

***Radiative corrections to quarkonium HFS
and η_b mass puzzle***

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Loops and Legs in Quantum Field Theory

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Topics covered

✓ QCD theory of quarkonium HFS

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- ✓ QCD theory of quarkonium HFS
- ✓ Discovery of pseudoscalar bottomonium state η_b

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- ✓ QCD theory of quarkonium HFS
- ✓ Discovery of pseudoscalar bottomonium state η_b
- ✓ $M^{\text{exp}}(\eta_b) - M^{\text{QCD}}(\eta_b) \sim 3\sigma \quad \Rightarrow \quad \text{“}\eta_b \text{ mass puzzle”}$

Heavy quarkonium experiment 2003-2008

Charmonium	Bottomonium
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$X(3872), X(3940)$	
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$Y(3940), Y(4260)$	
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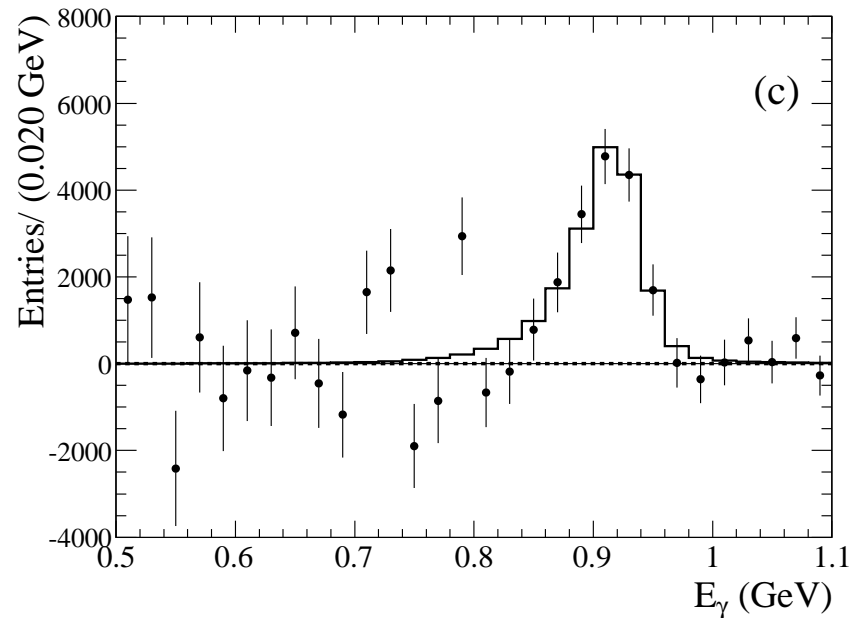
$Z(3930)$	
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<i>nothing</i>

Discovery of η_b - June 2008

BABAR Collaboration, Phys.Rev.Lett. 101, 071801 (2008)

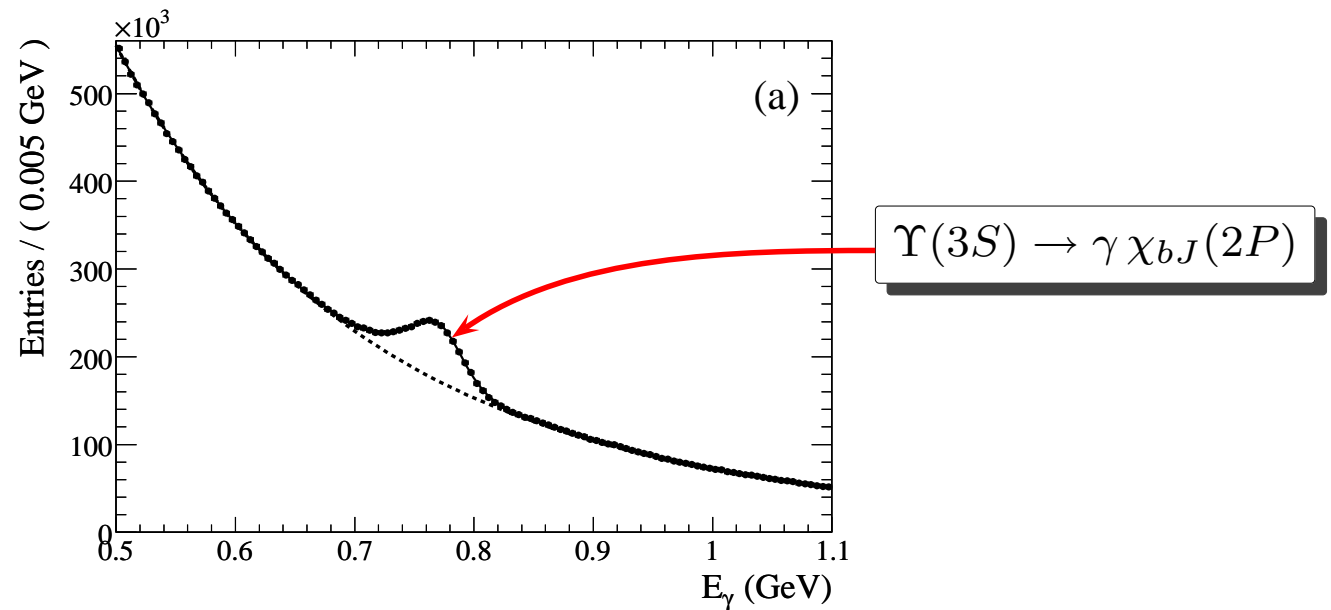
Single photon spectrum at $s = M(\Upsilon(3S))$



Peak at $E_\gamma = 921.2_{-2.8}^{+2.1}\text{stat} \pm 2.4\text{syst} \text{ MeV} \Rightarrow \Upsilon(3S) \rightarrow \gamma\eta_b$

Discovery of η_b - June 2008

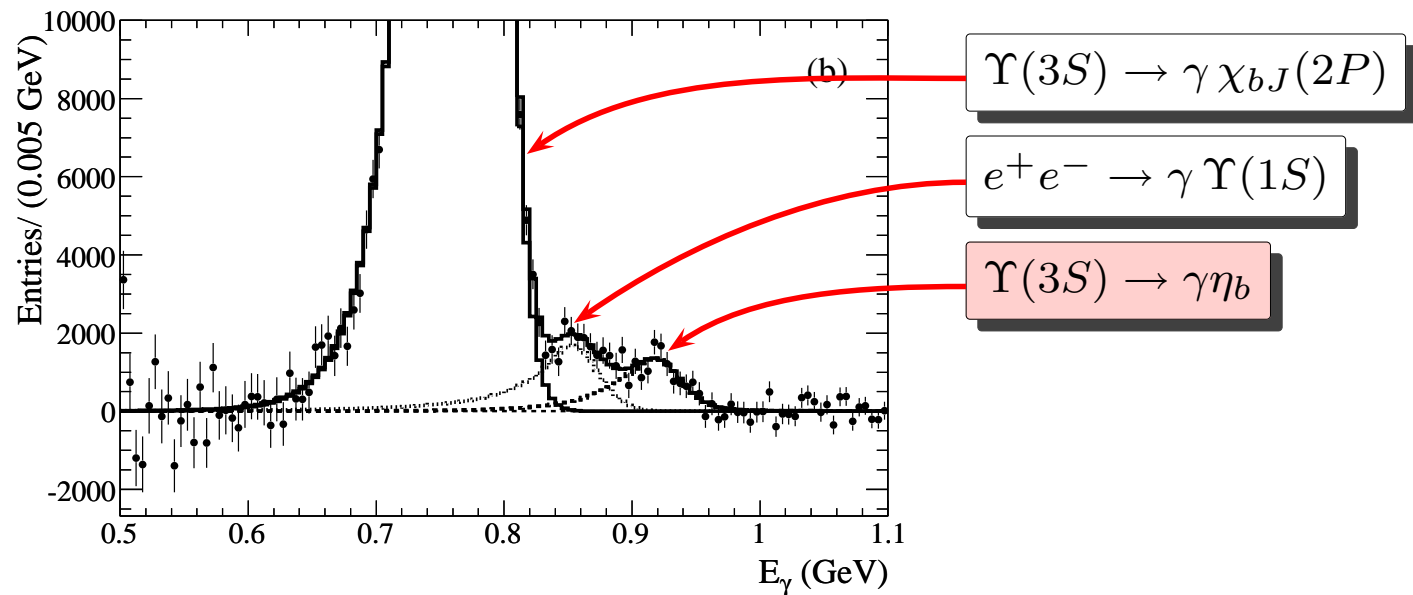
BABAR Collaboration, Phys.Rev.Lett. 101, 071801 (2008)



Inclusive spectrum

Discovery of η_b - June 2008

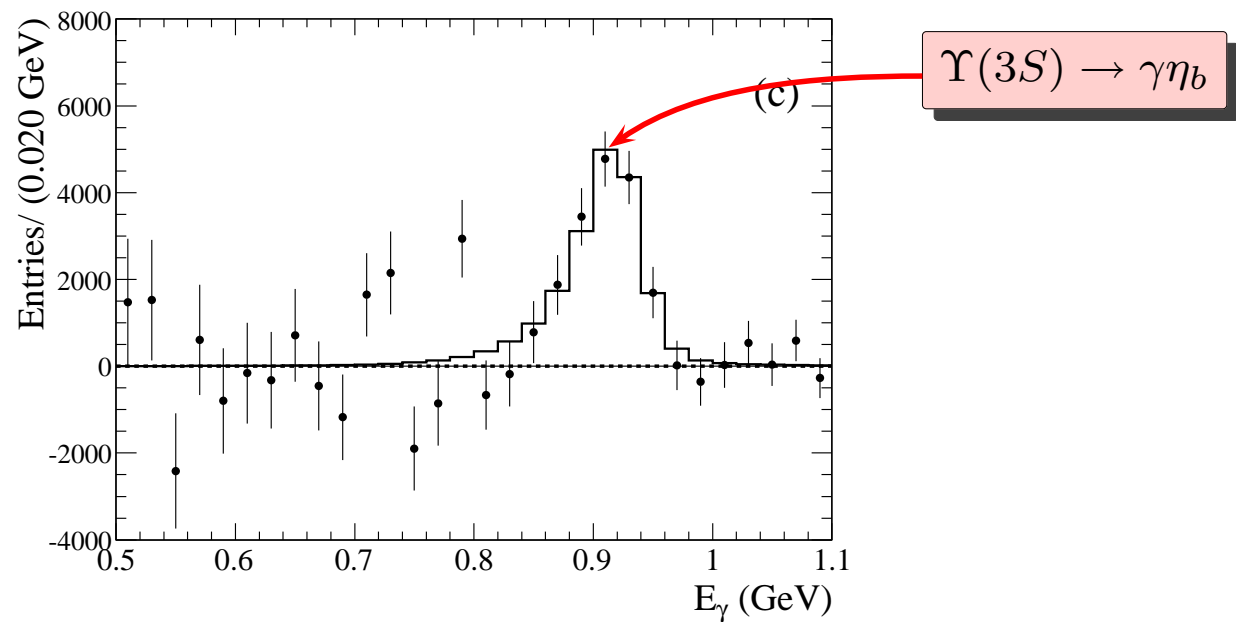
BABAR Collaboration, Phys.Rev.Lett. 101, 071801 (2008)



Non-peak background subtracted

Discovery of η_b - June 2008

BABAR Collaboration, Phys.Rev.Lett. 101, 071801 (2008)



All background subtracted

Hyperfine Splitting $E_{\text{hfs}} = M(\Upsilon_{1S}) - M(\eta_b)$

$$E_{\text{hfs}}^{\text{exp}} = 71.4 \pm 2.7 \text{ (syst)} \begin{matrix} +2.3 \\ -3.1 \end{matrix} \text{ (stat) MeV}$$

$\Upsilon(3S) \rightarrow \gamma\eta_b$

$$E_{\text{hfs}}^{\text{exp}} = 67.4 \pm 2.0 \text{ (syst)} \begin{matrix} +4.8 \\ -4.6 \end{matrix} \text{ (stat) MeV}$$

$\Upsilon(2S) \rightarrow \gamma\eta_b$

Theory of Hyperfine Splitting

✓ Perturbative QCD

✓ Lattice QCD

✗ Potential models

Perturbative QCD

- NRQCD, pNRQCD \Leftrightarrow systematic expansion in α_s
(Caswell, Lepage; Pineda, Soto; ...)
- Dimensional regularization \Leftrightarrow loops in pNRQCD
Threshold expansion
(Pineda, Soto; Beneke, Smirnov; ...)
- Log resummation \Leftrightarrow reduced scale dependence
(Luke, Manohar, Rothstein; Pineda; ...)
- Nonperturbative contribution $\mathcal{O}(v^2)$ \Leftrightarrow NNLO
(Voloshin; Leutwiller)

NLO corrections to HFS

Spin-flip potential:

$$V_S(\mathbf{q}^2) = \frac{4\pi C_F \alpha_s \mathbf{S}^2}{3m_q^2}$$

NLO corrections to HFS

Spin-flip potential:

$$V_S(\mathbf{q}^2) = \frac{4\pi C_F \alpha_s \mathbf{S}^2}{3m_q^2} \left\{ 1 + \frac{\alpha}{\pi} \left[\left(1 + \frac{7}{8} \left(\frac{1}{\epsilon} + \ln \frac{\mu^2}{m_q^2} \right) \right) C_A - \frac{1}{2} C_F + \frac{6 - 6 \ln 2 + i3\pi}{4} T_F \right] \right\}$$

hard contribution

QCD on-shell on-threshold
amplitude

NLO corrections to HFS

Spin-flip potential:

$$V_S(\mathbf{q}^2) = \frac{4\pi C_F \alpha_s \mathbf{S}^2}{3m_q^2} \left\{ 1 + \frac{\alpha}{\pi} \left[\left(1 + \frac{7}{8} \left(\frac{1}{\epsilon} + \ln \frac{\mu^2}{m_q^2} \right) \right) C_A - \frac{1}{2} C_F + \frac{6 - 6 \ln 2 + i3\pi}{4} T_F \right] + \frac{\alpha}{\pi} \left[\left(-\frac{7}{18} - \frac{7}{8} \left(\frac{1}{\epsilon} + \ln \frac{\mu^2}{q^2} \right) \right) C_A - \frac{5}{9} T_F n_l \right] \right\}$$

hard contribution

QCD on-shell on-threshold
amplitude

soft contribution

NQCD, static heavy
quark propagator

Coulomb potential:

$$V_C(\mathbf{q}^2) = -\frac{4\pi C_F \alpha_s(\mathbf{q}^2)}{q^2} \left[1 + \frac{\alpha_s}{\pi} \left(\frac{31}{36} C_A - \frac{5}{9} T_F n_l \right) \right]$$

NLO corrections to HFS

Quantum Mechanical PT (*potential contribution*)

$$\delta^{NLO} E_{\text{hfs}} = \langle \psi^{\text{Coulomb}} | V_S^{1\text{-loop}} | \psi^{\text{Coulomb}} \rangle + 2 \langle \psi^{1\text{-loop}} | V_S^{\text{tree}} | \psi^{\text{Coulomb}} \rangle$$

NLO corrections to HFS

Quantum Mechanical PT (*potential contribution*)

$$\delta^{NLO} E_{\text{hfs}} = \langle \psi^{\text{Coulomb}} | V_S^{1\text{-loop}} | \psi^{\text{Coulomb}} \rangle + 2 \langle \psi^{1\text{-loop}} | V_S^{\text{tree}} | \psi^{\text{Coulomb}} \rangle$$

Result for general n

$$E_{\text{hfs}}^{NLO}(n) = \frac{1}{3} \frac{C_F^4 \alpha_s^4}{n^3} \left\{ 1 + \frac{\alpha_s}{\pi} \left[\frac{7 C_A}{4} \ln \left(\frac{C_F \alpha_s}{n} \right) - \frac{C_F}{2} + \frac{-15 - 11n + 12n^2 \Psi_2(n)}{9n} \right] \right. \\ \left. \times n_f T_F + \frac{393 + 150n + 126 \gamma_E n + 126n \Psi_1(n) - 264n^2 \Psi_2(n)}{72n} C_A \right\}$$

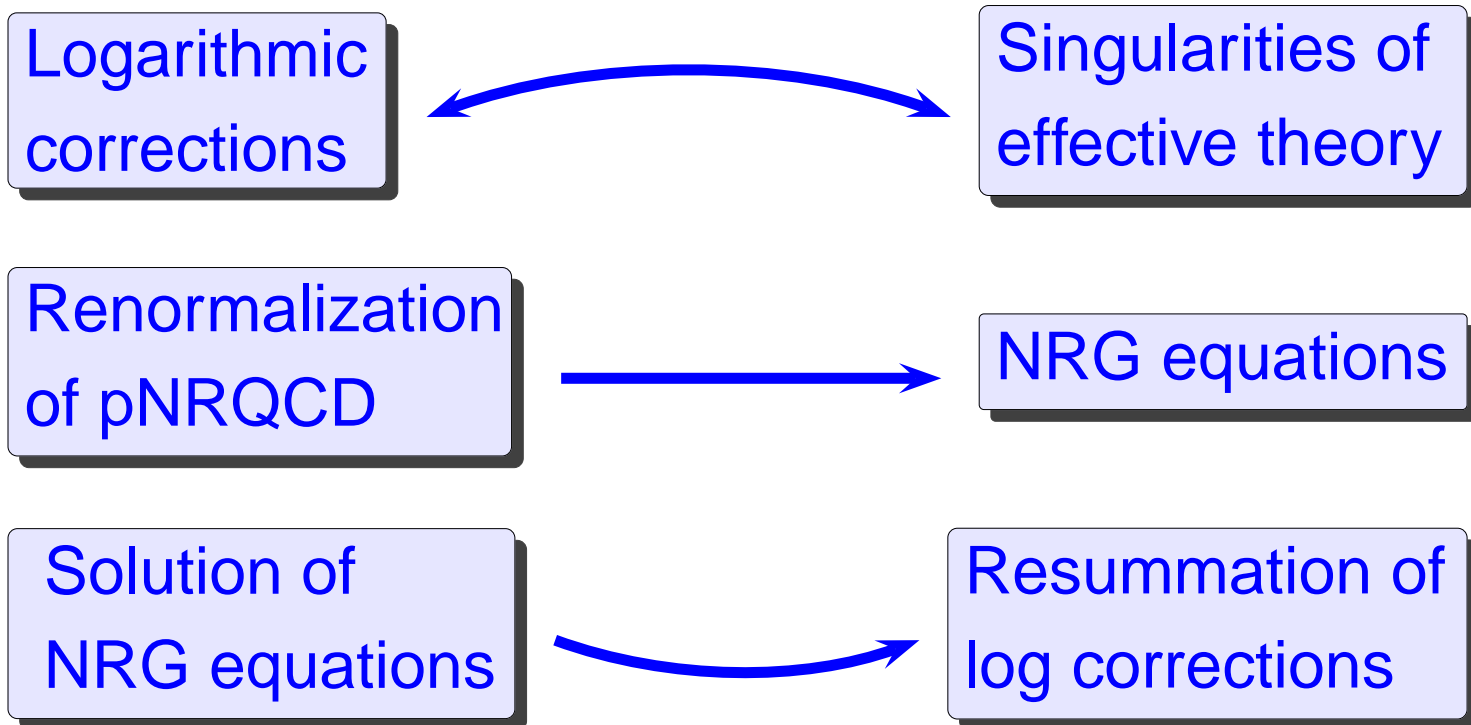
Penin, Steinhauser, Phys.Lett. B538, 335 (2002)

Nonrelativistic Renormalization Group

- Several scales: $m_q, m_q v, m_q v^2$
- Logarithmic integrals between the scales $\Leftrightarrow \ln v \Leftrightarrow \ln \alpha_s$

Nonrelativistic Renormalization Group

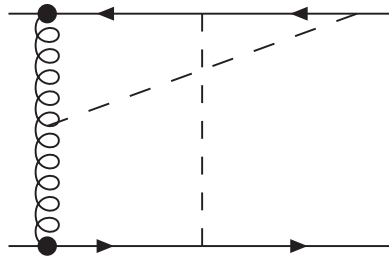
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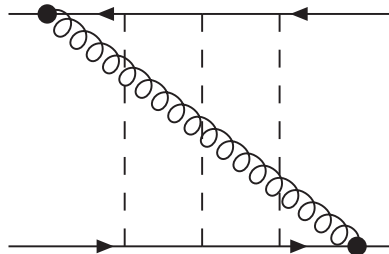
NRG running of spin-flip potential

● NLL running $\Rightarrow \mathcal{O}(\alpha_s^m \ln^{m-1} \alpha_s)$

● Two-loop soft running:



● Three-loop ultrasoft-potential running



NRG running of spin-flip potential

- NLL soft running

- *NRG equation*

$$\mu_s \frac{d}{d\mu_s} D_{S^2, s}^{(2)} = \alpha_s(\mu_s) c_F^2(\mu_s) \gamma_s(\alpha_s)$$

- *LL Fermi coupling* $c_F(\mu) = (\alpha_s(\mu)/\alpha_s(m_q))^{-C_A/\beta_0}$

- *Two-loop anomalous dimension* $\gamma_s = \gamma_s^{(1)} \frac{\alpha_s}{\pi} + \dots$

$$\gamma_s^{(2)} = \frac{1}{216} [C_A^2 (5 - 36 \pi^2) + 88 C_A n_l T_F + 4 n_l T_F (27 C_F - 40 n_l T_F)]$$

Penin, Pineda, Smirnov, Steinhauser, Phys.Lett. B593, 124 (2004)

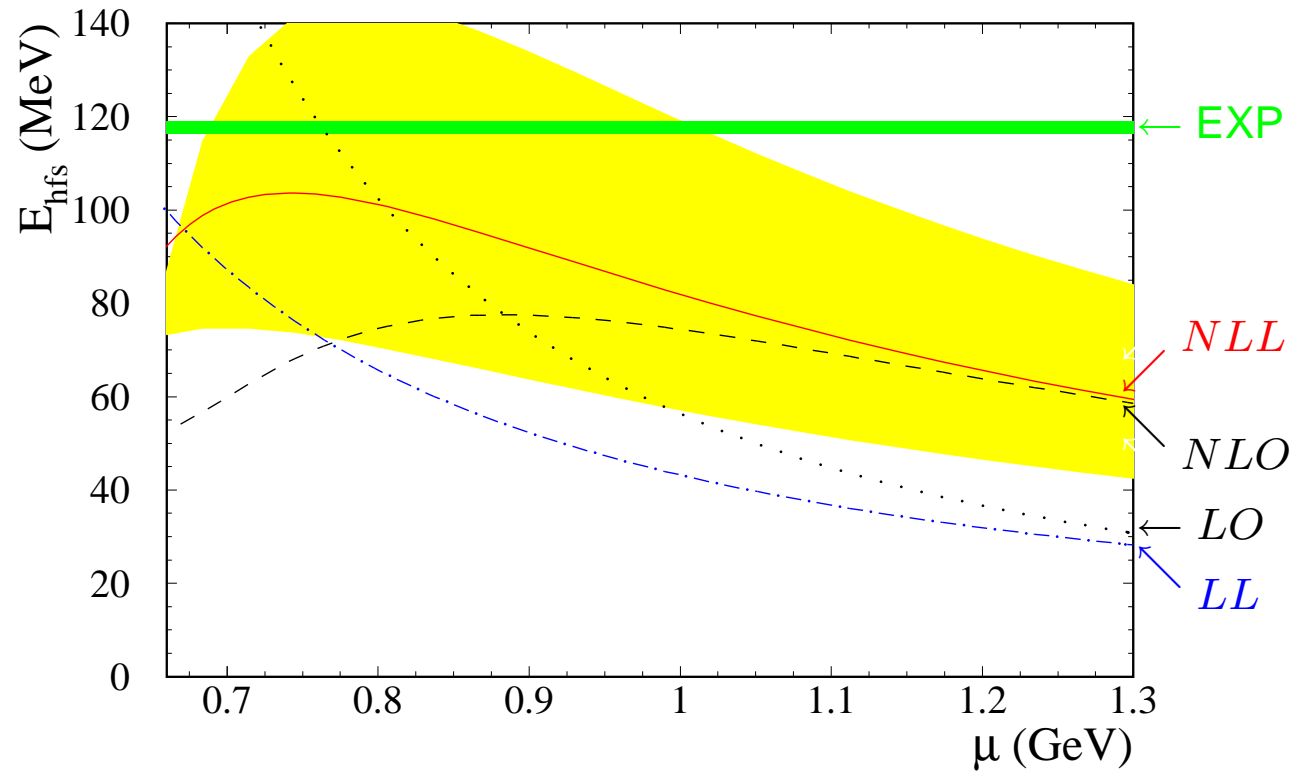
Bottomonium HFS in NLL approximation

Kniehl, Penin, Pineda, Smirnov, Steinhauser, Phys.Rev.Lett. 92, 242001 (2004)

$$\begin{aligned}
 M_{\Upsilon(1S)} - M_{\eta_b} = & \frac{C_F^4 \alpha_s^4(\nu) m_b}{3} \left\{ \frac{27}{14} y^{-1} - \frac{13}{14} y^{-\frac{18}{25}} + \frac{\alpha_s(m_b)}{\pi} \left[\left(\frac{1037}{224} + \frac{405086361761 \pi^2}{25617160800} \right. \right. \right. \\
 & \left. \left. - \frac{3}{4} \ln 2 \right) \times y^{-1} - \frac{1024 \pi^2}{143} y^{-\frac{39}{50}} - \left(\frac{102973}{26250} + \frac{184336 \pi^2}{25725} \right) y^{-\frac{18}{25}} + \frac{1024 \pi^2}{675} y^{-\frac{1}{2}} + \frac{671 \pi^2}{1029} y^{-\frac{1}{2}} \right. \\
 & \left. - \frac{3 \pi^2}{23} y^{-\frac{2}{25}} + \left(-\frac{13427921}{1260000} + \frac{88057 \pi^2}{151200} \right) y^{\frac{7}{25}} + \frac{4 \pi^2}{41} y^{\frac{16}{25}} + \frac{1377}{56} - \frac{1253587 \pi^2}{227500} - \frac{629 \pi^2}{7500} y \right. \\
 & \left. - \frac{2873 \pi^2}{7182} y^{\frac{32}{25}} {}_2F_1 \left(\frac{57}{25}, 1; \frac{82}{25}; \frac{y}{2} \right) + \frac{2873 \pi^2}{3591} y^{-1} {}_2F_1 \left(1, 1; \frac{82}{25}; -1 \right) + \left(\frac{675}{28} - \frac{533}{42} y^{\frac{7}{25}} \right) \right. \\
 & \left. \times \ln \left(\frac{\mu}{C_F \alpha_s(\mu) m_b} \right) + \frac{85248 \pi^2}{30625} y^{-1} \ln y + \left(-\frac{45834}{4375} y^{-1} + \frac{21216}{4375} - \frac{2873}{1575} y^{\frac{7}{25}} + \frac{243}{1250} y \right) \right. \\
 & \left. \times \pi^2 \ln(2 - y) \right\},
 \end{aligned}$$

$$y = \frac{\alpha_s(\mu)}{\alpha_s(m_b)}$$

HFS in charmonium



$$E_{\text{hfs}} = M(J/\Psi) - M(\eta_c)$$

$$n = 1$$

- QCD NLL approximation:

$$E_{\text{hfs}}^{\text{th}} = 112 \text{ MeV}$$

- Experiment:

$$E_{\text{hfs}}^{\text{exp}} = 117.7 \pm 1.3 \text{ MeV}$$

$$E_{\text{hfs}} = M(J/\Psi) - M(\eta_c)$$

$$n = 2$$

- QCD NLL approximation:

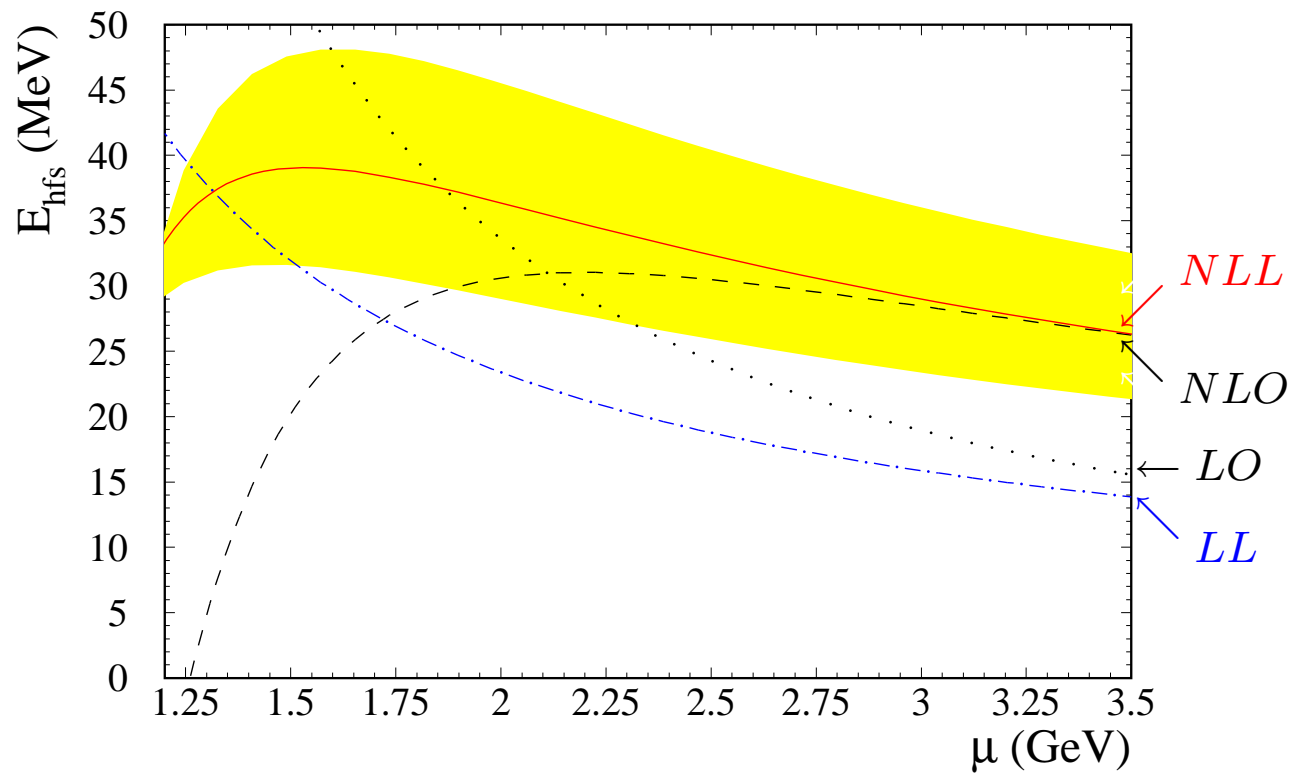
$$E_{\text{hfs}}^{\text{th}} = 41 \text{ MeV}$$

- Experiment:

$$E_{\text{hfs}}^{\text{exp}} = 41 \pm 3 \text{ MeV}$$

HFS in bottomonium

(Kniehl, Penin, Pineda, Smirnov, Steinhauser)



$$E_{\text{hfs}} = M(\Upsilon_{1S}) - M(\eta_b)$$

- QCD NLL approximation:

$$E_{\text{hfs}}^{\text{th}} = 41 \pm 11 \text{ (th)} \begin{matrix} +9 \\ -8 \end{matrix} (\delta\alpha_s) \text{ MeV}$$

- Experiment:

$$E_{\text{hfs}}^{\text{exp}} = 71.4 \pm 2.7 \text{ (syst)} \begin{matrix} +2.3 \\ -3.1 \end{matrix} \text{ (stat)} \text{ MeV}$$



Why perturbative QCD *works* for charmonium but *fails* for bottomonium HFS?

- *Nonperturbative correction scales as $1/m_q^4$*
Perturbative correction scales as $\alpha_s(\alpha_s m_q)$
- ➔ *Cancellation of huge perturbative and nonperturbative corrections for charmonium, no cancellation for bottomonium?*
- *Hard to believe since leading $\mathcal{O}(v^2)$ nonperturbative contribution is positive \Leftrightarrow bottomonium HFS would be even smaller!*

Lattice QCD

- “Full” lattice NRQCD simulations:

$$E_{\text{hfs}}^{\text{lat}} = 61 \pm 14 \text{ MeV}$$

A. Gray *et al.*, Phys.Rev. D72, 094507 (2005)

- Experiment:

$$E_{\text{hfs}}^{\text{exp}} = 71.4 \pm 2.7 \text{ (syst)} \begin{matrix} +2.3 \\ -3.1 \end{matrix} \text{ (stat) MeV}$$

Lattice QCD

- Hard cutoff $1/a \sim 2 \text{ GeV}$ from spin-average spectrum
- Logarithmic contribution of the hard modes:

$$\delta^{\text{hard}} E_{\text{hfs}} = -\frac{\alpha_s}{\pi} \frac{7 C_A}{4} \ln(am_b) E_{\text{hfs}} \approx -20 \text{ MeV}$$

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$$E_{\text{hfs}}^{\text{lat}} + \delta^{\text{hard}} E_{\text{hfs}} \approx 40 \text{ MeV}$$

➔ *in perfect agreement with perturbative QCD!*

η_b mass puzzle

● Origin of discrepancy



Underestimate of theoretical uncertainty



Underestimate of experimental uncertainty



Exotic e.g. light CP-odd Higgs Boson

Domingo *et al.*, Phys.Rev.Lett. 103, 111802 (2009)

η_b mass puzzle

- Origin of discrepancy



Underestimate of theoretical uncertainty



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Exotic e.g. light CP-odd Higgs Boson

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- BaBar guys are very angry

➔ better to solve the problem now!