# A new simple PDF parametrization: improved description of the HERA data

#### M. Bonvini, <u>F. Giuli</u>

#### xFitter External Workshop, Minsk (Belarus) 18/03/2019





#### **Theoretical motivations**

Previous fit based on default HERAPDF parameterisation:

$$xf(x,\mu_0^2) = A x^B (1-x)^C \left[1 + Dx + Ex^2\right] - A' x^{B'} (1-x)^{C'}$$

Negative term (only for gluon)

- Is this parameterisation flexible enough to describe small-x behaviour?
- ➤ Limited structure at small-x → shape strongly dominated by asymptotic behaviour of x<sup>B</sup>
- > More flexibility in the small-x regime is needed
- > It is also very important in the light of future higher-energy colliders:
  - Large Hadron-electron Collider (LHeC)
  - Future Circular electron-hadron or hadron-hadron Colliders (FCC-eh and FCC-hh)

## New proposed parametrization

> New general parametrization:

Multiplicative option:

$$xf(x,\mu_0^2) = A x^B (1-x)^C \left[1 + Dx + Ex^2\right] \left[1 + F \log x + G \log^2 x + H \log^3 x\right]$$
  
> Additive option:

$$xf(x,\mu_0^2) = A x^B (1-x)^C \left[ 1 + Dx + Ex^2 + F \log x + G \log^2 x + H \log^3 x \right]$$

 $\succ$  Difference between the two options in the medium-x region e.g. x ~ 0.1

$$\begin{split} xg(x,\mu_0^2) &= A_g \, x^{B_g} (1-x)^{C_g} \Big[ 1 + F_g \log x + G_g \log^2 x \Big] \\ xu_v(x,\mu_0^2) &= A_{u_v} \, x^{B_{u_v}} (1-x)^{C_{u_v}} \Big[ 1 + E_{u_v} x^2 + F_{u_v} \log x + G_{u_v} \log^2 x \Big] \\ xd_v(x,\mu_0^2) &= A_{d_v} \, x^{B_{d_v}} (1-x)^{C_{d_v}} \\ x\bar{u}(x,\mu_0^2) &= A_{\bar{u}} \, x^{B_{\bar{u}}} (1-x)^{C_a} \Big[ 1 + D_{\bar{u}} x + F_{\bar{u}} \log x \Big] \\ x\bar{d}(x,\mu_0^2) &= A_{\bar{d}} \, x^{B_{\bar{u}}} (1-x)^{C_{\bar{d}}} \Big[ 1 + D_{\bar{d}} x + F_{\bar{d}} \log x \Big], \end{split}$$
 Minimal parametrization 
$$xs(x,\mu_0^2) &= x\bar{s}(x,\mu_0^2) = r_s \, x\bar{d}(x,\mu_0^2) \qquad r_s = \frac{f_s}{1-f_s} \quad \text{with} \ \underline{f_s} = 0.4 \text{ fixed} \end{split}$$

## New proposed parametrization

- > Our new PDF parametrization:
  - Depends on 18 free parameters that must be fitted
  - > This is to be compared with HERAPDF2.0 (14 free parameters)
  - $\succ$  Two extra parameters for  $u_V$
  - $\succ$  Two extra parameters for  $\bar{u}$  and  $\bar{d}$
  - Major improvement comes from the gluon PDF (same number of free parameters)

$$\begin{split} xg(x,\mu_0^2) &= A_g \, x^{B_g} (1-x)^{C_g} \Big[ 1+F_g \log x + G_g \log^2 x \Big] \\ xu_v(x,\mu_0^2) &= A_{u_v} \, x^{B_{u_v}} (1-x)^{C_{u_v}} \Big[ 1+E_{u_v} x^2 + F_{u_v} \log x + G_{u_v} \log^2 x \Big] \\ xd_v(x,\mu_0^2) &= A_{d_v} \, x^{B_{d_v}} (1-x)^{C_{d_v}} \\ x\bar{u}(x,\mu_0^2) &= A_{\bar{u}} \, x^{B_{\bar{u}}} (1-x)^{C_a} \Big[ 1+D_{\bar{u}} x + F_{\bar{u}} \log x \Big] \\ x\bar{d}(x,\mu_0^2) &= A_{\bar{d}} \, x^{B_{\bar{u}}} (1-x)^{C_{\bar{d}}} \Big[ 1+D_{\bar{d}} x + F_{\bar{u}} \log x \Big], \end{split}$$
 Minimal parametrization 
$$xs(x,\mu_0^2) &= x\bar{s}(x,\mu_0^2) = r_s \, x\bar{d}(x,\mu_0^2) \qquad r_s = \frac{f_s}{1-f_s} \quad \text{with} \ \underline{f_s} = 0.4 \text{ fixed} \end{split}$$

### **Comparison to HERAPDF2.0**

Contribution to $\chi^2$	HERAPDF2.0	Our fit (new parametrization)
subset NC $e^+$ 920 $\bar{\chi}^2/{\rm n.d.p.}$	444/377	403/377
subset NC $e^+$ 820 $\tilde{\chi}^2$ /n.d.p.	66/70	74/70
subset NC $e^+$ 575 $\tilde{\chi}^2/\text{n.d.p.}$	219/254	221/254
subset NC $e^+$ 460 $\tilde{\chi}^2$ /n.d.p.	217/204	222/204
subset NC $e^- \tilde{\chi}^2$ /n.d.p.	219/159	220/159
subset CC $e^+ \bar{\chi}^2/\text{n.d.p.}$	45/39	38/39
subset CC $e^- \tilde{\chi}^2$ /n.d.p.	56/42	50/42
correlation term $+ \log$ term	91 + 5	75-3
Total $\chi^2/d.o.f.$	1363/1131	1301/1127
$\chi^{2} = \sum_{i} \frac{\left[D_{i} - T_{i}\left(1 - \sum_{j}\gamma_{i}\right) - \frac{1}{\delta_{i,\text{uncor}}^{2}T_{i}^{2} + \delta_{i,\text{stat}}^{2}\right]}{\delta_{i,\text{uncor}}^{2}T_{i}^{2} + \delta_{i,\text{stat}}^{2}}$ Exp. term	$\frac{\left[i_{j}b_{j}\right]^{2}}{D_{i}T_{i}} + \sum_{j}b_{j}^{2} + \sum_{j}Corr. \text{ term}$	$+\sum_{i} \log \frac{\delta_{i,\text{uncor}}^2 T_i^2 + \delta_{i,\text{stat}}^2 D_i T_i}{\delta_{i,\text{uncor}}^2 D_i^2 + \delta_{i,\text{stat}}^2 D_i^2}$ Log term





- Richer structure at medium-/high-x than HERAPDF2.0
- Gluon decreases more rapidly for x ~ 10<sup>-2</sup> and starts rising again for x < 10<sup>-4</sup>
- Up-valence rather different
- Down-valence is identical (same parametrization as in HERAPDF2.0)
- If compared to NNPDF3.0 (HERA data only), qualitatively same behavior

6



How could the fit quality improve so much?

- In most of the cases the agreement is at the same level
- Exception for low-Q<sup>2</sup> and low-x data, where a clear improvement of the theoretical description is manifest
- > This region is where the impact of log(1/x) terms is expected to be largest
- >  $\chi^2$  improvement of the same size as the one found in our small-x resummation paper Eur. Phys. J. C78 (2018) 621 is resummation really needed?

## From TR to FONLL

#### Various variations studied

First of all, migration from TR scheme to FONLL (to include small-x resummation in a later stage) – as done in <u>Eur. Phys. J. C78 (2018) 621</u>

#### Differences in the fit setup

heavy flavour scheme	TR	FONLL
initial scale $\mu_0$	$1.38  {\rm GeV}$	$1.6 \mathrm{GeV}$
charm matching scale $\mu_c$	$m_c$	$1.12m_{c}$
charm mass $m_c$	$1.43  {\rm GeV}$	$1.46  {\rm GeV}$

- > Raising the initial scale from the HERAPDF2.0 value ( $Q_0^2 = 1.9 \text{ GeV}^2$ ) to  $Q_0^2 = 2.56 \text{ GeV}^2$
- > FONLL scheme prefers  $m_c = 1.46 \text{ GeV}$  (while  $m_c^{HERA} = 1.43 \pm 0.06 \text{ GeV}$ )
- > The charm PDF must be generated perturbatively at a matching scale  $\mu_c > \mu_0 > m_c$  which needs to be larger than the default value  $\mu_c = m_c$
- > So  $\mu_c = 1.12 m_c$  (adopted also in Eur. Phys. J. C78 (2018) 621)

#### From TR to FONLL

Contribution to $\chi^2$	Old parametrization	New parametrization
subset NC $e^+$ 920 $\tilde{\chi}^2/{\rm n.d.p.}$	451/377	406/377
subset NC $e^+$ 820 $\tilde{\chi}^2/\text{n.d.p.}$	68/70	74/70
subset NC $e^+$ 575 $\tilde{\chi}^2$ /n.d.p.	220/254	222/254
subset NC $e^+$ 460 $\tilde{\chi}^2/\text{n.d.p.}$	218/204	225/204
subset NC $e^- \tilde{\chi}^2$ /n.d.p.	215/159	217/159
subset CC $e^+ \tilde{\chi}^2$ /n.d.p.	44/39	37/39
subset CC $e^- \tilde{\chi}^2$ /n.d.p.	57/42	50/42
correlation term $+ \log$ term	100 + 15	79 + 2
Total $\chi^2$ /d.o.f.	1388/1131	1312/1127

> Worse than the  $\chi^2$  presented in Slide 4 but:

- $\succ$  Old parameterisation  $\rightarrow$  deterioration of **25 units** wrt TR
- $\succ$  New parameterisation  $\rightarrow$  deterioration of just 11 units wrt TR
- PDFs largely compatible (backup)

### Stability of our fit

 $\succ$  First, we consider variations of the fit scale:

- >  $\mu_0 = 1.38 \text{ GeV}$  and  $\mu_c/m_c = 1.12 (\mu_c = 1.46 \text{ GeV}) \text{Down variation}$
- $\rightarrow$   $\mu_0$  = 1.60 GeV and  $\mu_c/m_c$  = 1.27 ( $\mu_c$  = 1.85 GeV) Intermediate step
- >  $\mu_0$  = 1.84 GeV and  $\mu_c/m_c$  = 1.27 ( $\mu_c$  = 1.85 GeV) Up variation
- Strange fraction variations:

>  $f_s = 0.5$  (up variation) and  $f_s = 0.3$  (down variation) – same as HERAPDF2.0

- Parametrization uncertainties addressed adding or removing parameters that do not change the fit quality. The ones giving the largest effect are:
  - > Adding  $F_{d_v}$
  - > Adding  $D_g$  (more flexibility at large-x)
  - > Adding  $H_g$  and removing  $F_g$  (possible effect at small-x)
- > The effect of all these variations is summarised in the next slide



- > The addition of the log term to  $d_v$  has the largest effect (negative for  $x \leq 10^{-3}$ )
- When D<sub>g</sub> is activated, large-x shape changes substantially, but in a region where the gluon PDF is very small and largely unconstrained
- > Effect of  $H_g$  (without  $F_g$ ) very mild
- Up/down variations of f<sub>s</sub> have a larger effect on the strange PDF (as expected)
- $\mu_0$  variations have small effects



#### Francesco Giuli - francesco.giuli@cern.ch

11

#### 18/03/2019



We combined all the uncertainties (exp+th+model) into a (symmetric) uncertainty band

Our final PDF set including the full uncertainty is largely compatible with NNPDF30

### Including small-x resummation

➢ Done using HELL (∨3.0 for the first time) – <u>1805.08785</u>, <u>1805.06460</u>, <u>1708.07510</u>

It provides resummed contributions to the DGLAP splitting functions, the heavy quark matching conditions and the DIS coefficient functions at NLLx

Contribution to $\chi^2$	HELL3.0 (NLL)	HELL3.0 $(LL')$	Hell2.0 $(LL')$
subset NC $e^+$ 920 $\tilde{\chi}^2$ /n.d.p.	402/377	403/377	403/377
subset NC $e^+$ 820 $\tilde{\chi}^2$ /n.d.p.	70/70	69/70	69/70
subset NC $e^+$ 575 $\tilde{\chi}^2$ /n.d.p.	219/254	219/254	218/254
subset NC $e^+$ 460 $\tilde{\chi}^2$ /n.d.p.	223/204	224/204	224/204
subset NC $e^- \tilde{\chi}^2/\text{n.d.p.}$	219/159	220/159	220/159
subset CC $e^+ \tilde{\chi}^2/\text{n.d.p.}$	38/39	38/39	38/39
subset CC $e^- \tilde{\chi}^2/\text{n.d.p.}$	49/42	49/42	49/42
correlation term $+ \log$ term	73 - 7	72 - 11	72 - 10
Total $\chi^2/{ m d.o.f.}$	1284/1127	1283/1127	1283/1127

If compared to NNLO fit, further reduction of ~30 units in  $\chi^2$ 

### Including small-x resummation

- The difference between two versions of HELL v3.0 is the introduction of a new default treatment of subleading logarithmic contributions
- > HELL v2.0 is the previous version of the HELL code



### Including small-x resummation

- The difference between two versions of HELL v3.0 is the introduction of a new default treatment of subleading logarithmic contributions
- > HELL v2.0 is the previous version of the HELL code



### **Conclusion and outlook**

- Paper announced on arXiv and submitted to EPJC <u>1902.11125</u>
- We proposed a new simple parametrization for the PDFs at the initial scale that includes a low degree polynominal in log(x) – more flexibility at low-x
- > Improvement of the fit quality (62 units reduction in  $\chi^2$  wrt HERAPDF2.0)
- > Accomplished using **18 parameters** (only 4 more than HERAPDF2.0)
- Stability of our fit tested upon several model and parametrization variations -> results very robust
- Flexibility of our parametrization allows for a more reliable determination of the uncertainties
- > The impact of supplementing the theoretical predictions with the resummation of small-x logarithms investigated  $\rightarrow$  gain of ~30 units in  $\chi^2$

#### > ADD THIS NEW PARAMETRIZATION IN THE NEW XFITTER RELEASE

LHAPDF grids available upon request

# **Backup Slides**

#### 18/03/2019

#### Francesco Giuli - francesco.giuli@cern.ch

17



Some differences are manifest (gluon/sea quarks)

> 1 $\sigma$  bands overlap or are very close to each other (apart from  $\overline{d} + \overline{s}$ )

## More sensitivity to the gluon PDF

- We also studied the inclusion of HERA Charm combined data (Eur.Phys.J. C78 (2018) no.6, 473)
- These data are directly sensitive to xg(x, Q<sup>2</sup>)





It is remarkable that the two FO fits are in agreement within uncertainties

#### A NLO fit

- We also tried a NLO fit (using both TR and FONLL-B) – preliminary
- FONLL-B provides a better description than TR
- At low-x, same structure in the gluon PDF





## Reduced cross section, F<sub>2</sub> and F<sub>L</sub>



- FO and resumed calculations very similar for the reduced xsec
- NNLO prediction for F<sub>2</sub> decreases at small x (softer gluon and quark singlet), while it rises steadily at resumed level (larger singlet)
- ➢ Resummed F<sub>L</sub> is quite flat in x and much larger than the NNLO one (x ≤ 10<sup>-3</sup>)
- Rise of F<sub>L</sub> due to the gluon PDF shape (rising for x ~ 10<sup>-4</sup>)

-01	

	$B_g$	$C_g$	$F_{g}$	$G_g$	$B_{u_v}$	$C_{u_v}$	$E_{u_v}$	$F_{u_v}$	$G_{u_v}$	$B_{dv}$	$C_{d_v}$	$C_{u}$	$D_u$	$A_d$	$B_d$	$C_{d}$	$D_d$	$F_d$
$B_{g}$	1.000	0.783	-0.508	-0.465	-0.055	0.055	0.074	-0.094	-0.098	0.000	-0.058	-0.176	0.043	-0.457	-0.525	-0.285		
$C_{g}$	0.783	1.000	-0.093	-0.070	-0.014	0.038	0.046	-0.100	-0.091	0.061	-0.061	-0.163	0.036	-0.352	-0.383	-0.345		
$F_{g}$	-0.508	-0.093	1.000	0.989	-0.072	0.117	0.142	-0.150	-0.119	0.051	0.063	0.011	-0.183	0.422	0.494	0.125		
$G_{g}$	-0.465	-0.070	0.989	1.000	-0.075	0.124	0.149	-0.157	-0.125	0.051	0.061	0.016	-0.218	0.488	0.558	0.110		
$B_{u_v}$	-0.055	-0.014	-0.072	-0.075	1.000	-0.202	-0.598	0.485	0.897	-0.226	-0.197	-0.127	-0.244	0.126	0.050	-0.634		
$C_{u_v}$	0.055	0.038	0.117	0.124	-0.202	1.000	0.846	-0.616	-0.381	-0.042	0.030	-0.535	-0.521	0.211	0.178	0.315		
$E_{uv}$	0.074	0.046	0.142	0.149	-0.598	0.846	1.000	-0.871	-0.777	0.184	0.248	-0.462	-0.443	0.164	0.157	0.646		
$F_{uv}$	-0.094	-0.100	-0.150	-0.157	0.485	-0.616	-0.871	1.000	0.806	-0.409	-0.445	0.356	0.523	-0.240	-0.206	-0.673		
$G_{u_v}$	-0.098	-0.091	-0.119	-0.125	0.897	-0.381	-0.777	0.806	1.000	-0.402	-0.384	0.002	0.048	-0.031	-0.064	-0.730		
$B_{d_v}$	0.000	0.061	0.051	0.051	-0.226	-0.042	0.184	-0.409	-0.402	1.000	0.940	0.390	0.133	0.075	0.069	0.383		
$C_{d_v}$	-0.058	-0.061	0.063	0.061	-0.197	0.030	0.248	-0.445	-0.384	0.940	1.000	0.262	0.013	0.123	0.112	0.437		
$C_{u}$	-0.176	-0.163	0.011	0.016	-0.127	-0.535	-0.462	0.356	0.002	0.390	0.262	1.000	0.721	0.005	0.056	-0.126		
$D_{a}$	0.043	0.036	-0.183	-0.218	-0.244	-0.521	-0.443	0.523	0.048	0.133	0.013	0.721	1.000	-0.595	-0.517	-0.083		
$A_{\bar{d}}$	-0.457	-0.352	0.422	0.488	0.126	0.211	0.164	-0.240	-0.031	0.075	0.123	0.005	-0.595	1.000	0.986	0.078		
$B_d$	-0.525	-0.383	0.494	0.558	0.050	0.178	0.157	-0.206	-0.064	0.069	0.112	0.056	-0.517	0.986	1.000	0.122		
$C_d$	-0.285	-0.345	0.125	0.110	-0.634	0.315	0.646	-0.673	-0.730	0.383	0.437	-0.126	-0.083	0.078	0.122	1.000		
$D_{\bar{d}}$	-0.042	0.022	0.029	-0.010	-0.665	0.241	0.571	-0.575	-0.706	0.188	0.215	-0.317	-0.048	-0.304	-0.252	0.752	1.000	
$F_{\bar{d}}$	0.562	0.363	-0.590	-0.652	0.013	-0.142	-0.142	0.166	0.086	-0.061	-0.095	-0.094	0.429	-0.941	-0.983	-0.147	0.200	1.000

- Correlation matrix between fit parameters
- Most of them are poorly correlated
- When present, F and G parameters strongly correlated (they probe the same x regime)
- > They are also correlated to B parameters (same reason as above)
- > Down-valence parameters highly correlated (same as for HERAPDF2.0)

#### Local minima

- While fitting data with fixed-order theory, we found a local minimum pretty far away from the global minimum presented in the paper
- > Main difference in the gluon PDF: global minimum with  $B_g < 0$ , while local minimum with  $B_g > 0$
- > The fit converged in the local minimum has an extra parameter in it: cubic logarithmic term in the gluon PDF ( $H_g$ )
- > Even though very significant differences in some parameters,  $\chi^2$  really similar
- > When transitioning from one minimum to the other in the parameter space, the  $\chi^2$  becomes much larger  $\rightarrow$  with a standard minimization routine it is highly unlikely that once the local minimum is found, it could converge to the global minimum
- > The physical expectation  $B_g < 0$  was crucial to guide us

23

				gluon
				NNLO fit, global min
				4
Ditted	NNLO (DONLL)	NNLO (DONLL)	NNLO INLL	Q <sup>2</sup> -3 GeV <sup>2</sup>
Fitted	NNLO (FONLL)	NNLO (FONLL)	NNLO+NLLx	3
parameter	local minimum	global minimum	HELL 3.0 (NLL)	
$B_{g}$	$0.34\pm0.07$	$-0.55\pm0.03$	$-0.52\pm0.04$	§ 1
$C_{g}$	$8.8 \pm 1.0$	$4.5\pm0.5$	$4.5\pm0.5$	× 1-
$F_{g}$	$0.76\pm0.04$	$0.230 \pm 0.003$	$0.217 \pm 0.005$	
$G_{g}$	$0.22\pm0.02$	$0.0131 \pm 0.0004$	$0.0112 \pm 0.0005$	0
$H_{g}$	$0.017 \pm 0.002$			
$B_{u_v}$	$0.85\pm0.06$	$0.83\pm0.06$	$0.76\pm0.06$	1
$C_{u_v}$	$4.5\pm0.1$	$4.6 \pm 0.2$	$4.6 \pm 0.1$	
$E_{uv}$	$1.7\pm0.8$	$1.9 \pm 1.0$	$2.6\pm1.1$	10 <sup>-4</sup> 10 <sup>-3</sup> 10 <sup>-2</sup> 10 <sup>-1</sup> 10 <sup>0</sup>
$F_{u_v}$	$0.38\pm0.04$	$0.37\pm0.05$	$0.35\pm0.04$	
$G_{u_v}$	$0.062\pm0.011$	$0.058 \pm 0.012$	$0.049 \pm 0.010$	NNLO fit, local min
$B_{d_v}$	$1.01\pm0.09$	$0.98\pm0.10$	$0.99\pm0.09$	15
$C_{d_v}$	$4.7\pm0.4$	$4.7\pm0.5$	$4.7\pm0.5$	$Q^2 = 3 \text{ GeV}^2$
$A_{\bar{d}}$	$0.070\pm0.008$	$0.13 \pm 0.02$	$0.14\pm0.02$	
$B_{ar{d}}$	$-0.45\pm0.02$	$-0.34\pm0.02$	$-0.33\pm0.02$	
$C_{ar{d}}$	$28 \pm 3$	$24\pm2$	$24\pm3$	***()
$D_{ar{d}}$	$76 \pm 17$	$40 \pm 12$	$38 \pm 10$	1 0.5 - · · · · · · · · · · · · · · · · · ·
$F_{ar{d}}$	$0.084 \pm 0.001$	$0.072 \pm 0.004$	$0.071 \pm 0.004$	45 X
$C_{ar{u}}$	$11 \pm 1$	$11 \pm 1$	$11 \pm 1$	
$D_{ar{u}}$	$33 \pm 6$	$20 \pm 4$	$18 \pm 4$	0
$\chi^2$ /d.o.f.	1314/1126	1312/1127	1284/1127	
•				-0.5 10 <sup>-4</sup> 10 <sup>-3</sup> 10 <sup>-2</sup> 10 <sup>-1</sup> 10 <sup>0</sup>