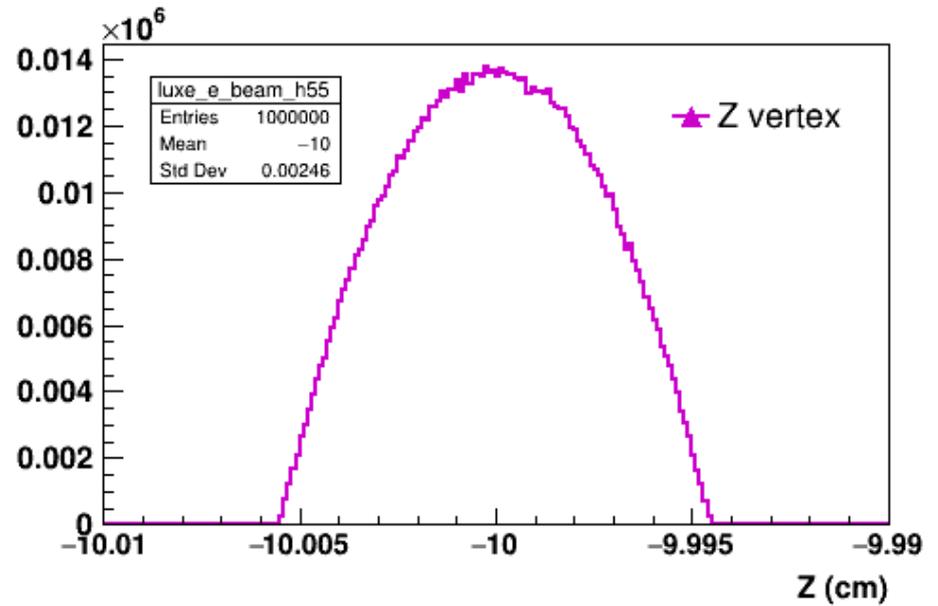
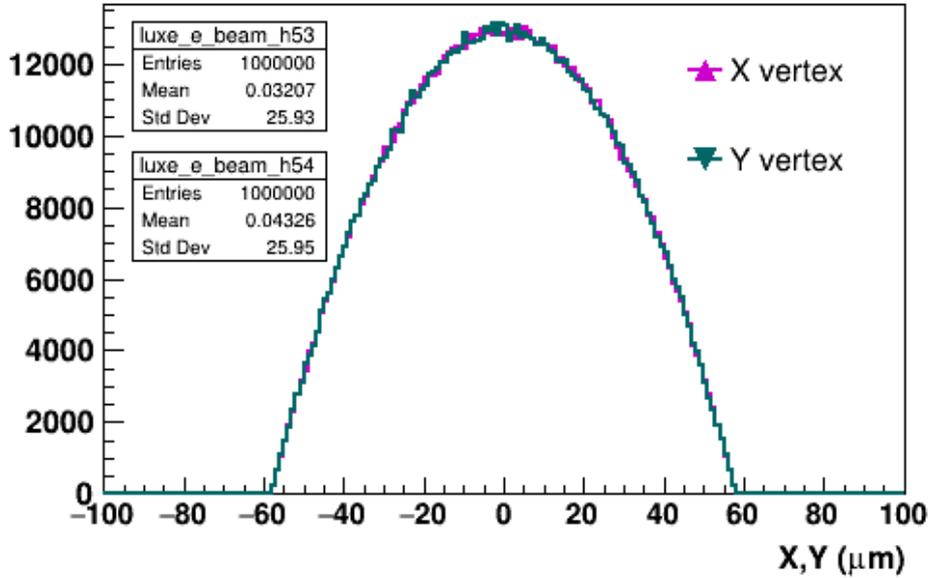


Bremsstrahlung simulation with Geant4

Oleksandr Borysov

LUXE Meeting
November 5, 2018

Particle Source



Geant4 General Particle Source:

```

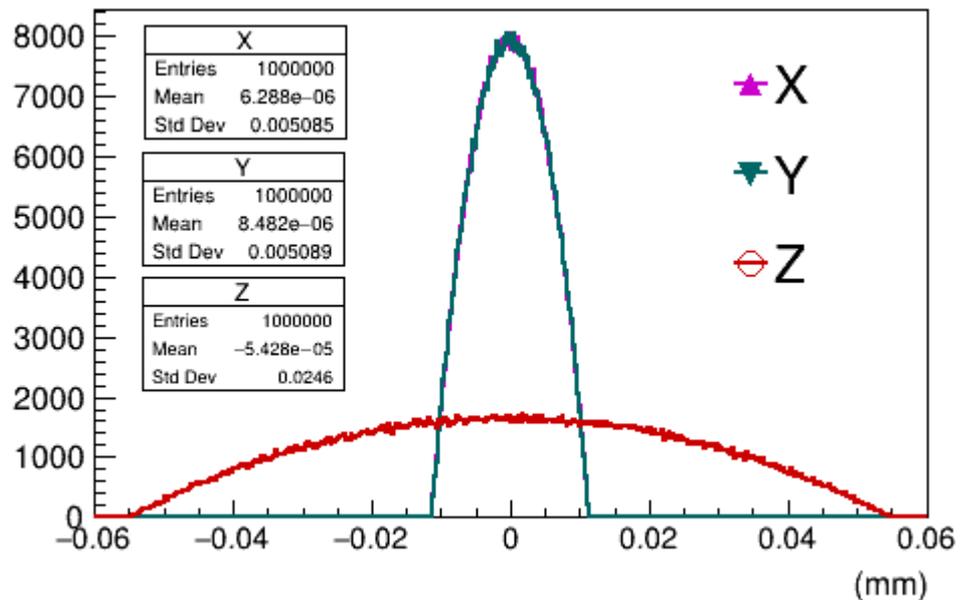
/gps/particle      e-
/gps/energy        17.5 GeV
/gps/position      (0.0, 0.0, -10.0) cm
/gps/number        1
/gps/pos/type      Volume
/gps/pos/shape     Ellipsoid
/gps/pos/centre   (0.0, 0.0, -10.0) cm
    
```

```

/gps/pos/halfx    58.0 um
/gps/pos/halfy    58.0 um
/gps/pos/halfz    55.0 um
    
```

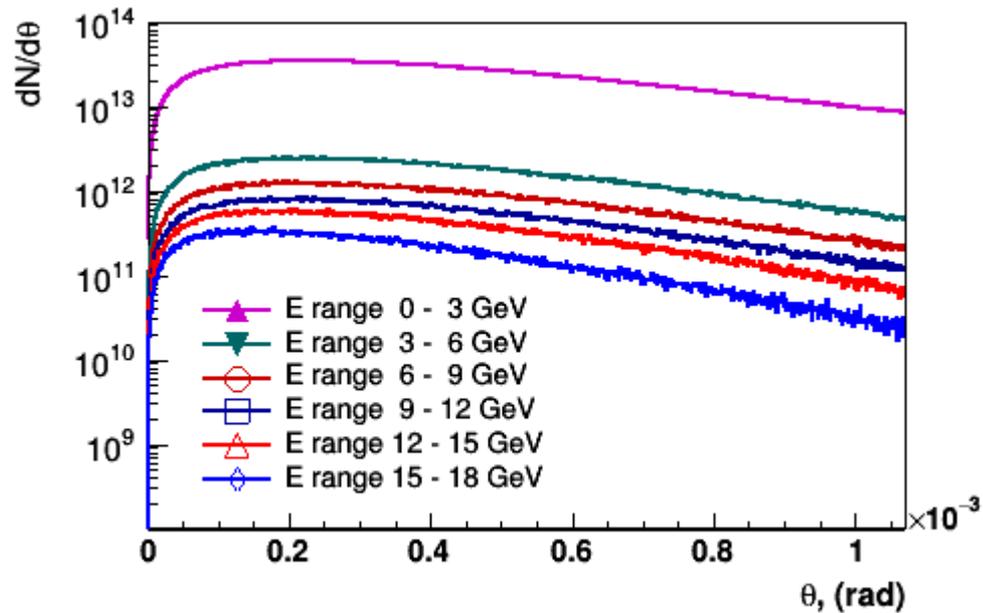
```

/gps/ang/type     focused
/gps/ang/focuspoint (0.0, 0.0, 500.0) cm
    
```

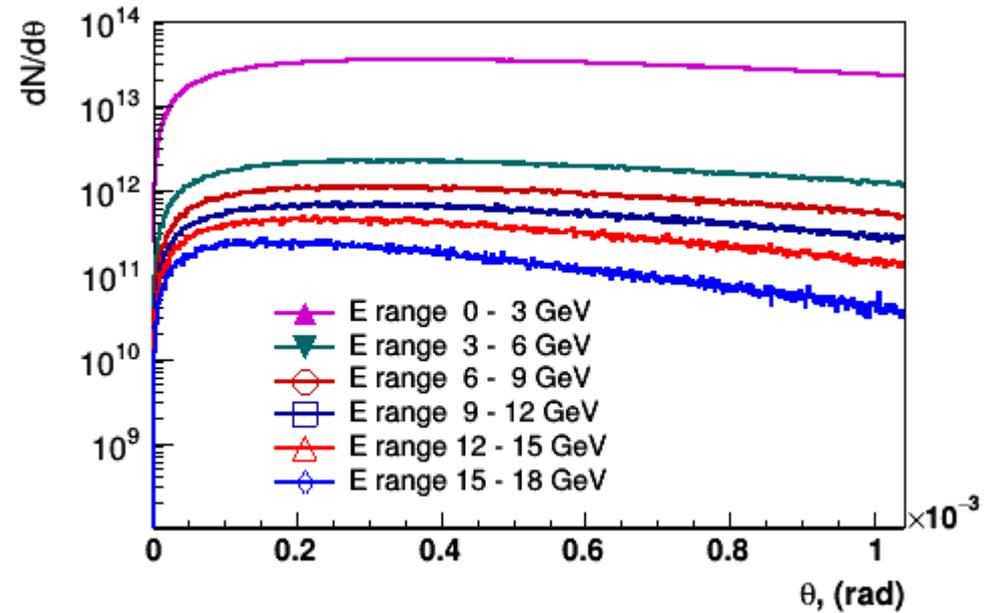


Bremsstrahlung on Al target

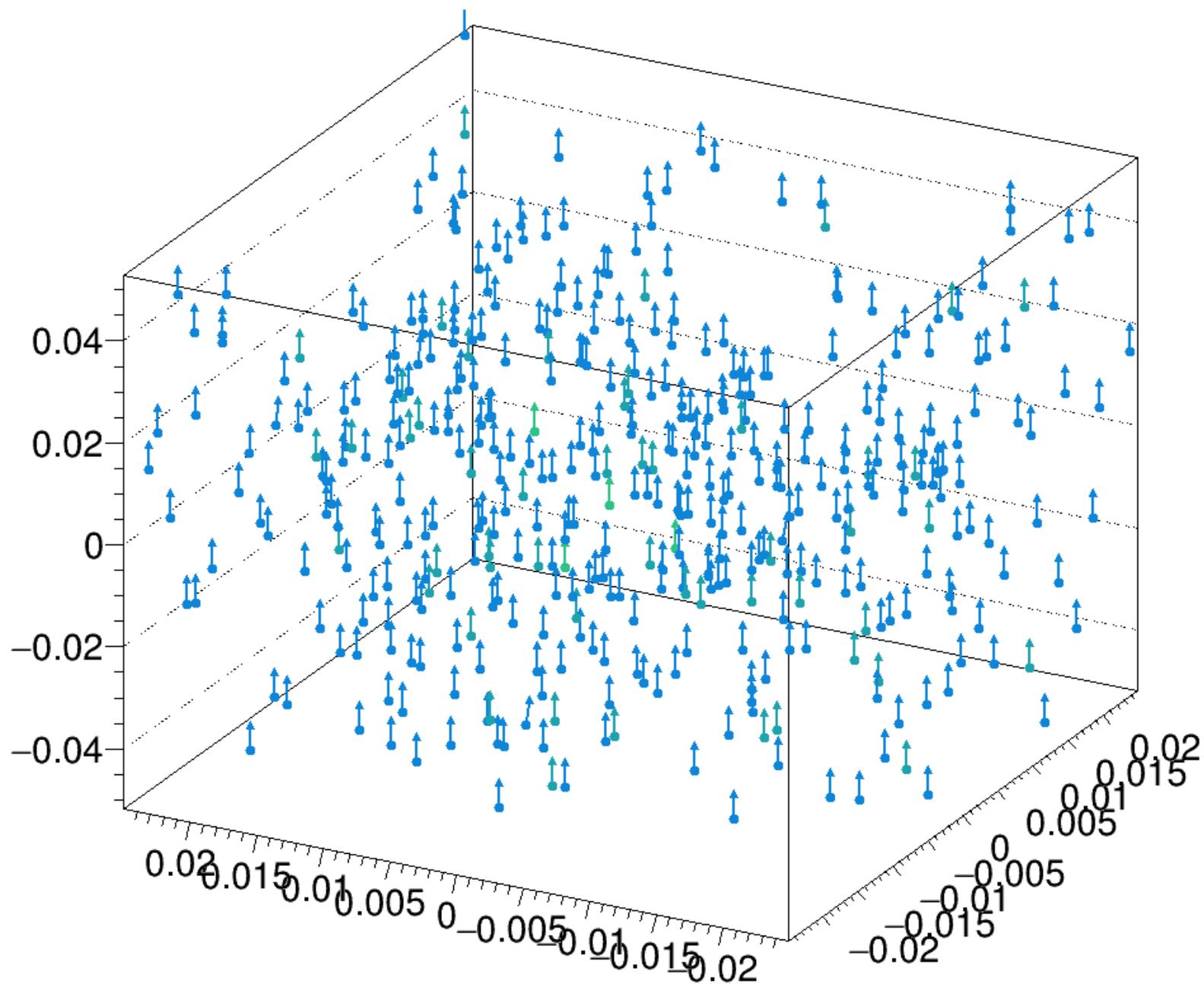
Al, 0.5X0 (44.5 mm)



Al, 1X0 (89 mm)



Photons at IP



Geant4 bremsstrahlung model

Seltzer-Berger bremsstrahlung model

In order to improve accuracy of the model described above a new model *G4SeltzerBergerModel* have been design which implementing cross section based on interpolation of published tables [SB85][SB86]. Single-differential cross section can be written as a sum of a contribution of bremsstrahlung produced in the field of the screened atomic nucleus $d\sigma_n/dk$, and the part $Z d\sigma_e/dk$ corresponding to bremsstrahlung produced in the field of the Z atomic electrons,

$$\frac{d\sigma}{dk} = \frac{d\sigma_n}{dk} + Z \frac{d\sigma_e}{dk}.$$

.... The angular dependence is

contained in the variable $u = E\theta m^{-1}$. For a given value of u the dependence of the shape of the function on Z , E and $\epsilon = k/E$ is very weak. Thus, the distribution can be approximated by a function

$$f(u) = C (ue^{-au} + due^{-3au})$$

where

$$C = \frac{9a^2}{9+d} \quad a = 0.625 \quad d = 27$$

Angular distribution of Bremsstrahlung photons and of positrons for calculations of terrestrial gamma-ray flashes and positron beams

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^a CWI, P.O. Box 94079, 1090GB Amsterdam, The Netherlands

^b Eindhoven University of Technology, P.O. Box 513, 5600MB Eindhoven, The Netherlands

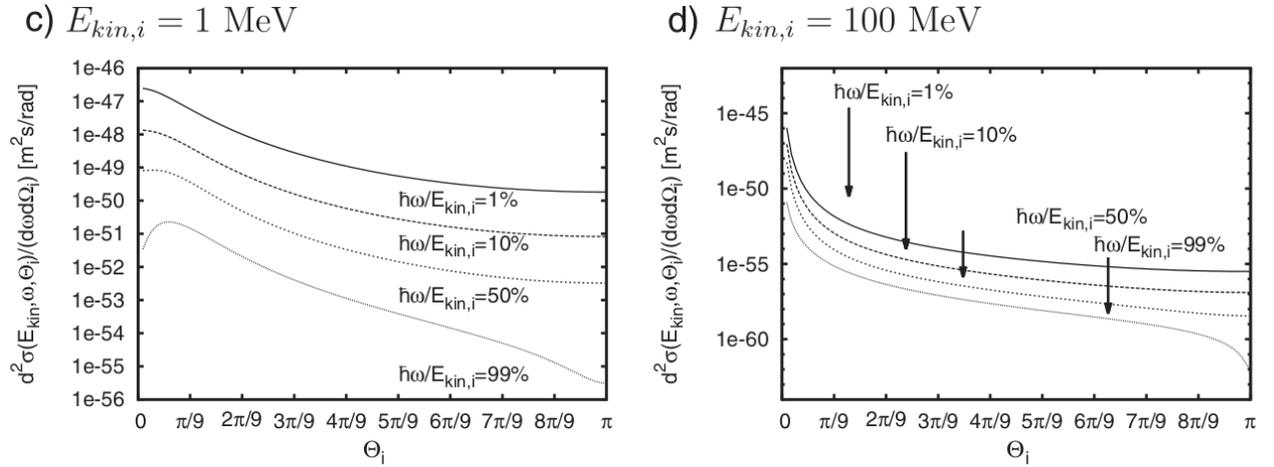


Fig. 3. The doubly differential cross section $d^2\sigma/(d\omega d\Omega_i)(E_{kin,i},\omega,\Theta_i)$ for Bremsstrahlung ($Z = 7$) versus the scattering angle Θ_i between emitted photon and incident electron. The electron energies are a) $E_{kin,i} = 10 \text{ keV}$, b) $E_{kin,i} = 150 \text{ keV}$, c) $E_{kin,i} = 1 \text{ MeV}$ and d) $E_{kin,i} = 100 \text{ MeV}$. In each plot the photon energy $\hbar\omega$ amounts to 1%, 10%, 50% and 95% of the kinetic energy of the incident electron.

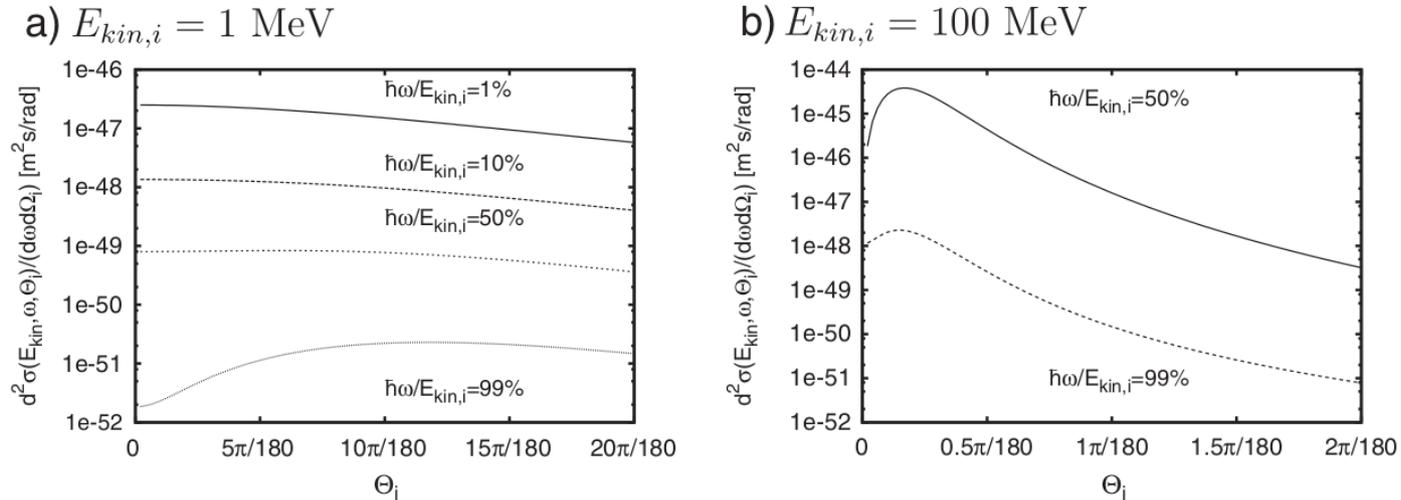
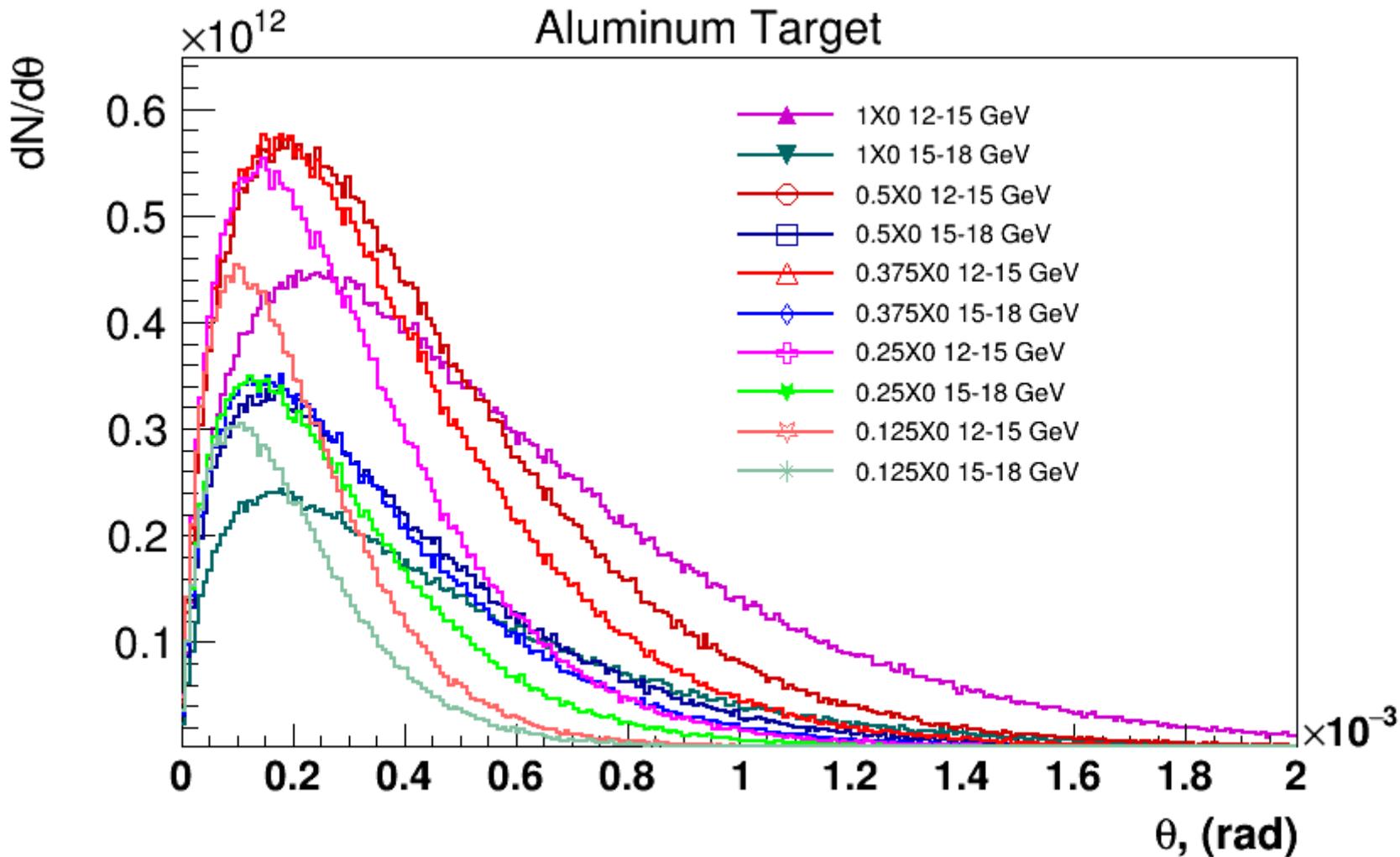


Fig. 4. The doubly differential cross section $d^2\sigma/(d\omega d\Omega_i)$ for Bremsstrahlung as in Fig. 3 for a smaller angular range.

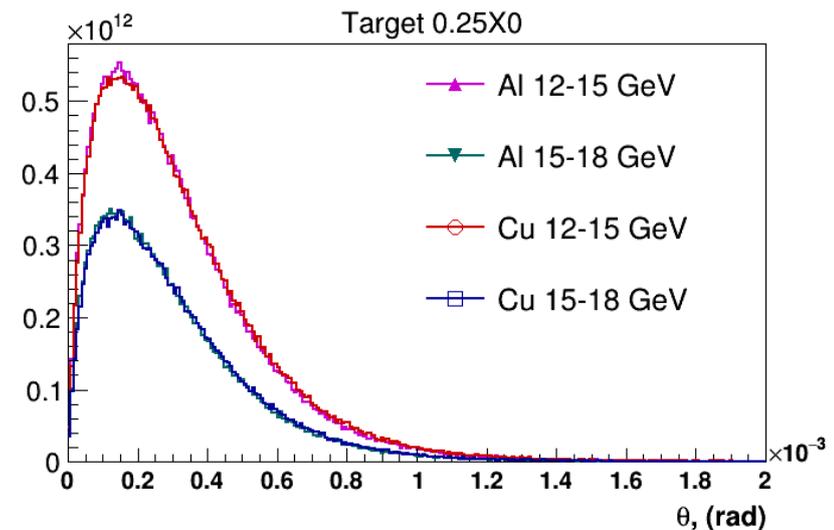
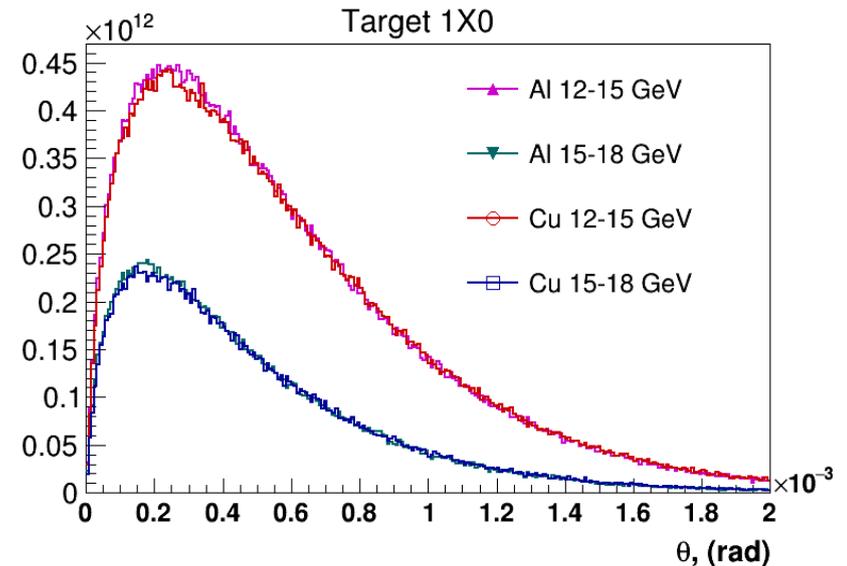
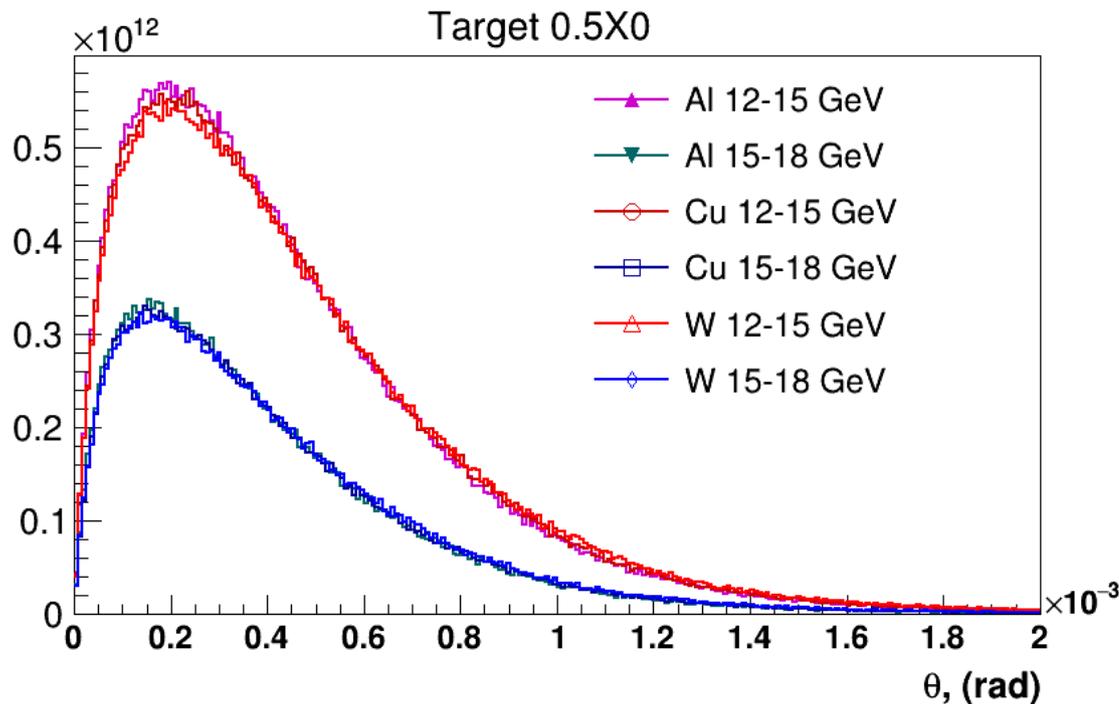
Target thickness

- 17.5 GeV e- beam, generated from single point along Z;
- Detector is 4 m away from the target.



Different target materials

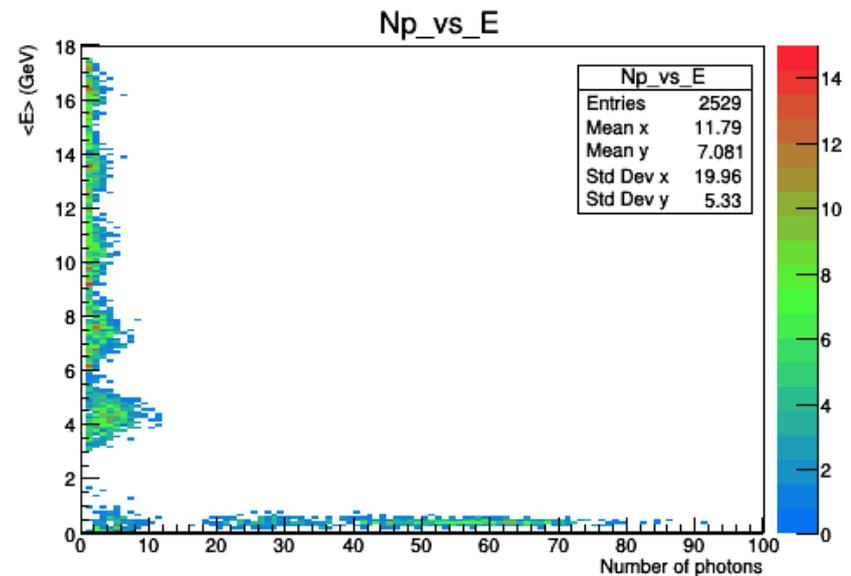
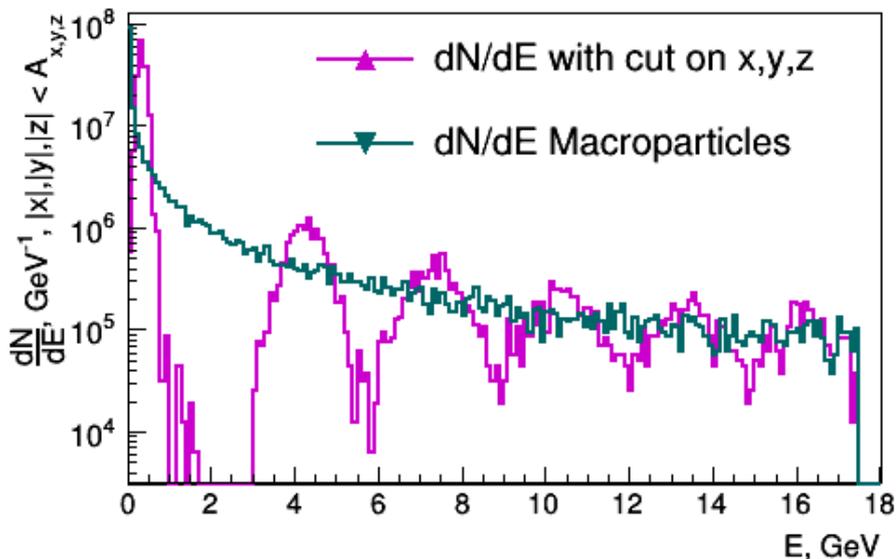
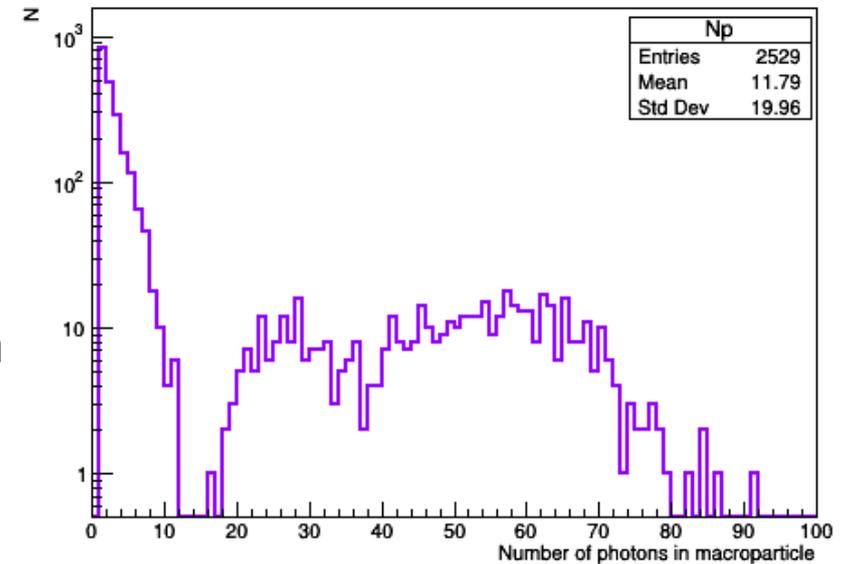
- 17.5 GeV e- beam, generated from single point along Z;
- Detector is 4 m away from the target.



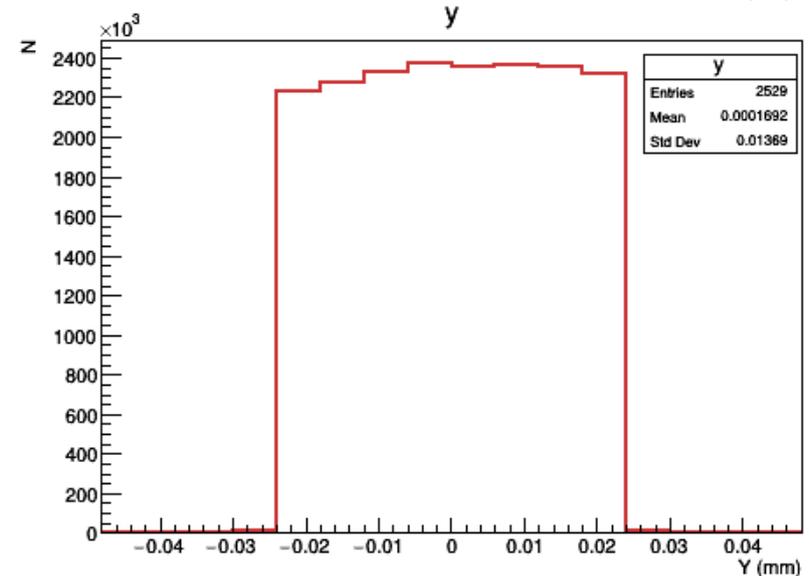
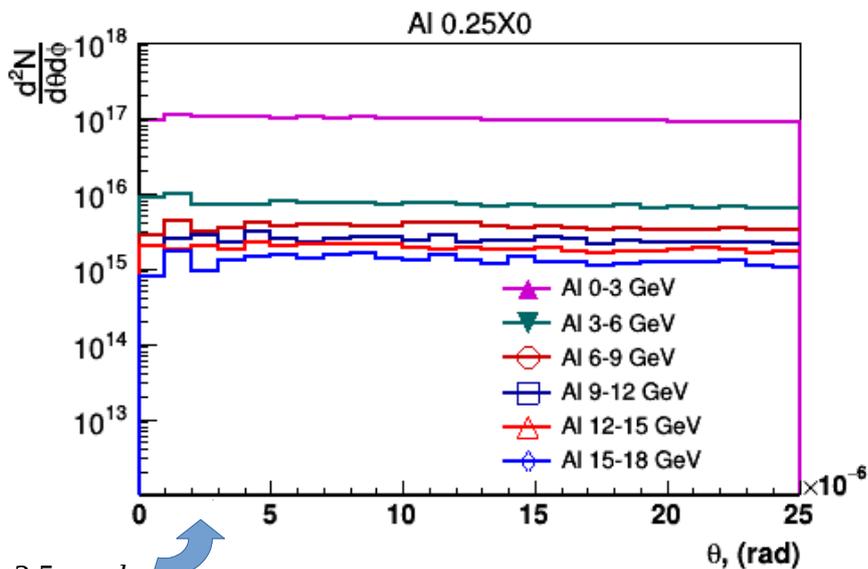
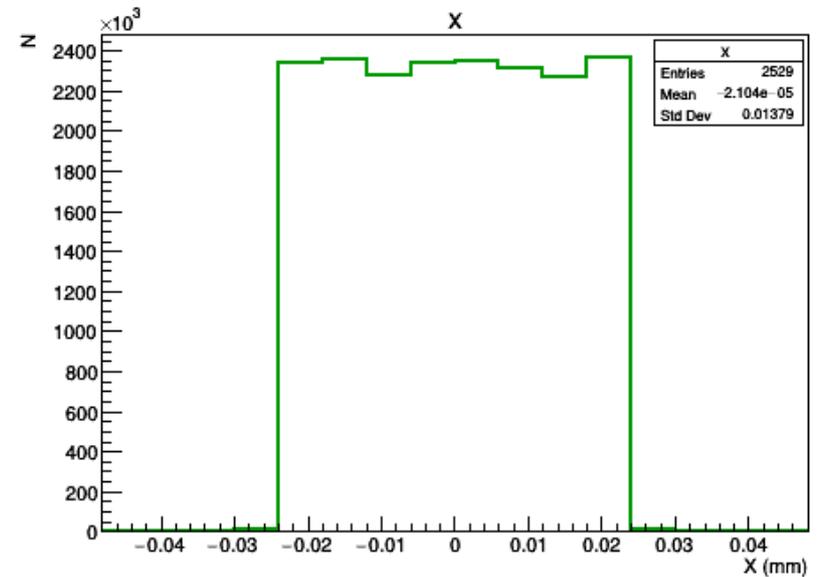
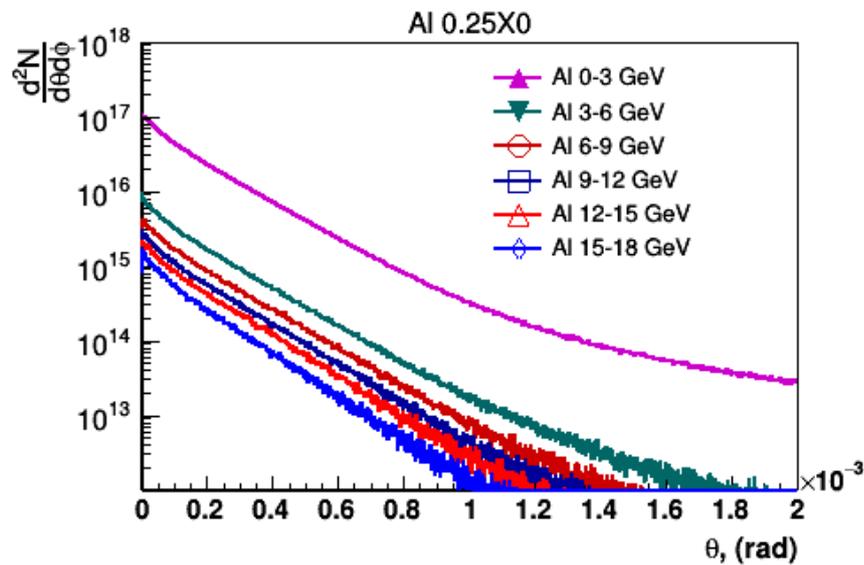
Macroparticles

E (GeV) x (um) y (um) z(um) beta_x beta_y beta_z PDG_NUM MP_Wgt
#-----

- Beam as on slide 2;
- Target: Al, 0.25X0, (22 mm).
 - 4D space (x, y, z, E) divided into cells;
 - Each cell is converted to macroparticle with $X = \langle x \rangle, \dots E = \langle E \rangle$;
 - Number of photons in the cell defines the weight.



Photon θ and Macroparticle position

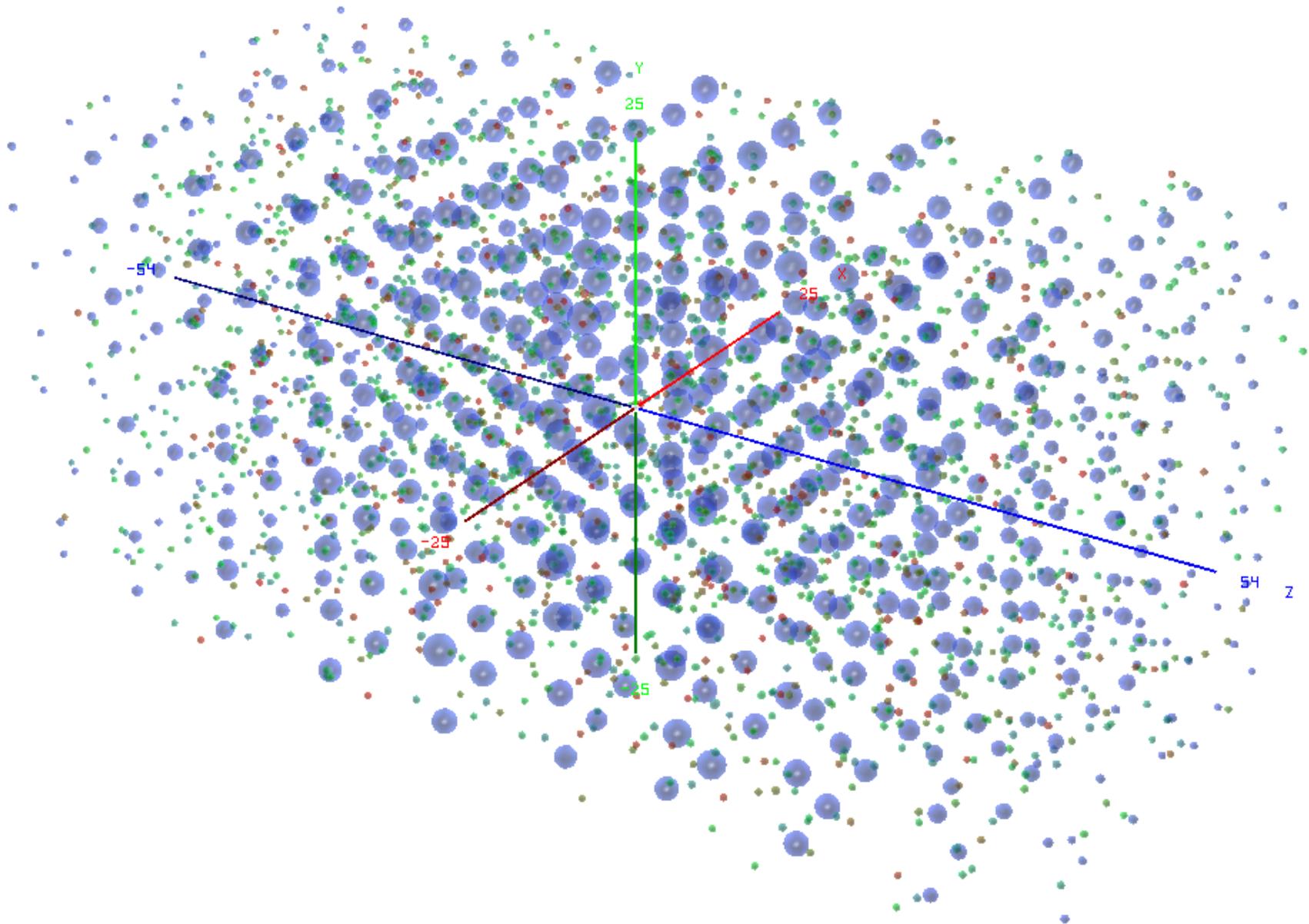


$$\frac{10 \mu m}{4 m} = 2.5 \mu rad$$

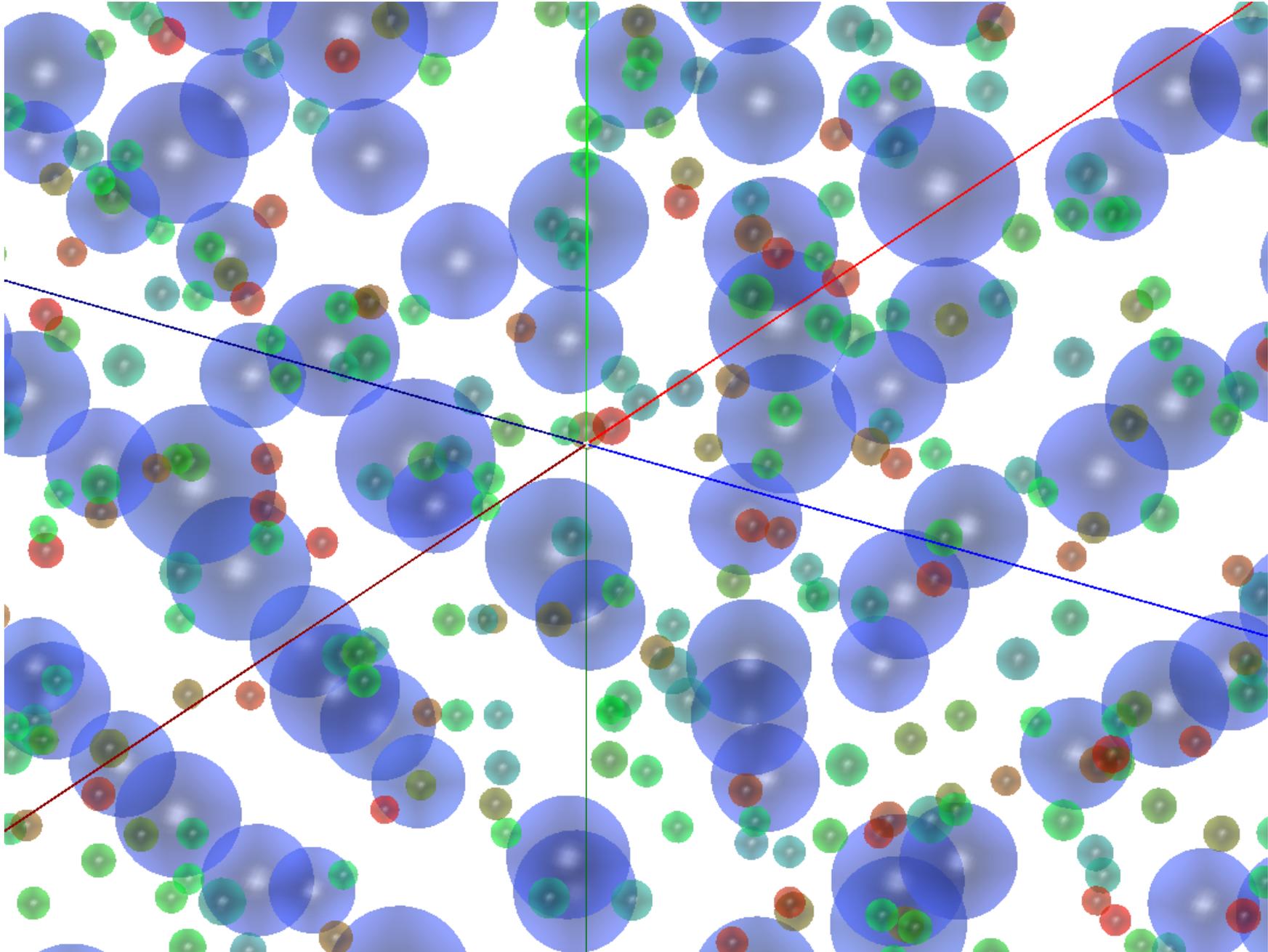


Macro particles Al, 0.25X0.

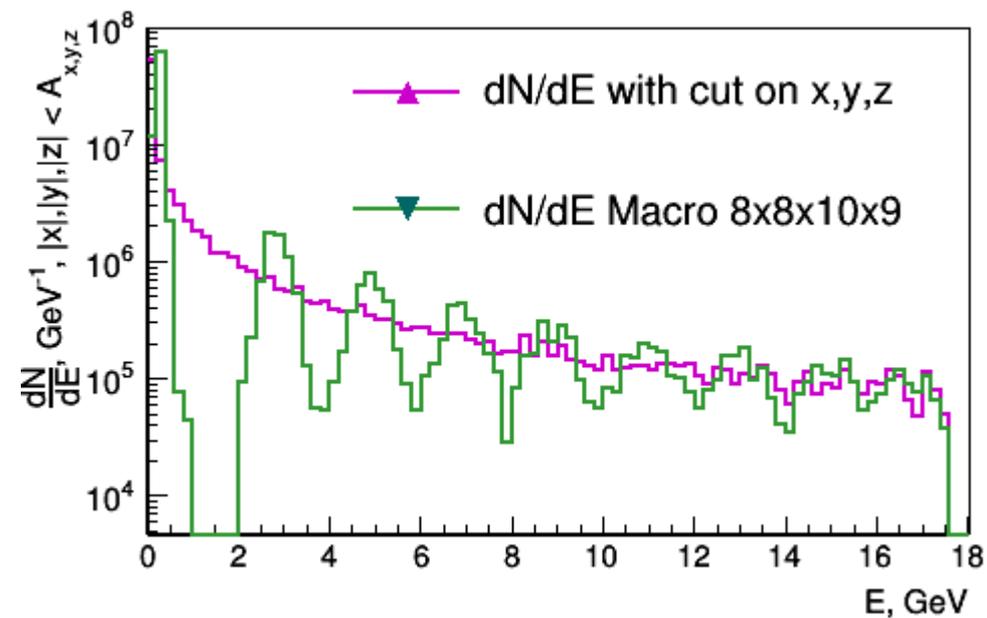
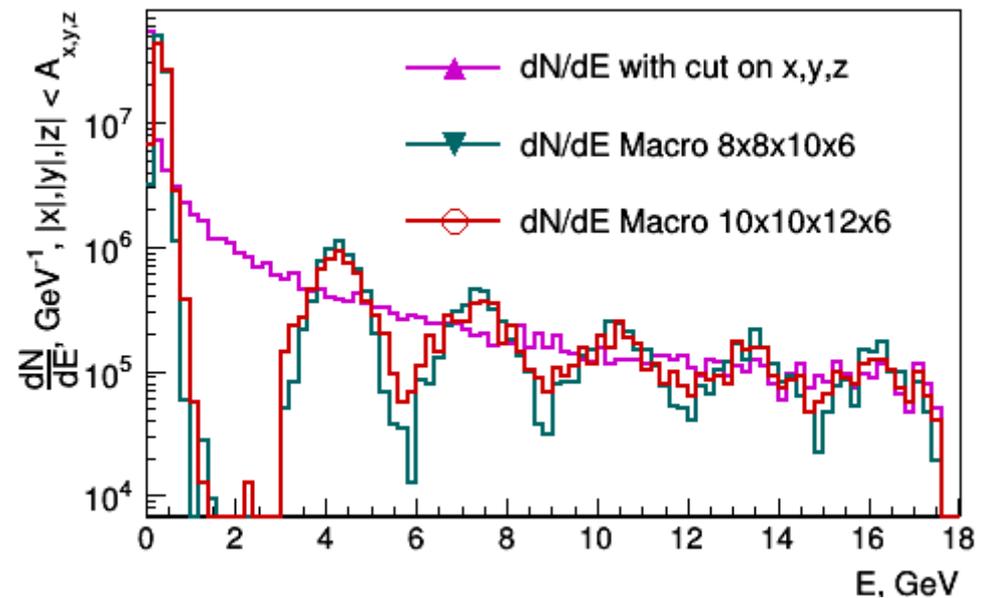
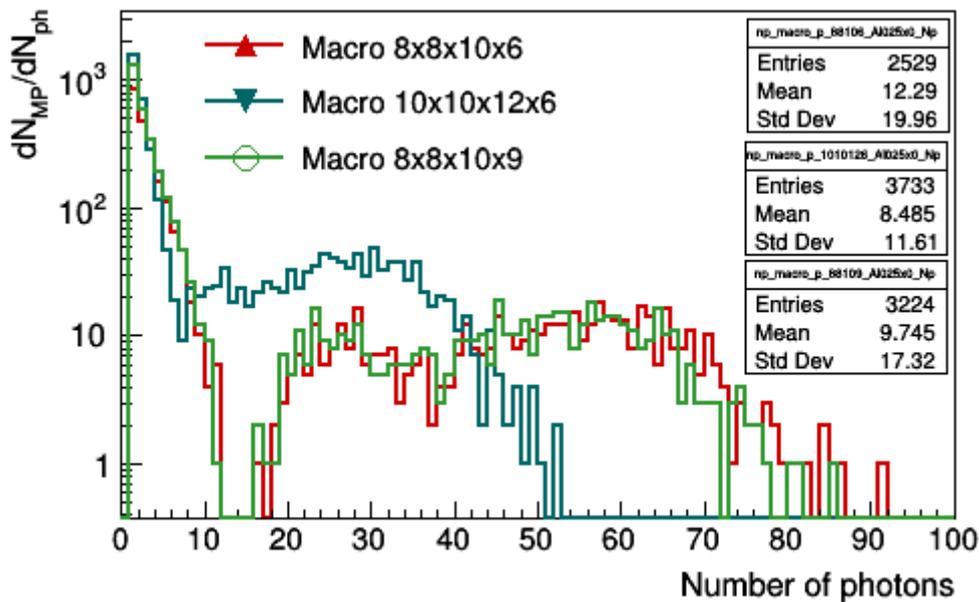
x/y/z/E: 8/8/10/6



Macro particles Al, 0.25X0.
x/y/z/E: 8/8/10/6



Macroparticles of different size



Macroparticles: Al 0.25X0 and 0.5X0

