# Seiberg-like Dualities for 2d Gauge Theories with a Boundary 

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## Outline

Introduction to Gauged Linear Sigma Models (GLSMs)

## Dualities of Non-Abelian GLSMs

Extending Dualities to GLSMs with a Boundary

An Example

Summary and Outlook

## Introduction to GLSMs

GLSMs are $\mathcal{N}=(2,2)$ gauge theories in 2 d with $U(1)$ R-symmetry [Witten '93]

Gauge group $\mathrm{G}=\left(U(1)^{\ell} \times K\right) / \Gamma$
( $K$ : Product of semi-simple groups, $\Gamma$ : Discrete group)
Spectrum : 1. Chiral Field $(\Phi)$ in a repr. of $G$
2. Vector Field $(V)$ in the adjoint repr. of $G$

GLSM data can be modelled s.t. the RG flow to the infrared results in an $\mathcal{N}=(2,2)$ SCFT.

## —Dualities of Non-Abelian GLSMs

## Dualities of Non-Abelian GLSMs

Amongst certain non-abelian GLSMs, there exist non-trivial dualities [Hori,Tong '07], [Hori '13]

Example: $\circ \mathrm{GLSM}$ with $\mathrm{G}=\operatorname{SU}(k)$ and $N$ fundamentals $\Phi_{i}$
Gauge invariant d.o.f.:

$$
B_{i_{1} \ldots i_{k}}=\epsilon_{\alpha_{1} \ldots \alpha_{k}} \Phi_{i_{1}}^{\alpha_{1}} \ldots \Phi_{i_{k}}^{\alpha_{k}} \quad \#:\binom{N}{k}
$$

- $\widetilde{\text { GLSM }}$ with $\mathrm{G}=\mathrm{SU}(N-k)$ and $N$ fundamentals $\widetilde{\Phi}_{i}$ Gauge invariant d.o.f.:

$$
\widetilde{B}_{i_{1} \ldots i_{N-k}}=\epsilon_{\alpha_{1} \ldots \alpha_{N-k}} \widetilde{\Phi}_{i_{1}}^{\alpha_{1}} \ldots \widetilde{\Phi}_{i_{N-k}}^{\alpha_{k}} \quad \#:\binom{N}{N-k}
$$

Check: Sphere partition function of dual theories coincide
[Benini, Cremonesi '14]
Remarks: 1. Similar to $\mathcal{N}=14 d$ Seiberg Duality
2. Dualities also exist for $\mathrm{O}(k)$ and $\mathrm{USp}(2 k)$ gauge groups

## A 'Geometric' Duality

We consider an example of a duality motivated by the target spaces of the GLSMs being geometric duals:

$$
\operatorname{Gr}(k, N) \simeq \operatorname{Gr}(N-k, N)
$$

Recall: $\operatorname{Gr}(k, N)$ is the space of complex $k$-planes in an $N$-dimensional ambient space; $\operatorname{dim}_{\mathbb{C}}(\operatorname{Gr}(k, N))=k \cdot(N-k)$

The duality is understood by associating a complex $k$-plane in $\mathbb{C}^{N}$ to its orthogonal complement, i.e. an $(N-k)$-plane.

Duality: $\circ \mathrm{G}=\mathrm{U}(k)$ and $N$ flavours $\Phi_{i}$ in the repr. $\square_{1}$ Target Space: $\operatorname{Gr}(k, N)$

- $\mathrm{G}=\mathrm{U}(N-k)$ and $N$ flavours $\widetilde{\Phi}_{i}$ in the repr. $\square_{1}$

Target Space: $\operatorname{Gr}(N-k, N)$

## Dualities of GLSMs with a Boundary

Duality discussion so far has centred on GLSMs without boundary. What is the effect of duality on boundaries of GLSMs?


GLSM on a Riemann surface with a boundary generically preserves only half of the susy (A- or B-type).
A supersymmetric action is typically invariant under a susy transformation upto a total derivative term $\Delta$.
$\Delta$ non-vanishing for theory with a boundary $\Rightarrow$ 'Warner Problem'

## The Warner Problem

$$
\Delta=\delta_{\text {susy }} \cdot \int_{D} d^{2} z d^{2} \theta W(\Phi)=\int_{\partial D} d \sigma d \theta W\left(\left.\Phi\right|_{\partial D}\right)
$$

To restore susy a term is added at the boundary of the form:

$$
S_{\partial D}=-\int_{\partial D} d \sigma d \theta\left(\pi \cdot J\left(\left.\Phi\right|_{\partial D}\right)\right)
$$

where $\pi$ are fermions at the bdy that satisfy: $\delta_{\text {susy }} \cdot \pi=E\left(\left.\Phi\right|_{\partial D}\right)$ and $J, E$ are functions s.t. $J(\Phi) \cdot E(\Phi)=W(\Phi)$.
Then

$$
\delta_{\text {susy }} S_{\partial D}=-\int_{\partial D} d \sigma d \theta W\left(\left.\Phi\right|_{\partial D}\right)=-\Delta \cdot \checkmark \text { (susy restored) }
$$

The boundary d.o.f. can be represented as:

$$
Q(\Phi)=J(\Phi) \pi+E(\Phi) \bar{\pi}, \text { s.t. } Q^{2}(\Phi)=J(\Phi) \cdot E(\Phi)=W(\Phi)
$$

## Defects between GLSMs

Check of boundary duality: Compare Hemisphere Partition functions:

$$
Z_{H^{2}}\left(\operatorname{GLSM}_{Q}\right)=Z_{H^{2}}\left({\widetilde{\operatorname{GLSM}_{\widetilde{Q}}}}_{\tilde{L}}\right)
$$

[Hori,Romo '13]
[Honda, Okuda '13]
$\rightarrow$ We aim for a stronger check: Categories of B-type boundary conditions on dual GLSMs must be equivalent.
In doing so, we first separate the two GLSMs with a 'defect'.


Corresponding to the defect, we construct an explicit map FM,

$$
\text { FM : Bdy } 2 \rightarrow \text { Bdy } 1
$$

This map is known as the Fourier-Mukai Kernel.

## The Fourier-Mukai Kernel

Fourier-Mukai Kernel from theory B to $A$ is a map s.t.,

$$
\text { FM : Bdys of } B \rightarrow \text { Bdys of } A
$$

$$
\mathscr{B} \mapsto \mathscr{A}
$$

More explicitly, consider the product of dual boundary theories and respective projections to individual theories:


Then the map FM corresponds to an object $\varphi_{\text {FM }} \in \operatorname{Bdys}$ of $(A \times B)$,

$$
p_{1 *}\left(\varphi_{\mathrm{FM}} \otimes \mathrm{p}_{2}^{*}(\mathscr{B})\right)=\mathscr{A} .
$$

$\Rightarrow F M$ acts as an integration kernel that maps bdys from $B$ to $A_{厄}$

## The Fourier-Mukai Kernel II: Application

How does the narrative of defects between dual GLSMs connect to their boundary d.o.f.?
Recall, boundary d.o.f of a GLSM $\equiv Q(x)$, with $Q^{2}(x)=W(x)$.


Since the defect is just a boundary in the product theory:

$\Rightarrow$ Defect d.o.f. between $\mathrm{GLSM}_{1}$ and $\mathrm{GLSM}_{2} \equiv Q\left(x_{1}, x_{2}\right)$, with

$$
Q^{2}\left(x_{1}, x_{2}\right)=W\left(x_{1}\right)-W_{2}\left(x_{2}\right)
$$

## Example: Grassmannian Duality

For the Grassmannian duality, $\operatorname{Gr}(k, N) \simeq \operatorname{Gr}(N-k, N)$, we consider the duality transformation on the defect:


| Matter | $(U(k) \times U(k))$ Rep. |
| :---: | :---: |
| $x_{i}$ | $\left(\square_{1}, 1_{0}\right)$ |
| $y_{i}$ | $\left(1_{0}, \square_{1}\right)$ |
| $\mathcal{M}$ | $\left(\square_{1}, \square_{-1}\right)$ |
| $\pi$ | $\left(\square_{-1}, 1_{0}\right)$ |


| Matter | $(U(k) \times U(N-k))$ Rep. |
| :---: | :---: |
| $x_{i}$ | $\left(\square_{1}, 1_{0}\right)$ |
| $\widetilde{y}_{i}$ | $\left(1_{0}, \square_{1}\right)$ |
| $\pi^{\prime}$ | $\left(\square_{-1}, \square_{-1}\right)$ |

$Q_{1}=(x+y \mathcal{M}) \pi$

$$
\begin{aligned}
Q_{2} & =(x \cdot \tilde{y}) \pi^{\prime} \\
Q_{2}^{2} & =0
\end{aligned}
$$

Categories of bdy conditions are shown to be equivalent.

## Summary and Outlook

Summary
GLSMs and Seiberg-like dualities thereof
Extension of dualities to GLSMs w/ boundary: Fourier-Mukai Kernel
Outlook
Analyse dualities in quiver gauge theories leading to cluster algebras. Strengthen the cluster algebra proposal by an extension to theories with a boundary.
[Benini et al '14]
Extend dualities to boundary GLSMs with other gauge groups that are Seiberg-like-dual, eg. USp(2k).
Realise derived equivalences of target spaces relevant from a string compactification point of view.

## Thank You

## Questions/Comments?

