

Strange quark distributions in the ABMP16 fit and beyond

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sa, Blümlein, Moch, Plačakytė PRD 96, 014011 (2017)
sa, Blümlein, Moch PLB 777, 134 (2018)

PDF fit framework

QCD evolution

massless NNLO, massive NLO OMEs
(*OPENQCDRAD*)

3-flavour PDFs

5-flavour PDFs

DIS inclusive

NNLO
(*OPENQCDRAD*)

Power corr.
(TMC+high-twist)

DIS heavy quark

NNLO(approx.)
(*OPENQCDRAD*)

Drell-Yan (W,Z, γ)

NNLO
(*FEWZ-grids*)

t-quark

(*Hathor, fasttop*)

Data used and fit quality

| Experiment | Process | Reference | NDP | χ^2 |
|-----------------------------------|--------------------------------------|-----------|-------|----------|
| <i>DIS</i> | | | | |
| HERA I + II | $e^\pm p \rightarrow e^\pm X$ | [4] | 1168 | 1510 |
| | $e^\pm p \rightarrow \nu^{(-)} X$ | | | |
| BCDMS | $\mu^+ p \rightarrow \mu^+ X$ | [61] | 351 | 411 |
| NMC | $\mu^+ p \rightarrow \mu^+ X$ | [60] | 245 | 343 |
| SLAC-49a | $e^- p \rightarrow e^- X$ | [54,62] | 38 | 59 |
| SLAC-49b | $e^- p \rightarrow e^- X$ | [54,62] | 154 | 171 |
| SLAC-87 | $e^- p \rightarrow e^- X$ | [54,62] | 109 | 103 |
| SLAC-89b | $e^- p \rightarrow e^- X$ | [56,62] | 90 | 79 |
| <i>DIS heavy-quark production</i> | | | | |
| HERA I + II | $e^\pm p \rightarrow e^\pm cX$ | [63] | 52 | 62 |
| H1 | $e^\pm p \rightarrow e^\pm bX$ | [15] | 12 | 5 |
| ZEUS | $e^\pm p \rightarrow e^\pm bX$ | [16] | 17 | 16 |
| • CCFR | $\nu^{(-)} N \rightarrow \mu^\pm cX$ | [64] | 89 | 62 |
| • CHORUS | $\nu N \rightarrow \mu^+ cX$ | [18] | 6 | 7.6 |
| • NOMAD | $\nu N \rightarrow \mu^+ cX$ | [17] | 48 | 59 |
| • NuTeV | $\nu^{(-)} N \rightarrow \mu^\pm cX$ | [64] | 89 | 49 |
| <i>DY</i> | | | | |
| FNAL-605 | $pCu \rightarrow \mu^+ \mu^- X$ | [68] | 119 | 165 |
| FNAL-866 | $pp \rightarrow \mu^+ \mu^- X$ | [69] | 39 | 53 |
| | $pD \rightarrow \mu^+ \mu^- X$ | | | |
| <i>Top-quark production</i> | | | | |
| ATLAS, CMS | $pp \rightarrow tqX$ | [27–32] | 10 | 2.3 |
| CDF&DØ | $\bar{p}p \rightarrow tbX$ | [53] | 2 | 1.1 |
| | $\bar{p}p \rightarrow tqX$ | | | |
| ATLAS, CMS | $pp \rightarrow t\bar{t}X$ | [33–52] | 23 | 13 |
| CDF&DØ | $\bar{p}p \rightarrow t\bar{t}X$ | [53] | 1 | 0.2 |

Direct constraint on strangeness

Collider DY in the next slide
(indirect constraint on strangeness)

DY data in the ABMP16 fit

| Experiment | ATLAS | | CMS | | DØ | | LHCb | | | |
|-------------------------|---|---|--|--|--|--|---|-------------------------|---|------------|
| \sqrt{s} (TeV) | 7 | 13 | 7 | 8 | 1.96 | | 7 | 8 | | |
| Final states | $W^+ \rightarrow l^+ \nu$ $W^- \rightarrow l^- \nu$ $Z \rightarrow l^+ l^-$ | $W^+ \rightarrow l^+ \nu$ $W^- \rightarrow l^- \nu$ $Z \rightarrow l^+ l^-$ | $W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$ (asym) | $W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$ | $W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$ (asym) | $W^+ \rightarrow e^+ \nu$ $W^- \rightarrow e^- \nu$ (asym) | $W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$ $Z \rightarrow \mu^+ \mu^-$ | $Z \rightarrow e^+ e^-$ | $W^+ \rightarrow \mu^+ \nu$ $W^- \rightarrow \mu^- \nu$ $Z \rightarrow \mu^+ \mu^-$ | |
| Cut on the lepton P_T | $P_T^l > 20$ GeV | $P_T^e > 25$ GeV | $P_T^\mu > 25$ GeV | $P_T^\mu > 25$ GeV | $P_T^\mu > 25$ GeV | $P_T^e > 25$ GeV | $P_T^\mu > 20$ GeV | $P_T^e > 20$ GeV | $P_T^\mu > 20$ GeV | |
| Luminosity (1/fb) | 0.035 | 0.081 | 4.7 | 18.8 | 7.3 | 9.7 | 1 | 2 | 2.9 | |
| NDP | 30 | 6 | 11 | 22 | 10 | 13 | 31(33) ^a | 17 | 32(34) | |
| | ABMP16 | 31.0 | 9.2 | 22.4 | 16.5 | 17.6 | 19.0 | 45.1(54.4) | 21.7 | 40.0(59.2) |
| | CJ15 | – | – | – | – | 20 | 29 | – | – | – |
| | CT14 | 42 | – | – ^b | – | – | 34.7 | – | – | – |
| | HERAFitter | – | – | – | – | 13 | 19 | – | – | – |
| | MMHT16 | 39 ^c | – | – | 21 | 21 ^c | 26 | (43) | 29 | (59) |
| | NNPDF3.1 | 29 | – | 19 | – | 16 | 35 | (59) | 19 | (47) |

^a The values of NDP and χ^2 correspond to the unfiltered samples.

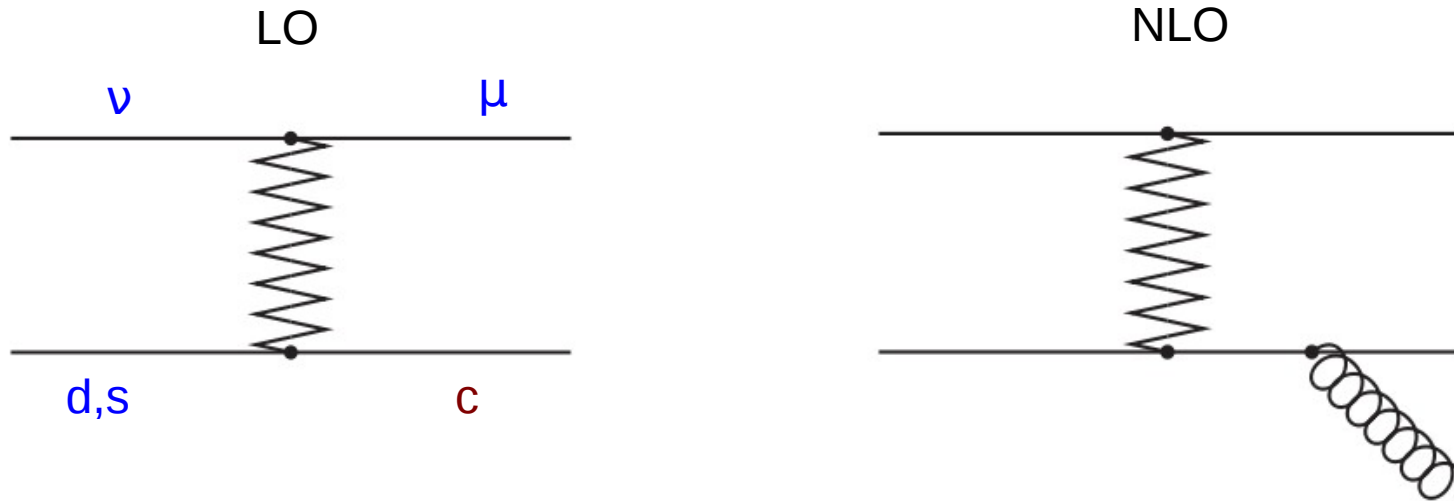
^b For the statistically less significant data with the cut of $P_T^\mu > 35$ GeV the value of $\chi^2 = 12.1$ was obtained.

^c The value obtained in MMHT14 fit.

| Experiment | NDP | χ^2 after the data sets excuded | | | | |
|------------|-----|--------------------------------------|-------|------|-------------|-------------|
| | | – | ATLAS | CMS | DØ | LHCb |
| ATLAS | 36 | 37.7 | – | 37.0 | 38.3 | 39.6 |
| CMS | 33 | 26.6 | 25.6 | – | 26.0 | 23.5 |
| DØ | 23 | 48.5 | 48.1 | 47.7 | – | 44.2 |
| LHCb | 80 | 98.2 | 100.2 | 97.4 | 78.8 | – |

- Good overall agreement in NNLO with some tension between DØ and LHCb data

Strange sea from the νN DIS



Two decay modes of c -quark are used: hadronic (emulsion experiments) and semi-leptonic (electronic experiments)

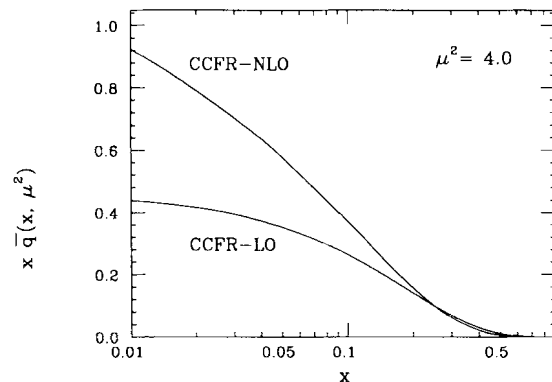


Fig. 3. The quark sea distribution $x \bar{q}(x, \mu^2 = 4.0 \text{ GeV}^2/c^2)$ determined at next-to-leading order and leading order

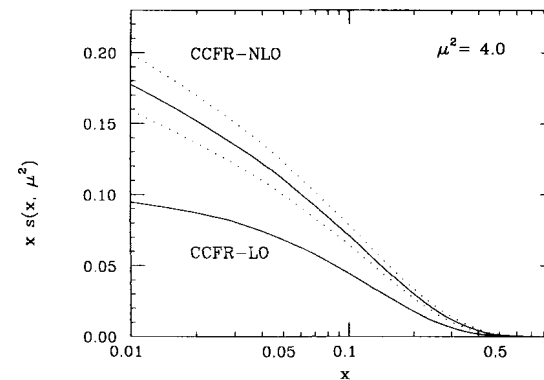
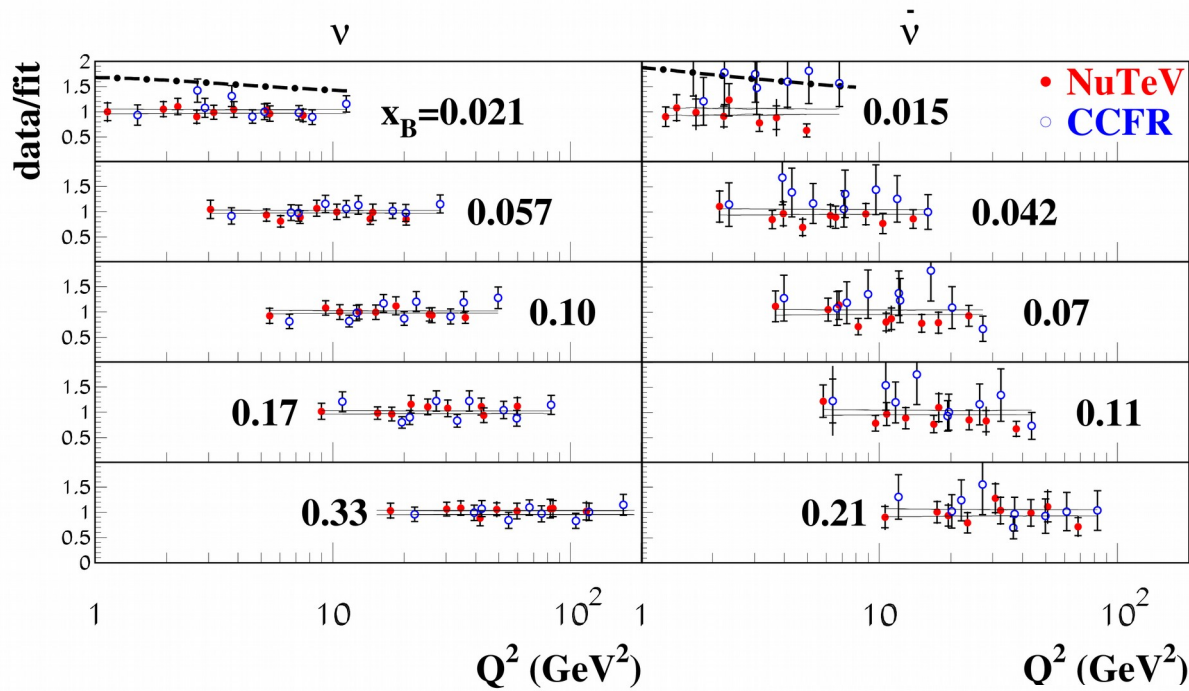


Fig. 4. The strange quark distribution $x s(x, \mu^2 = 4.0 \text{ GeV}^2/c^2)$ determined at next-to-leading order (described in section 4.1) and leading order. The band around the NLO curve indicates the $\pm 1\sigma$ uncertainty in the distribution

CCFR ZPC 65, 189 (1995)

Primary source for the strange sea was for a long time neutrino-induced charm production measured by CCFR/NuTeV at Fermilab preferring a suppression of ~ 0.5 w.r.t. non-strange sea

NuTeV/CCFR data in the PDF fit framework

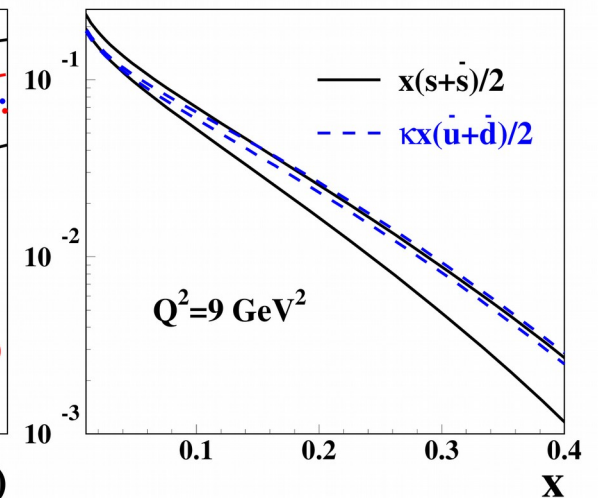
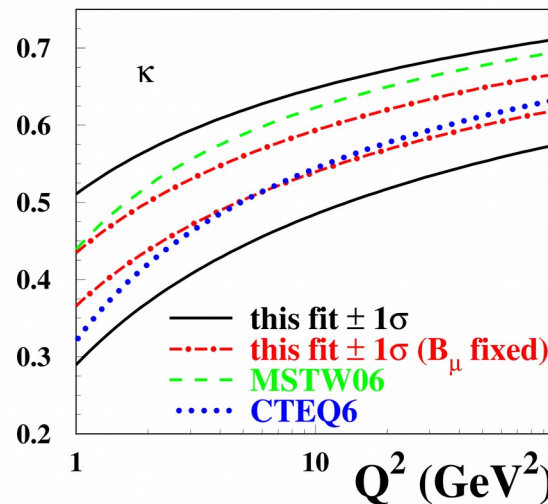


- CCFR and NuTeV are in a good agreement
- Charge asymmetry in the strange sea is consistent with 0 within uncertainties

sa, Kulagin, Petti PLB 675, 433 (2009)

$$\kappa_s(\mu^2) = \frac{\int_0^1 x[s(x, \mu^2) + \bar{s}(x, \mu^2)]dx}{\int_0^1 x[\bar{u}(x, \mu^2) + \bar{d}(x, \mu^2)]dx},$$

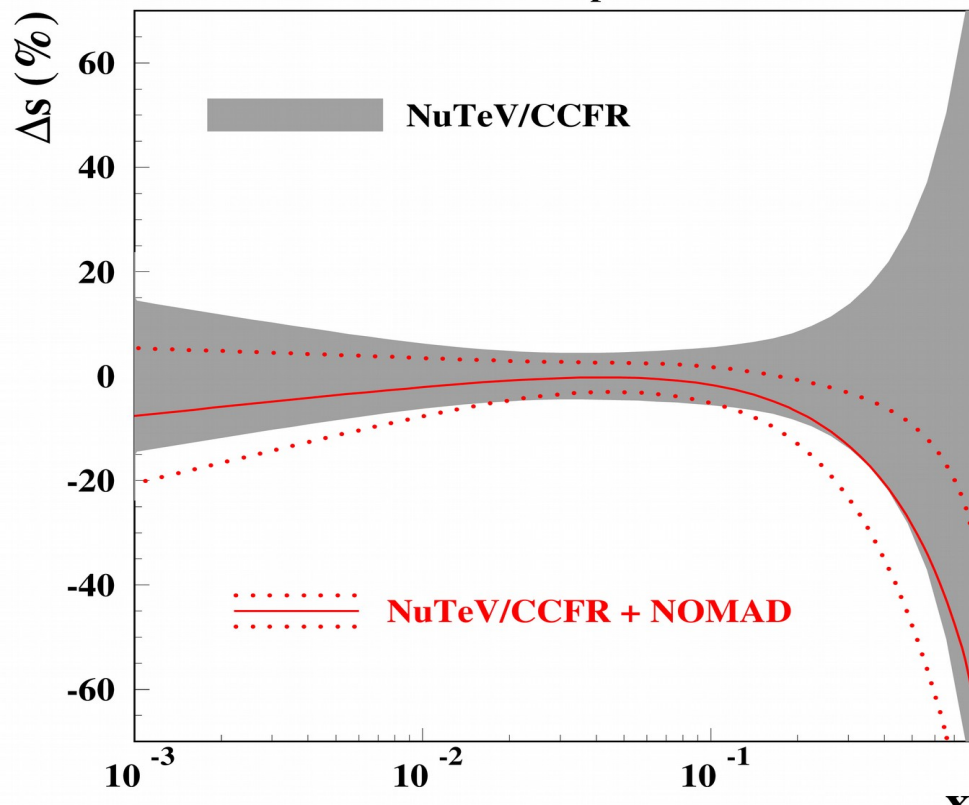
Integral suppression factor
 $K_s(20 \text{ GeV}^2) = 0.62 \pm 0.04$ is obtained



NOMAD charm data

$\mu=3 \text{ GeV}, n_f=3$

NOMAD NPB 876, 339 (2013)



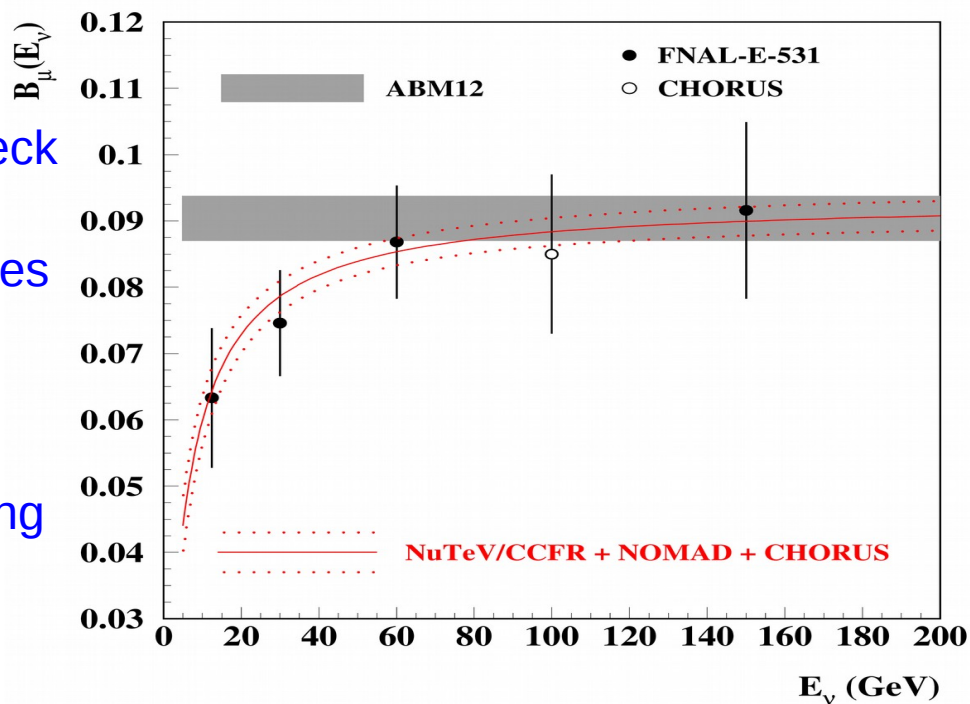
- The data on ratio $2\mu/\text{incl. CC ratio}$ with the 2μ statistics of 15000 events (much bigger than in earlier CCFR and NuTeV samples).
- Systematics, nuclear corrections, etc. cancel in the ratio
- Pull down strange quarks at $x>0.1$ with a sizable uncertainty reduction

The semi-leptonic branching ratio B_μ is a bottleneck

– weighted average of the charmed-hadron rates

$$B_\mu(E_\nu) = \sum_h r^h(E_\nu) B^h = a/(1+b/E_\nu)$$

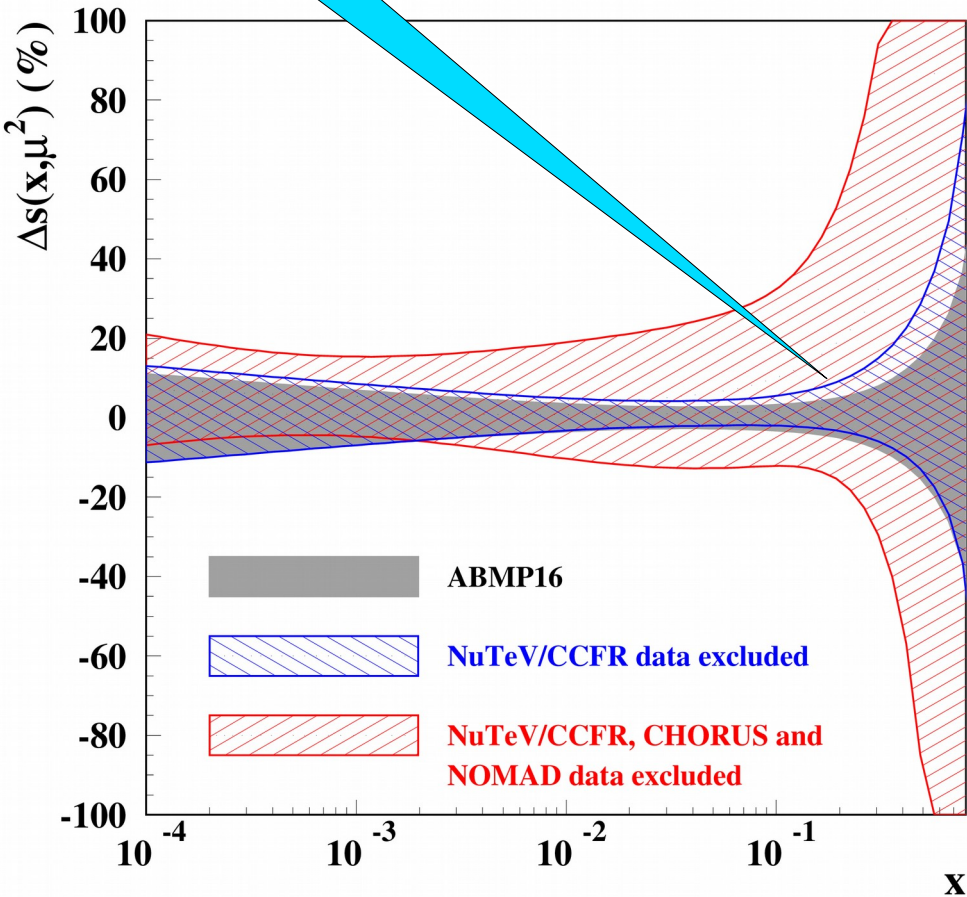
– fitted simultaneously with the PDFs, etc. using the constraint from the emulsion data



Constraints on strange sea

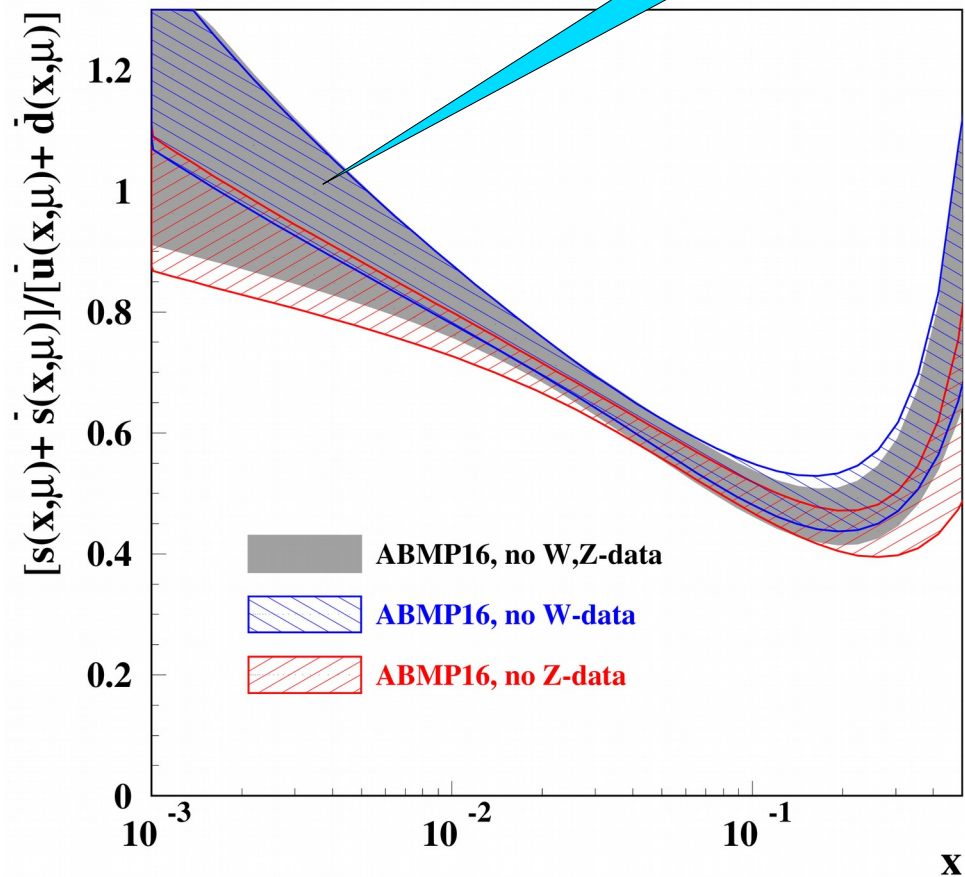
Controlled by
NOMAD

$\mu=3 \text{ GeV}, N_F=3$



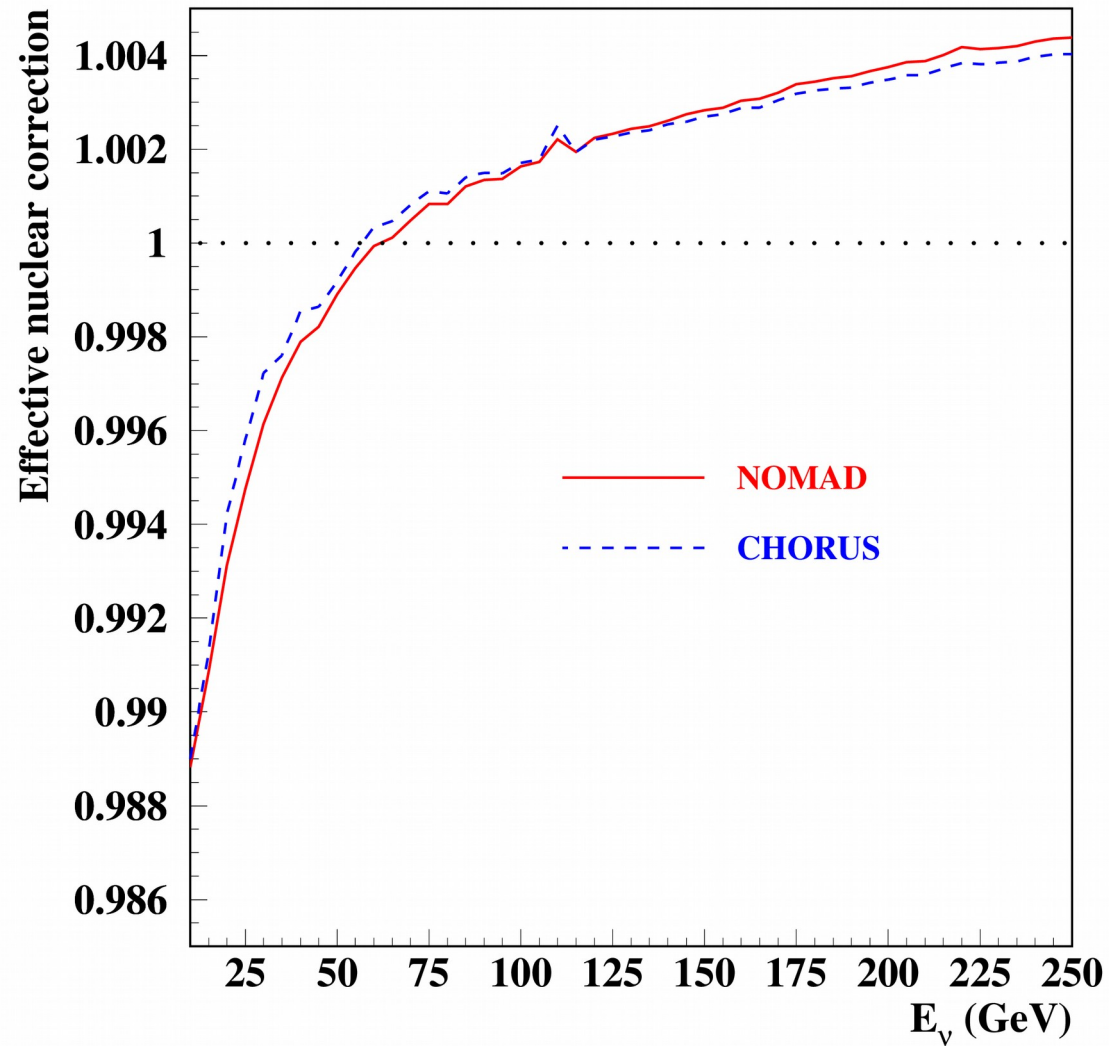
Controlled by
DY&DIS(incl.)

$\mu=3 \text{ GeV}, N_F=3$



- Uncertainty of $\sim 5\%$ is achieved at x around 0.1
- NuTeV/CCFR data play no essential role \rightarrow impact of the nuclear corrections is greatly reduced (NOMAD and CHORUS give the ratio CC/incl.)

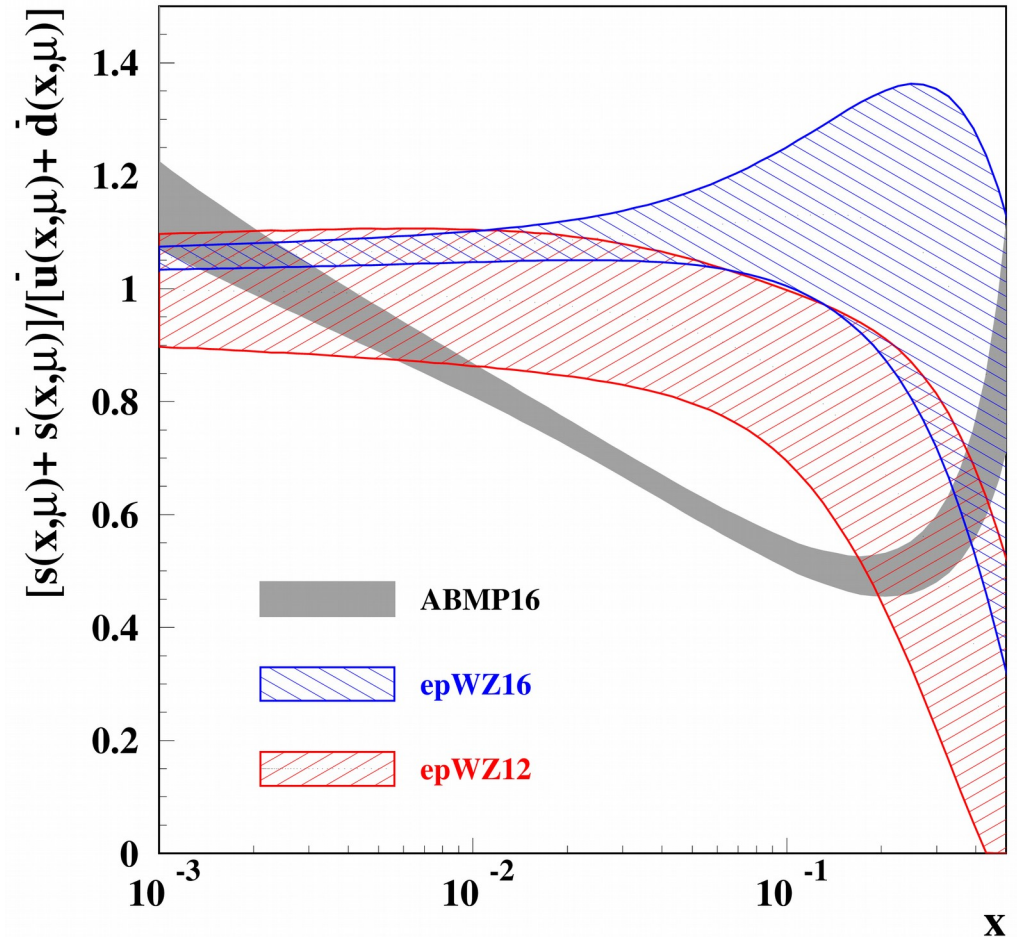
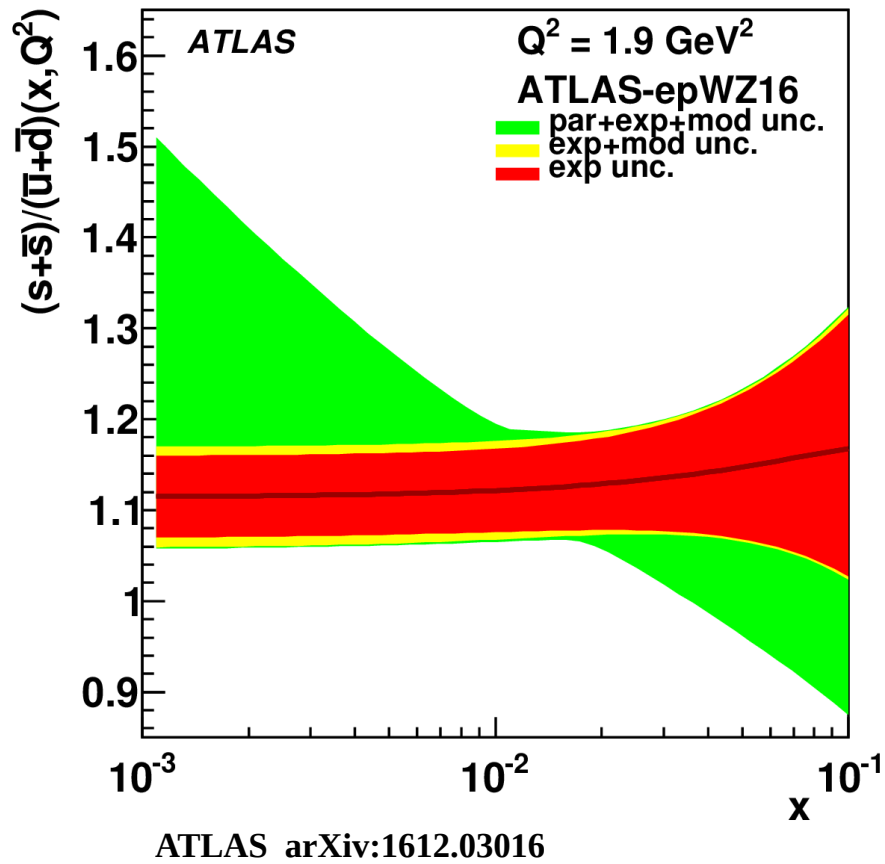
Nuclear corrections in NOMAD data



Nuclear corrections cancel in the ratio

ATLAS strange enhancement

$\mu^2 = 20 \text{ GeV}^2, N_F = 5$



The epWZ16 strange-sea determined from analysis of the combined HERA-ATLAS data is enhanced as compared to other (earlier) determinations

ABM strange sea determination is in particular based on the dimuon neutrino-nucleon DIS production (NuTeV/CCFR and NOMAD) that gives a strange sea suppression ~ 0.5 at $x \sim 0.2$

- Disentangling d - and s - contribution?
- Impact of the nuclear corrections?
-?

Test-fit data set (epWZ16 and CJ15 studies)

sa, Blümlein, Moch PLB 777, 134 (2018)
sa, Blümlein, Kulagin, Moch, Petti hep-ph/1808.06871

| Experiment | Process | <i>NDP</i> |
|------------|---------|------------|
|------------|---------|------------|

DIS

| | | |
|---------------------------------|--|------|
| HERA I+II | $e^\pm p \rightarrow e^\pm X$ $e^\pm p \rightarrow \bar{\nu} X$ | 1168 |
| Fixed-target (BCDMS, NMC, SLAC) | $l^\pm p \rightarrow l^\pm X$ | 1935 |

DIS heavy-quark production

| | | |
|---|---------------------------------------|-----|
| HERA I+II | $e^\pm p \rightarrow e^\pm c X$ | 52 |
| H1, ZEUS | $e^\pm p \rightarrow e^\pm b X$ | 29 |
| Fixed-target (CCFR, CHORUS, NOMAD, NuTeV) | $\bar{\nu} N \rightarrow \mu^\pm c X$ | 232 |

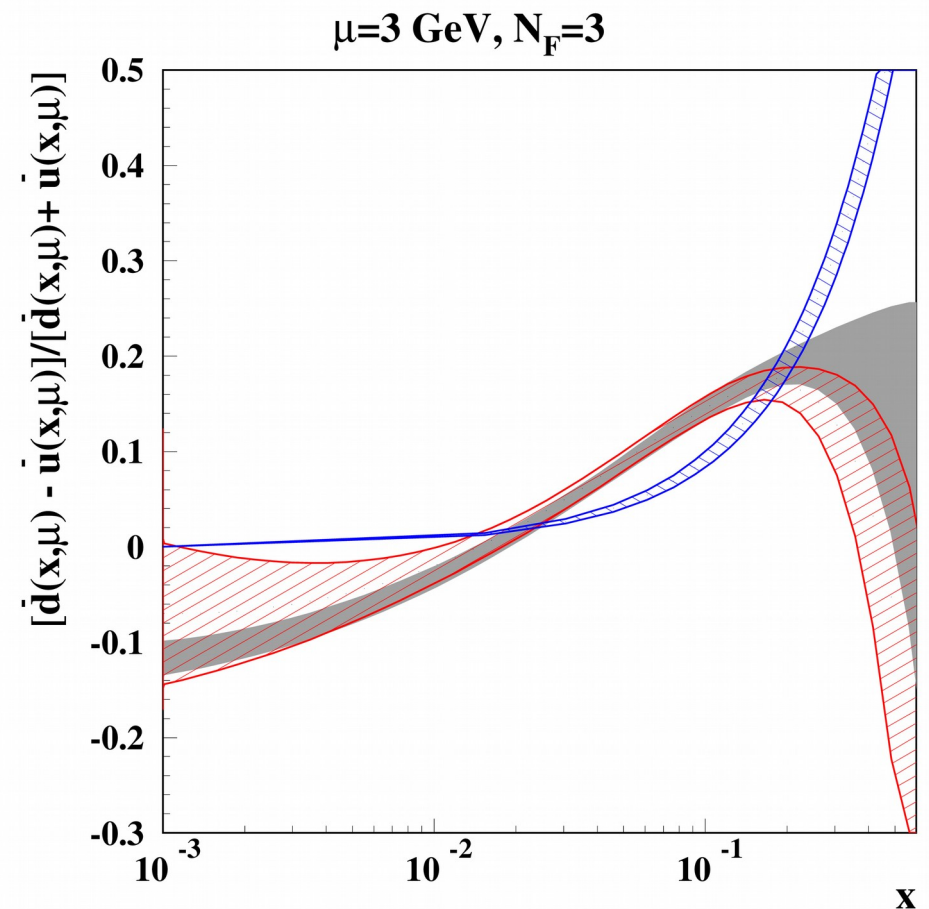
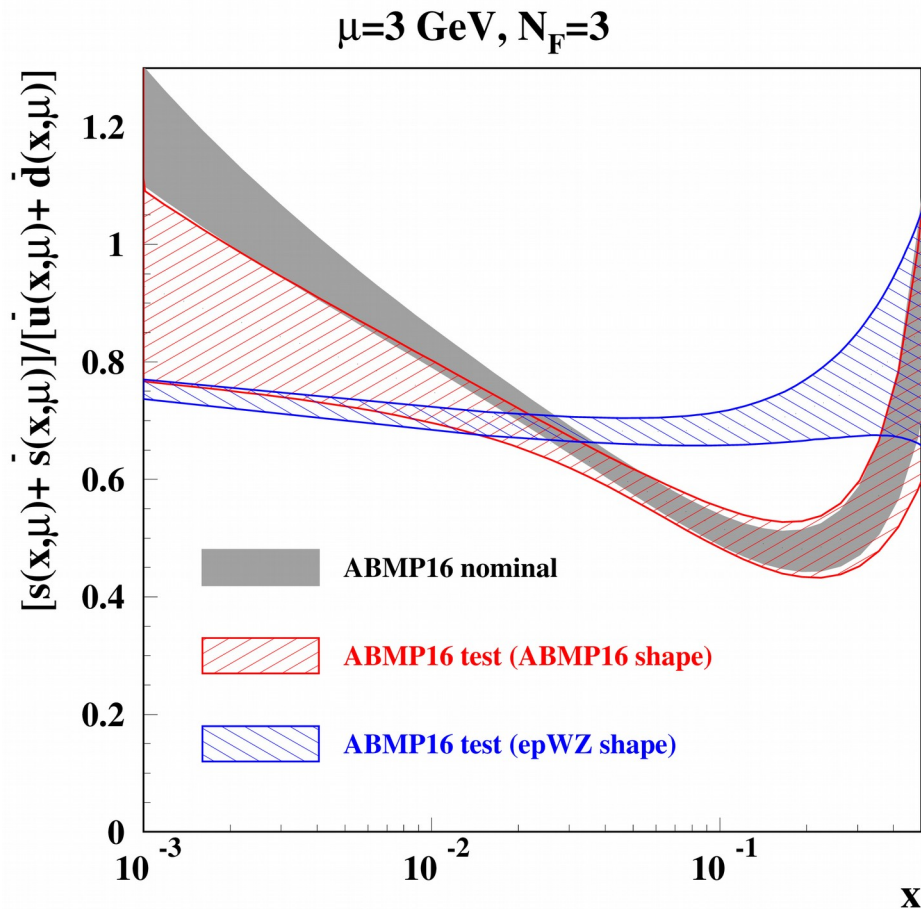
DY

| | | |
|-----------------------------------|--------------------------------|-----|
| Fixed-target (FNAL-605, FNAL-866) | $pN \rightarrow \mu^+ \mu^- X$ | 158 |
|-----------------------------------|--------------------------------|-----|

The ABMP16 framework with:

- DY data replaced by the deuteron ones \Rightarrow comparable quark disentangling at moderate and large x
- t-quark data excluded (no relevance for present study)

Test fit (the PDF shape comparison)



The strange sea is enhanced for the epWZ shape despite the ATLAS data are not used. However, the dimuon data description is not deteriorated: $\chi^2=167$ versus 161 for the ABMP shape \Rightarrow enhancement is achieved by the price of the d-quark sea suppression

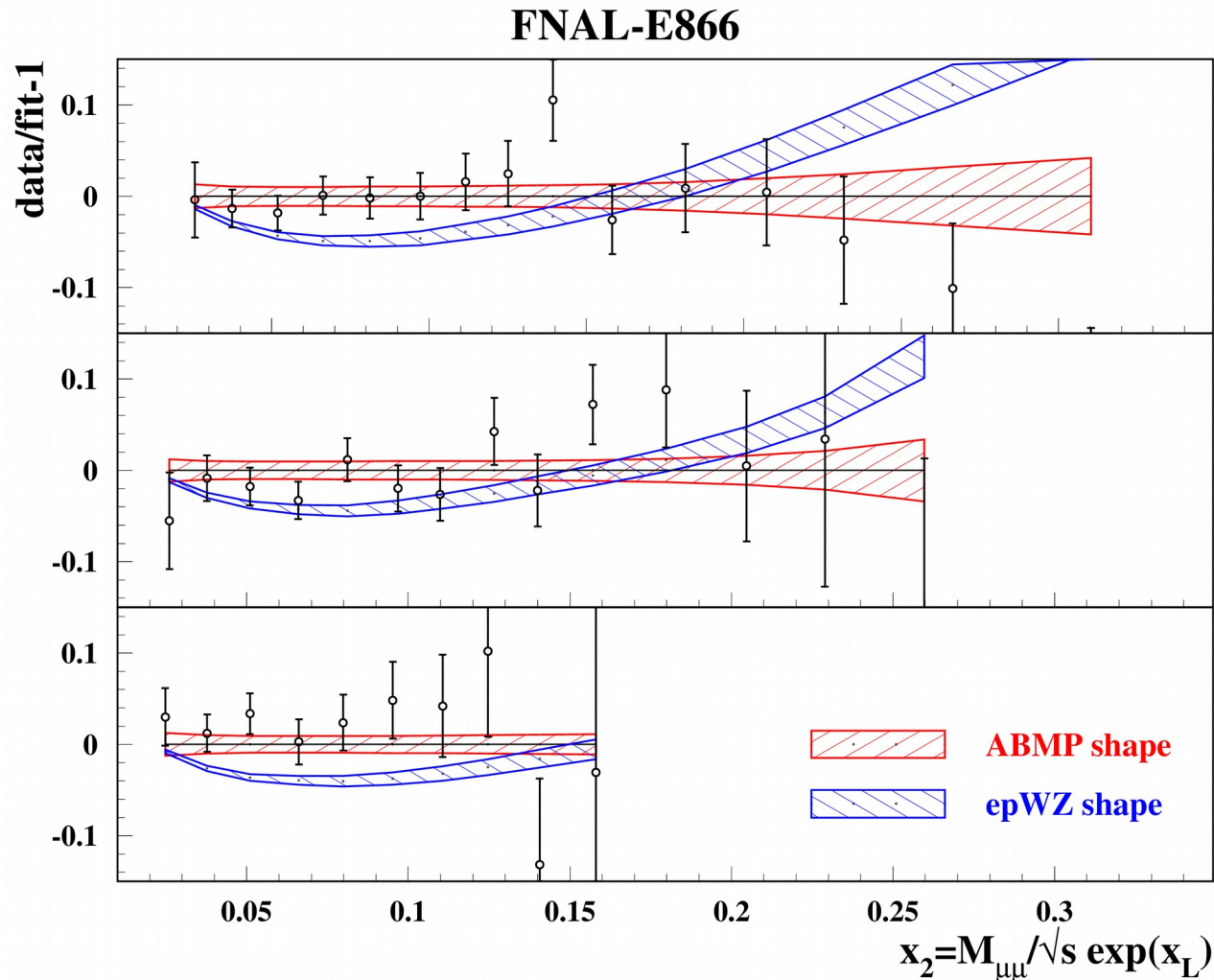
sa, Blümlein, Caminada, Lipka, Lohwasser,
Moch, Petti, Plačákytė PRD 91, 094002 (2015)

Checking styles of PDF shape

| | ABMP16 | CJ15 | CT10 | CT14 | epWZ16 | MMHT14 |
|-------------------------------|--|---|--|--|---|--------------------------------------|
| N_{PDF} | 28 | 21 | 26 | 26 | 14 | 31 |
| μ_0^2 (GeV ²) | 9 | 1.69 | 1.69 | 1.69 | 1.9 | 1 |
| χ^2 | 4065 | 4108 | 4148 | 4153 | 4336 | 4048 |
| PDF shape | $x^\alpha(1-x)^\beta \exp[P(x, \ln(x))]$ | $x^\alpha(1-x)^\beta P(x, \sqrt{x})$ | $x^\alpha(1-x)^\beta \exp[P(x, \sqrt{x})]$ | $x^\alpha(1-x)^\beta \exp[P(x, \sqrt{x})]$ | $x^\alpha(1-x)^\beta P(x, \sqrt{x})$ | $x^\alpha(1-x)^\beta P(x, \sqrt{x})$ |
| Constraints | | $\bar{u}=\bar{d} \quad (x \rightarrow 0)$ | $\alpha_{uv}=\alpha_{dv}$ $\alpha_{\bar{u}}=\alpha_{\bar{d}}=\alpha_s$ $\bar{u}=\bar{d} \quad (x \rightarrow 0)$ | $\alpha_{uv}=\alpha_{dv}$ $\beta_{uv}=\beta_{dv}$ $\alpha_{\bar{u}}=\alpha_{\bar{d}}=\alpha_s$ | $\alpha_{\bar{u}}=\alpha_{\bar{d}}=\alpha_s$ $\bar{u}=\bar{d} \quad (x \rightarrow 0)$ | |
| $\alpha_s(M_Z)$ | 0.1153 | 0.1147 | 0.1150 | 0.1160 | 0.1162 | 0.1158 |

- Various PDF-shape modifications provide comparable description with $N_{\text{PDF}} \sim 30$
- Some deterioration, which happens in cases is apparently due to constraints on large(small)-x exponents

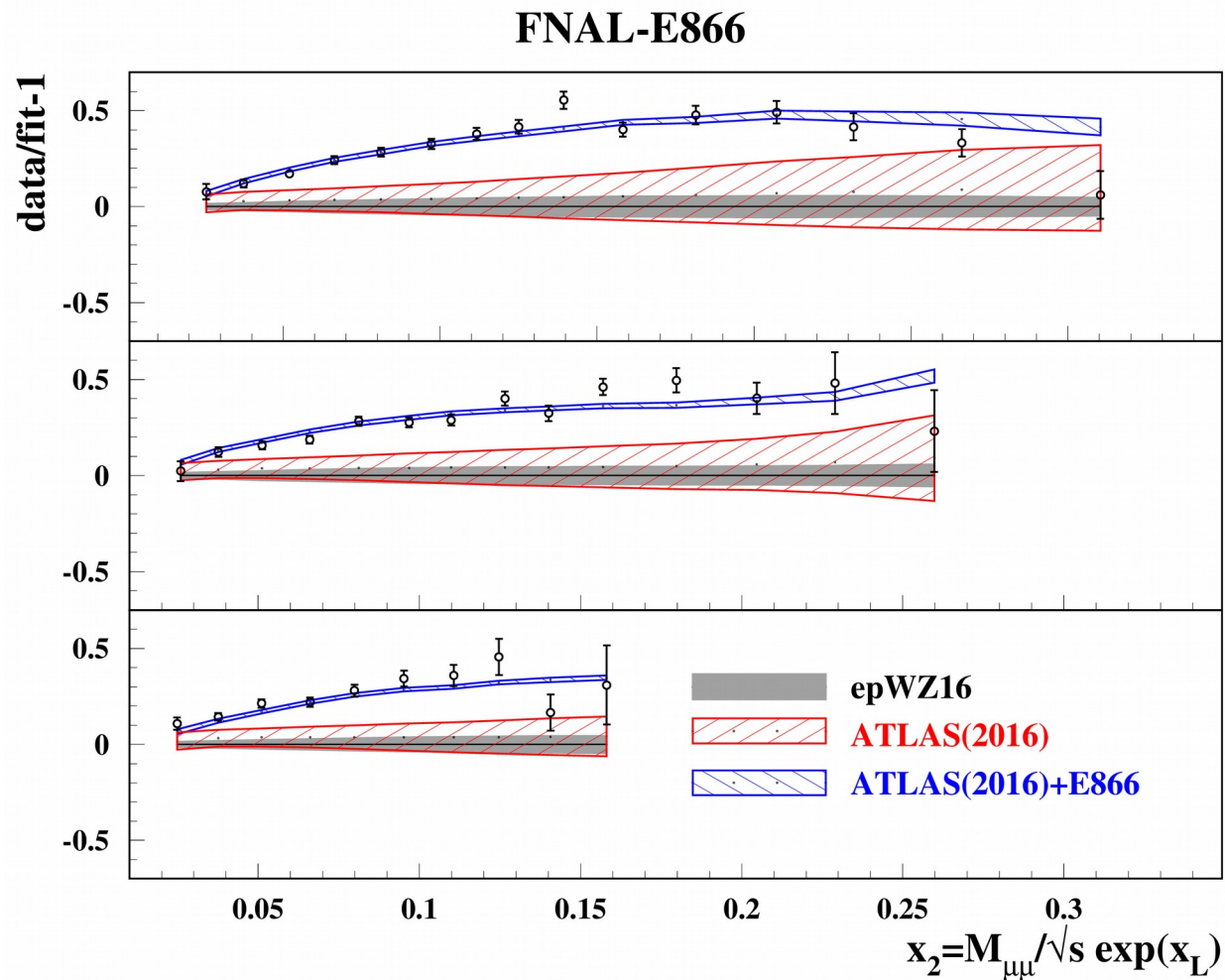
E866 data in the test fit



The E866 data on p/d DY cross sections are sensitive to the iso-spin sea asymmetry

The epWZ shape does not allow to accommodate E866 data: $\chi^2/NDP=96/39$ versus $49/39$ for the ABMP shape; the errors in epWZ predictions are suppressed at small x , evidently due to over-constrained PDF shape at small x

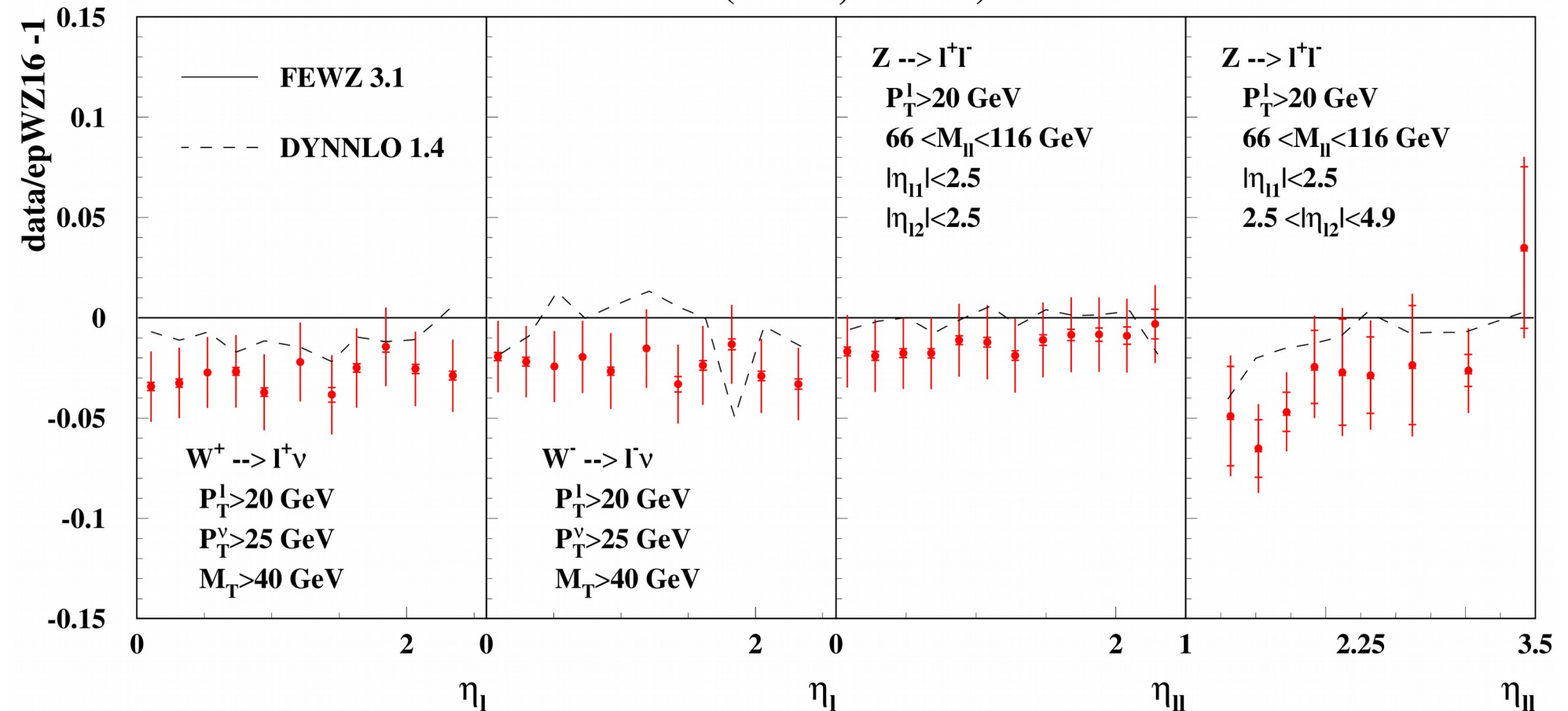
Consistency of ATLAS and E866 data



- The uncertainties in epWZ predictions are quite narrow and several σ off the E866 data \Rightarrow E866 cannot be accommodated into the fit ATLAS
- The ABMP16 shape gives much wider error band \Rightarrow E866 data are well accommodated: $\chi^2/NDP=48/39$ and $40/34$ for the E866 and ATLAS, respectively

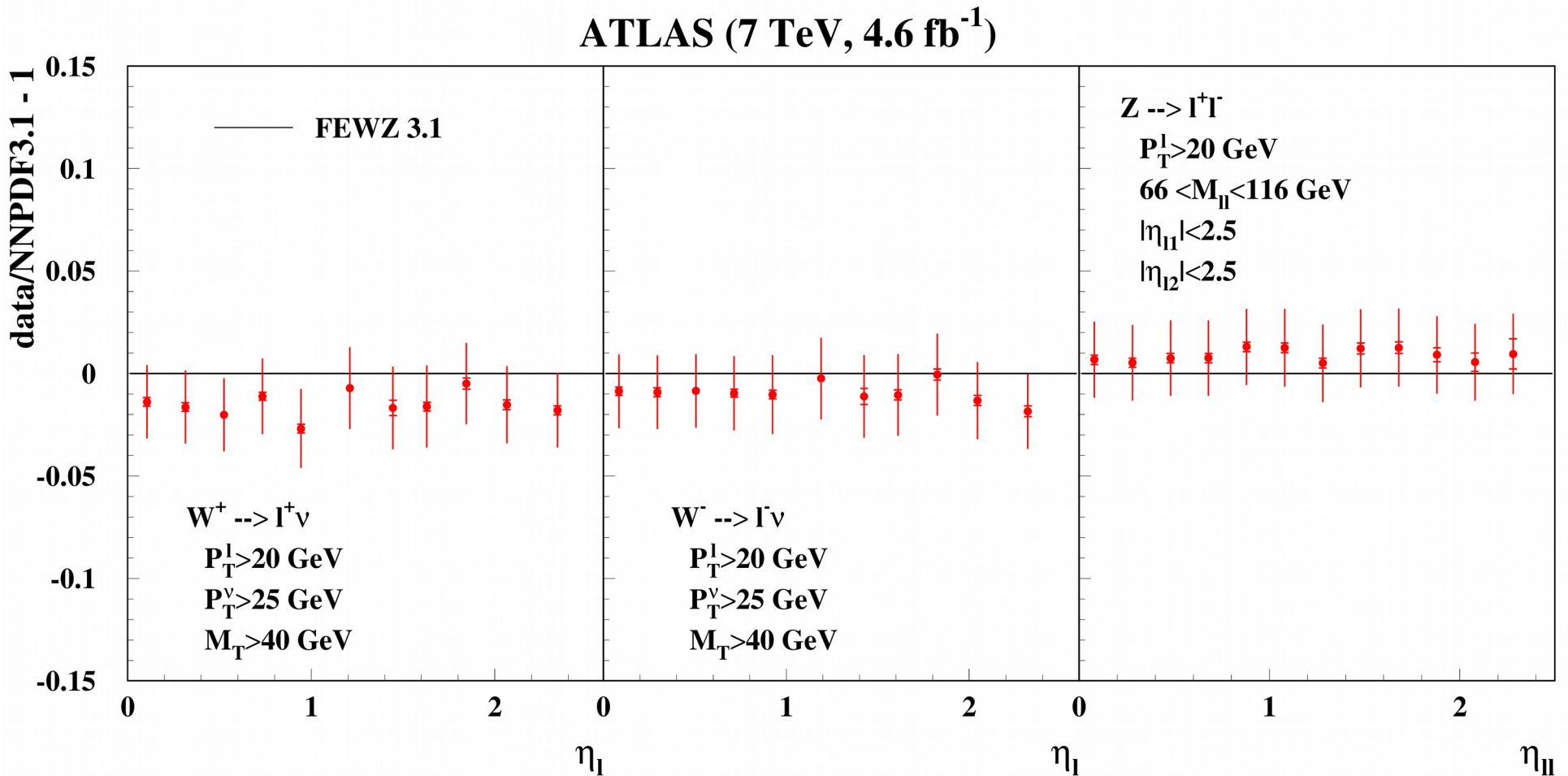
Closure test of the NN31 fit

ATLAS (7 TeV, 4.6 fb⁻¹)



The epWZ16 predictions go systematically above the ATLAS data \Rightarrow either statistical bias or inaccurate theory predictions (epWZ16 fit uses combination of the NLO calculations with the NNLO K-factors)

Closure test of the NNPDF3.1 fit

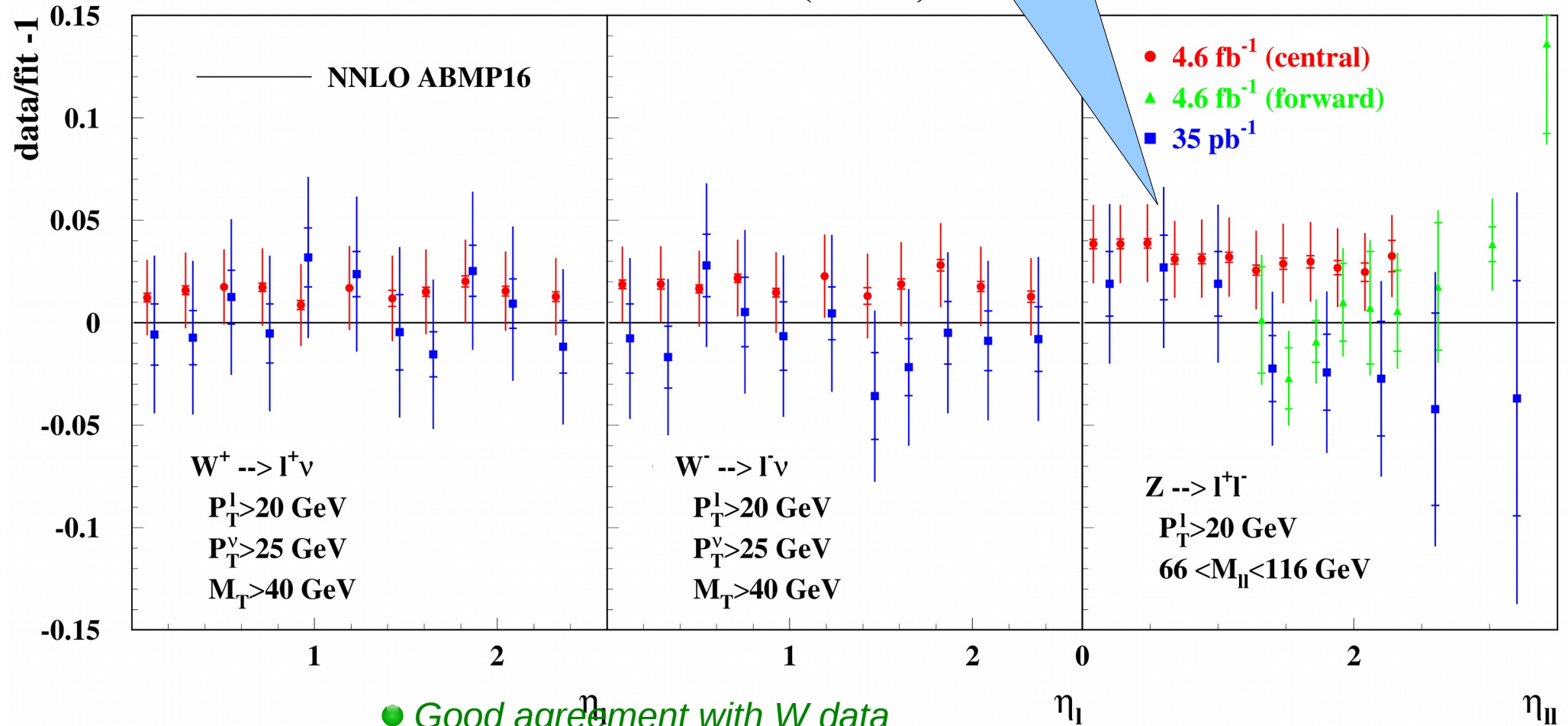


- Different trend for W and Z data $\Rightarrow \chi^2/NDP = 400/34$; problems with the flavor disentangling
- Suppressed (fitted) charm distribution requires corresponding enhancement of strangeness due to constraint from W data

New input: ATLAS at 7 TeV

Signal of strangeness enhancement?

ATLAS (7 TeV)

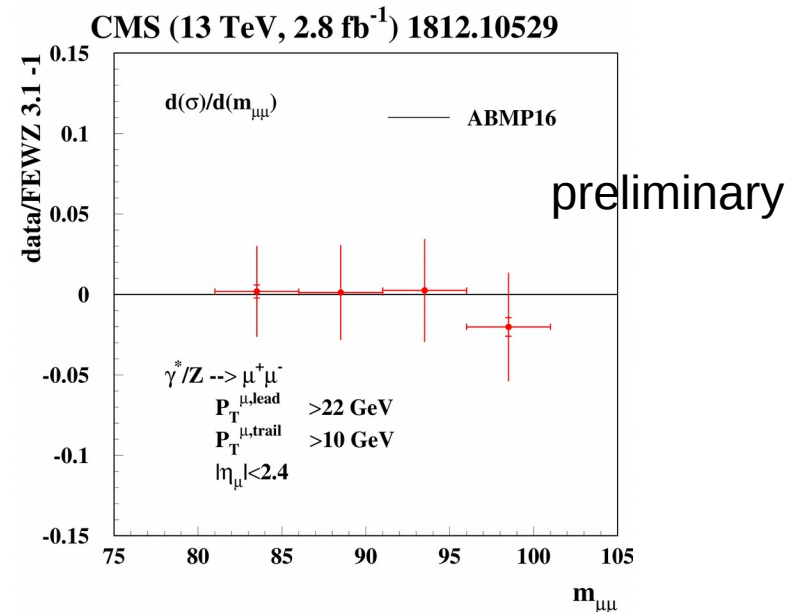
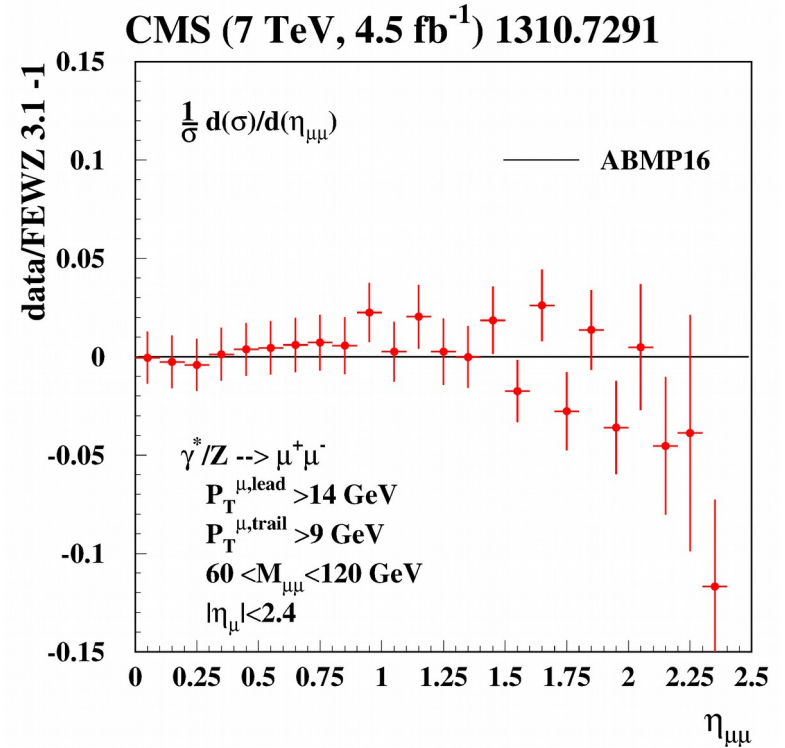
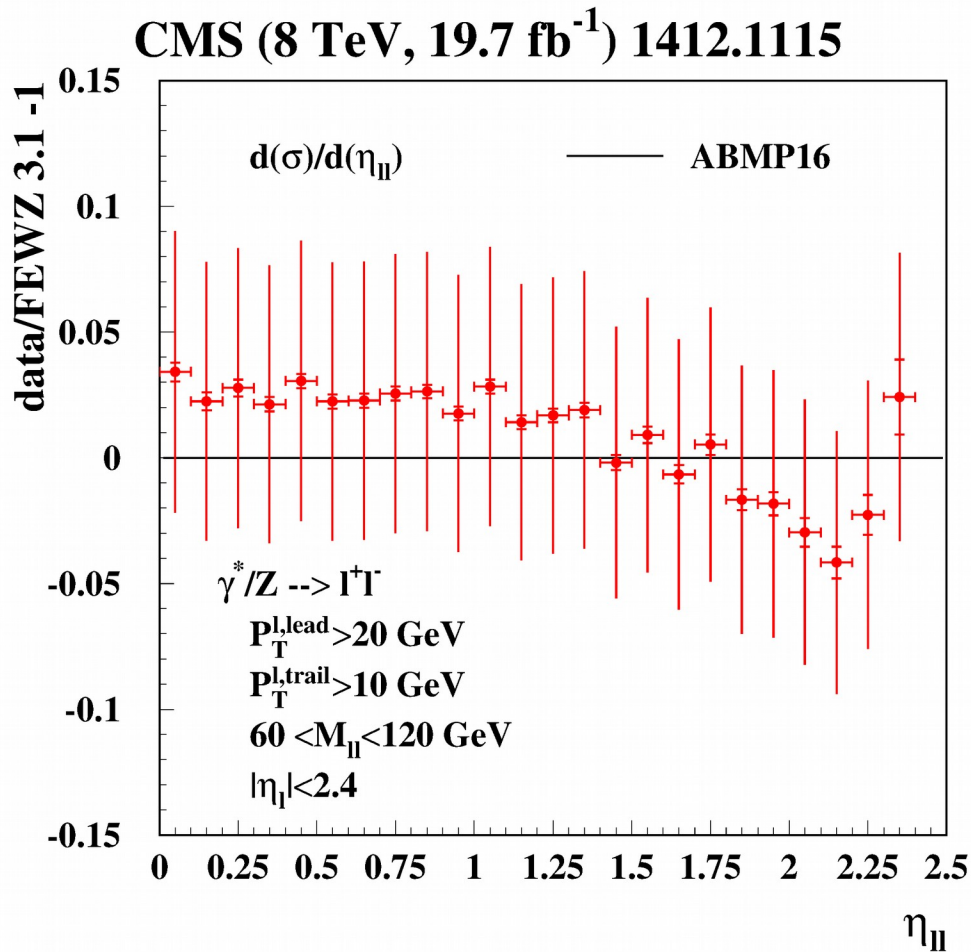


● Good agreement with W data

● Undershooting Z-boson data

● Different trends for the central and forward Z-boson data

CMS data on Z-boson production



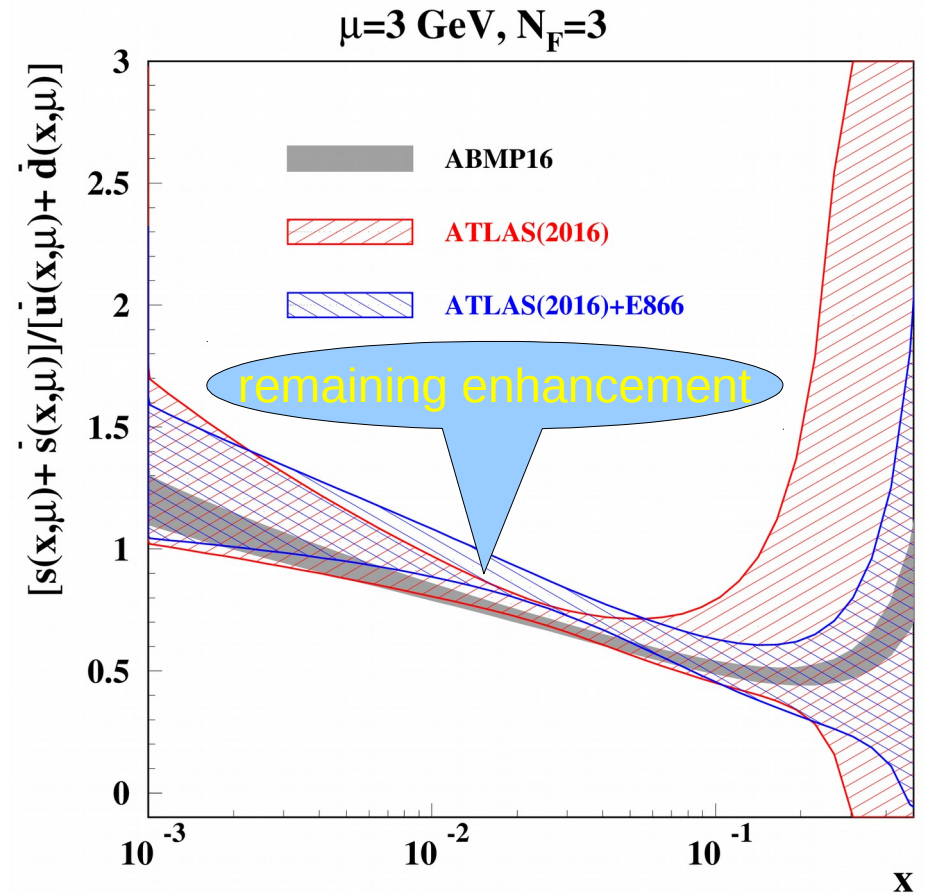
The CMS data go somewhat lower than the ATLAS ones and the trend is different at large rapidity; further clarification is necessary

Impact of ATLAS data with flexible PDF shape

| | $\kappa_s(\mu^2=20 \text{ GeV}^2)$ |
|---------------------|------------------------------------|
| HERA+ATLAS | 0.81(18) |
| HERA+ATLAS+E866 | 0.72(8) |
| ABMP16(incl. NOMAD) | 0.66(3) |

κ_s is integral strange sea suppression factor:

$$\kappa_s(\mu^2) = \frac{\int_0^1 x[s(x, \mu^2) + \bar{s}(x, \mu^2)]dx}{\int_0^1 x[\bar{u}(x, \mu^2) + \bar{d}(x, \mu^2)]dx},$$



- For the flexible PDF shape the strangeness is in a broad agreement with the one extracted from the dimuon data
- The E866 data are consistent with the ATLAS(2016) set: $\chi^2/\text{NDP}=48/39$ and $40/34$, respectively.

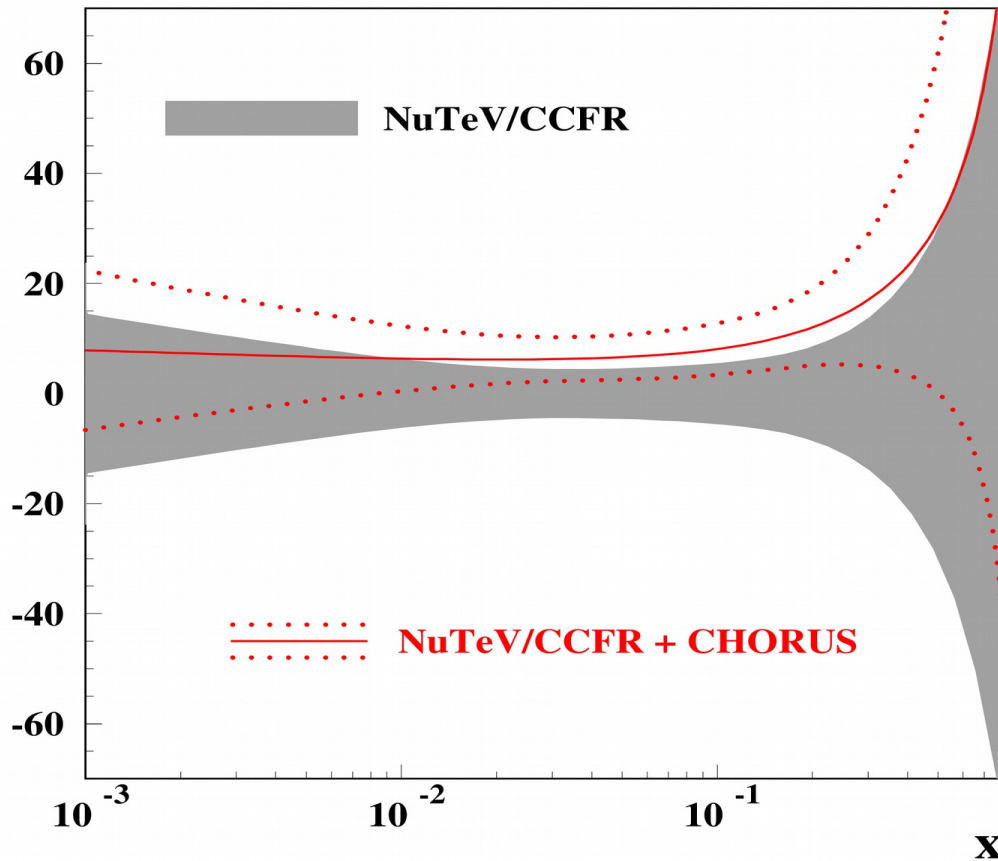
Summary

- The strange sea suppression observed in the early νN DIS experiments is confirmed by recent precise measurements (NOMAD, CHORUS)
- These data sets can be accommodated into the global PDF fit with a consistent treatment of the fixed-target and collider Drell-Yan data
- The ATLAS analysis based on the combination of Drell-Yan and HERA DIS data demonstrates strange sea enhancement by the price of disagreement with the Fermilab fixed-target Drell-Yan data (E-866, E-906) and overconstrained PDF shape at small x
- A refined comparison with recent CMS measurements is desirable in order to confirm small strange-sea enhancement at $x \sim 0.01$ driven by the recent ATLAS Drell-Yan data

EXTRAS

CHORUS charm data

$\mu=3 \text{ GeV}, n_f=3$



CHORUS data pull strangeness up, however the statistical significance of the effect is poor

sa, Blümlein, Caminada, Lipka, Lohwasser, Moch, Petti, Placakyte hep-ph/1404.6469

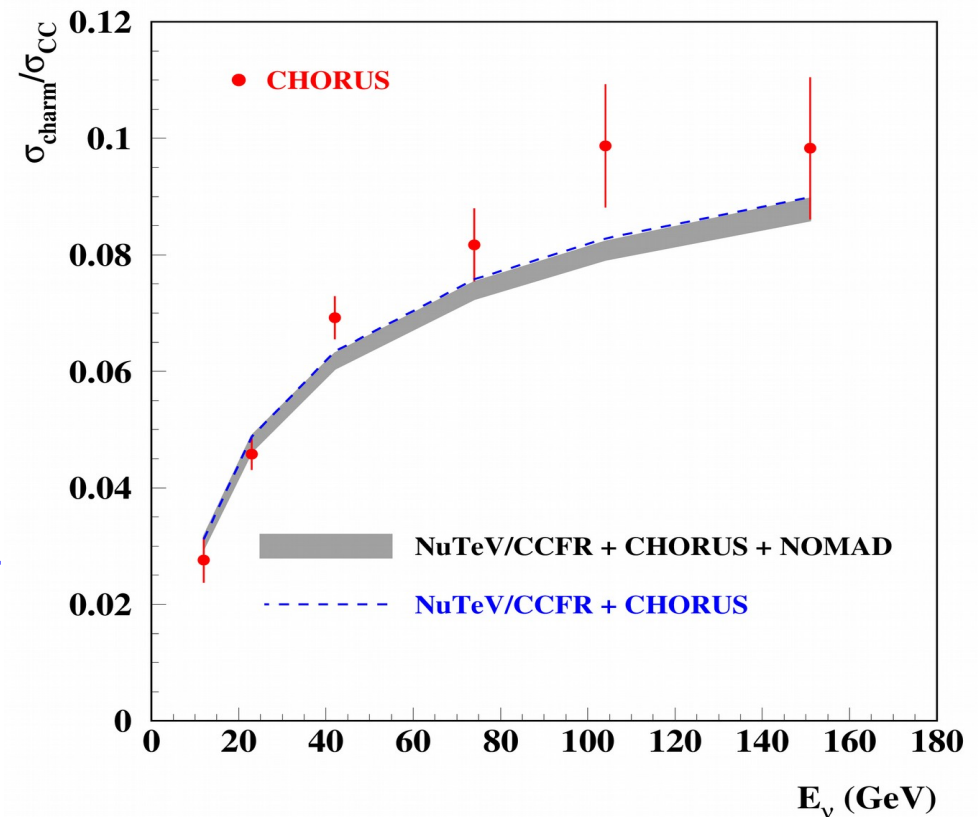
Emulsion data on charm/CC ratio with the charmed hadron vertex measured

CHORUS NJP 13, 093002 (2011)

– full phase space measurements

– no sensitivity to B_μ

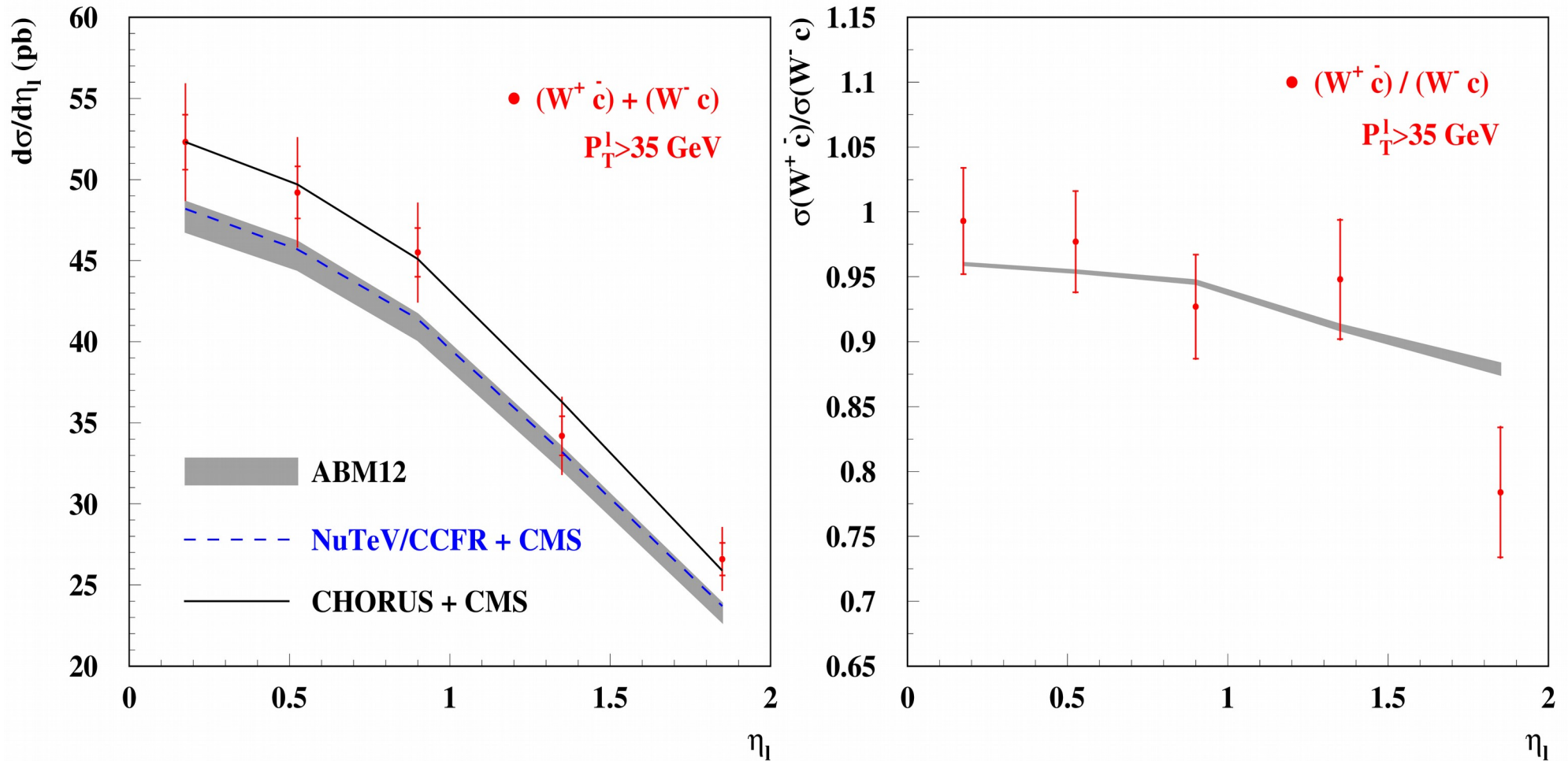
– low statistics (2013 events)



CMS W +charm data

CMS Collaboration JHEP 02, 013 (2014)

CMS (7 TeV, 5 1/fb)

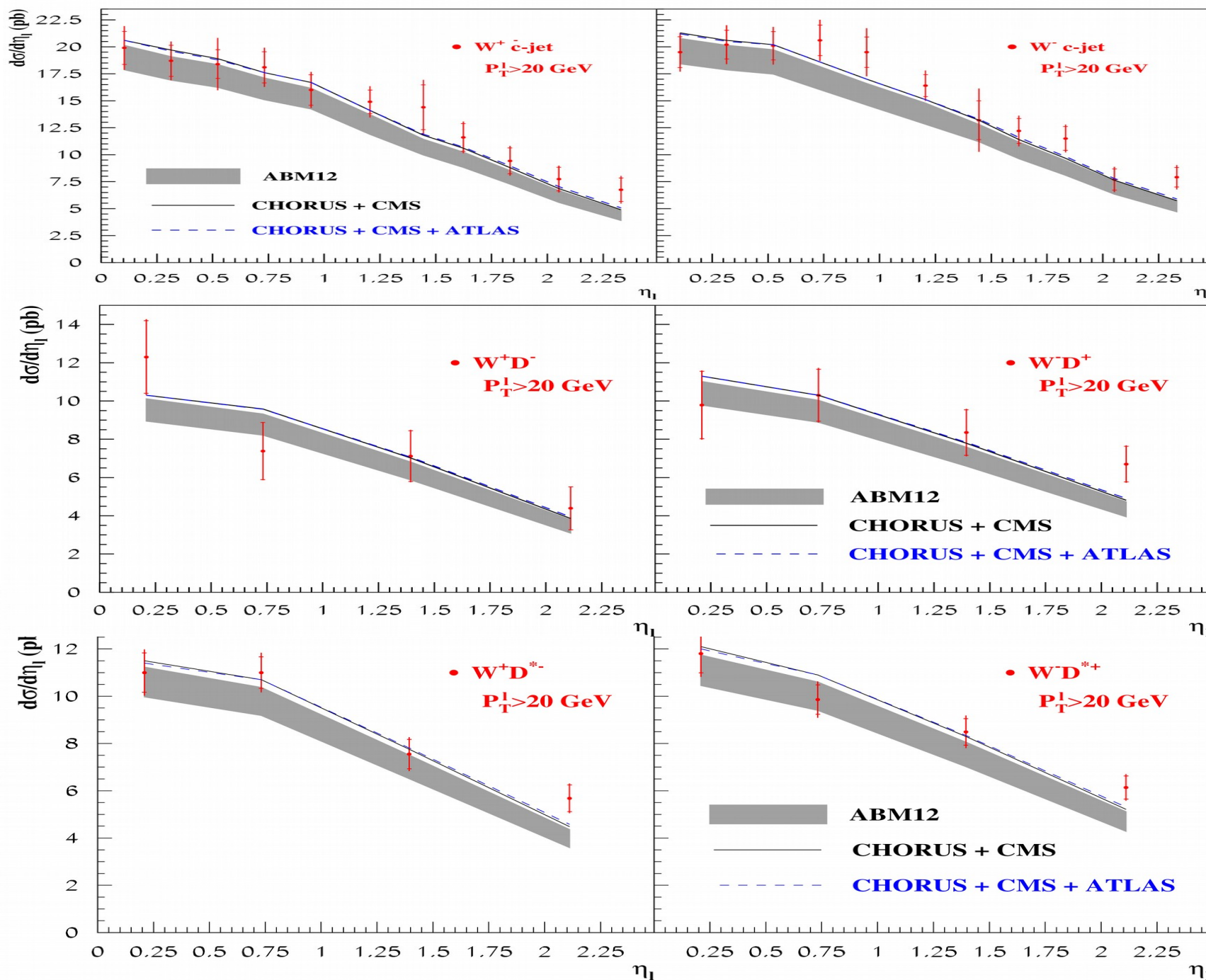


- CMS data go above the NuTeV/CCFR by 1σ ; little impact on the strange sea
- The charge asymmetry is in a good agreement with the charge-symmetric strange sea
- Good agreement with the CHORUS data

ATLAS W+charm data

ATLAS arXiv:1402.6263

ATLAS (7 TeV, 4.6 1/fb)

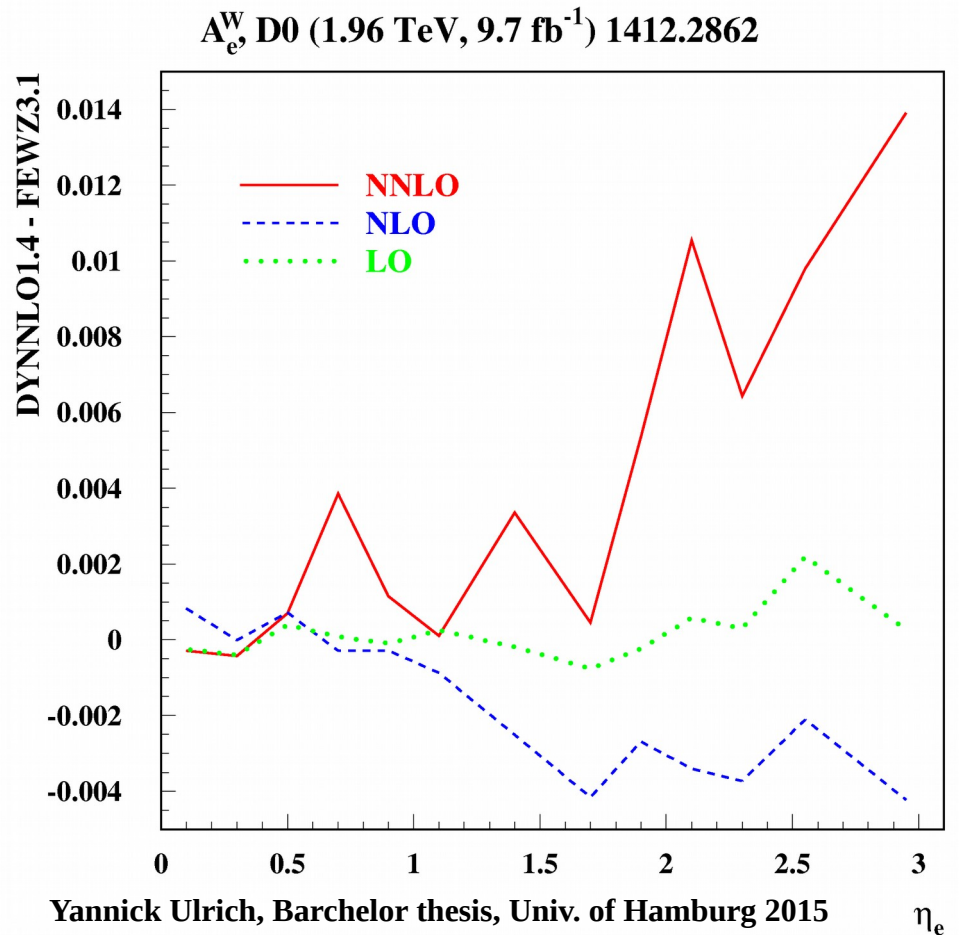
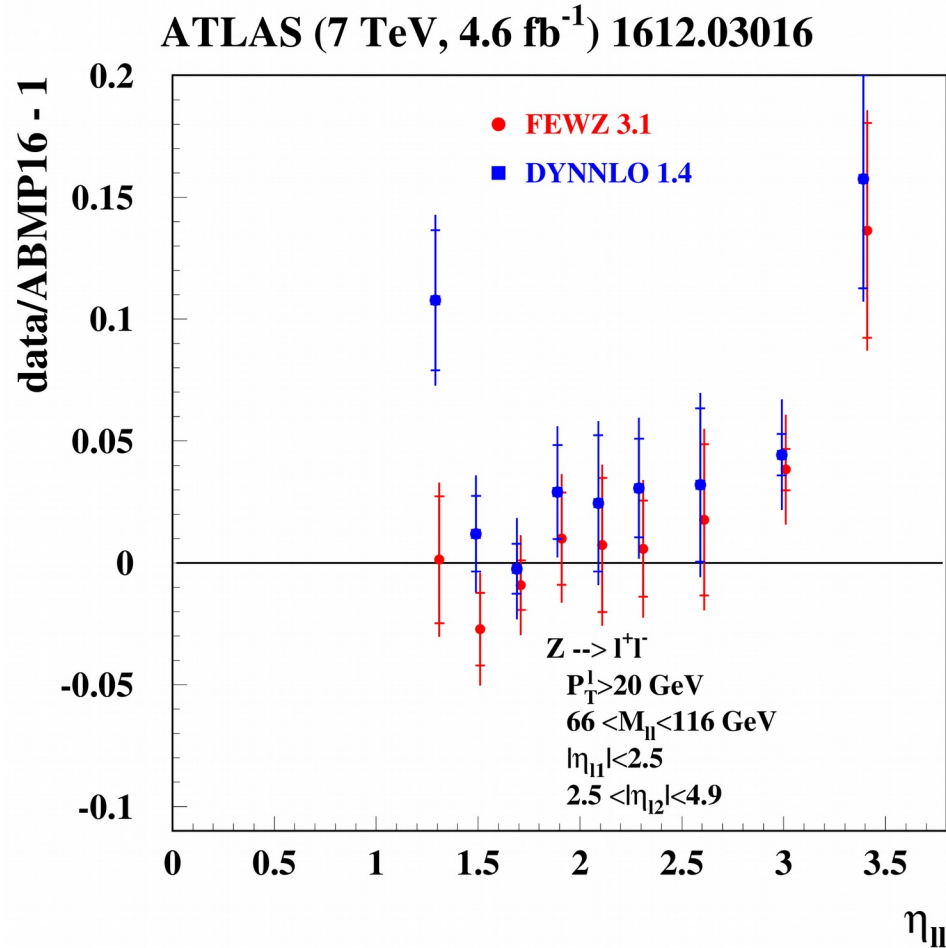


Details of the epWZ and ABMP16 fits

| | epWZ16 | ABMP16 |
|-----------|--|---|
| Data | HERA, ATLAS W&Z | HERA, LHC and Tevatron W&Z, fixed-target DIS and charm production, fixed-target DY, |
| PDF shape | $xu_v(x, \mu_0^2) = A_{u_v} x^{B_{u_v}} (1-x)^{C_{u_v}} (1 + E_{u_v} x^2),$ $xd_v(x, \mu_0^2) = A_{d_v} x^{B_{d_v}} (1-x)^{C_{d_v}},$ $x\bar{u}(x, \mu_0^2) = A_{\bar{u}} x^{B_{\bar{u}}} (1-x)^{C_{\bar{u}}},$ $x\bar{d}(x, \mu_0^2) = A_{\bar{d}} x^{B_{\bar{d}}} (1-x)^{C_{\bar{d}}},$ $xg(x, \mu_0^2) = A_g x^{B_g} (1-x)^{C_g} - A'_g x^{B'_g} (1-x)^{C'_g},$ $x\bar{s}(x, \mu_0^2) = A_{\bar{s}} x^{B_{\bar{s}}} (1-x)^{C_{\bar{s}}},$ <p style="text-align: center;">15 free parameters</p> | $xq_v(x, \mu_0^2) = \frac{2\delta_{qu} + \delta_{qd}}{N_q^v} (1-x)^{b_{qv}} x^{a_{qv}} P_{qv}(x),$ $xq_s(x, \mu_0^2) = A_{qs} (1-x)^{b_{qs}} x^{a_{qs}} P_{qs}(x),$ $xg(x, \mu_0^2) = A_g (1-x)^{b_g} x^{a_g} P_g(x),$ $P_p(x) = (1 + \gamma_{-1,p} \ln x) (1 + \gamma_{1,p} x + \gamma_{2,p} x^2 + \gamma_{3,p} x^3),$ <p style="text-align: center;">25 free parameters</p> |

ABMP16 PDFs are selected more flexible in order to accommodate more data as compared to the EpWZ16 fit, which was evolved form the HERA data analysis

NNLO tools benchmarking



DYNNLO-FEWZ difference not fully understood; further benchmarking is needed

LHC data on Z-boson production

