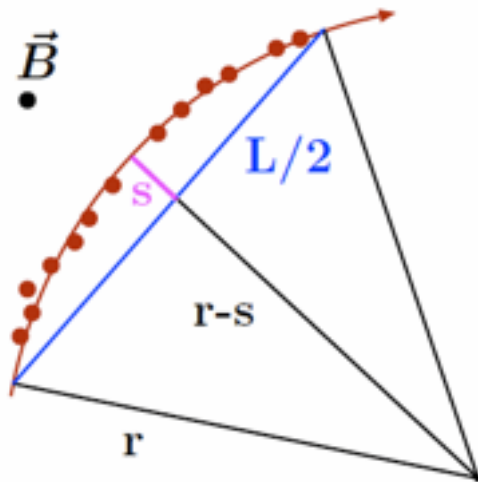
● Figure of merit: sagitta

Segment of a circle: $s = r - \sqrt{r^2 - \frac{L^2}{4}}$

$$\Rightarrow r = \frac{s}{2} + \frac{L^2}{8s} \approx \frac{L^2}{8s} \quad (s \ll L)$$

With the radius-momentum-B-field relation: $r = \frac{p_T}{0.3 B} \Rightarrow s = \frac{0.3 B L^2}{8 p_T}$

DETERMINATION OF THE MOMENTUM IN MAGNETIC FIELD II



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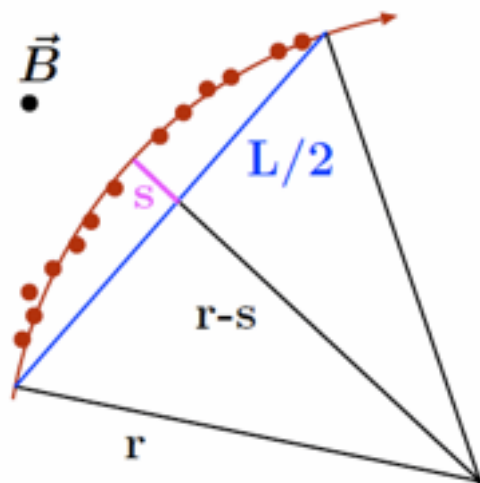
Momentum resolution due to position measurement:

Gluckstern $\frac{\sigma_{p_T}}{p_T} = \frac{\sigma_s}{s} = \sqrt{\frac{720}{n+4}} \frac{\sigma_y p_T}{0.3 B L^2}$

NIM, 24, P381, 1963



DETERMINATION OF THE MOMENTUM IN MAGNETIC FIELD II



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NIM, 24, P381, 1963

➡ The larger the magnetic field **B**, the length **L** and the number of measurement points **n**, and the better the spatial resolution, the better is the momentum resolution

IMPULS RESOLUTION: SPATIAL RESOLUTION AND MULTIPLE SCATTERING

- More components are influencing the momentum resolution $\sigma(p_T)/p_T$ of a tracking system:

- Inaccuracy of the tracking detector: $\sigma(p_T) \propto p_T$

- Influence of the particle due to MS: $\sigma(x)_{MS} \propto \frac{1}{p}$

- The angular resolution of the detector

Multiple scattering angle:
 $\theta \propto \frac{1}{p}$

$$\left(\frac{\sigma_{p_T}}{p_T}\right)^2 = \left(\sqrt{\frac{720}{n+4} \frac{\sigma_y p_T}{0.3BL^2}}\right)^2 + \left(\frac{52.3 \times 10^{-3}}{\beta B \sqrt{LL_y \sin \theta}}\right)^2 + (\cot \theta \sigma_\theta)^2$$

Position resolution

Multiple scattering

Angular resolution

- p_T resolution improves as $1/B$ and depends on p as $1/L^2$ or $1/L^{-1/2}$
- For low momentum ($\beta \rightarrow 0$), MS will dominate the momentum resolution.
- Improving the spatial resolution (σ_y) only improves momentum resolution if the first term is dominate.

TRACKER: IMPORTANT PARAMETER

● **Detector efficiency** ϵ : probability to detect a transversing particle

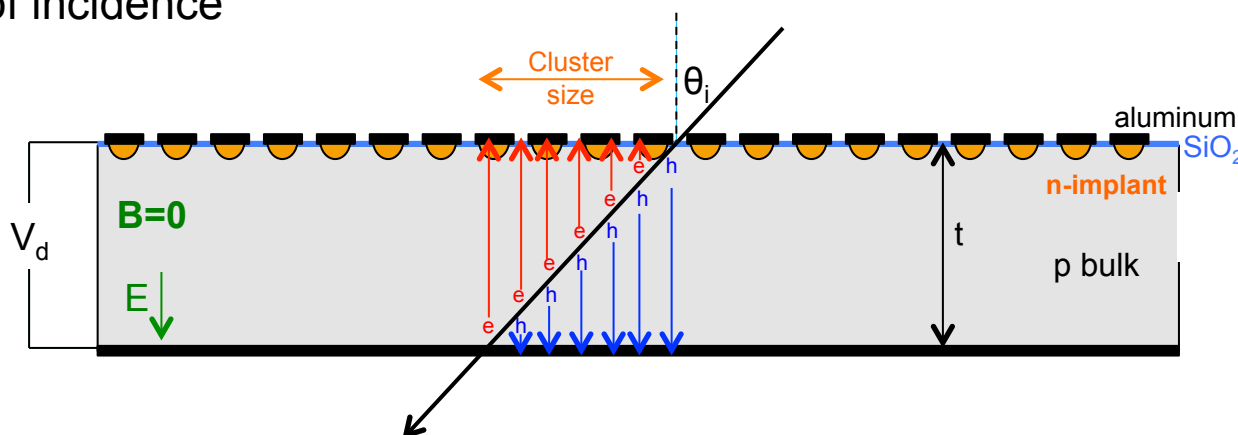
- should be as close to 100% as possible
- i.e. 12 layer silicon detector with 98% efficiency per layer -> overall tracking efficiency is only 78%
- needs to be measured in test beam

$$\epsilon_{\text{track}} = (\epsilon_{\text{layer}})^n$$

n = number of layer is tracking system

● **Cluster size** : number of hit pixels/strips belonging to one track

- usually given in unit of strips or pixels
- depending on angle of incidence



Optimally measured in test beam

TRACKER: IMPORTANT PARAMETER

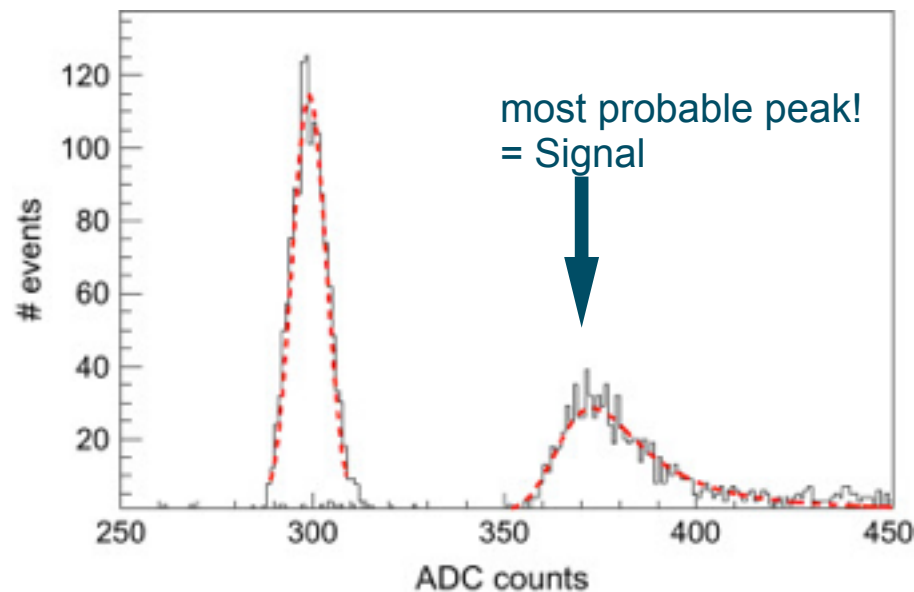
- **Signal/noise ratio**: signal size for a certain input signal over the intrinsic noise of the detector
 - parameter for analog signals
 - good understanding of **electrical noise charge** needed
 - leakage current (ENC_I)
 - detector capacity (ENC_C)
 - det. parallel resistor (ENC_{Rp})
 - det. series resistor (ENC_{Rs})
- signal induced by source or laser (or test beam particles)
- optimal S/N for a MiP is larger than 20

$$ENC = \sqrt{ENC_C^2 + ENC_I^2 + ENC_{Rp}^2 + ENC_{Rs}^2}$$

example for silicon detector

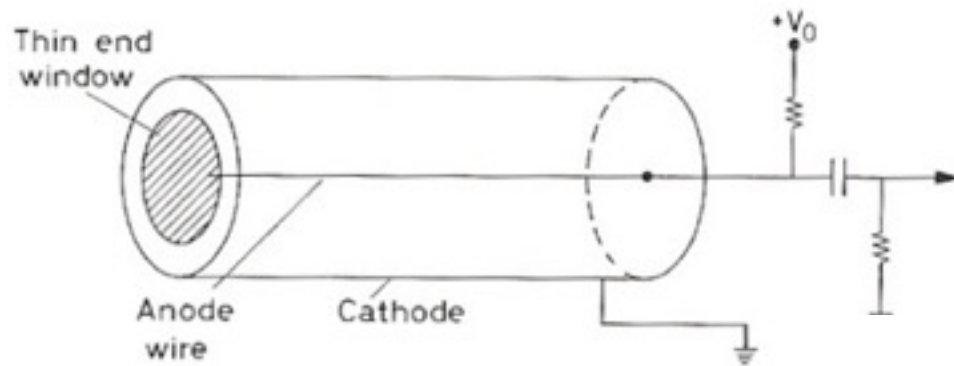
With analog readout:

Gaussian distributed “non-signal”
= sigma -> noise

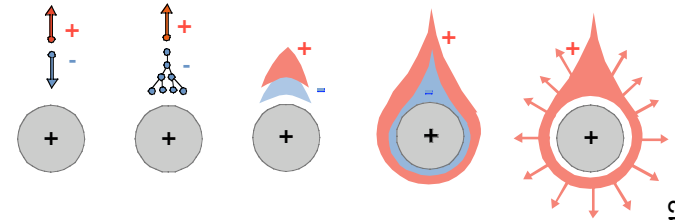


IV.A GAS-DETECTORS

ANOTHER CLASSIC: IONISATION CHAMBER

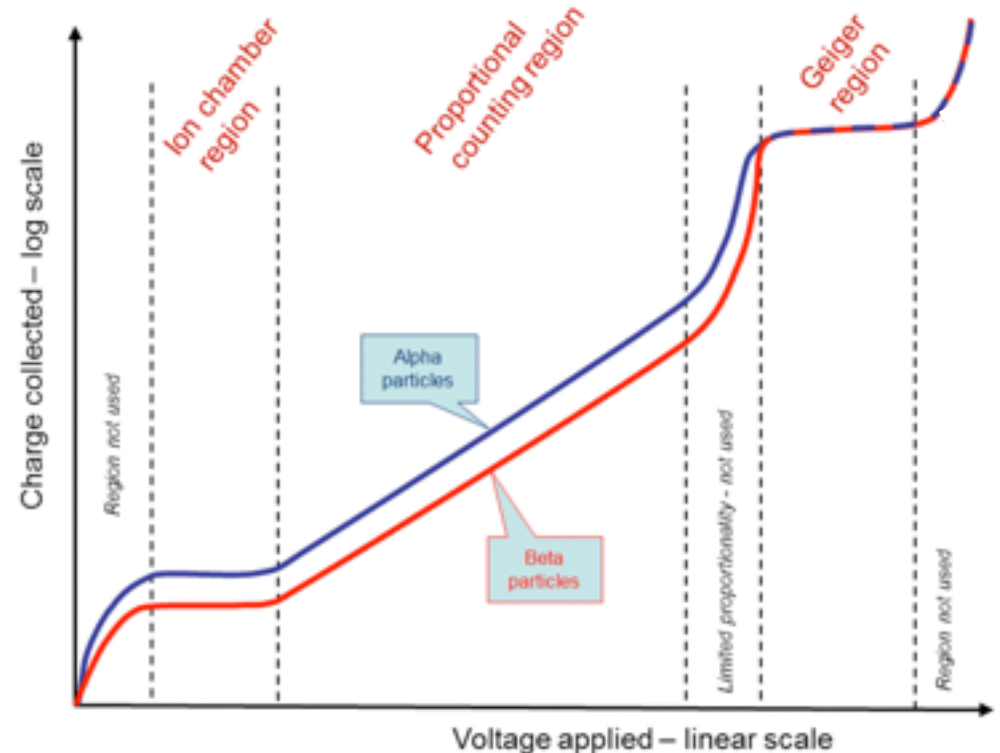


Signal



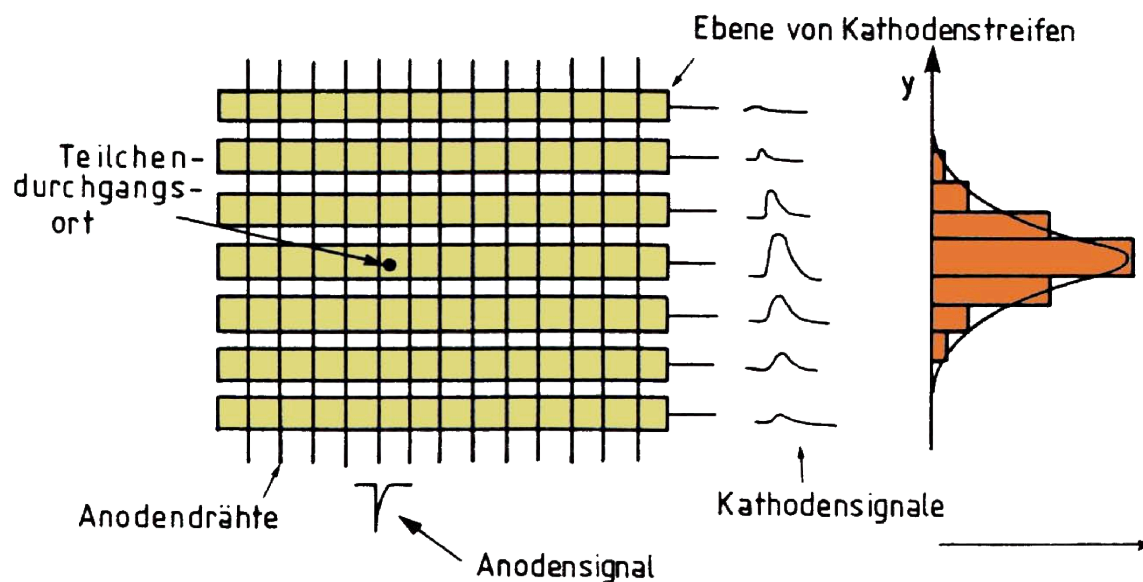
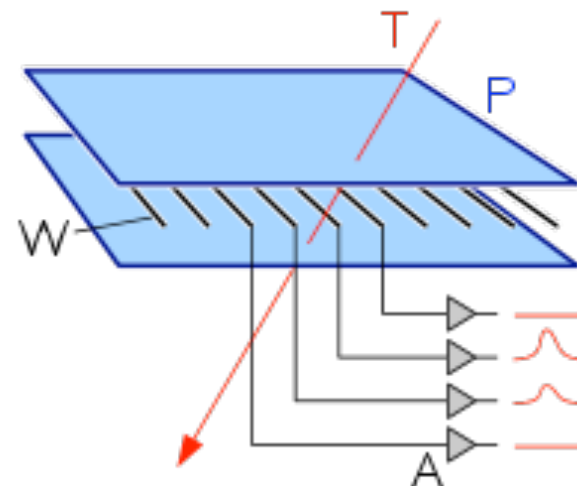
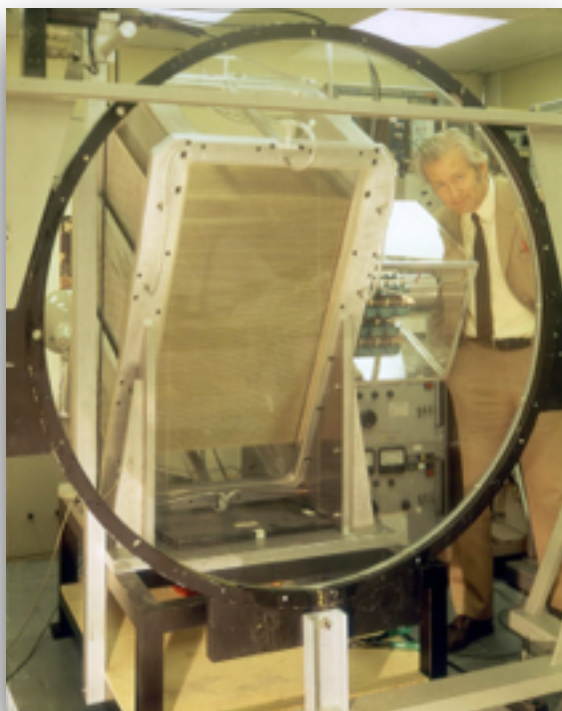
Variation of ion pair charge with applied voltage

- Passage of particles creates within the gas volume electron-ion pair (ionisation)
- Electrons are accelerated in a strong electric field \rightarrow amplification
- The signal is proportional to the original deposited charge or is saturated (depending on the voltage)

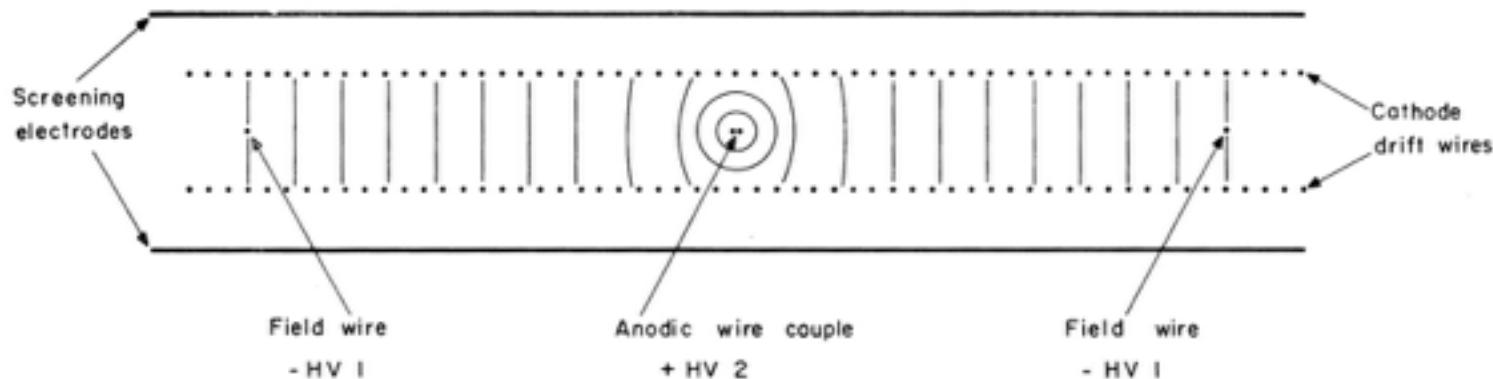


CONTINUATION OF IONISATION CHAMBERS

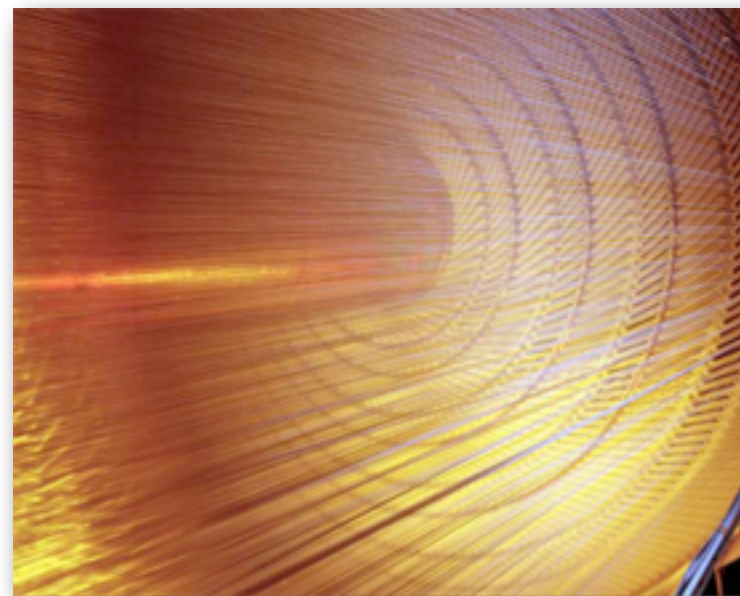
- Extreme successful approach to improve the spatial resolution of gas detectors
- Multi wire proportional chamber (MWPC)
- Gas-filled box with a large number of parallel detector wires, each connected to individual amplifiers
- G. Charpak 1968
(Nobel-prize 1992)



ADDING THE TIME: DRIFT CHAMBER

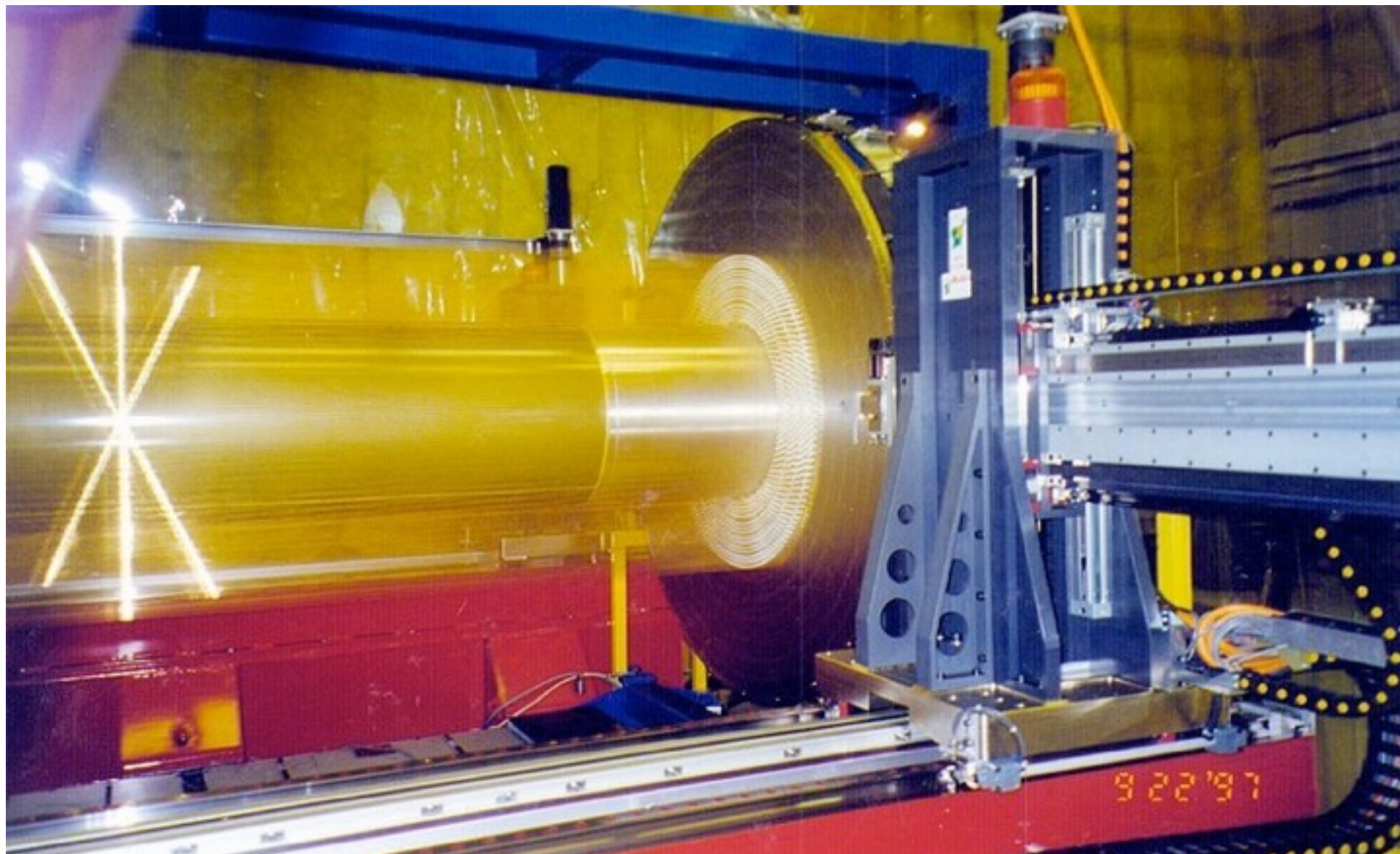


- Electric field is designed in a way that electrons drift with a constant velocity and only amplify very close to the wire
- If time of arrival of a particle is known (trigger), one can derive from the signal arrival time at the anode the position of the track
- Condition: the HV field distribution and therefore the drift velocity within the gas is well known

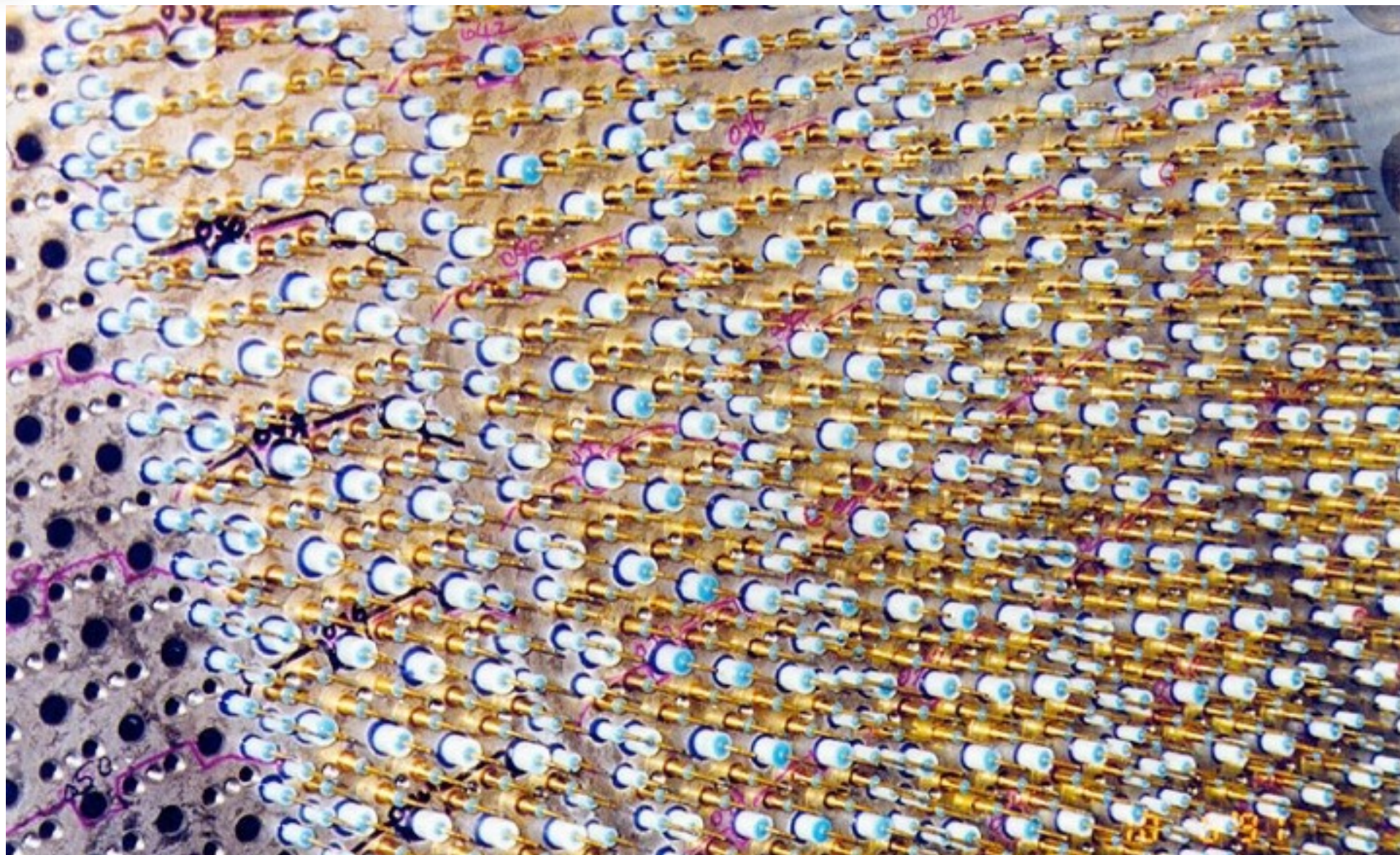


Wire chamber CDF (@Tevatron)

WIRE STRINGING IN PROGRESS

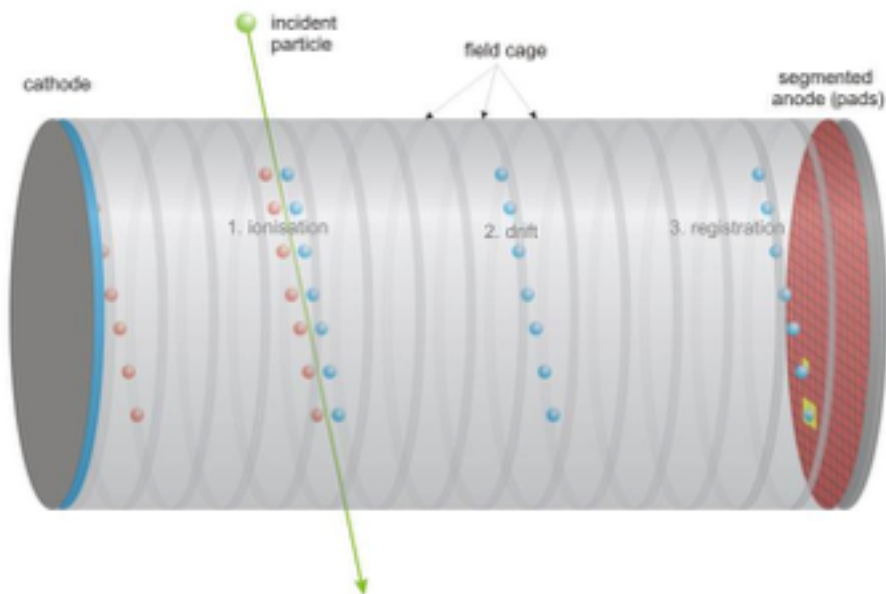


END PLATE CLOSE UP



TPC- TIME PROJECTION CHAMBER: 3D

- Combination of the the 2D track information and the time results in a real 3D point

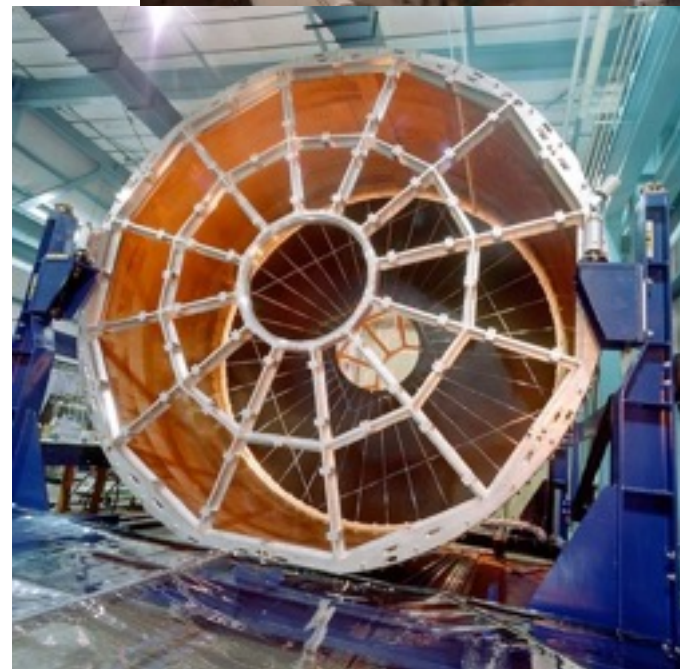


Pic: O. Schäfer



Pic: DESY

- Readout of the anode usually with multi wire projection chambers
- Nowadays new developments under way.



Pic: ALICE Collaboration

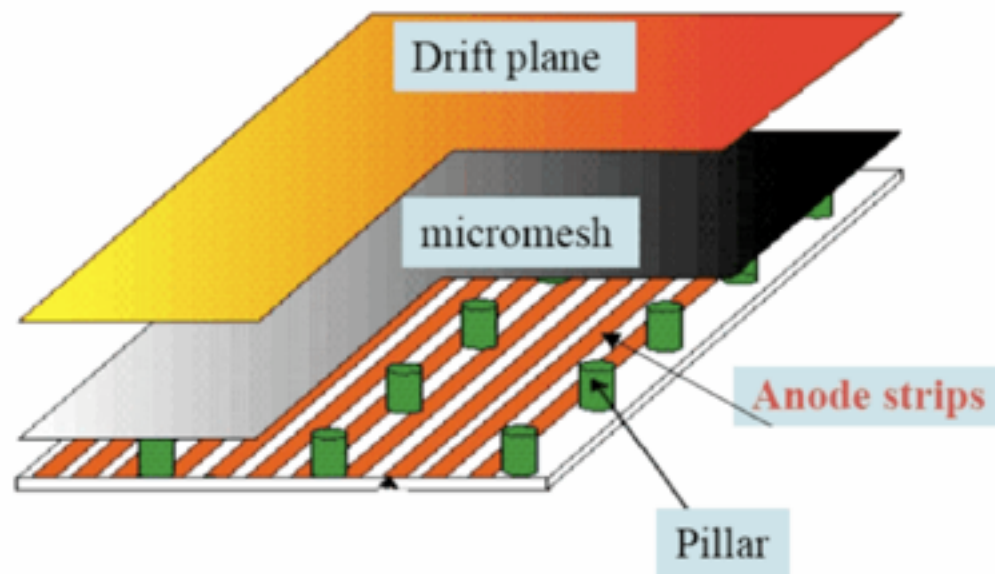
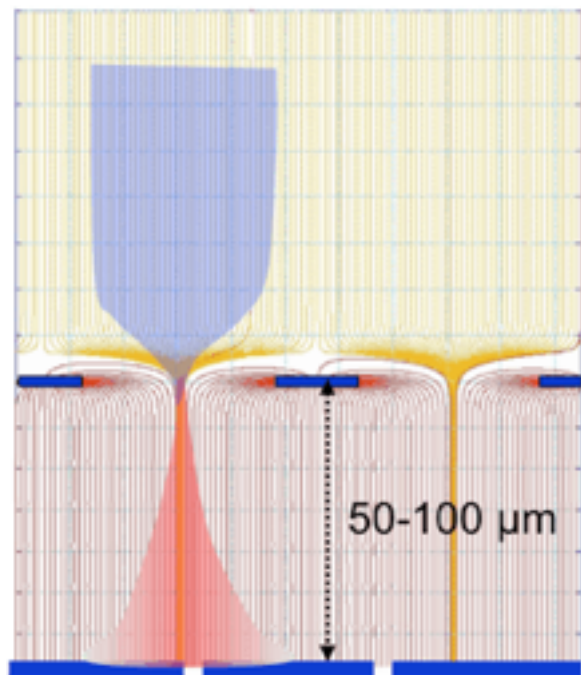
NEW DEVELOPMENTS

- Largely improved spacial resolution and higher particle rates:

Micro-Pattern Gas Detectors

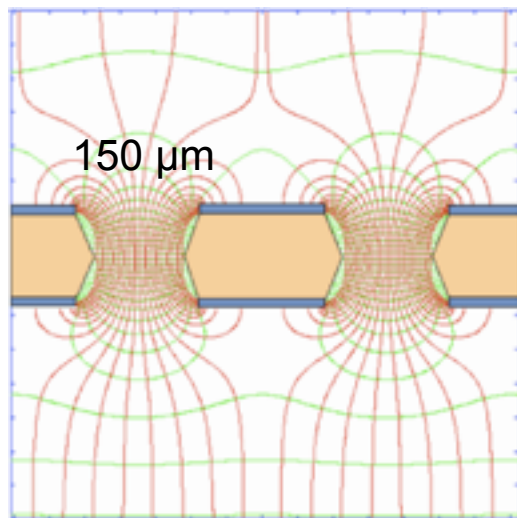
- a number of developments were started, some with a lot of problems
- two technologies are currently the most successful: GEMs and MicroMegas
- MicroMegas: Avalanche amplification in a small gap

Y. Giomataris et al, NIM A376, 29(1996)

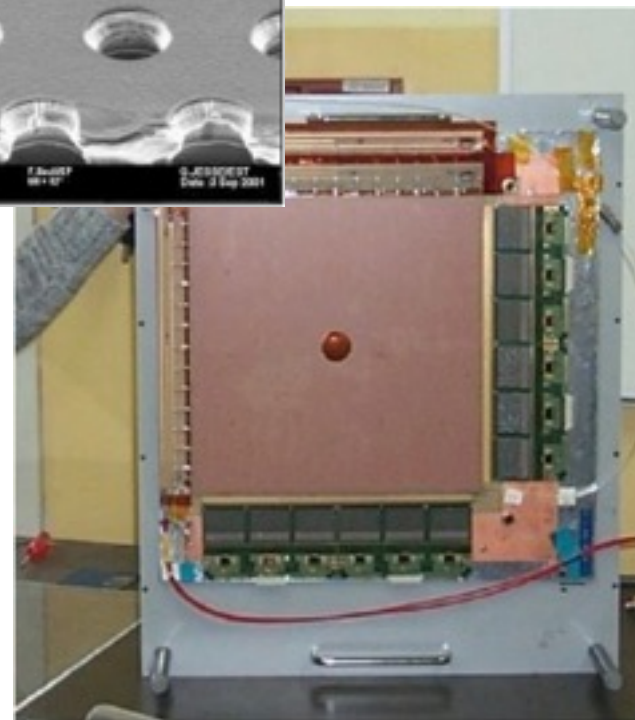
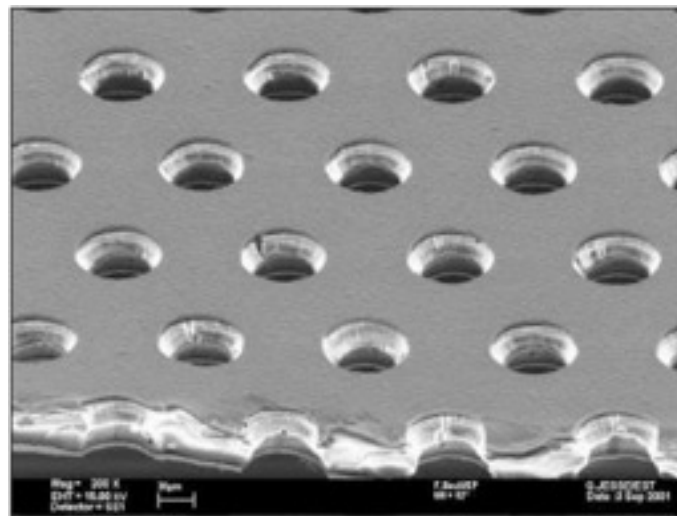


NEW DEVELOPMENTS

- GEM: Gas Electron Multiplier: Gas amplification in small holes in a special foil



50 μm

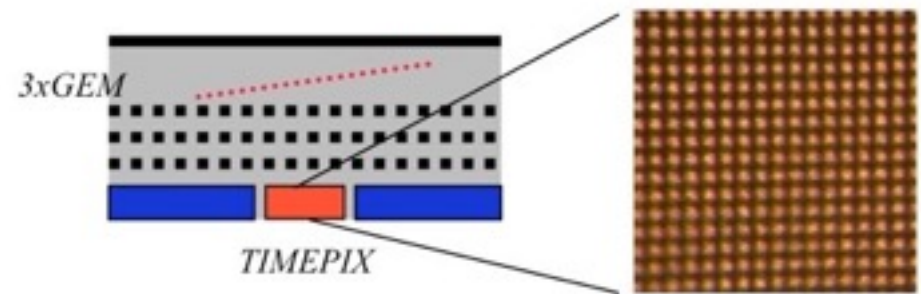


Charge collection on two separate levels: 2D structure possible: separation of amplification and read out

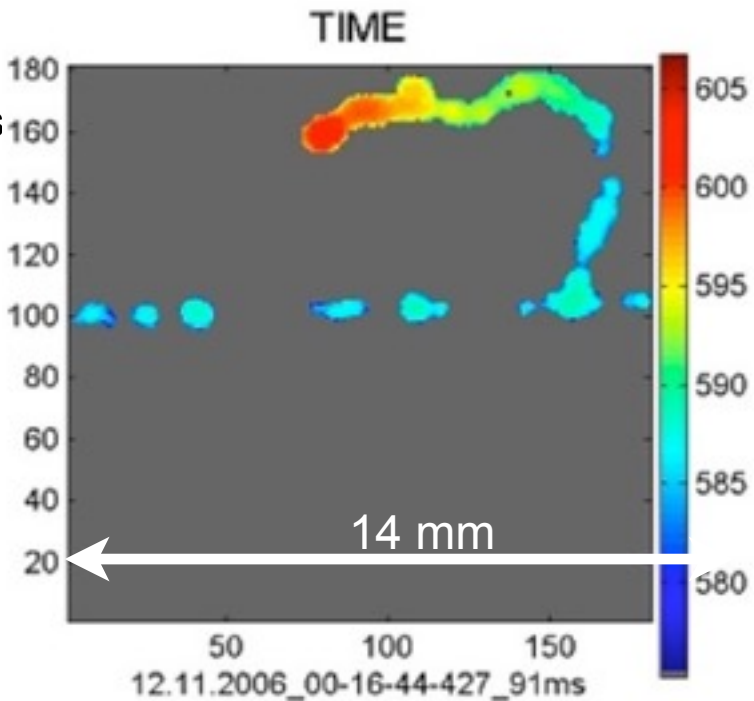
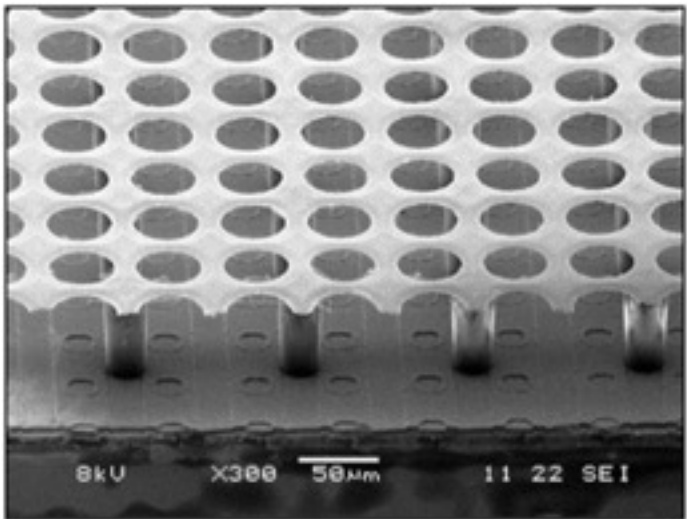
Both technologies, MicroMegas and GEMs are used in experiments. Typical spatial resolution: $\sim 70 \mu\text{m}$

MPGDs AS NEXT GENERATION DETECTOR

- Combination of gas detectors and Silicon
- Integration of MPGDs with pixel read out chips

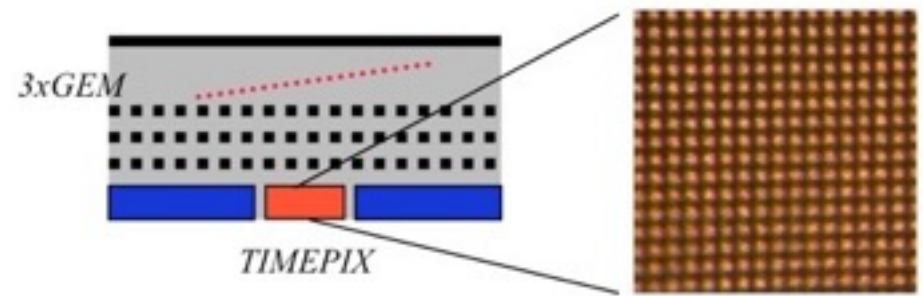


- Amplification and read out made of silicon

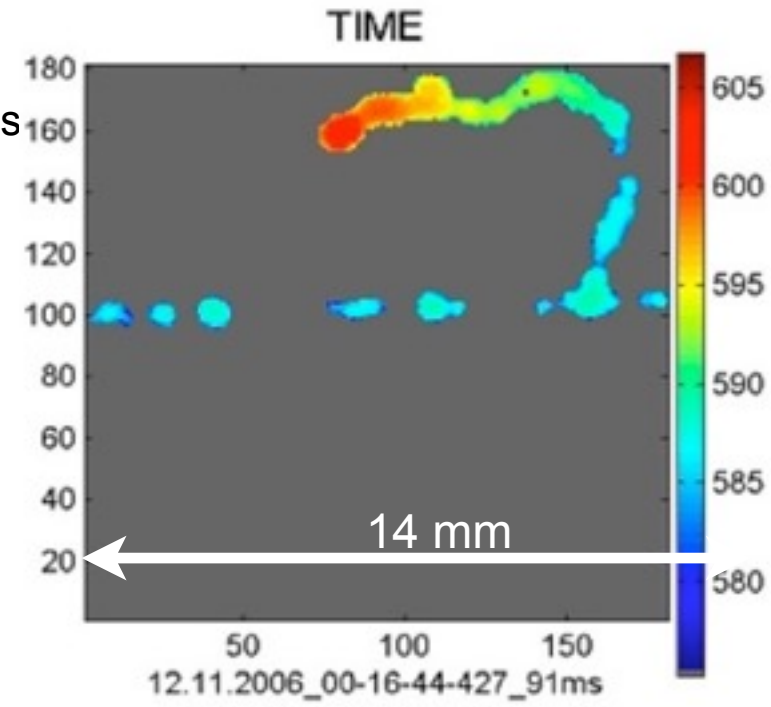
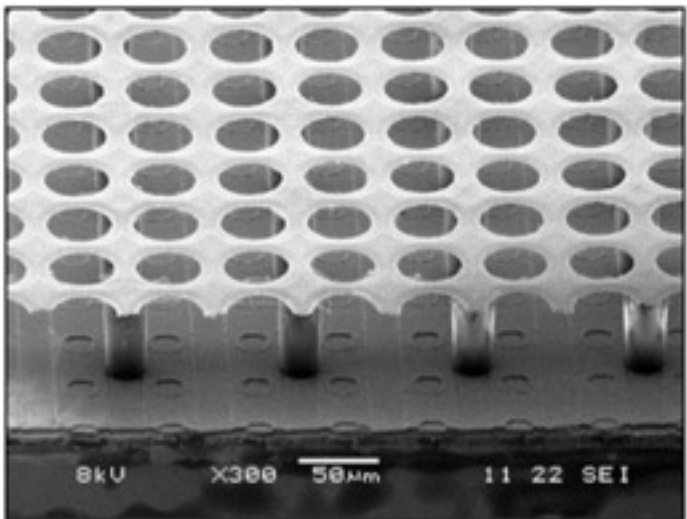


MPGDs AS NEXT GENERATION DETECTOR

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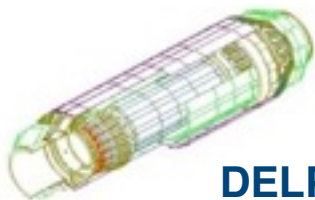


- Advantages of gas detectors:
- Low radiation length
 - Gas can be replaced regularly: Reduction of radiation damages!

V.B SEMICONDUCTOR-DETECTORS

LARGE SILICON SYSTEMS

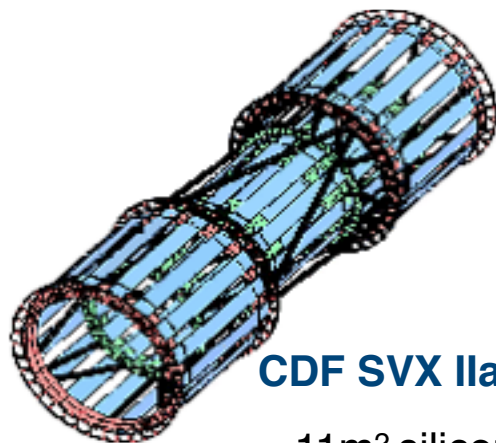
Since ~ 30 years: Semiconductor detectors for precise position measurements.



DELPHI (1996)

~ 1.8m² silicon area

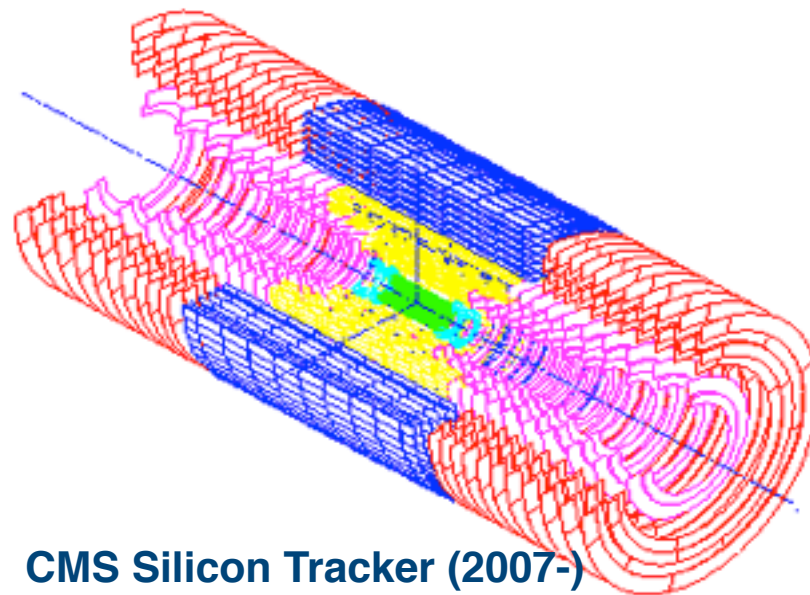
~ 175 000 readout channels



CDF SVX IIa (2001-2012)

~ 11m² silicon area

~ 750 000 readout channels



CMS Silicon Tracker (2007-)

~12,000 modules

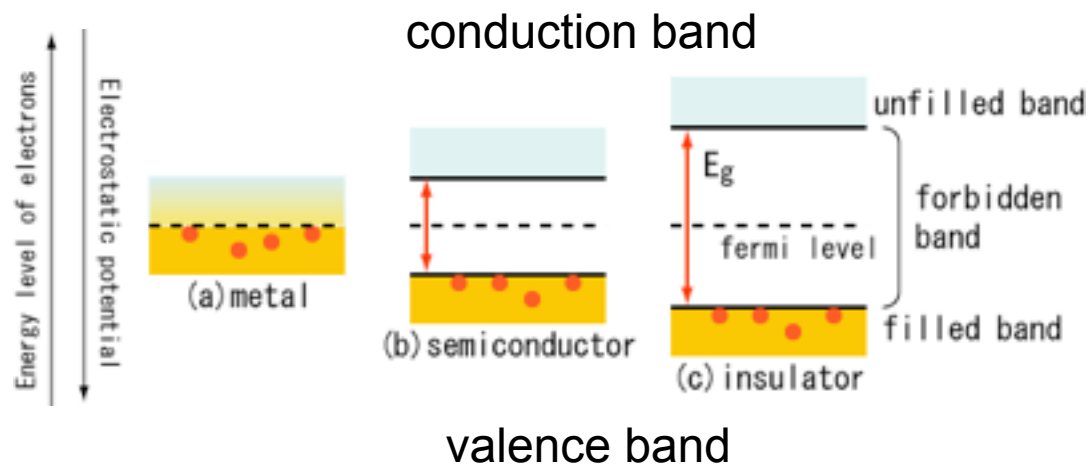
~ 223 m² silicon area

~25,000 silicon wafers

~ 10M readout channels

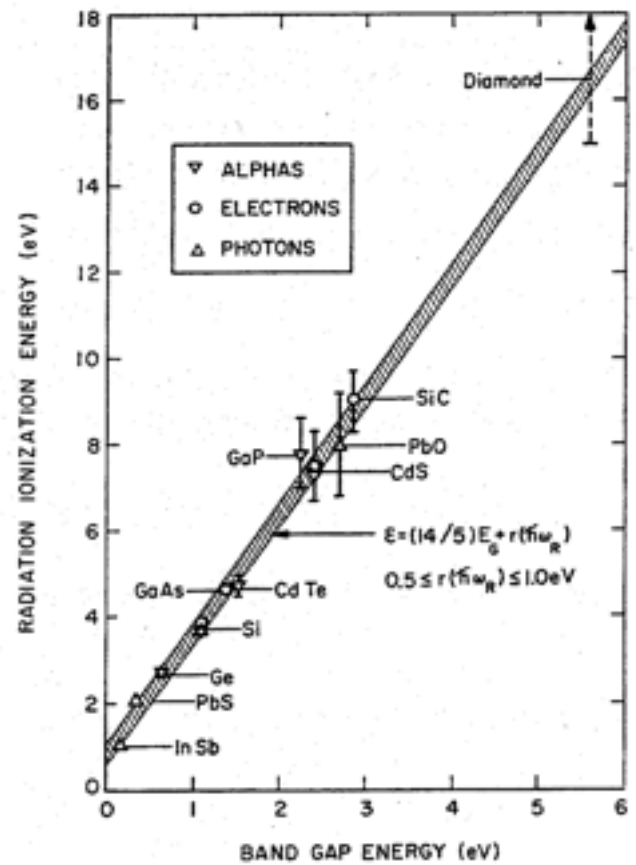
SEMICONDUCTOR BASICS I

- In free atoms the electron energy levels are discrete.
- In a solid, energy levels split and form a nearly-continuous band.



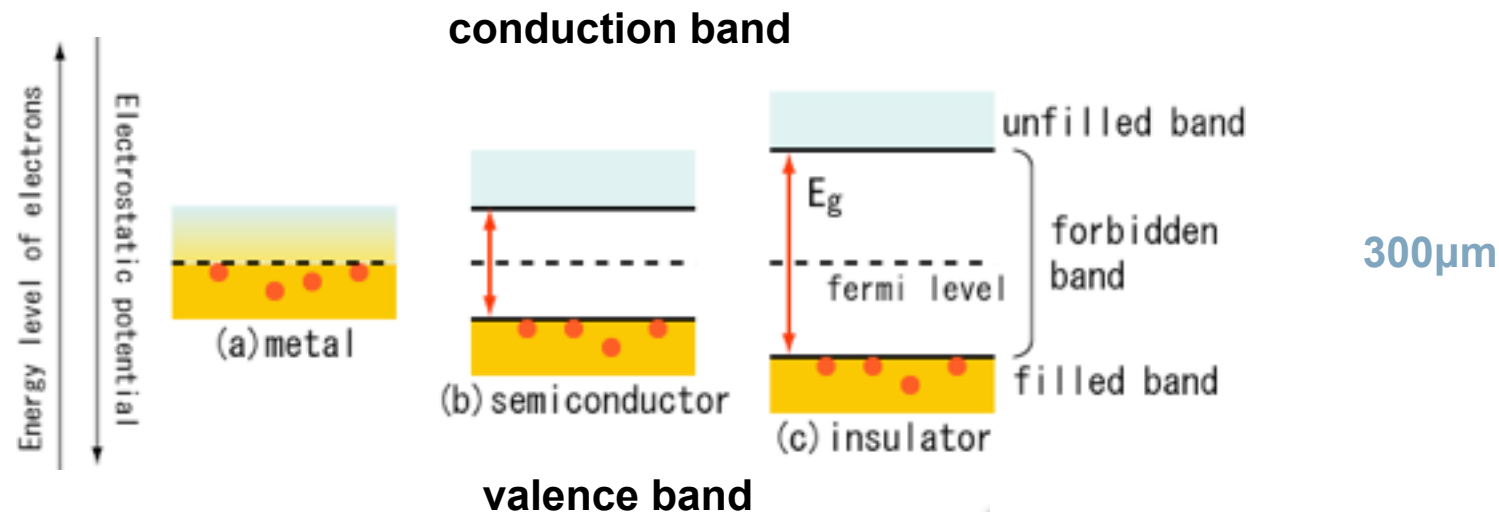
- Large gap: the solid is an insulator.
- No gap: it is a conductor.
- Small band gap: semiconductor
- For silicon, the band gap is 1.1 eV, but it takes 3.6 eV to ionize an atom -> rest of the energy goes to phonon excitations (heat).

C.A. Klein, J. Applied Physics 39 (1968) 2029



SEMICONDUCTOR BASICS

- In free atoms electron energy levels are discrete.
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- Large gap: the solid is an insulator.
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- Small band gap: semiconductor

Signal of charged particle:

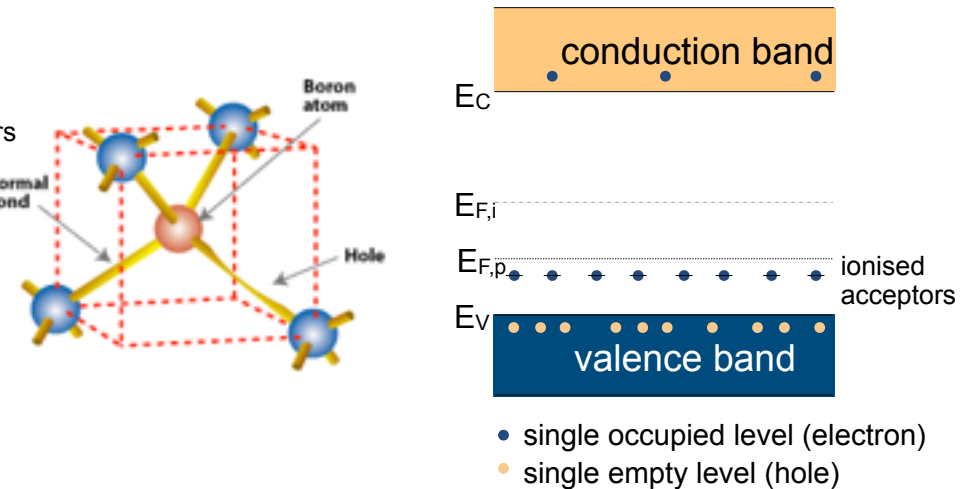
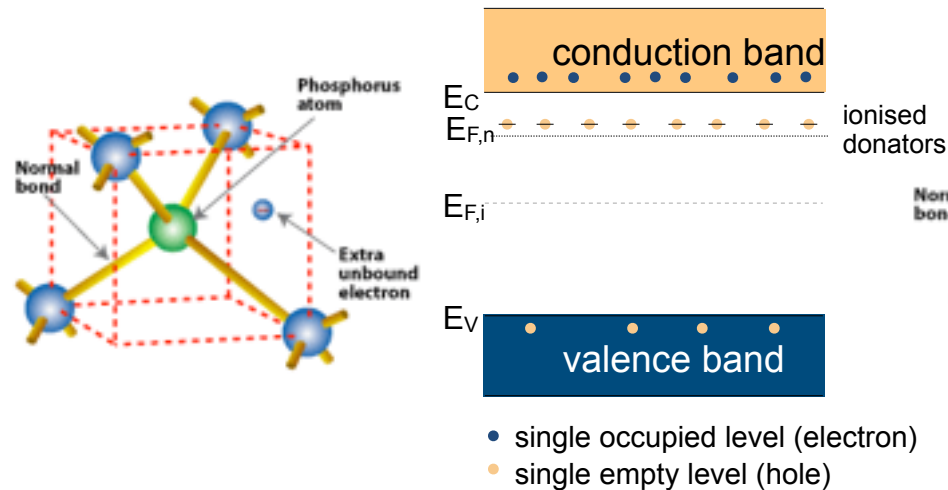
~ 3×10^4 electron/hole pairs

Intrinsic charge carrier

~ 4×10^8 electron/hole pairs

- For silicon, the band gap is 1.1 eV, but it takes 3.6 eV to ionise an atom
 - Remaining energy goes to phonon excitations (heat).

DOPING SILICON



n type semiconductor:

- ⊙ Negative charge carriers (electrons) by adding impurities of donor ions (e.g. Phosphorus (type V))
- ⊙ **Donors** introduce energy levels close to conduction band thus almost fully ionised (E_F closest to CB)

Electrons are the majority carriers.

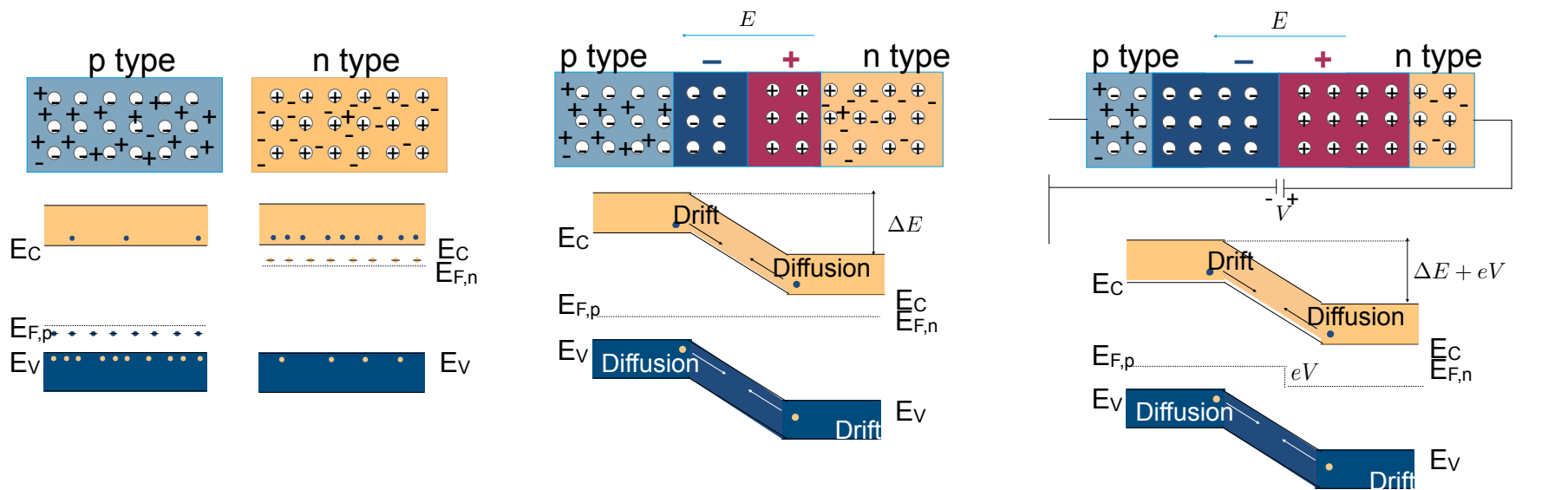
p type semiconductor:

- ⊙ Positive charge carriers (holes) by adding impurities of acceptor ions (e.g. Boron (type III)).
- ⊙ Acceptors introduce energy levels close to valence band thus 'absorb' electrons from VB, creating holes (E_F closest to VB).

Holes are the majority carriers.

BASIS OF SILICON DETECTOR: PN JUNCTION

- At interface of p type and n type semiconductor difference in the Fermi levels causes diffusion of excessive carriers to the other material until thermal equilibrium is reached.
- Stable space charge region free of charge carriers is called **depletion zone**.
- Typical current-voltage of a p-n junction: exponential current increase in forward bias, small saturation in reverse bias.



Applying an external voltage V with the cathode to p and the anode to n (reverse biasing), e-h pairs are pulled out of the depletion zone. → **larger depletion zone** → **suppress current across the junction**

PRINCIPLE OF SEMICONDUCTOR

- Creation of electric field: voltage to deplete thickness d

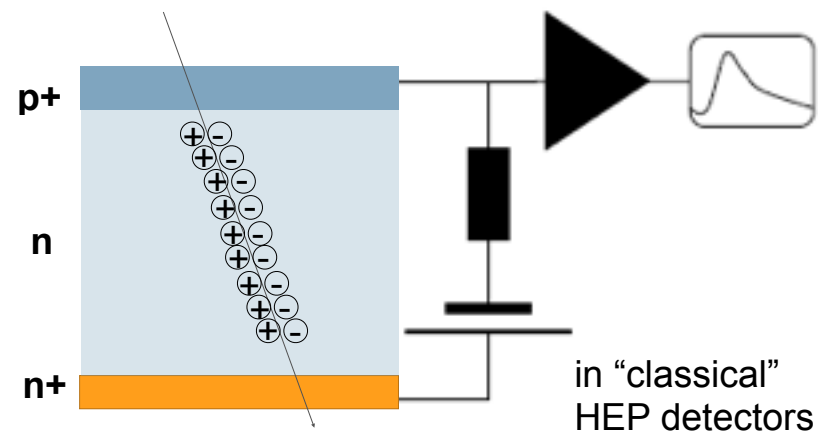
with $n_A \gg n_D$

$$d = \sqrt{\frac{2\epsilon\epsilon_0 V_{dep}}{en_D}}$$

for $d = 300\mu m$ $V_{dep} \approx 160V$

- Passage of a charged particle: Electron-hole pairs formed in the depletion zone
 - Drift under the influence of the electric field
 - Signal depends on width of depletion zone

The signal is induced by the motion of charge after incident radiation (not when the charge reaches the electrodes).



Typical numbers

Doping concentration

$$n_A \approx 10^{19} cm^{-3}$$

Acceptors

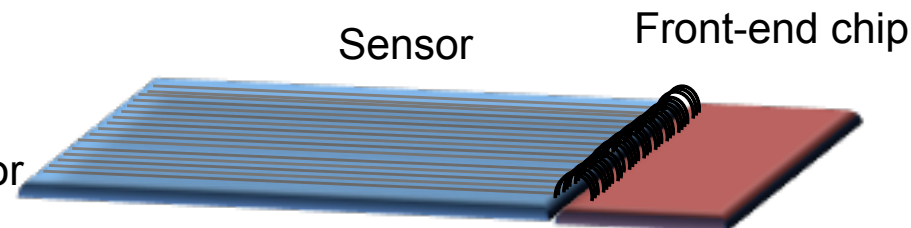
$$n_D \approx 2 \cdot 10^{12} cm^{-3}$$

Donators

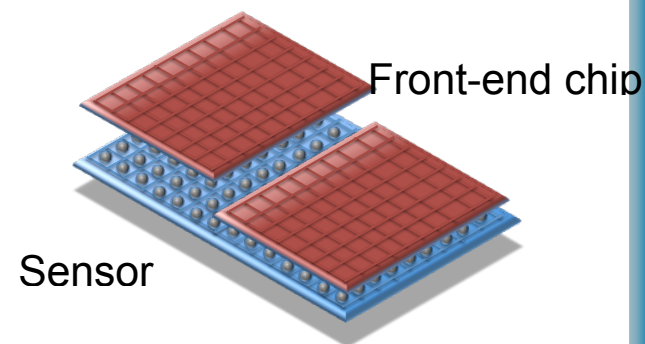
in "classical" HEP detectors

STRIPS AND PIXELS

- **Strips detector:** charge sensed by long narrow strips
1D information (typically 20 - 100 μ m)
- 2D information by double sided processing or adding back to back second layer slightly rotated (stereo angle)
- In regions with higher track density one dimensional measurements can lead to ambiguities.
- **Pixel detector:** charge sensed by small pixels on one side of sensor
 - Hybrid pixels: sensors and readout joined via bump bonds
 - Monolithic pixels: sensor and readout on one substrate



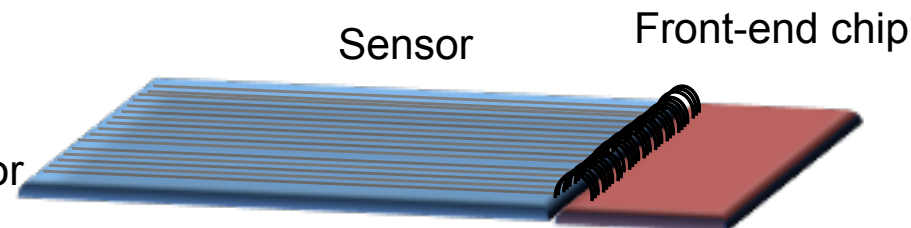
Microstrips detector



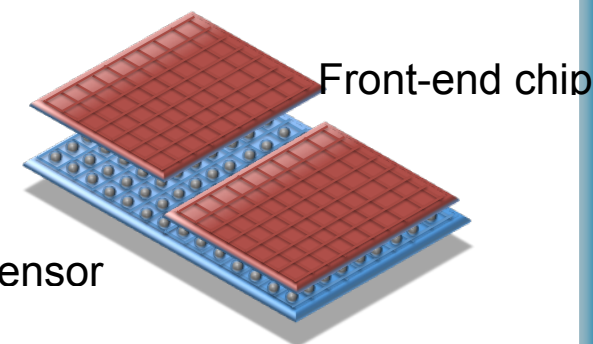
Hybrid pixel detector

STRIPS AND PIXELS

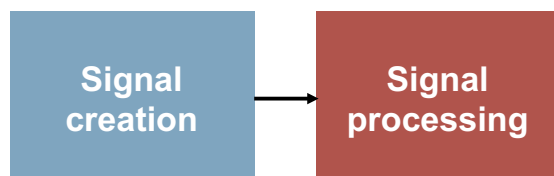
- **Strips detector:** charge sensed by long narrow strips
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Microstrips detector



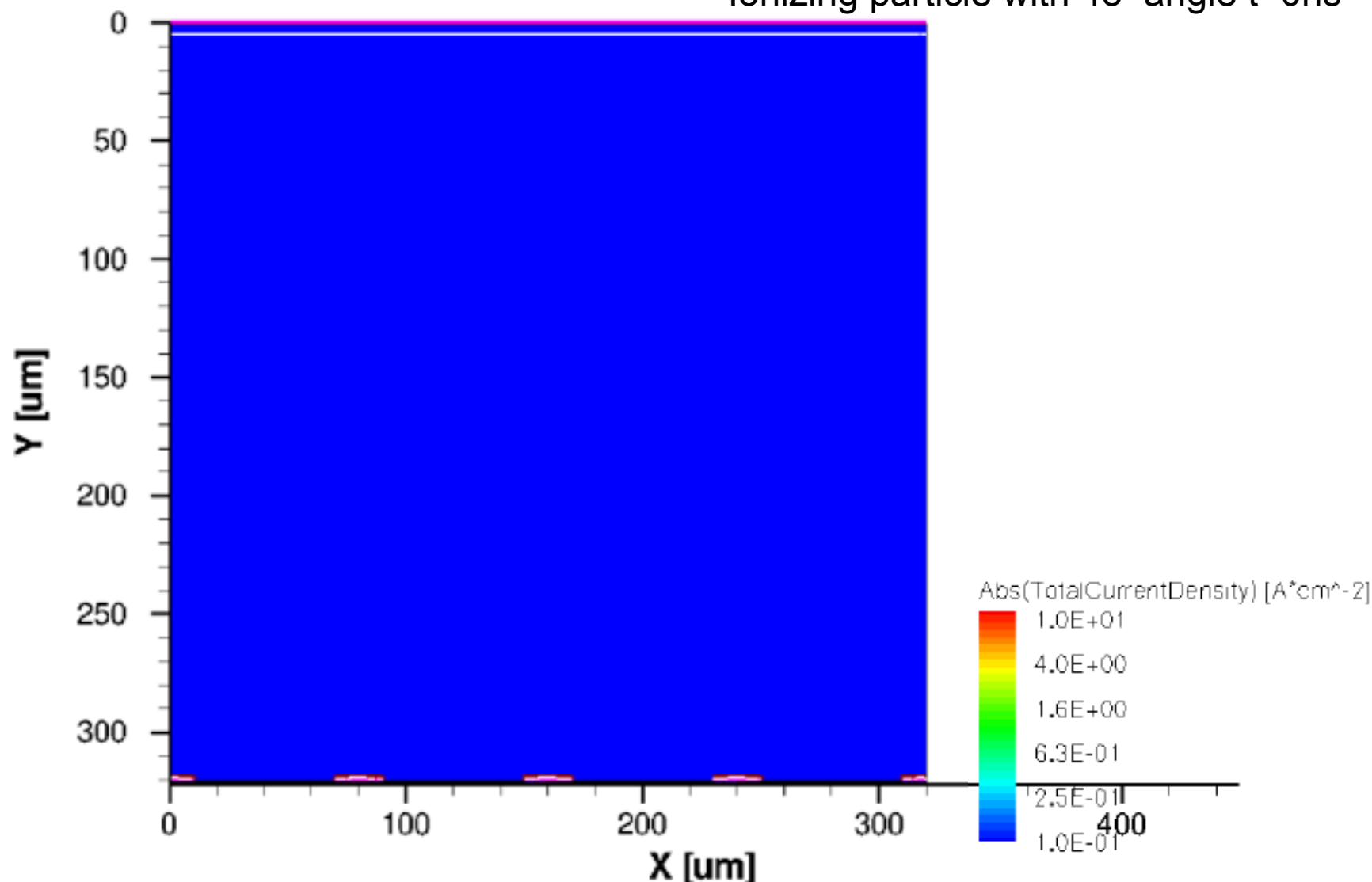
Hybrid pixel detector



- Signals created in silicon by charged particle
 - Very small signals (fC): need amplification
 - Measurement of amplitude/hit and/or time
 - Several thousands to millions of channels

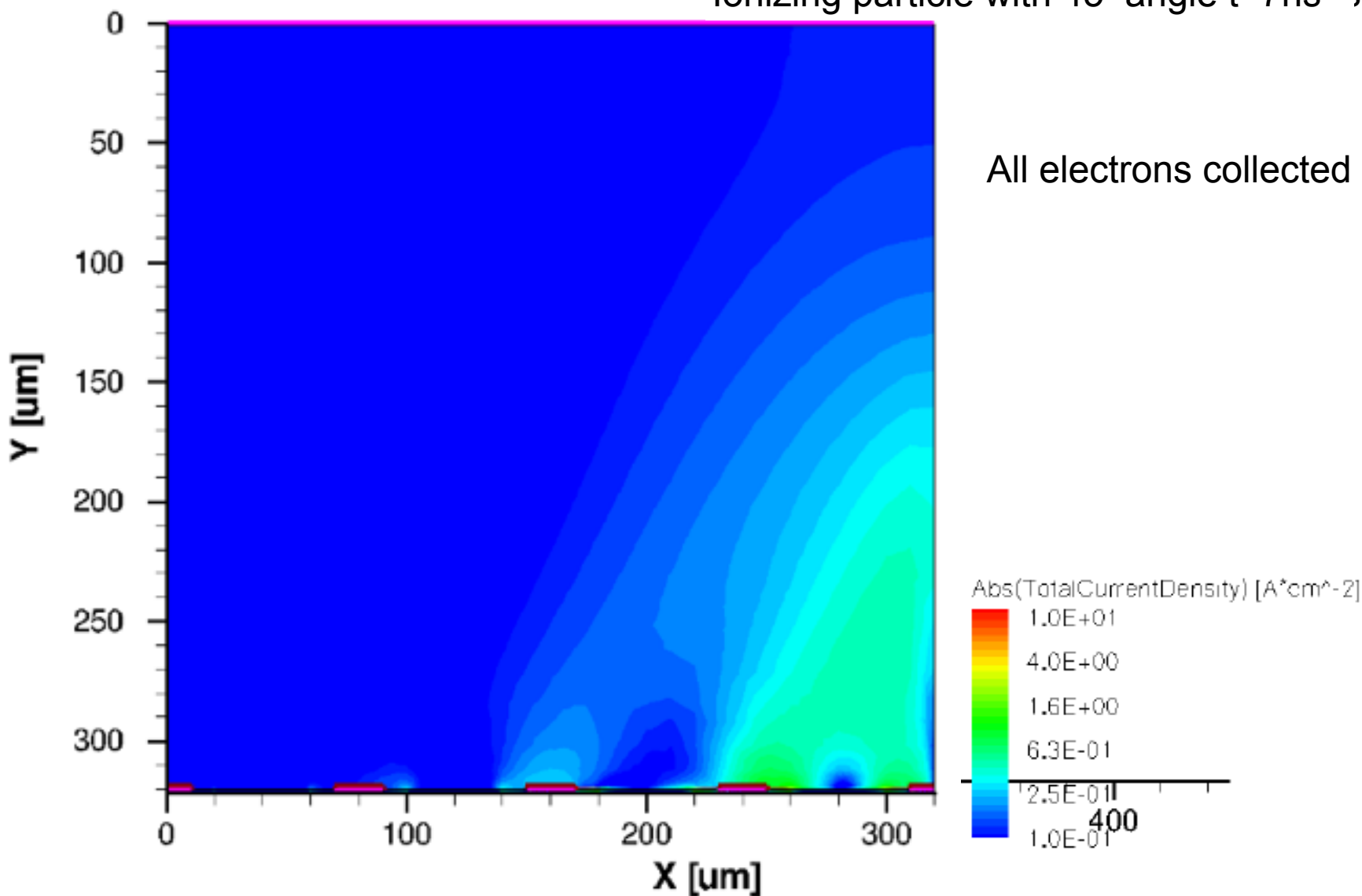
CURRENT DENSITY

Ionizing particle with 45° angle $t=0\text{ns}$



CURRENT DENSITY

Ionizing particle with 45° angle $t=7\text{ns}$;

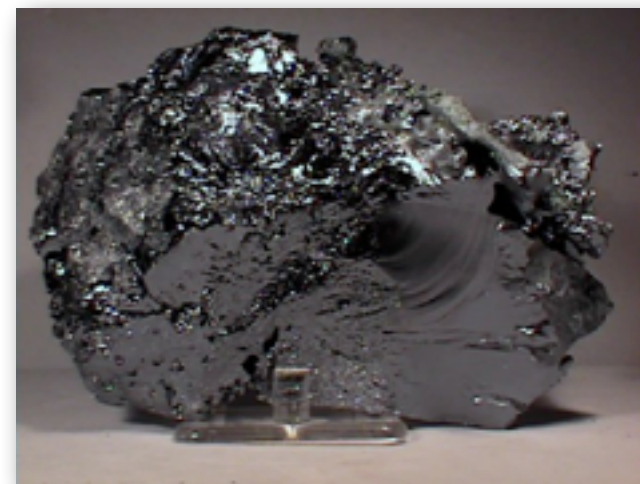


MATERIAL PROPERTIES

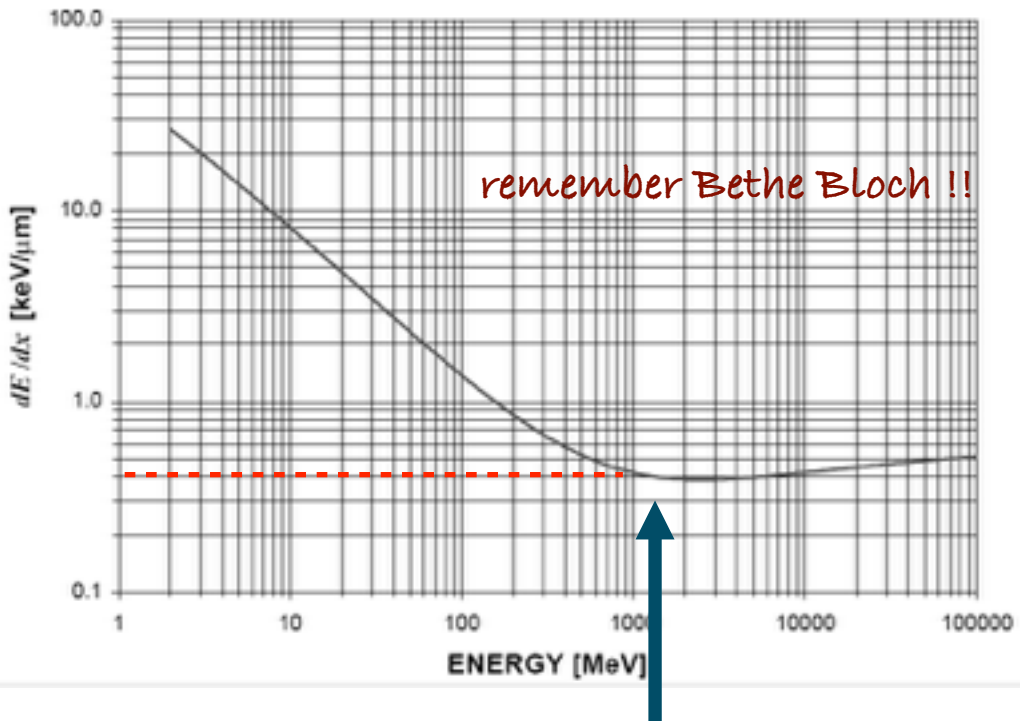
	Si	Ge	GaAs	CdTe	Diamond	SiC
band gap	1.12	0.67	1.42	1.56	5.48	2.99
energy for e-p pair [eV]	3.6	2.9	4.2	4.7	13.1	6.9
e- for MIP (300 μ m)	24000	50000	35000	35000	9300	19000
Z	14	32	31+33	48+52	6	14+6

Why is silicon used more often ?

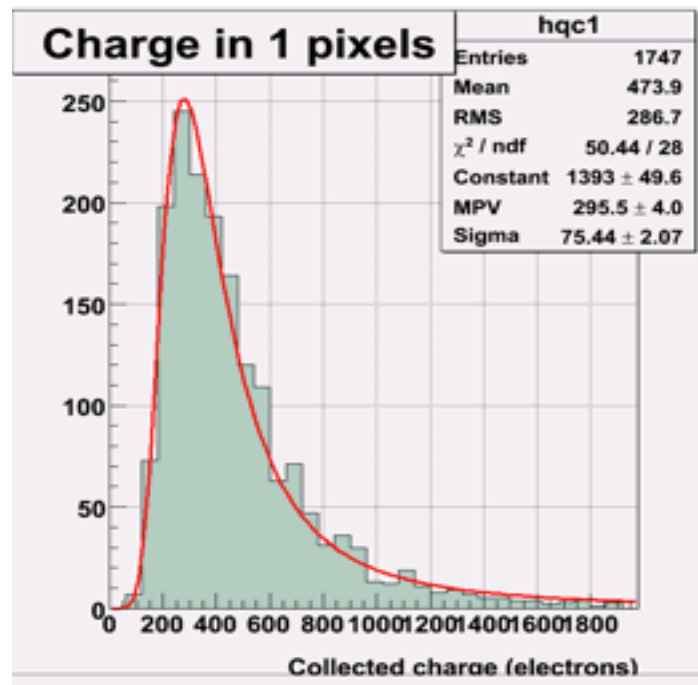
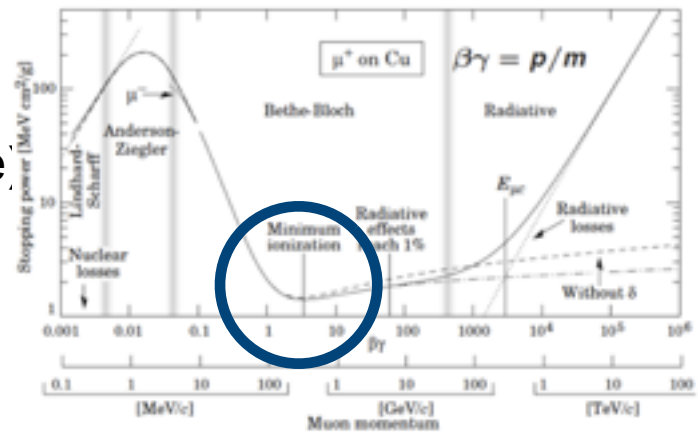
- Silicon is the only material which can be produced in larger areas in high quality
- compare to $kT = 0.026$ eV at room temperature -> dark current under control
- high density compared to gases: $\rho=2.33\text{g/cm}^3$
- good mechanical stability -> possible to produce mechanically stable layers
- large charge carrier mobility
- fast charge collection $\delta t \sim 10\text{ns}$
- well understood -> radiation tolerant



PROTONS IN SILICON

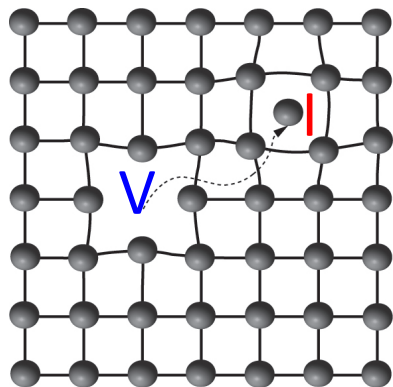


ue



PROBLEM: RADIATION DAMAGE

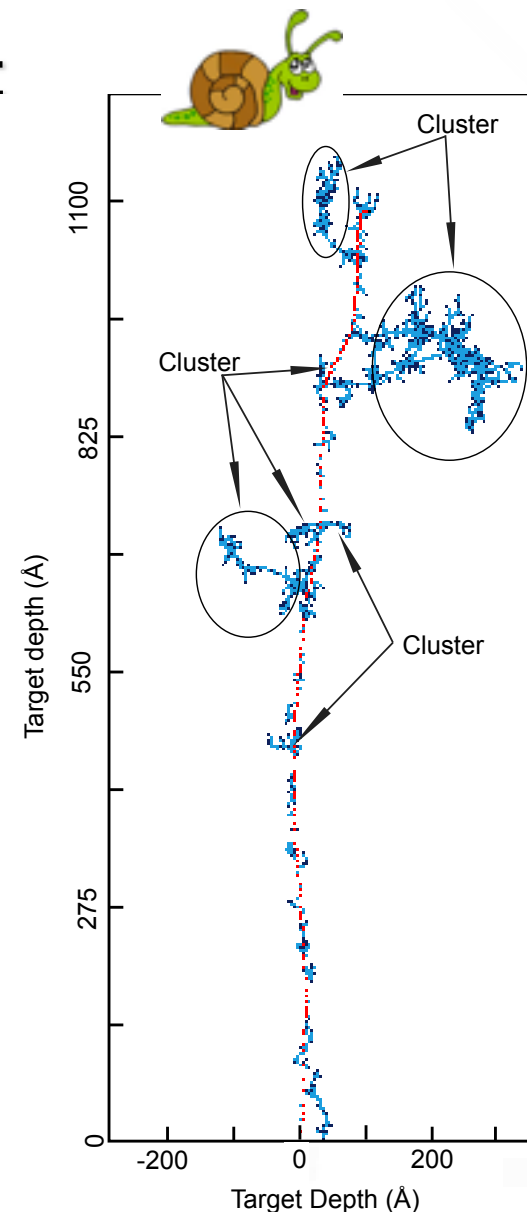
- Radiation damage the silicon on atomic level significantly leading to macroscopic effect.
- **Bulk effects:** displacement damage and built up of crystal defects due to Non Ionising Energy Loss (NIEL) (**main problem for sensors**).
unit: 1MeV equivalent n/cm² (up to 10¹⁵ n_{eq}/cm²)



Defects composed of:
V acancies and I nterstitials

Compound defects with impurities possible!

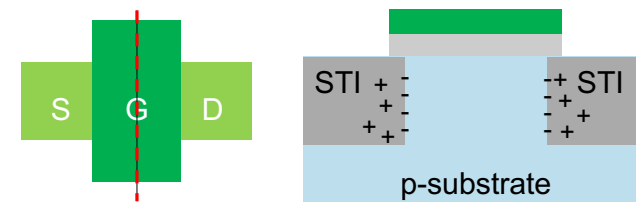
- **Surface effects:** Generation of charge traps due to ionizing energy loss (Total ionising dose, TID)
(**main problem for electronics**).
unit: Rad (up to 100 MRad)



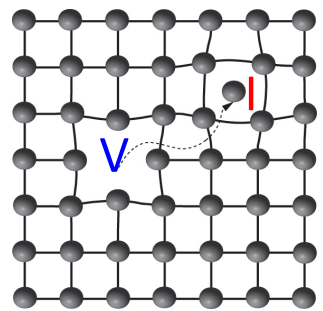
Simulation of 50 keV PKA damage cascade (1 MeV n)

PROBLEM: RADIATION DAMAGE

- Radiation damages the silicon on atomic level significantly leading to macroscopic effects.
- **Surface effects:** Generation of charge traps due to ionising energy loss — Total ionising dose, TID **(problem for sensors and readout electronics)**.
 - Cumulative long term trapping of positive charge
 - Increase of leakage current and oxide breakdown
- **Bulk effects:** displacement damage and build up of crystal defects due to non ionising energy loss (NIEL) **(main problem for sensors)**.
 - Unit: 1MeV equivalent n/cm²



STI = shallow trench interface



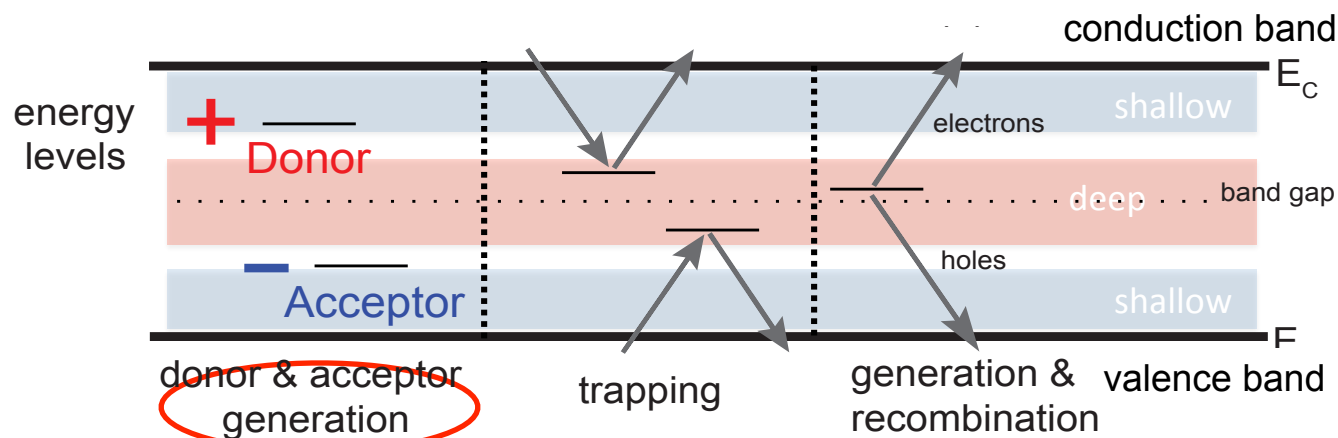
Defects composed of:
Vacancies and **I**nterstitials

Compound defects with
impurities possible!

Detector	NIEL [1MeV n _{eq} /cm ²]	TID [Mrad]
ALICE ITS	10 ¹³	<1
Belle II (per year)	1.2x10 ¹³	1.9
ATLAS/CMS Outer Tracker	10 ¹⁵	<100
ATLAS/CMS Inner Tracker	10 ¹⁶	1000

RADIATION DAMAGE: MACROSCOPIC EFFECTS

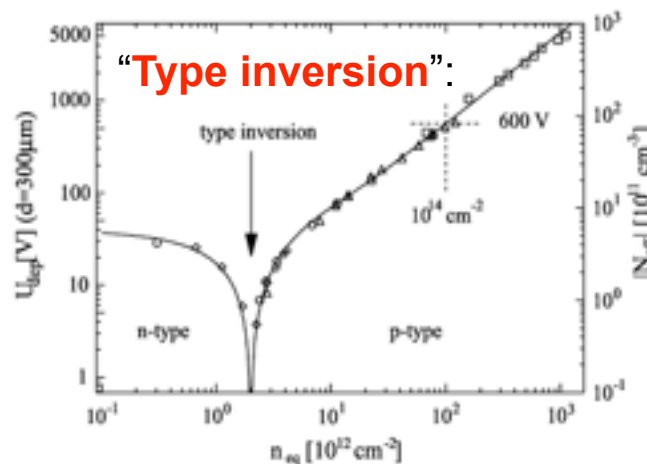
- Impact of defects on detector properties depends on defect level in band gap



Donor&acceptor generation:
Change of effective doping concentration (N_{eff})

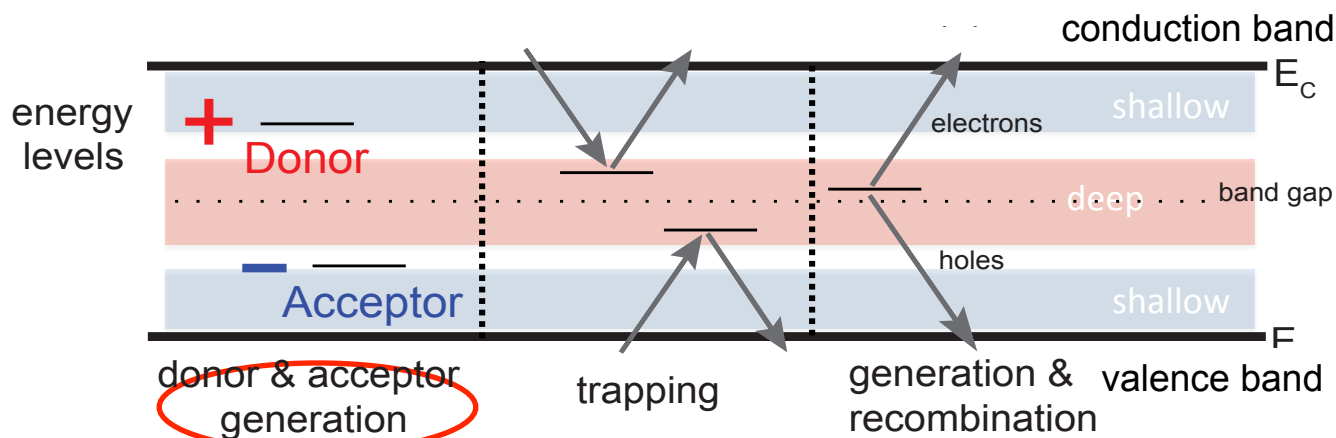
$$V_{dep} = d^2 N_{eff} \frac{q}{e \epsilon \epsilon_0}$$

- Increase of depletion voltage
- Under-depleted operation



RADIATION DAMAGE: MACROSCOPIC EFFECTS

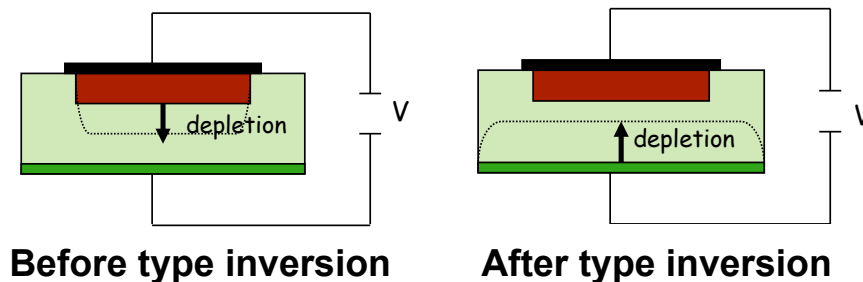
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Donor&acceptor generation:
Change of effective doping concentration (N_{eff})

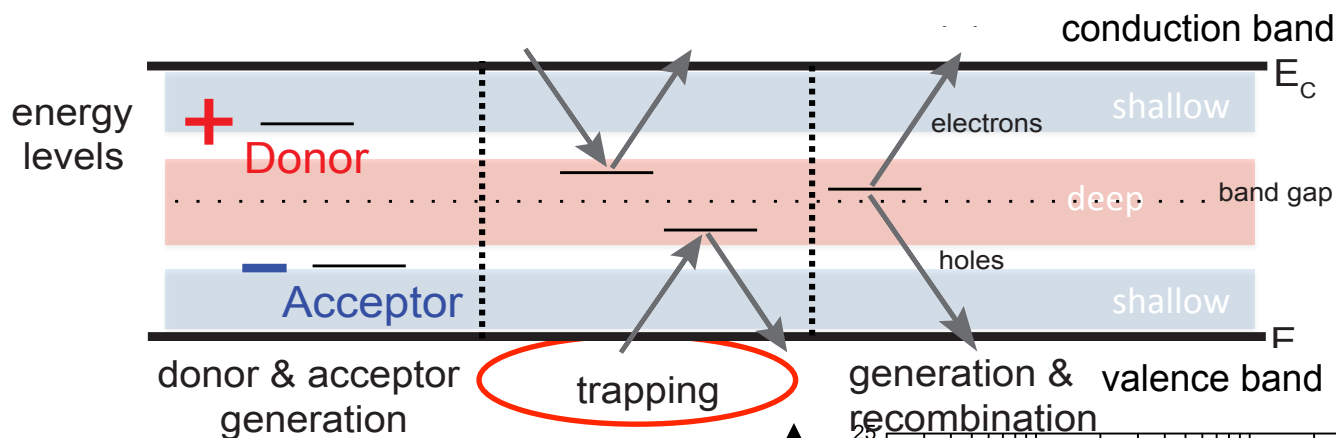
$$V_{dep} = d^2 N_{eff} \frac{q}{e \epsilon \epsilon_0}$$

- Increase of depletion voltage
- Under-depleted operation



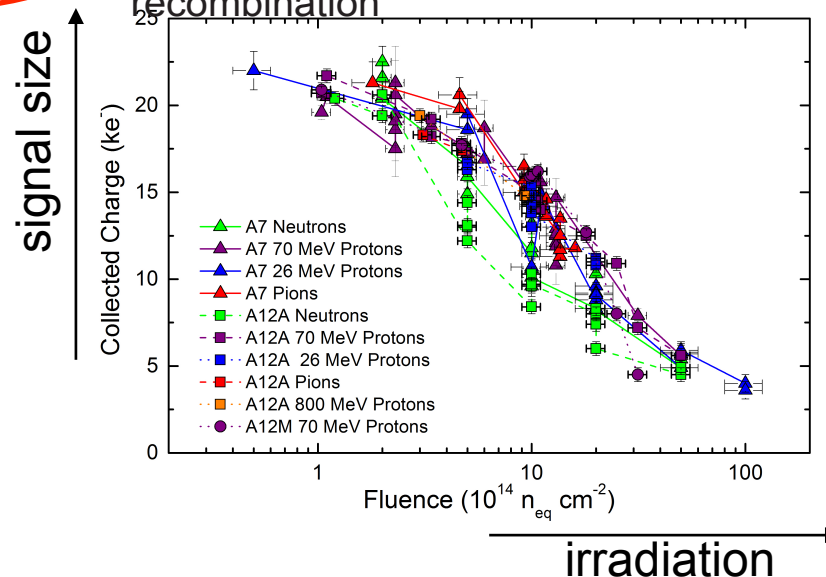
RADIATION DAMAGE: MACROSCOPIC EFFECTS

- Impact of defects on detector properties depends on defect level in band gap



Increased charge trapping

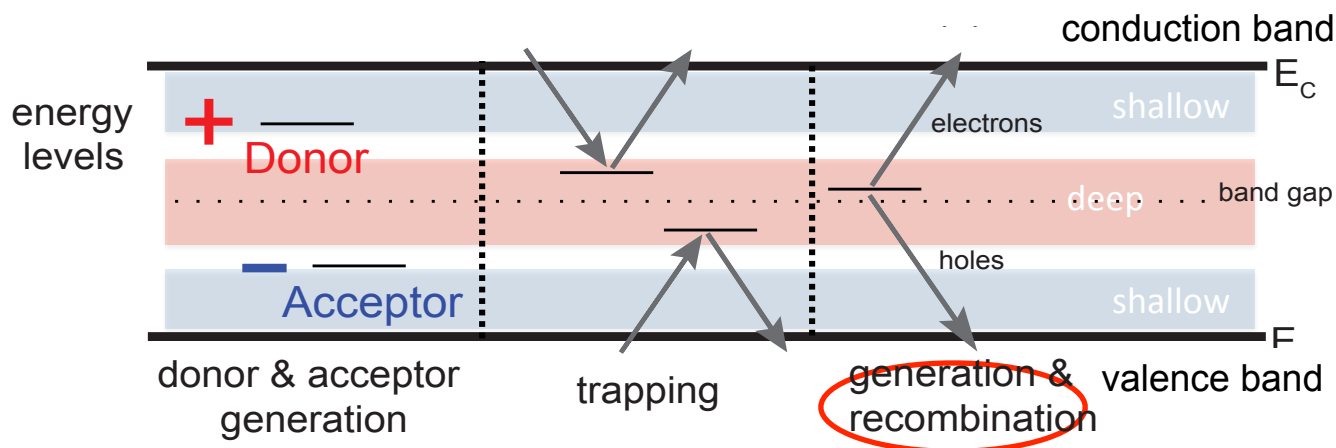
Lower signal (less charge)
Reduced charge collection efficiency



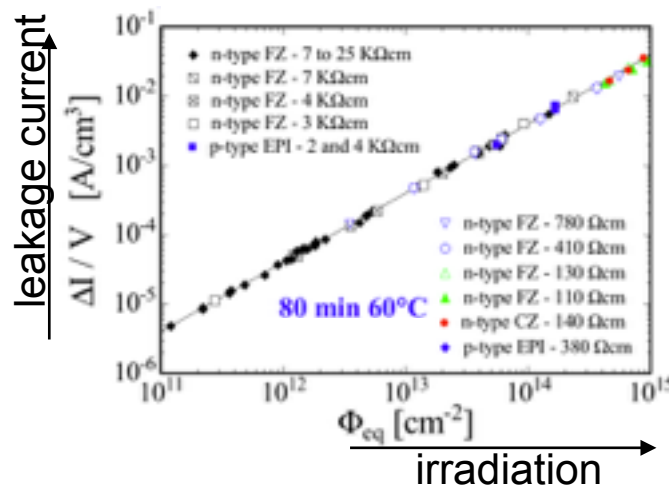
ATLAS ITk Strips TDR, April 2017

RADIATION DAMAGE: MACROSCOPIC EFFECTS

- Impact of defects on detector properties depends on defect level in band gap



Increase of leakage current
 higher shot noise; thermal runaway
 → Cooling during operation helps
 (leakage current depends on T)

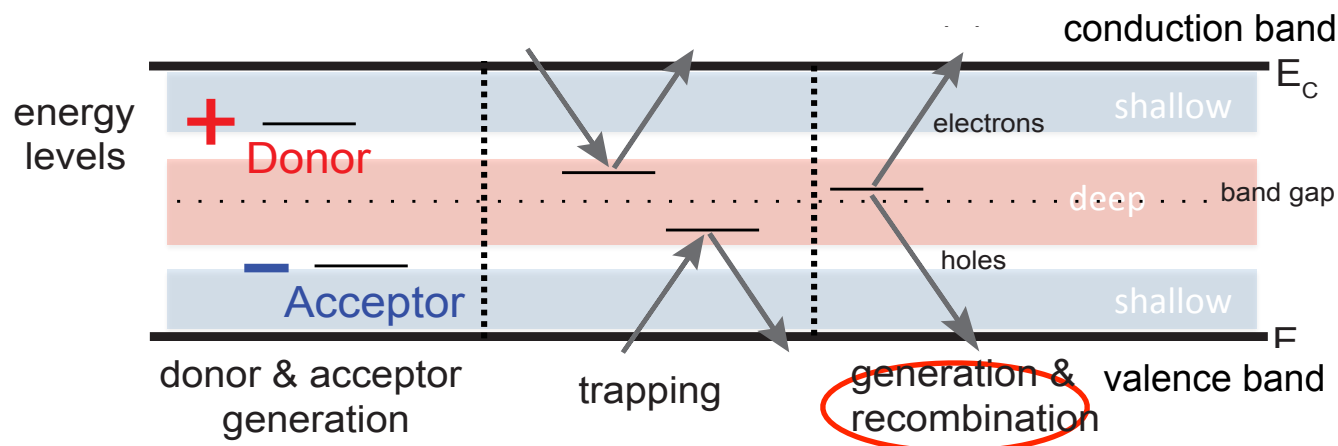


$$\alpha = \frac{\Delta I}{V \Phi_e}$$

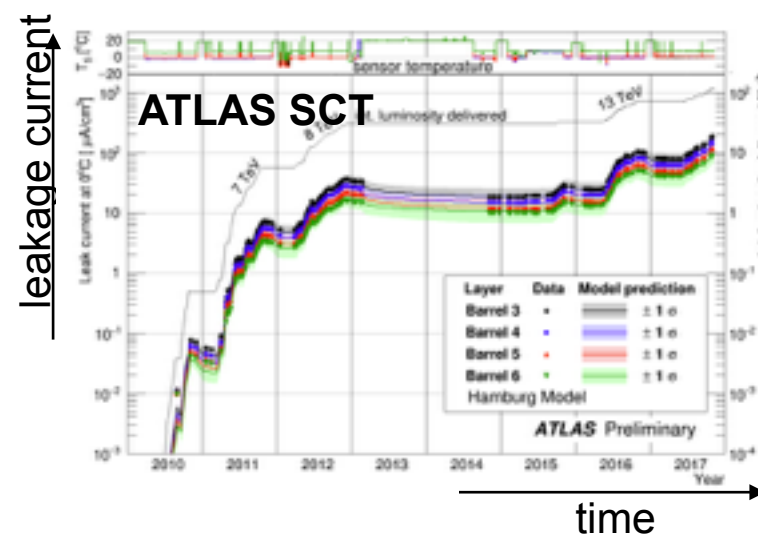
α is constant over
 several orders of
 fluence and
 independent of

RADIATION DAMAGE: MACROSCOPIC EFFECTS

- Impact of defects on detector properties depends on defect level in band gap

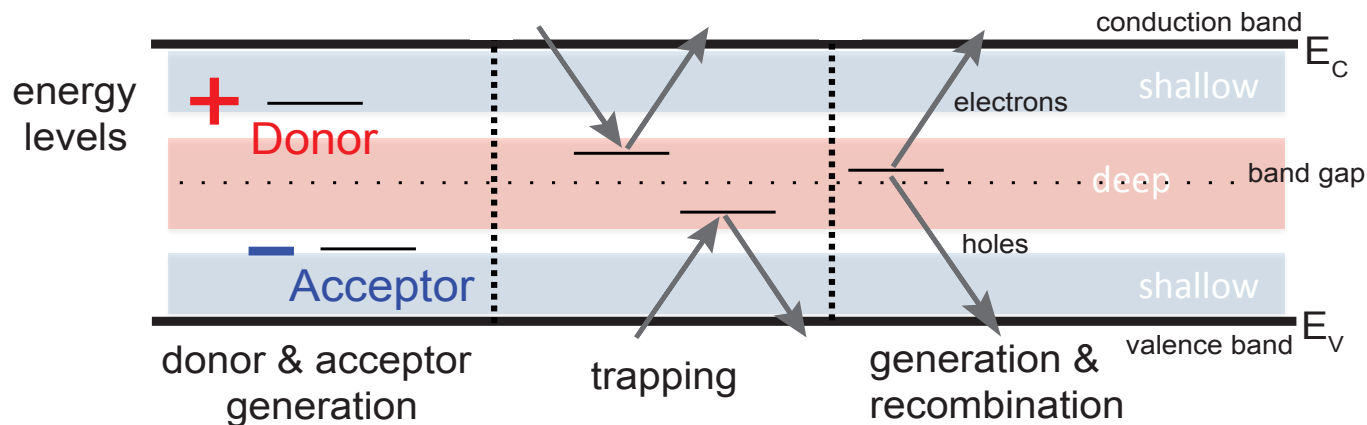


Increase of leakage current
 higher shot noise; thermal runaway
 → Cooling during operation helps
 (leakage current depends on T)



RADIATION DAMAGE: BULK DEFECTS

- Impact of defects on detector properties depends on defect level in band gap



Change of effective doping concentration (N_{eff})

Can contribute to space charge:

- increase of depletion voltage
- under-depleted operation

Increased charge trapping

Loss of signal (reduced charge collection efficiency)

Increase of leakage current

higher shot noise
thermal runaway

COFFEE BREAK ??

