

Isometric Tensor Network States

Frank Pollmann

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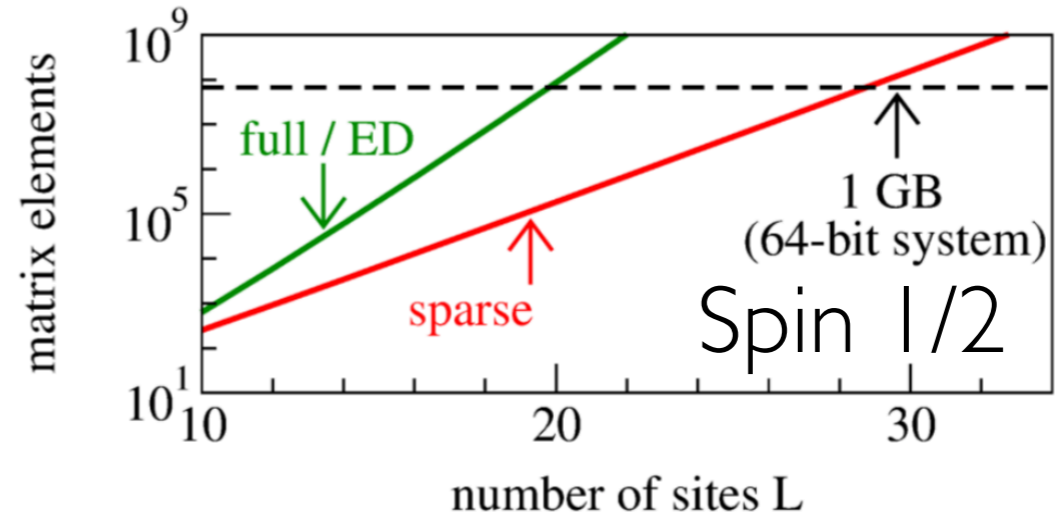
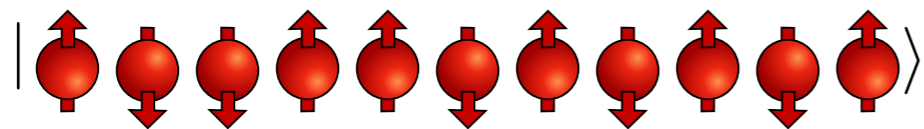


Zeuthen, Mar 4 2019

Complexity of a quantum many-body problem

Diagonalize a Hamiltonian in the full many-body Hilbert space

$$|\psi\rangle = \sum_{j_1, j_2, \dots, j_L} \psi_{j_1, j_2, \dots, j_L} |j_1\rangle |j_2\rangle \dots |j_L\rangle, \quad j_n = 1 \dots d$$



- ➔ Full diagonalization up to ~ 20 sites
- ➔ Sparse methods up to ~ 30 sites

Outlook

Efficient representation of quantum many-body states

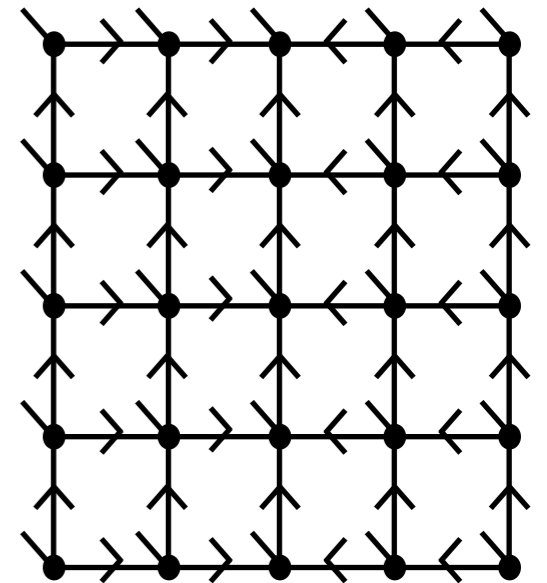
- ▶ Brief review of **Matrix-Product States**



- ▶ **Isometric Tensor Network States in 2D:**

Tensor-network state ansatz that allows for efficient contractions

[Zaletel and FP; arXiv:1902.05100]



Entanglement

Generic quantum state has a d^L dimensional Hilbert space

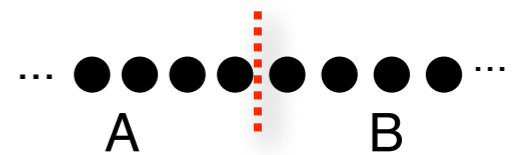
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Decompose a pure state into a superposition of product states (**Schmidt decomposition**)



$$|\psi\rangle = \sum_{i,j} C_{i,j} |i\rangle_A \otimes |j\rangle_B = \sum_{\alpha} \Lambda_{\alpha} |\alpha\rangle_A \otimes |\alpha\rangle_B$$

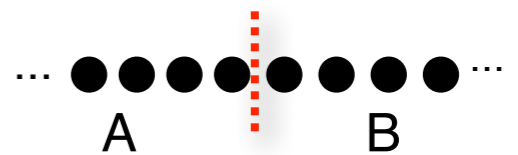
with $\langle \alpha | \alpha' \rangle = \delta_{\alpha \alpha'}$

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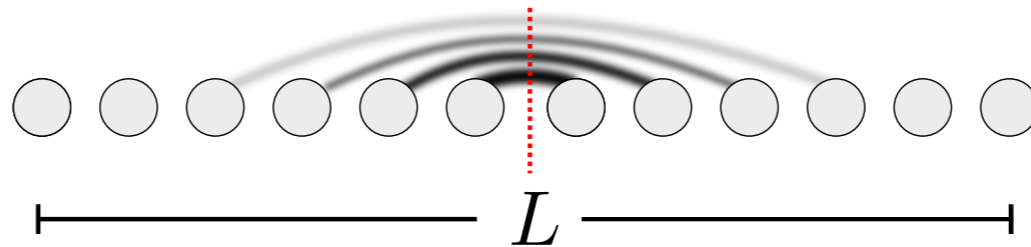
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Entanglement entropy as a measure for the amount of entanglement $S = - \sum_{\alpha} \Lambda_{\alpha}^2 \log \Lambda_{\alpha}^2$

Entanglement

Area law for ground states of local (gapped) Hamiltonians
in 1D systems

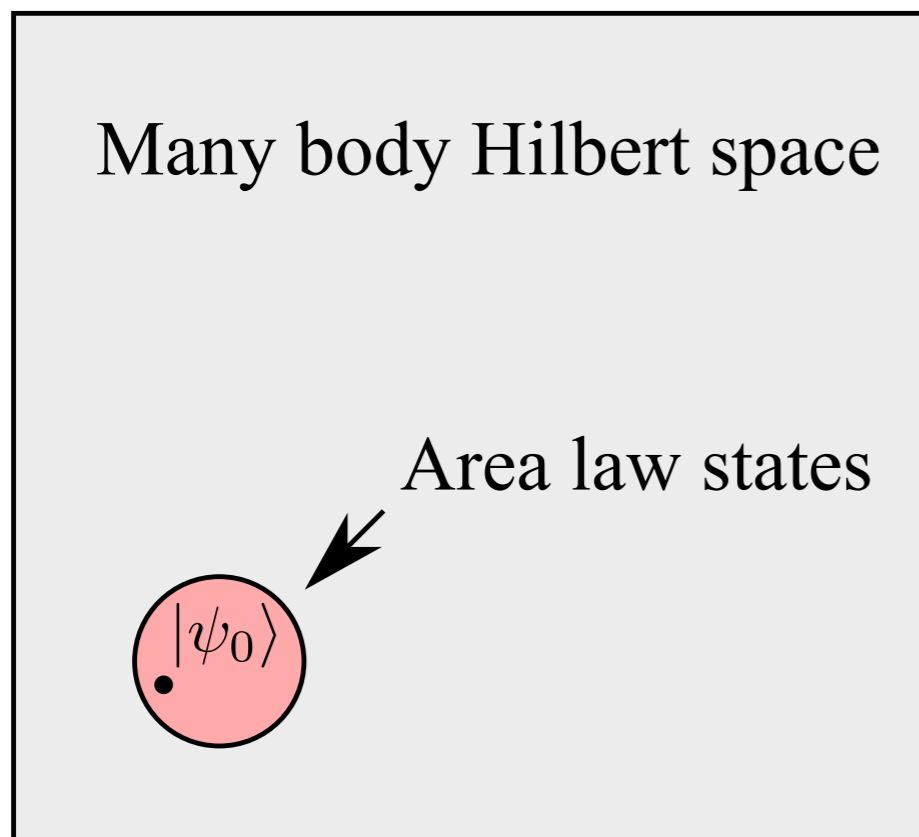
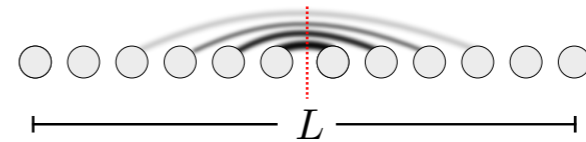
$$S(L) = \text{const.} \quad [\text{Srednicki '93, Hastings '07}]$$



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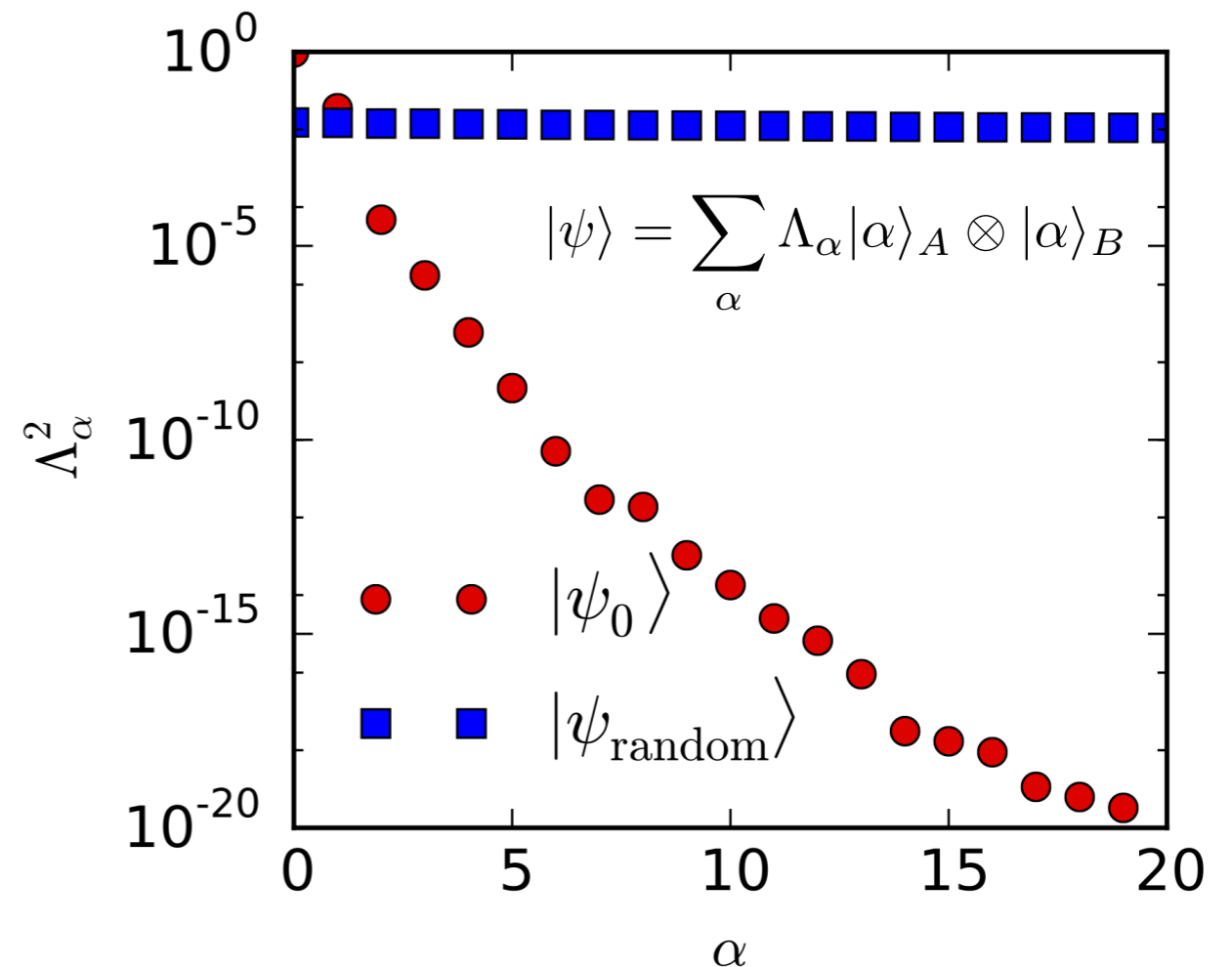
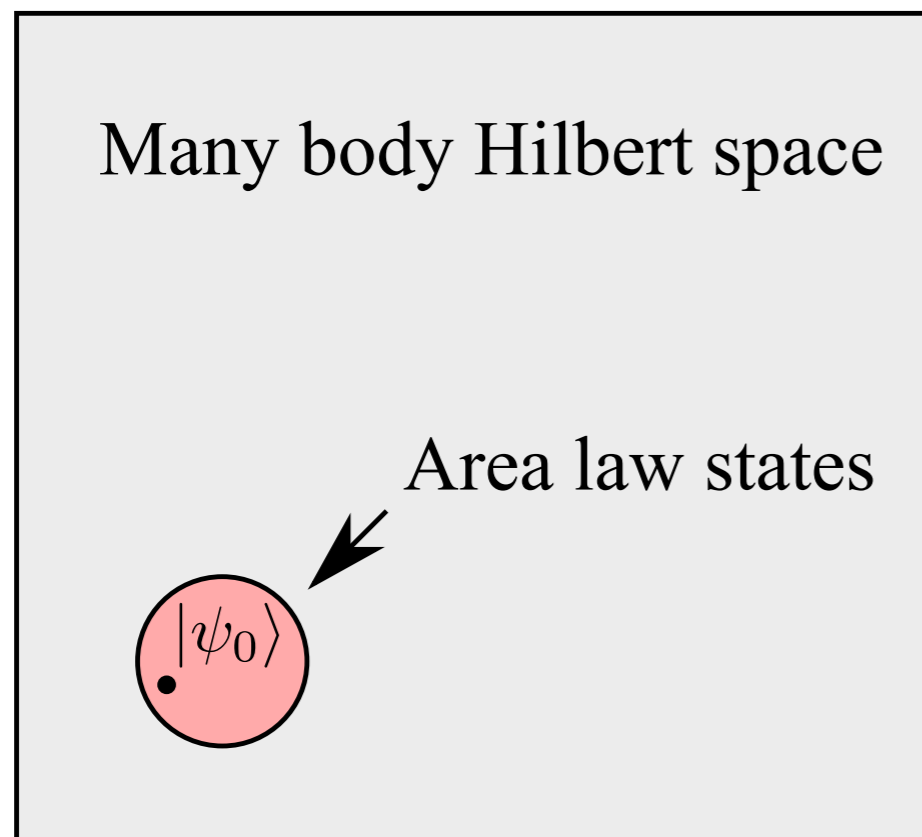
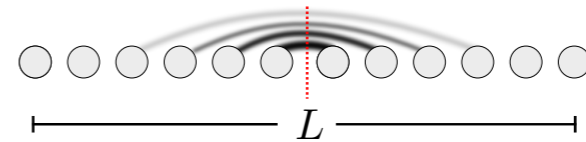


All ground states live in a tiny corner of the Hilbert space!

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Matrix-Product States

Matrix-product states (MPS): Reduction of the number of variables: $d^L \rightarrow Ld\chi^2$ [M. Fannes et al. 92]

$$\psi_{j_1, j_2, j_3, j_4, j_5} = \begin{array}{c} M^{[1]} \quad M^{[2]} \quad M^{[3]} \quad M^{[4]} \quad M^{[5]} \\ \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet \\ | \quad | \quad | \quad | \quad | \end{array}$$

$$M_{\alpha, \beta}^j = \begin{array}{c} M \\ \alpha \text{---} \bullet \text{---} \beta \\ | \\ j \end{array}$$

$\alpha, \beta = 1 \dots \chi$
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Canonical form: Use the gauge degree of freedom ($A^j = XM^jX^{-1}$) to find a convenient representation

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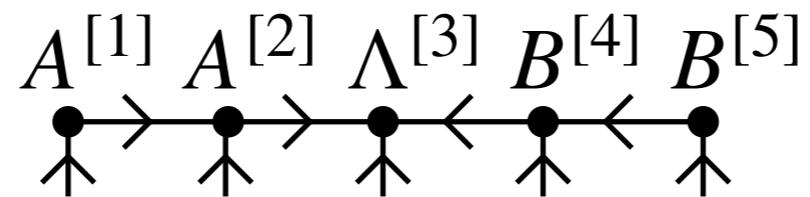
(Isometries)

Center matrix Λ represents wave function

$$|\psi\rangle = \sum_{\alpha, \beta, j} \Lambda_{\alpha, \beta}^j |\alpha\rangle |j\rangle |\beta\rangle \quad (\text{orthogonal states } |j\rangle, |\alpha\rangle, |\beta\rangle)$$

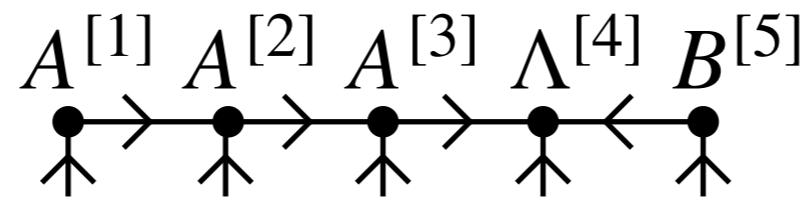
Density Matrix Renormalization Group

Moving the center matrix: $\Lambda^\ell B^{[\ell+1]} = A^{[\ell]} \Lambda^{[\ell+1]}$ accomplished by an **orthogonal factorization** (e.g. QR or SVD)



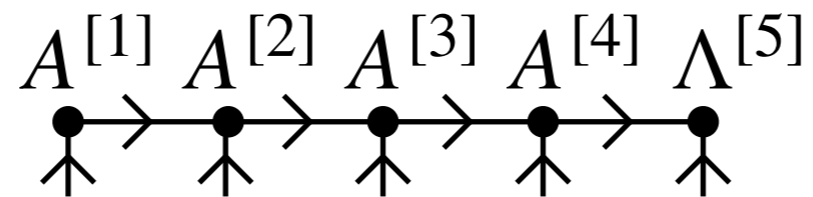
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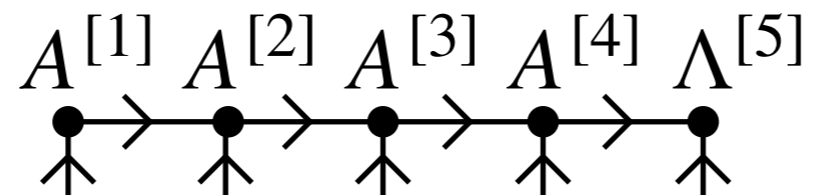
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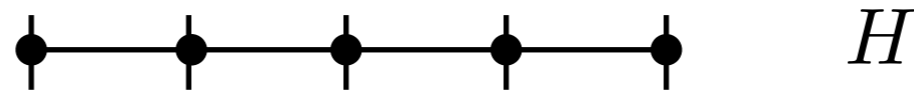


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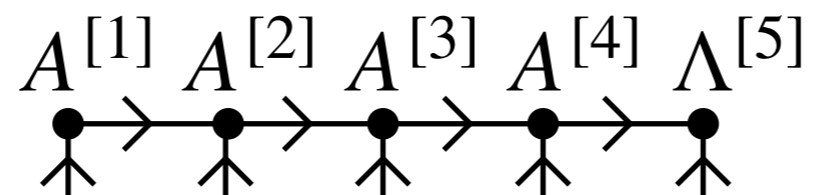


Find the **ground state** iteratively

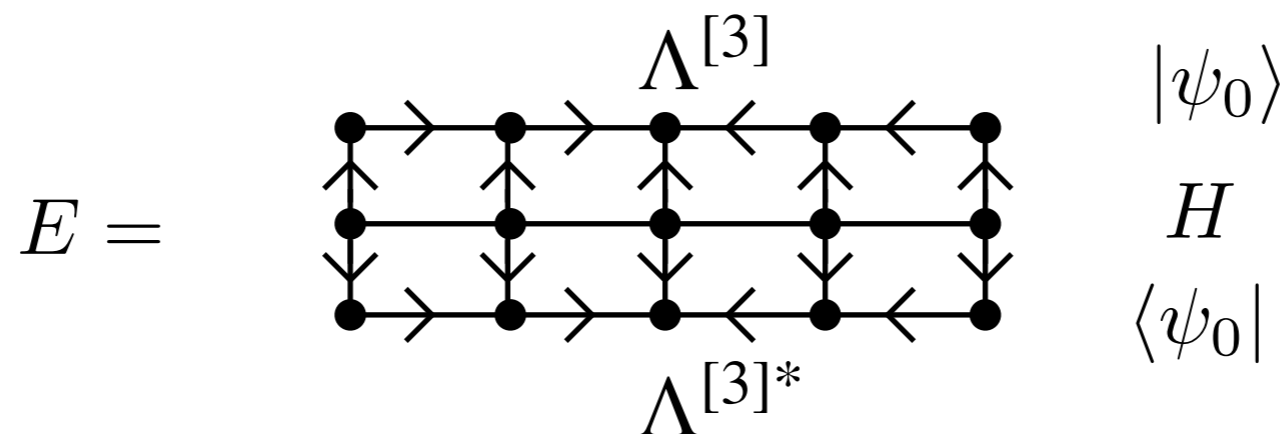


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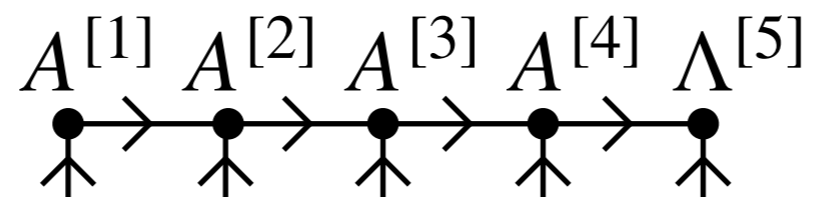


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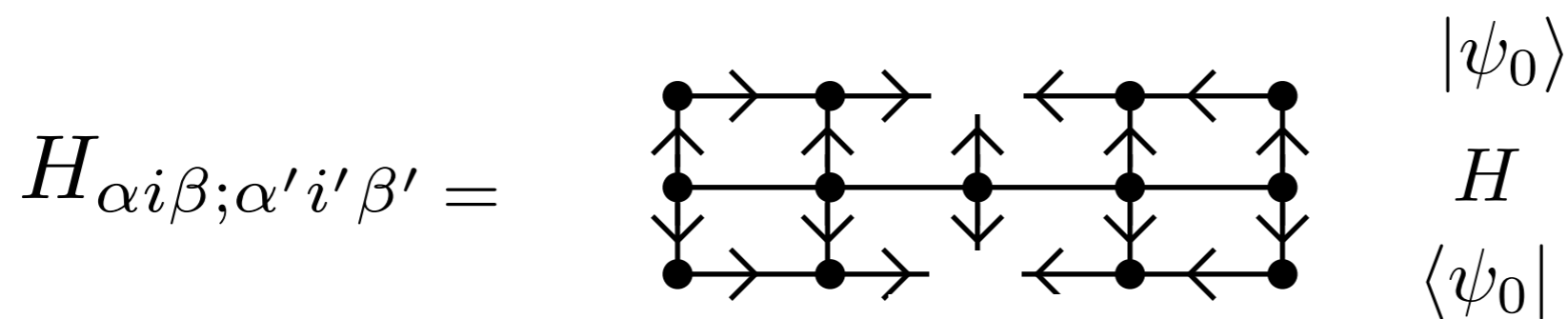


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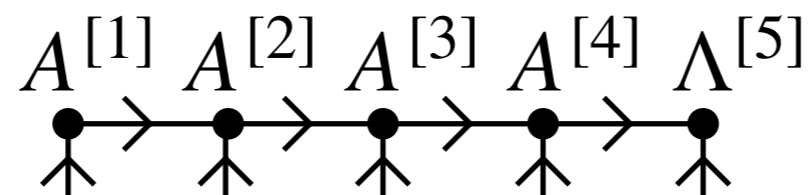


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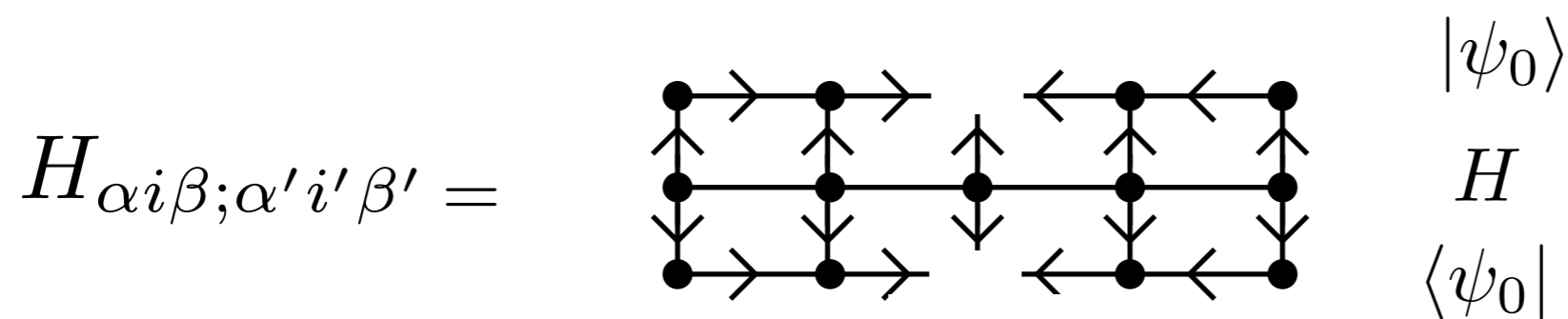


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Find the **ground state** iteratively



Locally minimize the energy of $H_{\alpha i \beta; \alpha' i' \beta'}$ (e.g., Lanczos)

Density matrix renormalization group (DMRG) [White '92, Schollwoeck '11]

Tensor Network States in 2D

MPS capture 1D area law \rightarrow Exponential scaling in 2D

$$\psi_{j_1, j_2, j_3, j_4, j_5} \approx \begin{array}{c} M^{[1]} \quad M^{[2]} \quad M^{[3]} \quad M^{[4]} \quad M^{[5]} \\ \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet \text{---} \bullet \\ | \quad | \quad | \quad | \quad | \end{array}$$

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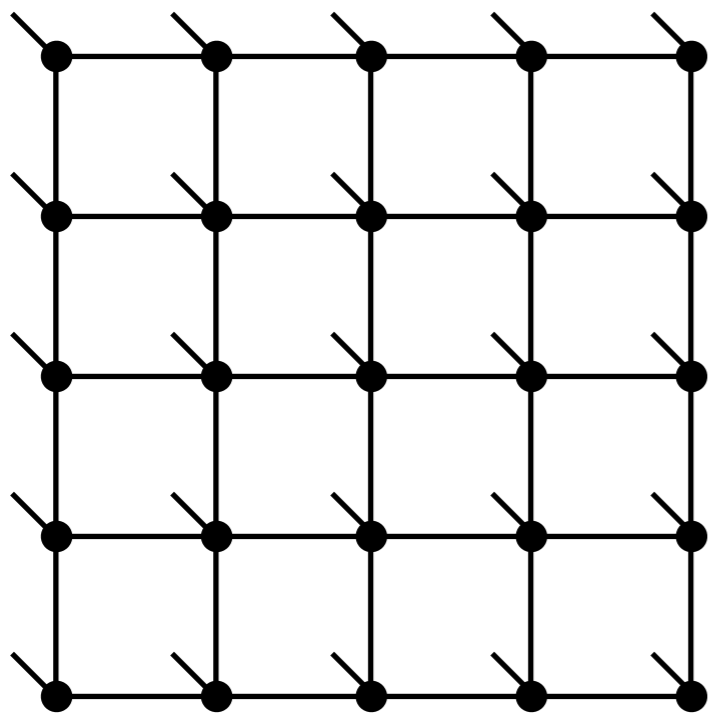
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How to generalize the MPS approach to 2D?



$$T_{\alpha, \beta, \gamma, \delta}^j = \begin{array}{c} \diagup \\ | \\ \bullet \\ | \\ \diagdown \end{array}$$

► **Tensor Network States (TNS)**

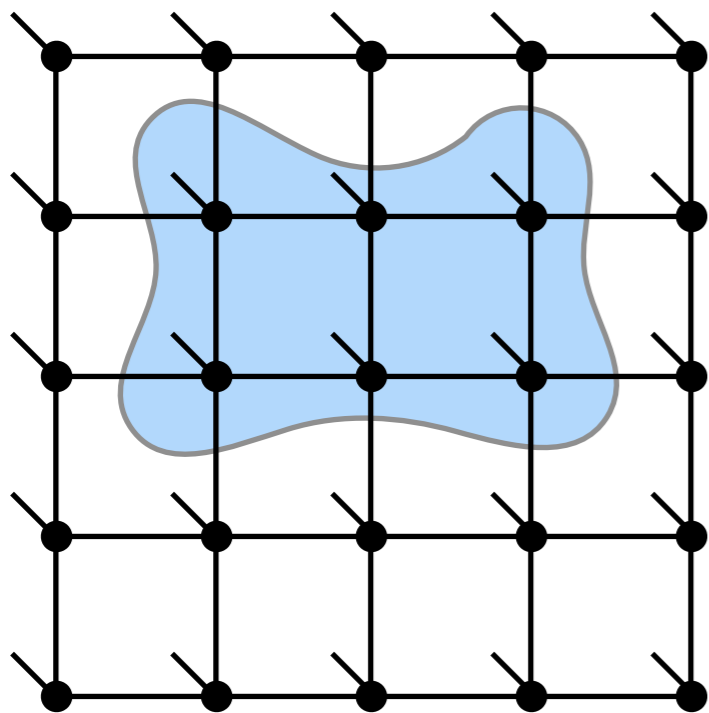
[Maeshima et al. '01, Verstraete and Cirac '04]

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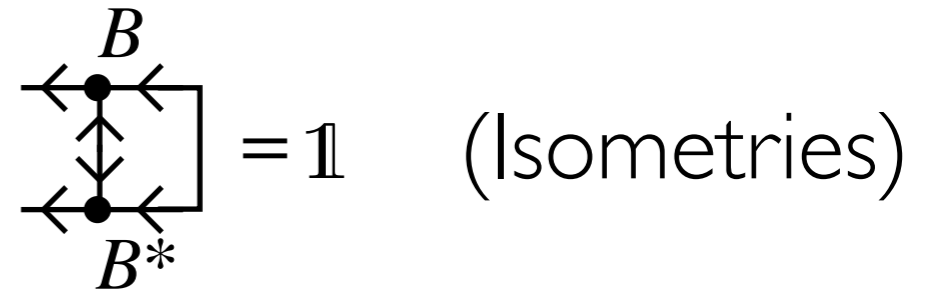
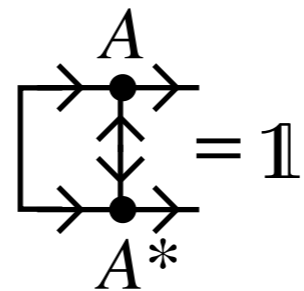
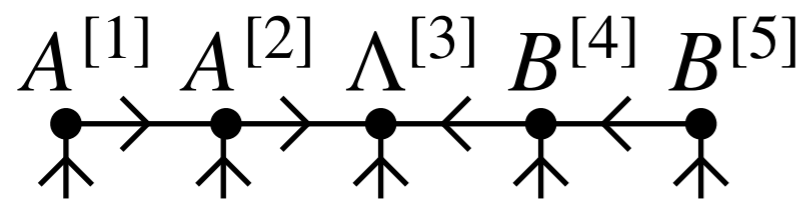
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▶ **Capture 2D area law** 😊

▶ Difficult to handle numerically:
Exact contraction of the 2D network
is still **exponentially hard** 😞

Isometric Tensor Network States in 2D

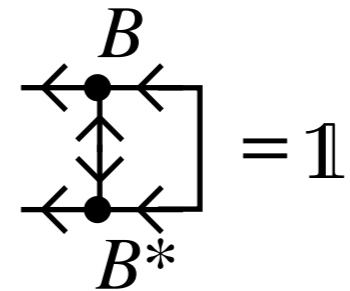
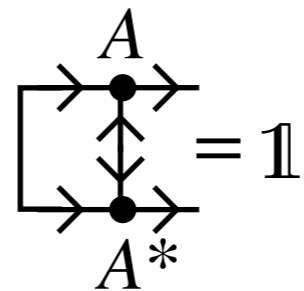
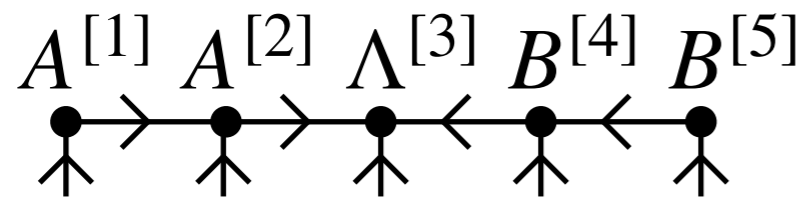
Recall: **Canonical form of 1D MPS**



(Isometries)

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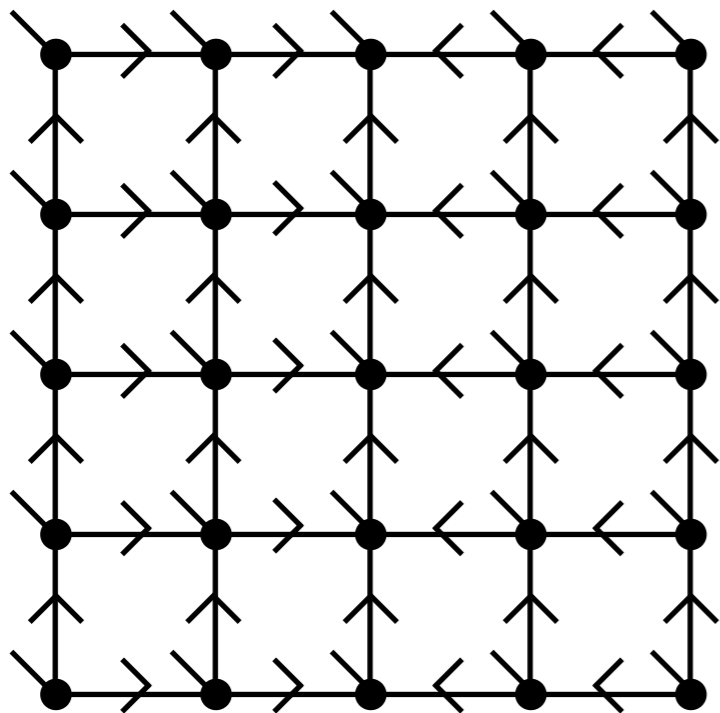
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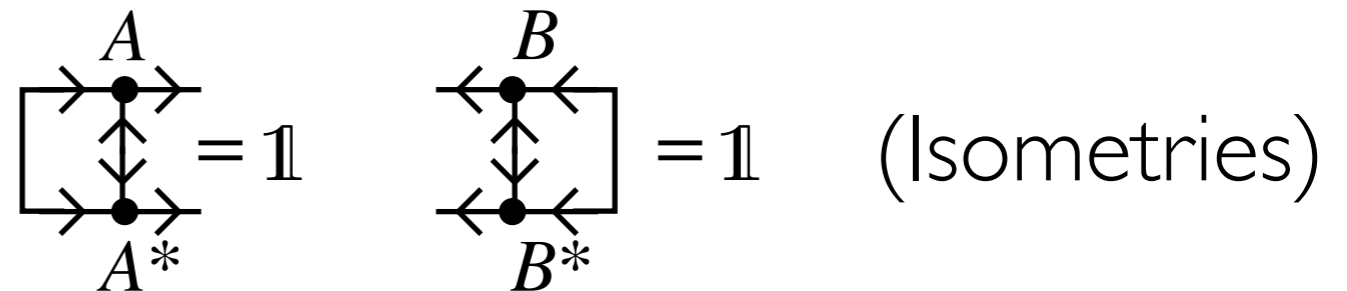
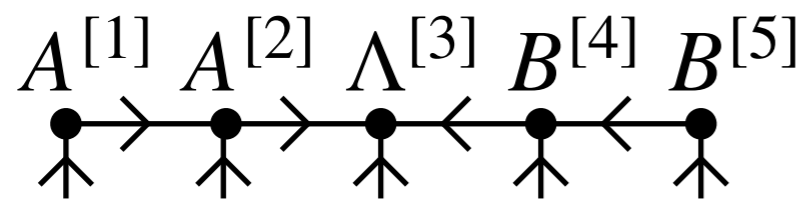
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$A^{[1]} A^{[2]} \Lambda^{[3]} B^{[4]} B^{[5]}$



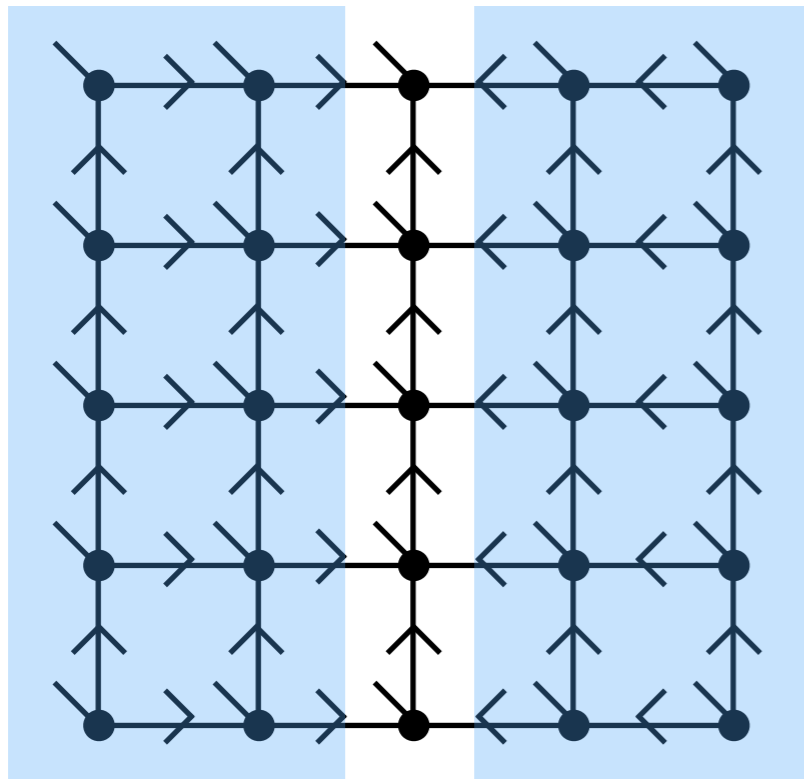
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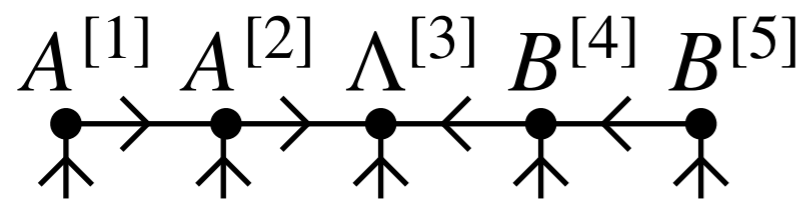
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- Isometric tensors are **efficiently contractable**

Isometric Tensor Network States in 2D

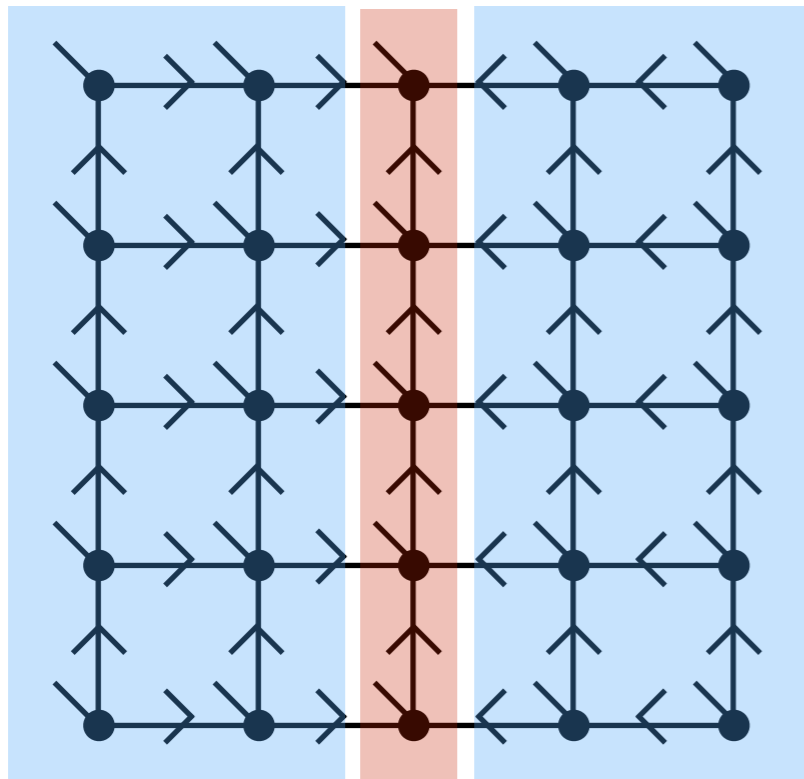
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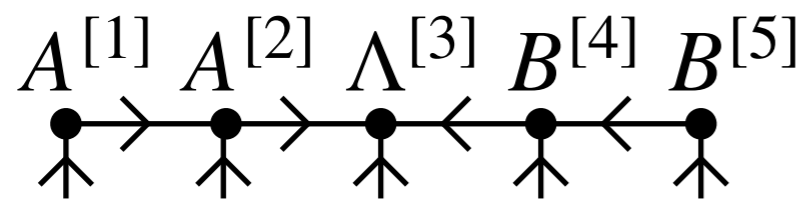
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- ▶ Isometric tensors are **efficiently contractable**
- ▶ Orthogonality center column is a **1D MPS**: Standard DMRG techniques

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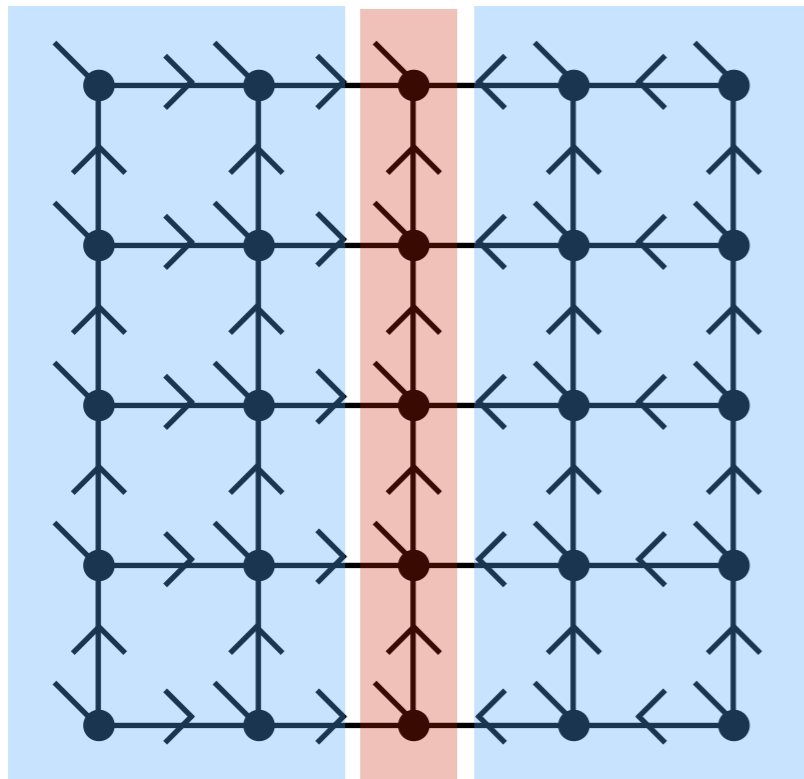
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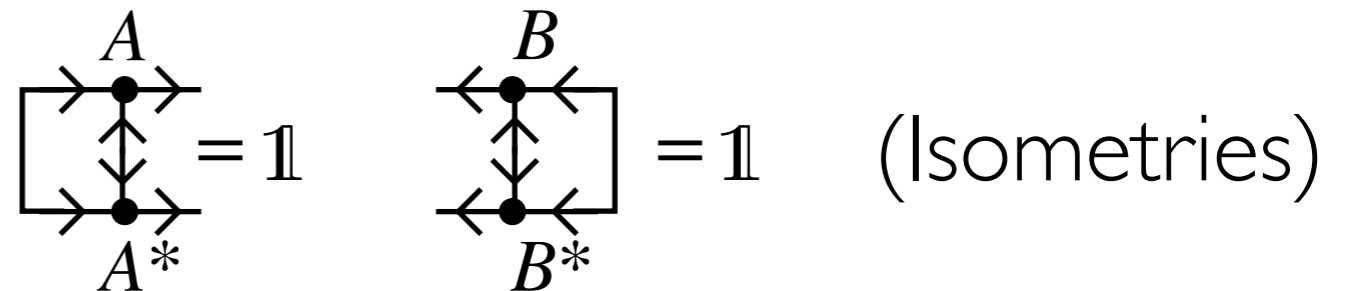
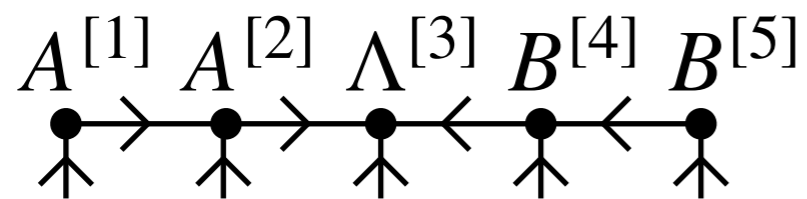
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- ▶ **Subset of TNS: Unclear what its variational power is!**

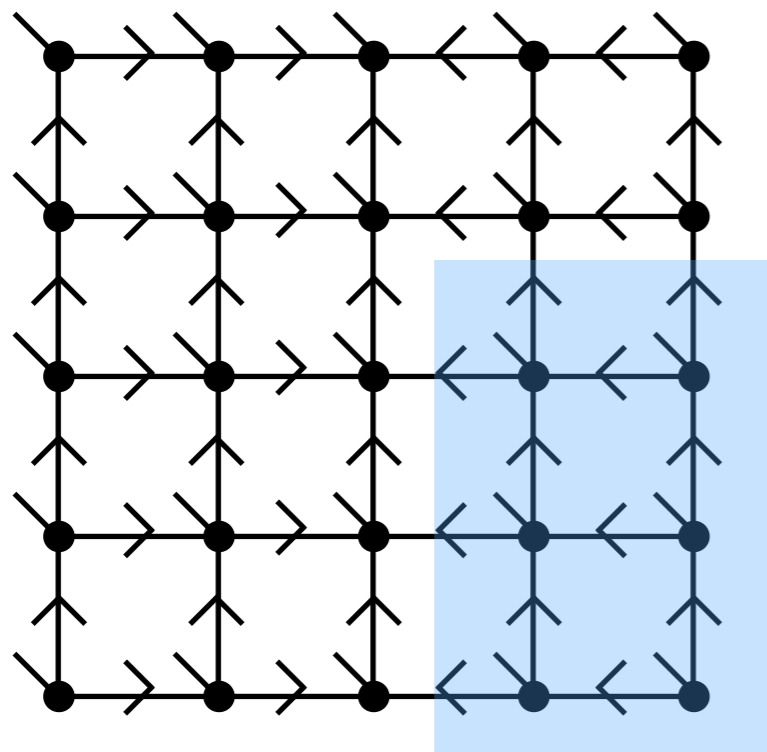
Isometric Tensor Network States in 2D

Recall: **Canonical form of 1D MPS**



Isometric TNS

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- ▶ **Subregions** with only outgoing arrows have isometric boundary maps
- ▶ **Causal structure**: time flows opposite to the direction of the arrows

Isometric Tensor Network States in 2D

How to shift the orthogonality center?

Recall: **ID MPS** $\Lambda^\ell B^{[\ell+1]} = A^{[\ell]} \Lambda^{[\ell+1]}$ solved by QR or SVD

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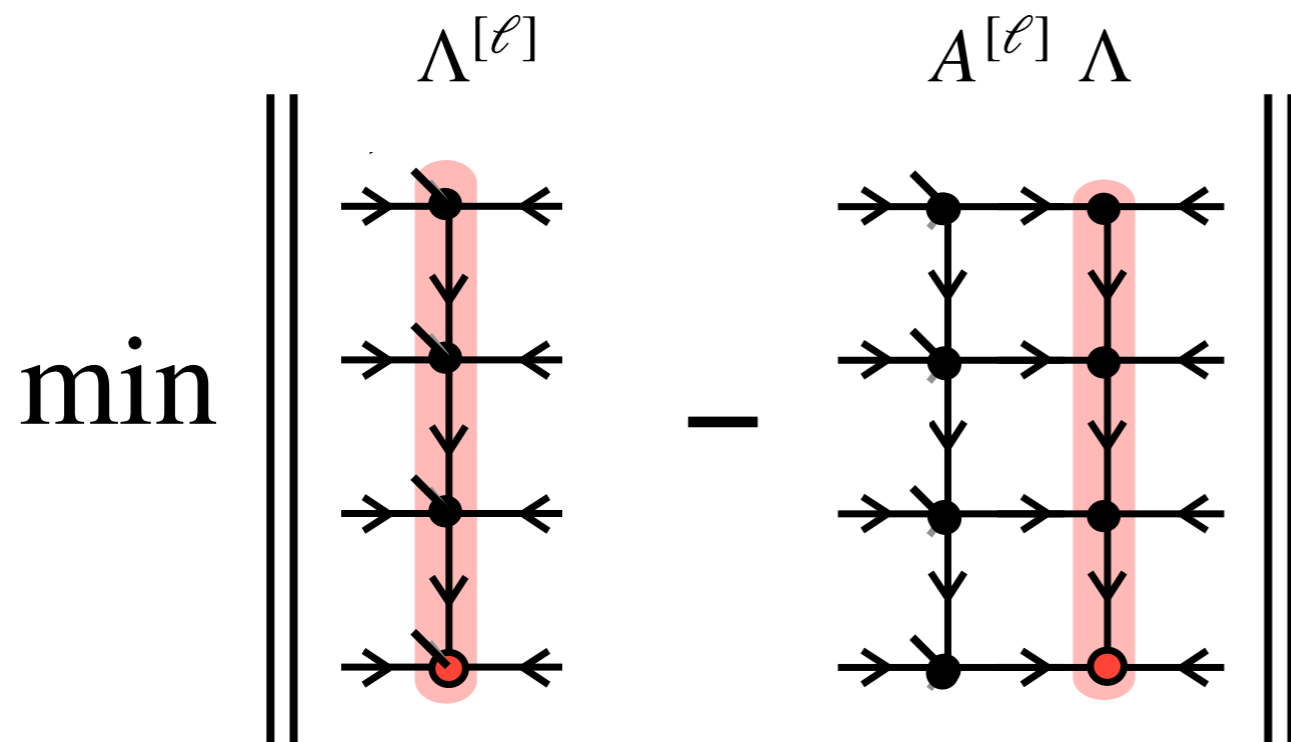
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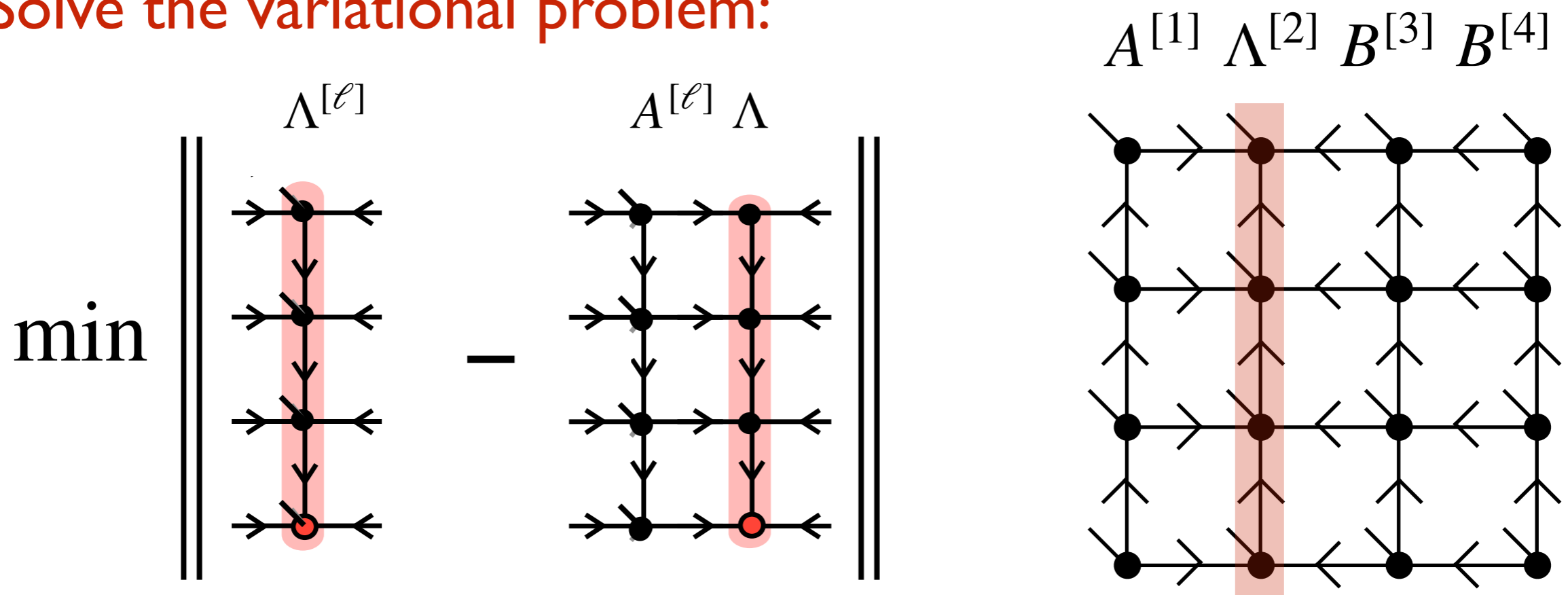
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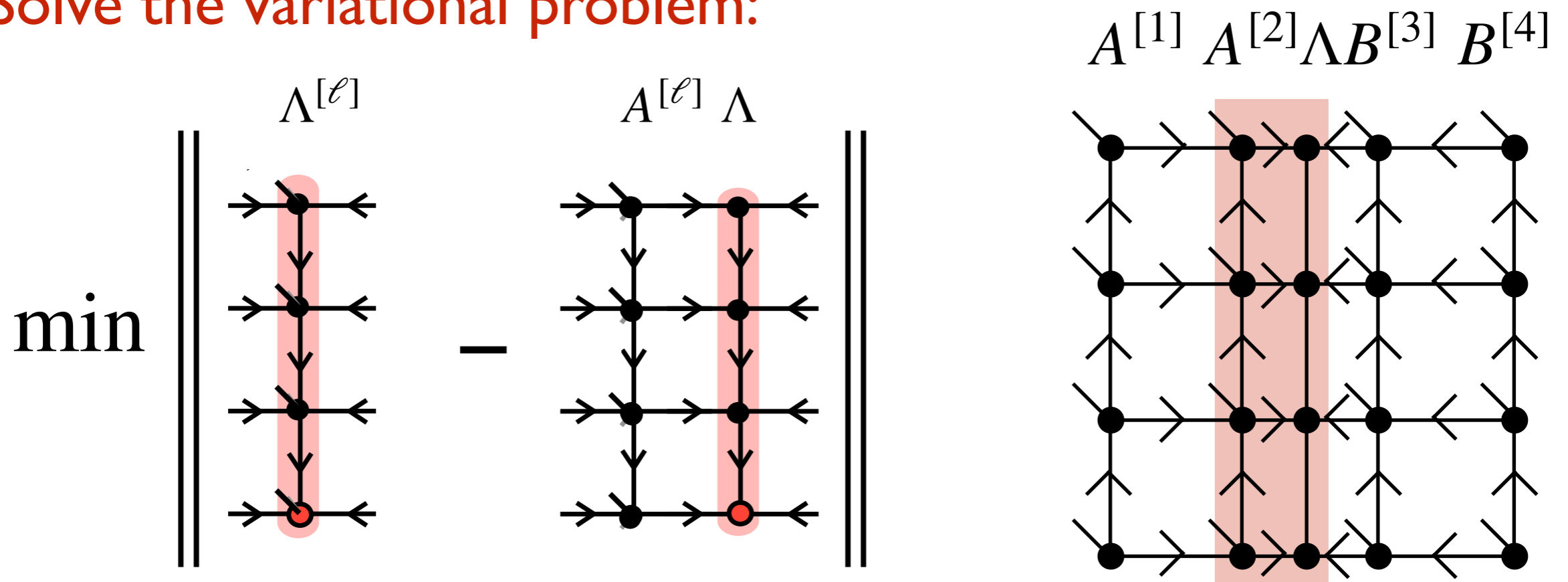
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Not possible for 2D TNS as it would destroy the locality of Λ

Solve the variational problem:



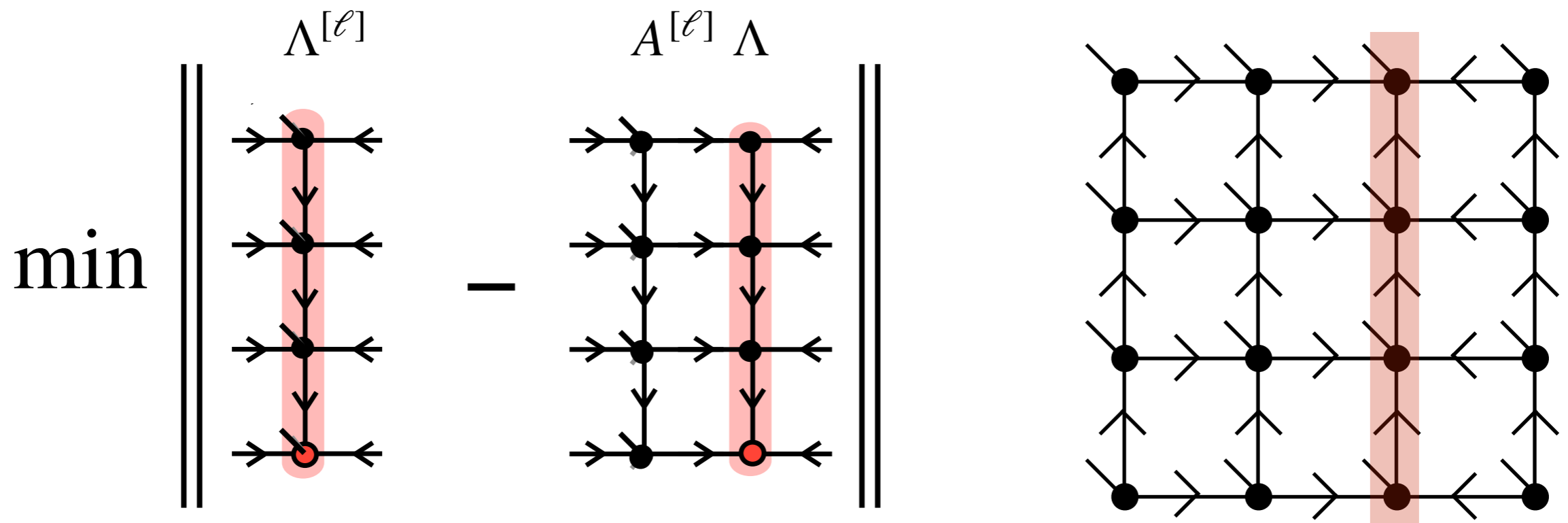
Isometric Tensor Network States in 2D

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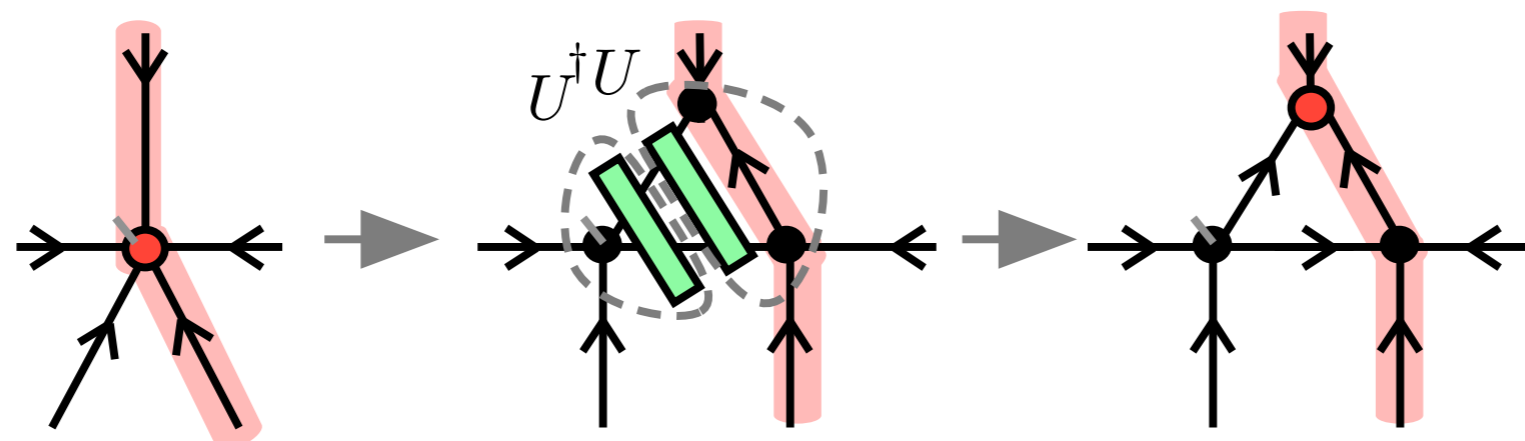
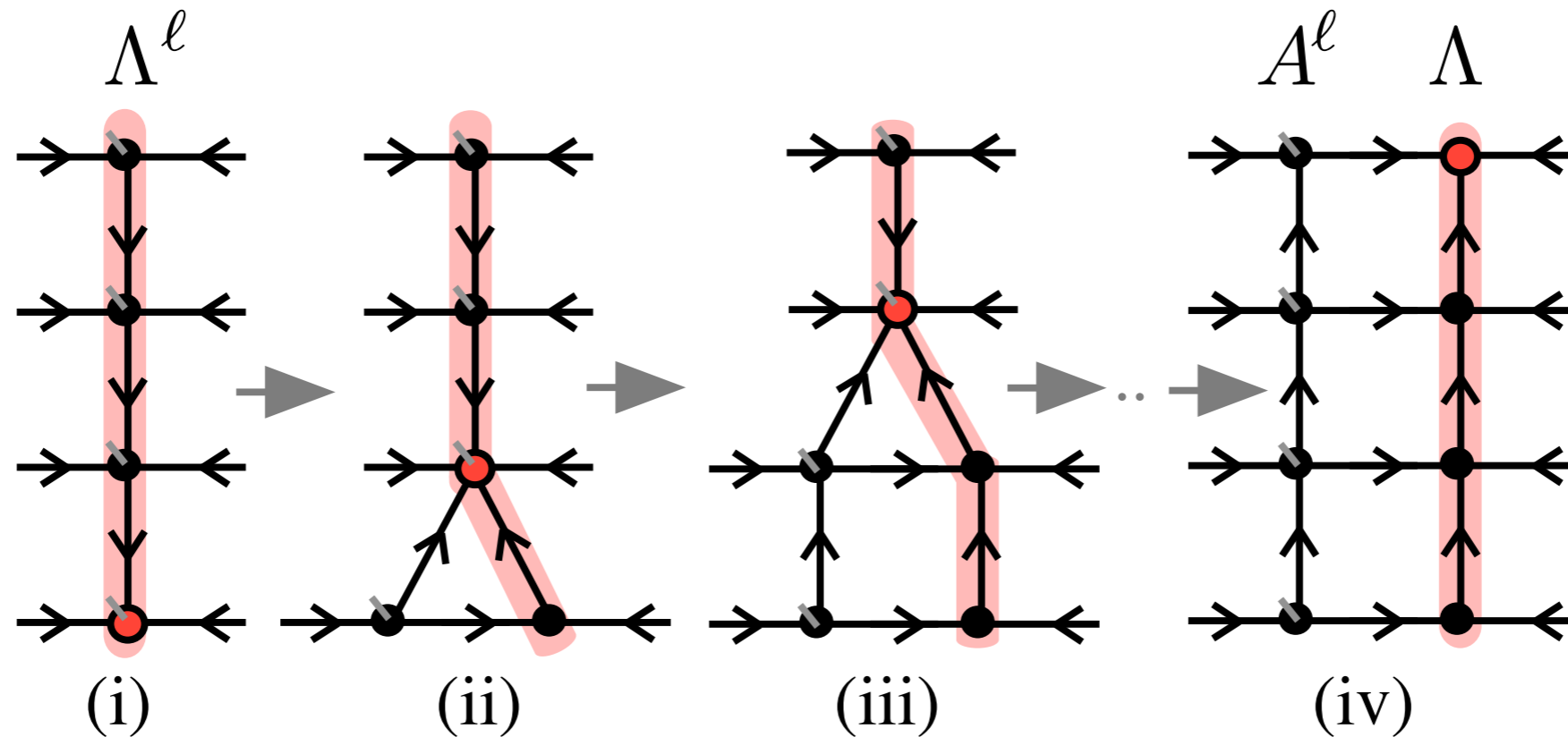
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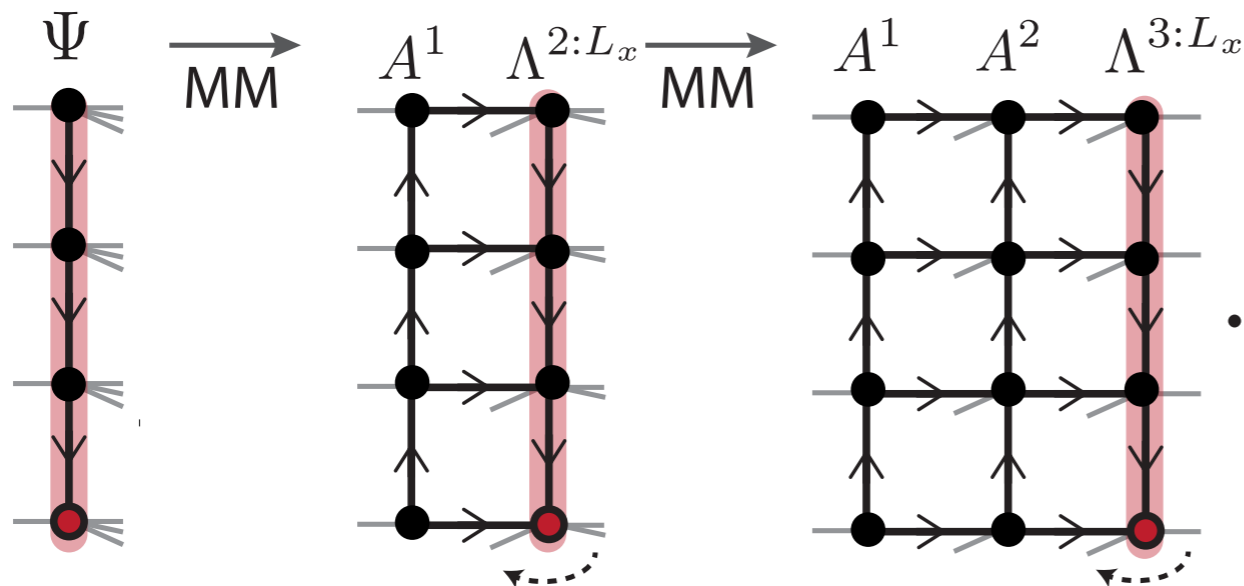
Isometric Tensor Network States in 2D

Sequential splitting based on disentangling: **“Moses Move” (MM)**



Convert quasi 1D MPS to isometric TNS

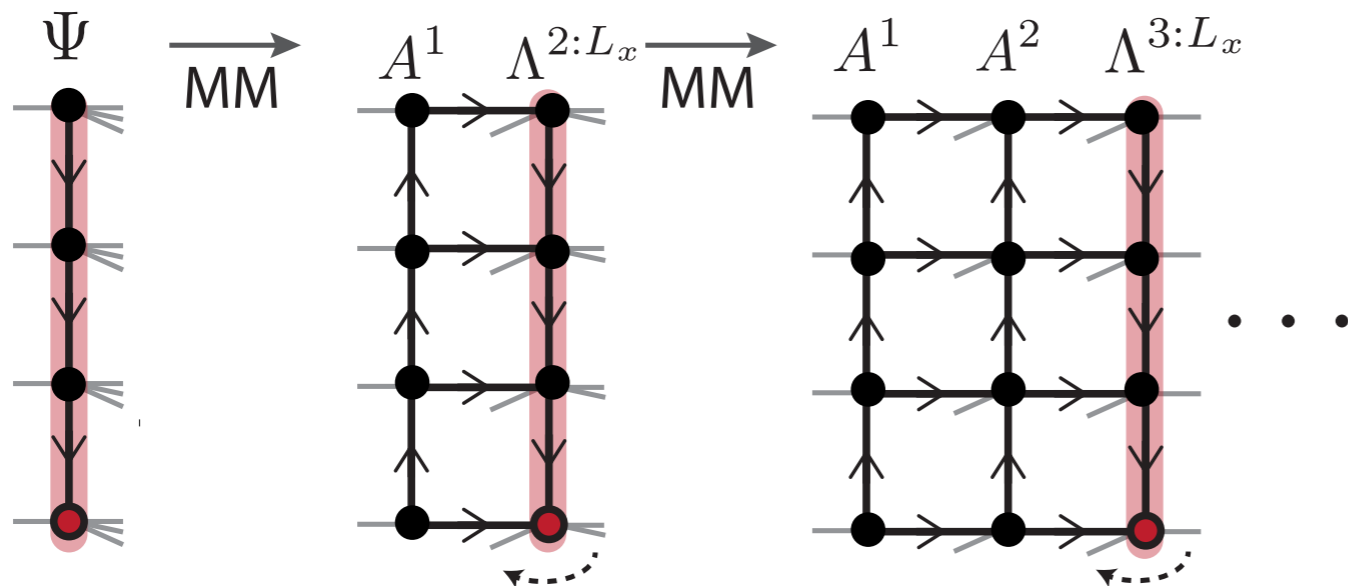
“Peel off” layers from MPS representation of 2D state



- ▶ **Sequentially disentangle the state**
- ▶ **Efficient compression**

Convert quasi 1D MPS to isometric TNS

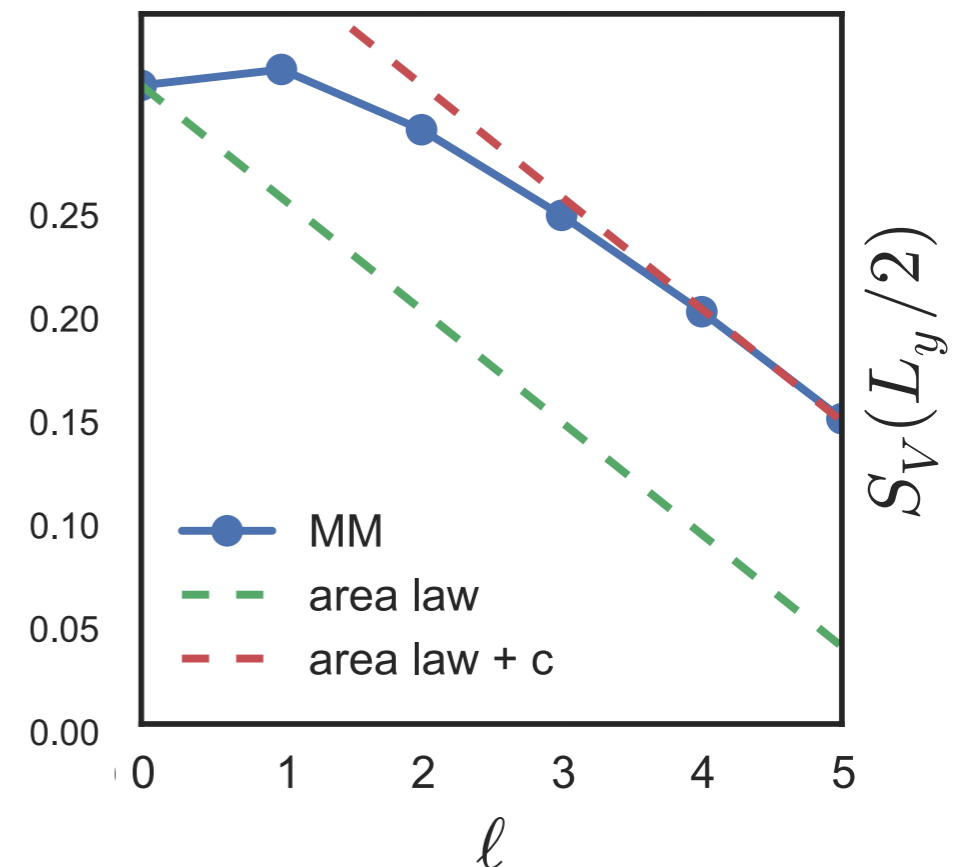
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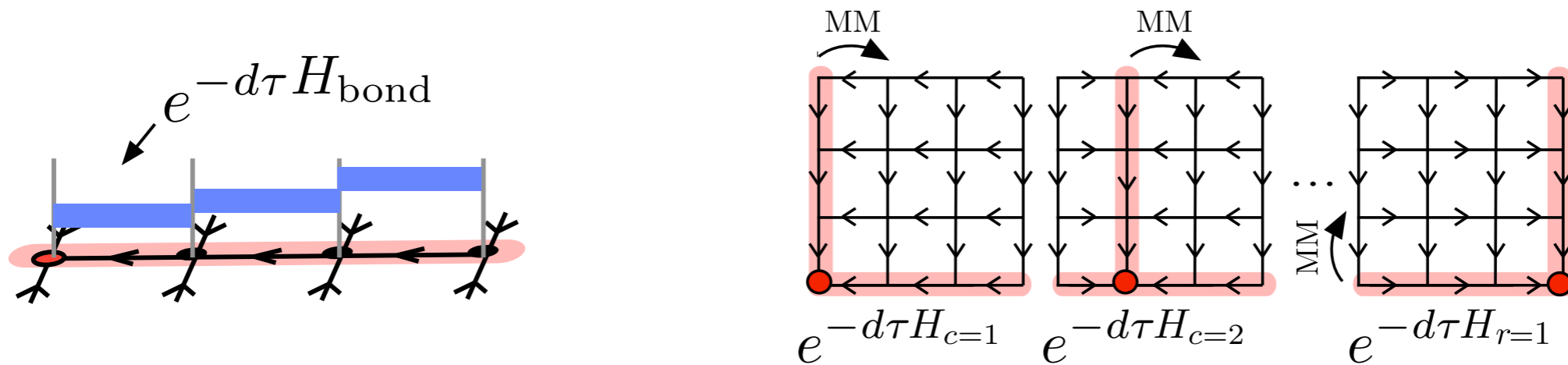
2D transverse field Ising Model ($g = 3.5$)

$$H = - \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z - g \sum_i \sigma_i^x$$



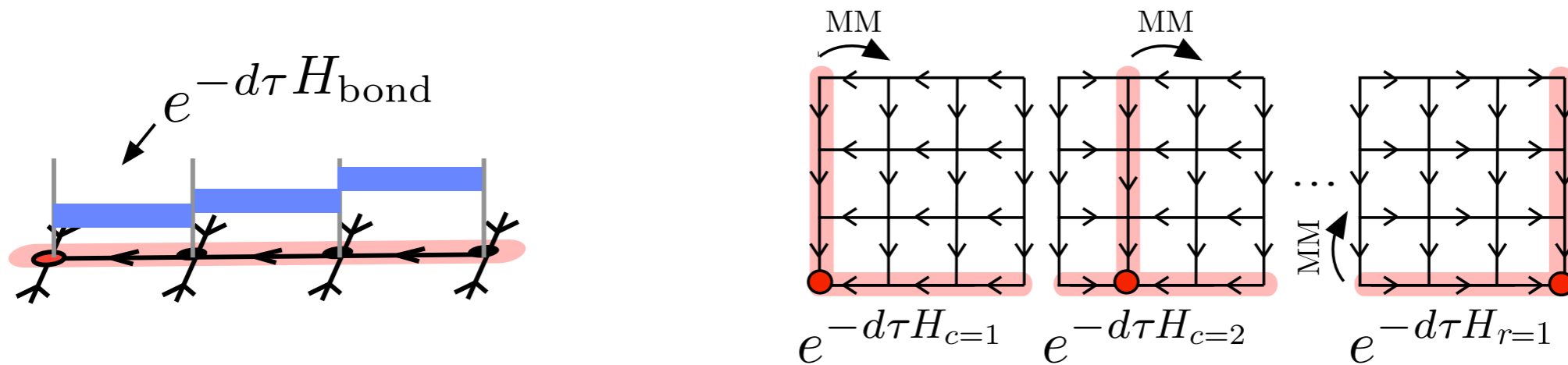
Ground states of 2D Hamiltonians

Sequentially apply **1D Time-Evolving Block Decimation (TEBD)** algorithm on the center columns/rows: 2nd order [Vidal '03]



Ground states of 2D Hamiltonians

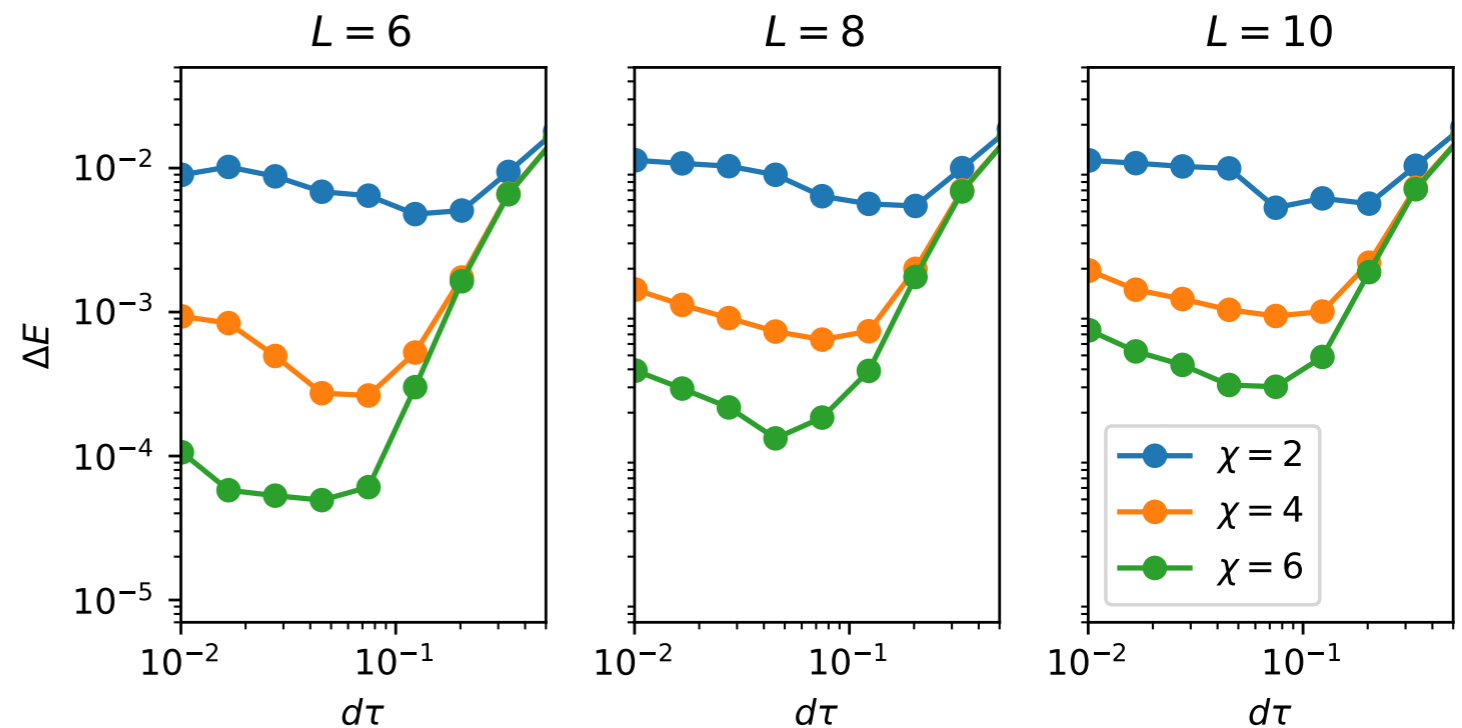
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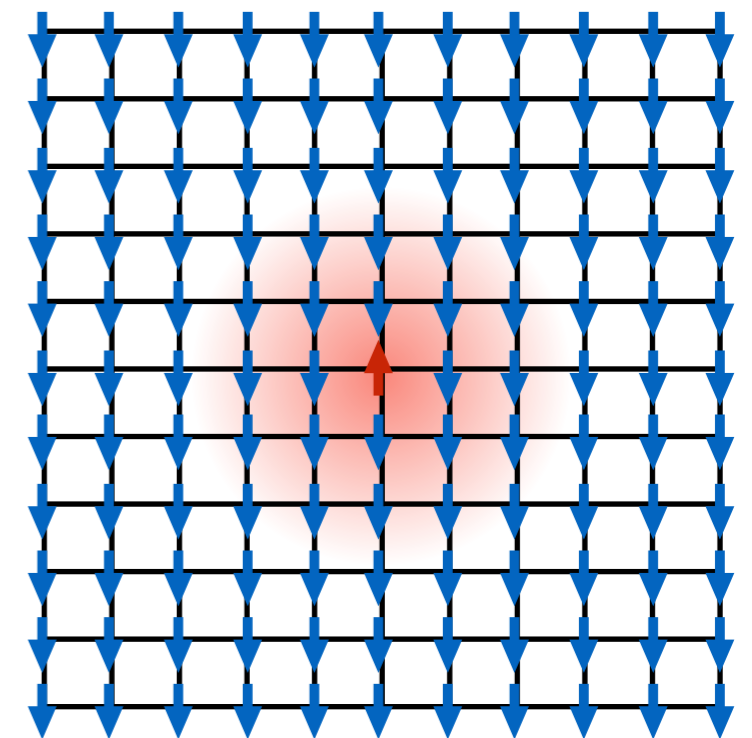
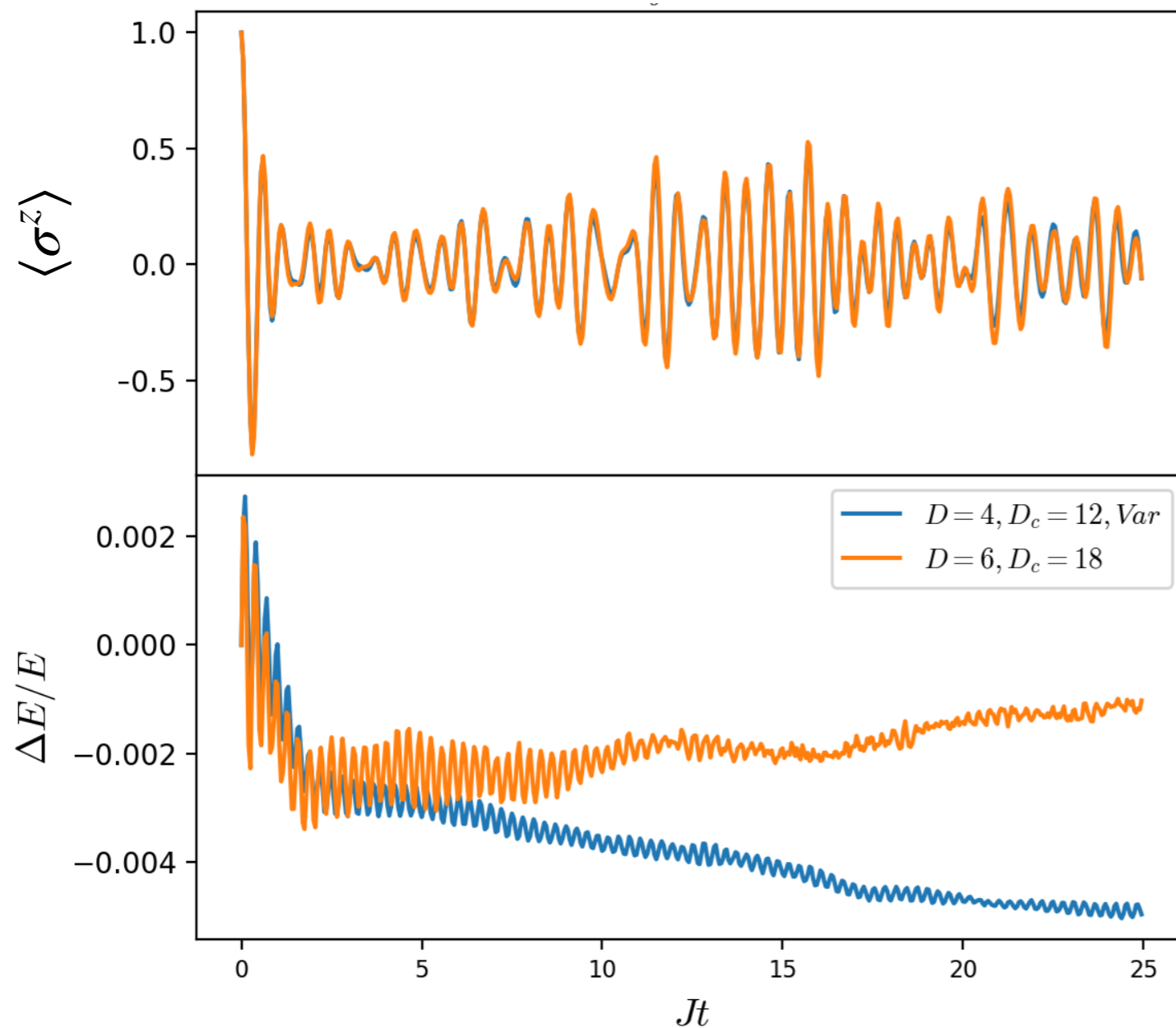
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Imaginary time evolution: $|\psi_0\rangle$



Real time evolution of 2D Hamiltonians

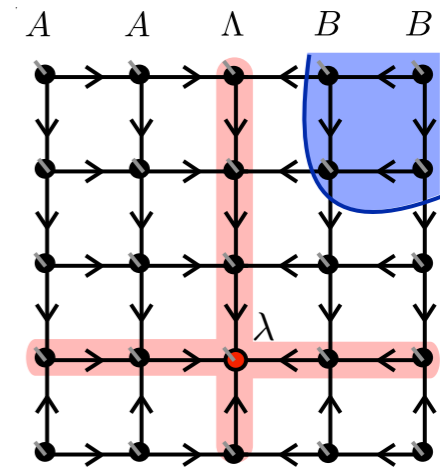
Real time evolution of $|\psi_0(t)\rangle = e^{-iHt} \sigma^+ |\psi_0\rangle$ for the transverse field Ising model (paramagnetic phase)



► **Good convergence at small bond dimension χ**

Summary

2D tensor-network state ansatz that allows for efficient contractions



- ▶ Subset of TNS: **Variational power?**
- ▶ Sequential splitting based on disentangling **Moses Move**
- ▶ **TEBD²** to obtain ground states and perform time evolution

[Zaletel and FP; arXiv:1902.05100]



Thank You!