

Asymptotically Safe Extensions of the MSSM

[Work in progress in collaboration with Gudrun Hiller and Daniel Litim]

DESY Theory Workshop 2019
Quantum field theory meets gravity

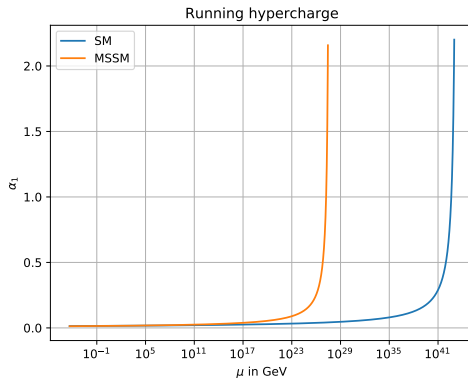
Kevin Moch

TU Dortmund

25.09.2019

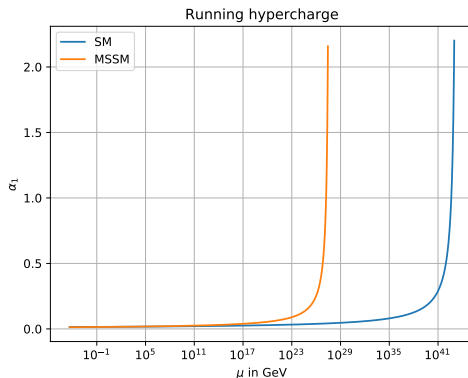
Motivation

- The Standard Model (SM) is not fundamental



Motivation

- The Standard Model (SM) is not fundamental

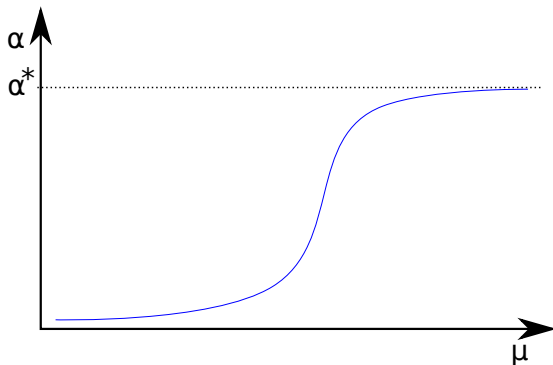


- Not only a perturbative artifact

[H. Gies and J. Jaeckel, 2004]

Motivation

- The Standard Model (SM) is not fundamental
- Asymptotically safe (AS) models are physical at all energies



Motivation

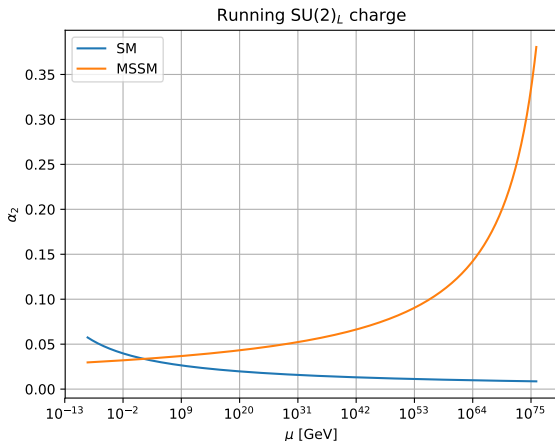
- The Standard Model (SM) is not fundamental
- Asymptotically safe (AS) models are physical at all energies
- The SM can be extended at $\sim \mathcal{O}(\text{TeV})$ to become AS without gravity
[A. Bond, G. Hiller, K. Kowalska, D. Litim, 2017]

Motivation

- The Standard Model (SM) is not fundamental
- Asymptotically safe (AS) models are physical at all energies
- The SM can be extended at $\sim \mathcal{O}(\text{TeV})$ to become AS without gravity
[A. Bond, G. Hiller, K. Kowalska, D. Litim, 2017]
- Also possible for the minimal supersymmetric SM (MSSM)?

Motivation

- The Standard Model (SM) is not fundamental
- Asymptotically safe (AS) models are physical at all energies
- The SM can be extended at $\sim \mathcal{O}(\text{TeV})$ to become AS without gravity
[A. Bond, G. Hiller, K. Kowalska, D. Litim, 2017]
- Also possible for the minimal supersymmetric SM (MSSM)?
- In the MSSM, also the $SU(2)_L$ coupling has a Landau pole



The MSSM is not AS, with or without R -parity violating terms.

The goal

Find an asymptotically safe MSSM extension (without gravity)

The goal

Find an asymptotically safe MSSM extension (without gravity)

Yukawa interactions are necessary for AS within perturbation theory

[A. Bond, D. Litim, 2016]

The goal

Find an asymptotically safe MSSM extension (without gravity)

Yukawa interactions are necessary for AS within perturbation theory

[A. Bond, D. Litim, 2016]

⇒ Focus on gauge-Yukawa models

The goal

Find an asymptotically safe MSSM extension (without gravity)

Yukawa interactions are necessary for AS within perturbation theory

[A. Bond, D. Litim, 2016]

⇒ Focus on gauge-Yukawa models

It was argued that perturbatively, AS SUSY models do not exist

[S. Martin, J. Wells, 2000]

The goal

Find an asymptotically safe MSSM extension (without gravity)

Yukawa interactions are necessary for AS within perturbation theory

[A. Bond, D. Litim, 2016]

⇒ Focus on gauge-Yukawa models

It was argued that perturbatively, AS SUSY models do not exist

[S. Martin, J. Wells, 2000]

Loophole: Semi-simple gauge groups → AS SUSY models found
(MSSM not yet included)

[A. Bond, D. Litim, 2017]

Framework

- $\mathcal{N} = 1$ supersymmetry in 4D

Framework

- $\mathcal{N} = 1$ supersymmetry in 4D
- Gauge-Yukawa models, superpotential $W = \frac{1}{6} Y^{ijk} \psi_i \psi_j \psi_k$

Framework

- $\mathcal{N} = 1$ supersymmetry in 4D
- Gauge-Yukawa models, superpotential $W = \frac{1}{6} Y^{ijk} \psi_i \psi_j \psi_k$
- Gauge anomaly absence

Framework

- $\mathcal{N} = 1$ supersymmetry in 4D
- Gauge-Yukawa models, superpotential $W = \frac{1}{6} Y^{ijk} \psi_i \psi_j \psi_k$
- Gauge anomaly absence
- Beta functions at 2-loop gauge and 1-loop Yukawa level

[S. Martin, M. Vaughn, 1994]

Framework

- $\mathcal{N} = 1$ supersymmetry in 4D
- Gauge-Yukawa models, superpotential $W = \frac{1}{6} Y^{ijk} \psi_i \psi_j \psi_k$
- Gauge anomaly absence
- Beta functions at 2-loop gauge and 1-loop Yukawa level

[S. Martin, M. Vaughn, 1994]

- Define couplings $\alpha_i = \frac{g_i^2}{(4\pi)^2}$, $\alpha_{Y^{ijk}} = \frac{|Y^{ijk}|^2}{(4\pi)^2}$, $\alpha = (\alpha_i, \alpha_{Y^{ijk}})$

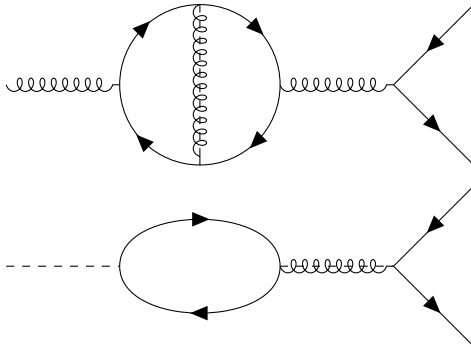
Framework

- $\mathcal{N} = 1$ supersymmetry in 4D
- Gauge-Yukawa models, superpotential $W = \frac{1}{6} Y^{ijk} \psi_i \psi_j \psi_k$
- Gauge anomaly absence
- Beta functions at 2-loop gauge and 1-loop Yukawa level
[S. Martin, M. Vaughn, 1994]
- Define couplings $\alpha_i = \frac{g_i^2}{(4\pi)^2}$, $\alpha_{Y^{ijk}} = \frac{|Y^{ijk}|^2}{(4\pi)^2}$, $\alpha = (\alpha_i, \alpha_{Y^{ijk}})$
- Fixed point (FP) α^* with $\beta_i(\alpha^*) = 0$ physical and perturbative:

$$0 \leq \alpha_i^* < 1 .$$

$$\beta_i = \alpha_i^2 \left[-B_i + C_i \alpha_i + \sum_{j \neq i} \underline{C_{ij}} \alpha_j - \sum_m \underline{D_{im}} \alpha_{y_m} \right],$$

$$\beta_{y_m} = \alpha_{y_m} \left[\underline{E_m} \alpha_{y_m} + \sum_{n \neq m} \underline{E_{mn}} \alpha_{y_n} - \sum_i \underline{F_{mi}} \alpha_i \right].$$



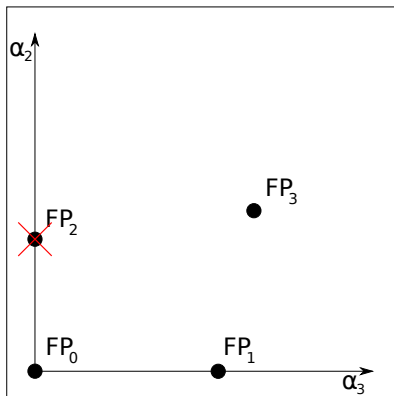
$$\beta_i = \alpha_i^2 \left[-\textcircled{B_i} + C_i \alpha_i + \sum_{j \neq i} \underline{C_{ij}} \alpha_j - \sum_m \underline{D_{im}} \alpha_{y_m} \right],$$

$$\beta_{y_m} = \alpha_{y_m} \left[\underline{E_m} \alpha_{y_m} + \sum_{n \neq m} \underline{E_{mn}} \alpha_{y_n} - \sum_i \underline{F_{mi}} \alpha_i \right].$$

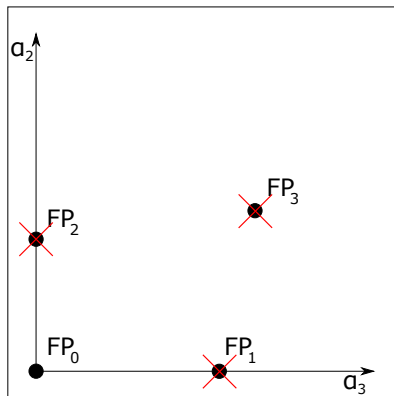
More particles: $\textcircled{B_i} \searrow$, $C_i \nearrow$

Gauge group $SU(3)_C \otimes SU(2)_L$

$B_3 > 0, B_2 < 0$:



$B_3 < 0, B_2 < 0$:



$$\mathrm{SU}(3)_C \otimes \mathrm{SU}(2)_L \otimes \mathrm{U}(1)_Y$$
$$B_3 > 0, B_2 < 0, B_1 < 0.$$

$$\text{SU}(3)_C \otimes \text{SU}(2)_L \otimes \text{U}(1)_Y$$

$$B_3 > 0, B_2 < 0, B_1 < 0.$$

For B_3 to stay positive, we may only add the colored fields

- 1) One $\mathbf{3}$ and one $\bar{\mathbf{3}}$,
- 2) Two $\mathbf{3}$ and two $\bar{\mathbf{3}}$.

MSSM extensions

The four possible candidates for AS MSSM extensions are

- 1) MSSM + two quark singlets

MSSM extensions

The four possible candidates for AS MSSM extensions are

- 1) MSSM + two quark singlets
- 2) MSSM + four quark singlets

MSSM extensions

The four possible candidates for AS MSSM extensions are

- 1) MSSM + two quark singlets
- 2) MSSM + four quark singlets
- 3) MSSM + two quark doublets

MSSM extensions

The four possible candidates for AS MSSM extensions are

- 1) MSSM + two quark singlets
- 2) MSSM + four quark singlets
- 3) MSSM + two quark doublets
- 4) MSSM + 4th generation (1 quark doublet and 2 quark singlets)

MSSM extensions

The four possible candidates for AS MSSM extensions are

- 1) **MSSM + two quark singlets** \leftarrow AS models found in scans
- 2) **MSSM + four quark singlets** \leftarrow No AS models found in scans
- 3) **MSSM + two quark doublets** \leftarrow No AS models found in scans
- 4) **MSSM + 4th generation** \leftarrow No AS models found in scans

On UV/IR attractive directions

$$\beta_i(g)|_{\alpha^*} \approx \sum_j M_{ij}(\alpha_j - \alpha_j^*) ,$$

with stability matrix

$$M_{ij} = \left. \frac{\partial \beta_i}{\partial \alpha_j} \right|_{\alpha^*} .$$

On UV/IR attractive directions

$$\beta_i(g)|_{\alpha^*} \approx \sum_j M_{ij}(\alpha_j - \alpha_j^*) ,$$

with stability matrix

$$M_{ij} = \left. \frac{\partial \beta_i}{\partial \alpha_j} \right|_{\alpha^*} .$$

Negative/Positive eigenvalue of $M \leftrightarrow$ UV/IR attractive.

On UV/IR attractive directions

$$\beta_i(g)|_{\alpha^*} \approx \sum_j M_{ij}(\alpha_j - \alpha_j^*) ,$$

with stability matrix

$$M_{ij} = \left. \frac{\partial \beta_i}{\partial \alpha_j} \right|_{\alpha^*} .$$

Negative/Positive eigenvalue of $M \leftrightarrow$ UV/IR attractive.

If e.g. $\alpha_2^* = 0$:

$$\beta_2|_{\alpha^*} = -B_{2,\text{eff}}^{\alpha^*} \alpha_2^2 + \mathcal{O}(\alpha_2^3) , \quad B_{2,\text{eff}}^{\alpha^*} = B_2 - C'_{23} \alpha_3^* - C'_{21} \alpha_1^* .$$

On UV/IR attractive directions

$$\beta_i(g)|_{\alpha^*} \approx \sum_j M_{ij}(\alpha_j - \alpha_j^*) ,$$

with stability matrix

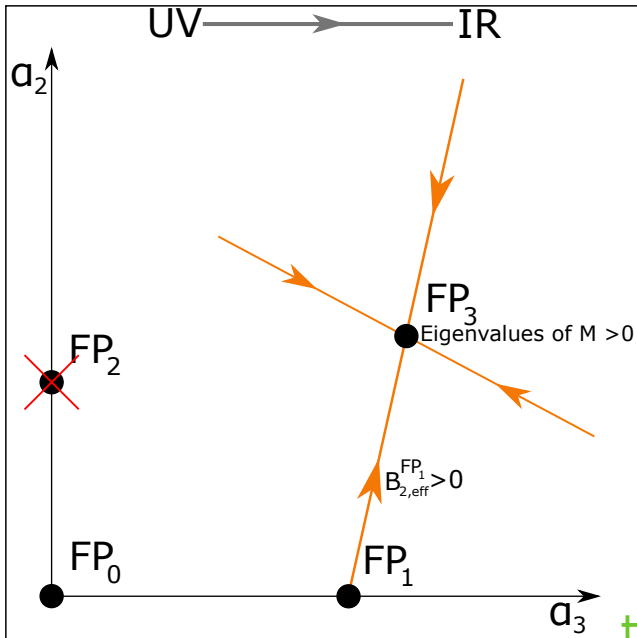
$$M_{ij} = \left. \frac{\partial \beta_i}{\partial \alpha_j} \right|_{\alpha^*} .$$

Negative/Positive eigenvalue of $M \leftrightarrow$ UV/IR attractive.

If e.g. $\alpha_2^* = 0$:

$$\beta_2|_{\alpha^*} = -B_{2,\text{eff}}^{\alpha^*} \alpha_2^2 + \mathcal{O}(\alpha_2^3) , \quad B_{2,\text{eff}}^{\alpha^*} = B_2 - C'_{23} \alpha_3^* - C'_{21} \alpha_1^* .$$

Negative/Positive $B_{2,\text{eff}}^{\alpha^*} \leftrightarrow$ IR/UV attractive.



MSSM + two quark singlets

Scanning $\sim 3.600.000$ models yields 281 AS models with FP_1 UV, FP_3 IR.

MSSM + two quark singlets

Scanning $\sim 3.600.000$ models yields 281 AS models with FP_1 UV, FP_3 IR.

AS Example:

Superfield	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	Multiplicity
quark doublet Q	3	2	$+\frac{1}{6}$	3
up-quark \bar{u}	$\bar{3}$	1	$-\frac{2}{3}$	3
down-quark \bar{d}	$\bar{3}$	1	$+\frac{1}{3}$	3
lepton doublet L	1	2	$-\frac{1}{2}$	3
lepton singlet \bar{e}	1	1	$+1$	3
up-Higgs H_u	1	2	$+\frac{1}{2}$	1
down-Higgs H_d	1	2	$-\frac{1}{2}$	1
BSM quark \bar{d}_4	$\bar{3}$	1	$+\frac{1}{3}$	1
BSM anti-quark d_4	3	1	$-\frac{1}{3}$	1
BSM lepton doublet $L_{4,5}$	1	2	$-\frac{1}{2}$	2
BSM anti-lepton doublet $\bar{L}_{1,2}$	1	$\bar{2}$	$+\frac{1}{2}$	2

MSSM + two quark singlets

Scanning $\sim 3.600.000$ models yields 281 AS models with FP_1 UV, FP_3 IR.

AS Example:

Superfield	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	Multiplicity
quark doublet Q	3	2	$+\frac{1}{6}$	3
up-quark \bar{u}	$\bar{3}$	1	$-\frac{2}{3}$	3
down-quark \bar{d}	$\bar{3}$	1	$+\frac{1}{3}$	3
lepton doublet L	1	2	$-\frac{1}{2}$	3
lepton singlet \bar{e}	1	1	$+\frac{1}{2}$	3
up-Higgs H_u	1	2	$+\frac{1}{2}$	1
down-Higgs H_d	1	2	$-\frac{1}{2}$	1
BSM quark \bar{d}_4	$\bar{3}$	1	$+\frac{1}{3}$	1
BSM anti-quark d_4	3	1	$-\frac{1}{3}$	1
BSM lepton doublet $L_{4,5}$	1	2	$-\frac{1}{2}$	2
BSM anti-lepton doublet $\bar{L}_{1,2}$	1	$\bar{2}$	$+\frac{1}{2}$	2

$$\begin{aligned}
 W = & y_1 \bar{d}_4 Q_1 L_1 + y_2 \bar{d}_4 Q_2 L_2 + y_3 \bar{d}_1 Q_1 L_4 + y_4 \bar{d}_2 Q_1 L_5 \\
 & + y_5 \bar{u}_2 Q_1 \bar{L}_1 + y_6 \bar{u}_1 Q_2 \bar{L}_2 + y_t \bar{u}_3 Q_3 H_u + y_b \bar{d}_3 Q_3 H_d
 \end{aligned}$$

MSSM + two quark singlets

Scanning $\sim 3.600.000$ models yields 281 AS models with FP_1 UV, FP_3 IR.

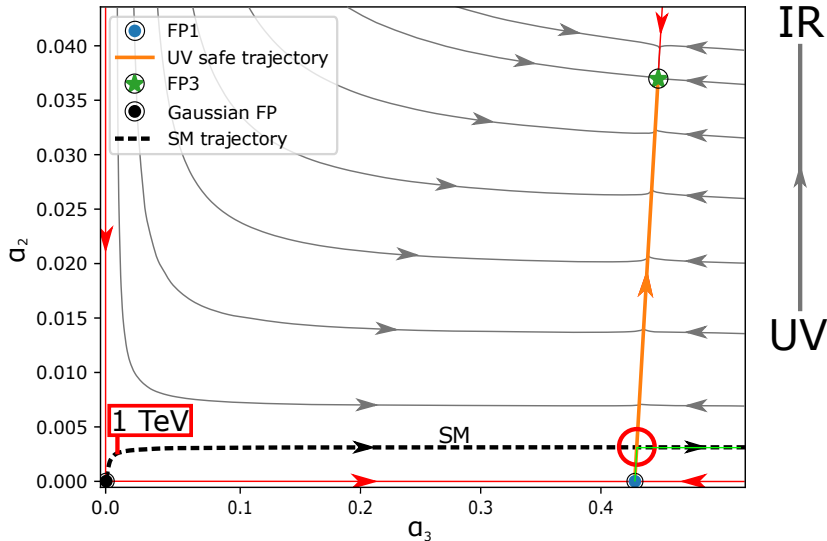
AS Example:

Superfield	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	Multiplicity
quark doublet Q	3	2	$+\frac{1}{6}$	3
up-quark \bar{u}	$\bar{3}$	1	$-\frac{2}{3}$	3
down-quark \bar{d}	$\bar{3}$	1	$+\frac{1}{3}$	3
lepton doublet L	1	2	$-\frac{1}{2}$	3
lepton singlet \bar{e}	1	1	$+\frac{1}{2}$	3
up-Higgs H_u	1	2	$+\frac{1}{2}$	1
down-Higgs H_d	1	2	$-\frac{1}{2}$	1
BSM quark \bar{d}_4	$\bar{3}$	1	$+\frac{1}{3}$	1
BSM anti-quark d_4	3	1	$-\frac{1}{3}$	1
BSM lepton doublet $L_{4,5}$	1	2	$-\frac{1}{2}$	2
BSM anti-lepton doublet $\bar{L}_{1,2}$	1	$\bar{2}$	$+\frac{1}{2}$	2

$$W = y_1 \bar{d}_4 Q_1 L_1 + y_2 \bar{d}_4 Q_2 L_2 + y_3 \bar{d}_1 Q_1 L_4 + y_4 \bar{d}_2 Q_1 L_5 \\ + y_5 \bar{u}_2 Q_1 \bar{L}_1 + y_6 \bar{u}_1 Q_2 \bar{L}_2 + y_t \bar{u}_3 Q_3 H_u + y_b \bar{d}_3 Q_3 H_d$$

→ R parity violation necessary!

General RG flow picture



MSSM + two quark singlets

$$3 \leq \# \text{ BSM Lepton doublets} \leq 13$$

MSSM + two quark singlets

$$3 \leq \# \text{ BSM Lepton doublets} \leq 13$$

$$\alpha_3^* > 0.43 \text{ in FP}_1$$

MSSM + two quark singlets

$$3 \leq \# \text{ BSM Lepton doublets} \leq 13$$

$$\alpha_3^* > 0.43 \text{ in FP}_1$$

\Rightarrow Matching onto the SM at $\mathcal{O}(\text{MeV})$.

MSSM + two quark singlets

$$3 \leq \# \text{ BSM Lepton doublets} \leq 13$$

$$\alpha_3^* > 0.43 \text{ in FP}_1$$

\Rightarrow Matching onto the SM at $\mathcal{O}(\text{MeV})$.

One can generally show that for all MSSM extensions

- $\bullet \text{FP}_1|_{\alpha_3} > 3/110 \approx 0.027$

MSSM + two quark singlets

$$3 \leq \# \text{ BSM Lepton doublets} \leq 13$$

$$\alpha_3^* > 0.43 \text{ in FP}_1 \\ \Rightarrow \text{Matching onto the SM at } \mathcal{O}(\text{MeV}).$$

One can generally show that for all MSSM extensions

- $\text{FP}_1|_{\alpha_3} > 3/110 \approx 0.027$
- $\text{FP}_3 \text{ exists} \Leftrightarrow \text{FP}_1 \text{ exists and UV.}$
 $\rightarrow \text{FP}_3|_{\alpha_3} > \text{FP}_1|_{\alpha_3}$

On the α_1 direction

All encountered FPs have $\alpha_1^* = 0$.

On the α_1 direction

All encountered FPs have $\alpha_1^* = 0$. Near such FPs the running of α_1 reads

$$\beta_1 = -B_{1,\text{eff}}\alpha_1^2 + \mathcal{O}(\alpha_1^3).$$

On the α_1 direction

All encountered FPs have $\alpha_1^* = 0$. Near such FPs the running of α_1 reads

$$\beta_1 = -B_{1,\text{eff}}\alpha_1^2 + \mathcal{O}(\alpha_1^3).$$

We always find $B_{1,\text{eff}} < 0$ for both FP_1 and FP_3
 $\Rightarrow \alpha_1 \equiv 0$ on UV-safe trajectories.

On the α_1 direction

All encountered FPs have $\alpha_1^* = 0$. Near such FPs the running of α_1 reads

$$\beta_1 = -B_{1,\text{eff}}\alpha_1^2 + \mathcal{O}(\alpha_1^3).$$

We always find $B_{1,\text{eff}} < 0$ for both FP_1 and FP_3
 $\Rightarrow \alpha_1 \equiv 0$ on UV-safe trajectories.

Matching onto the SM not
possible at all!

Non- perturbative checks

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

- Fixed point in SUSY = superconformal field theory (SCFT)

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

- Fixed point in SUSY = superconformal field theory (SCFT)
- Central charge a of CFT in the UV bigger than in the IR:

$$a_{UV} > a_{IR} \quad [\text{J. Cardy, 1988}]$$

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

- Fixed point in SUSY = superconformal field theory (SCFT)
- Central charge a of CFT in the UV bigger than in the IR:

$$a_{UV} > a_{IR} \quad [\text{J. Cardy, 1988}]$$

- Central charge a is a function of R -charges of global $U(1)_R$ -group:

$$a = a(R)$$

[D. Anselmi, D. Freedman, M. Grisaruc, A. Johansen, 1998]

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

- Fixed point in SUSY = superconformal field theory (SCFT)
- Central charge a of CFT in the UV bigger than in the IR:
$$a_{UV} > a_{IR} \quad [\text{J. Cardy, 1988}]$$
- Central charge a is a function of R -charges of global $U(1)_R$ -group:
$$a = a(R)$$

[D. Anselmi, D. Freedman, M. Grisaru, A. Johansen, 1998]
- R maximizes a in a SCFT [K. Intriligator, B. Wecht, 2003]

Are FP_1/FP_3 physical and UV/IR beyond perturbation theory?

- Fixed point in SUSY = superconformal field theory (SCFT)
- Central charge a of CFT in the UV bigger than in the IR:
$$a_{UV} > a_{IR} \quad [\text{J. Cardy, 1988}]$$
- Central charge a is a function of R -charges of global $U(1)_R$ -group:
$$a = a(R)$$

[D. Anselmi, D. Freedman, M. Grisaruc, A. Johansen, 1998]
- R maximizes a in a SCFT [K. Intriligator, B. Wecht, 2003]
- \rightarrow Benchmarks are in agreement with exact SCFT relations

Conclusion

- AS plus SUSY very restrictive

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks
- AS benchmarks in agreement with exact relations from SCFT

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks
- AS benchmarks in agreement with exact relations from SCFT
- Two challenges remain:
 - 1) Matching scale onto SM too low, $\mathcal{O}(\text{MeV})$

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks
- AS benchmarks in agreement with exact relations from SCFT
- Two challenges remain:
 - 1) Matching scale onto SM too low, $\mathcal{O}(\text{MeV})$
 - 2) For all models, UV-safe trajectories have $\alpha_1 = 0$

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks
- AS benchmarks in agreement with exact relations from SCFT
- Two challenges remain:
 - 1) Matching scale onto SM too low, $\mathcal{O}(\text{MeV})$
 - 2) For all models, UV-safe trajectories have $\alpha_1 = 0$
- Next step: Include non-abelian gauge factor to circumvent restrictions from the MSSM

Conclusion

- AS plus SUSY very restrictive
- AS MSSM extensions found → **exactly** two additional quarks
- AS benchmarks in agreement with exact relations from SCFT
- Two challenges remain:
 - 1) Matching scale onto SM too low, $\mathcal{O}(\text{MeV})$
 - 2) For all models, UV-safe trajectories have $\alpha_1 = 0$
- Next step: Include non-abelian gauge factor to circumvent restrictions from the MSSM

Stay tuned!

- [1] H. Gies and J. Jaeckel, Phys. Rev. Lett. **93** (2004) 110405 doi:10.1103/PhysRevLett.93.110405 [hep-ph/0405183].
- [2] A. D. Bond, G. Hiller, K. Kowalska and D. F. Litim, “Directions for model building from asymptotic safety,” JHEP **1708** (2017) 004 doi:10.1007/JHEP08(2017)004 [arXiv:1702.01727 [hep-ph]].
- [3] K. Kowalska, A. Bond, G. Hiller and D. Litim, “Towards an asymptotically safe completion of the Standard Model,” PoS EPS -HEP2017 (2017) 542. doi:10.22323/1.314.0542
- [4] A. D. Bond and D. F. Litim, “Theorems for Asymptotic Safety of Gauge Theories,” Eur. Phys. J. C **77** (2017) no.6, 429 Erratum: [Eur. Phys. J. C **77** (2017) no.8, 525] doi:10.1140/epjc/s10052-017-4976-5, 10.1140/epjc/s10052-017-5034-z [arXiv:1608.00519 [hep-th]].
- [5] S. P. Martin and J. D. Wells, Phys. Rev. D **64** (2001) 036010 doi:10.1103/PhysRevD.64.036010 [hep-ph/0011382].

- [6] A. D. Bond and D. F. Litim, “Asymptotic safety guaranteed in supersymmetry,” *Phys. Rev. Lett.* **119** (2017) no.21, 211601 doi:10.1103/PhysRevLett.119.211601 [arXiv:1709.06953 [hep-th]].
- [7] S. P. Martin and M. T. Vaughn, “Two loop renormalization group equations for soft supersymmetry breaking couplings,” *Phys. Rev. D* **50** (1994) 2282 Erratum: [*Phys. Rev. D* **78** (2008) 039903] doi:10.1103/PhysRevD.50.2282, 10.1103/PhysRevD.78.039903 [arXiv:9311340 [hep-ph]].
- [8] J. L. Cardy, “Is There a c Theorem in Four-Dimensions?,” *Phys. Lett. B* **215** (1988) 749. doi:10.1016/0370-2693(88)90054-8
- [9] D. Anselmi, D. Z. Freedman, M. T. Grisaru and A. A. Johansen, “Nonperturbative formulas for central functions of supersymmetric gauge theories,” *Nucl. Phys. B* **526** (1998) 543 doi:10.1016/S0550-3213(98)00278-8 [arXiv:9708042 [hep-th]].

- [10] K. A. Intriligator and B. Wecht, “The Exact superconformal R symmetry maximizes a ,” Nucl. Phys. B **667** (2003) 183
doi:10.1016/S0550-3213(03)00459-0 [arXiv:0304128 [hep-th]].