

QCD axions off the beaten tracks

Based on: “Astrophobic Axions” Phys.Rev.Lett. 120 (2018) no.26, 261803 [arXiv:1712.04940]

In collaboration with [L. Di Luzio](#) (IPPP, Durham), [F. Mescia](#) (Barcelona U.), [P. Panci](#) (CERN) and [R. Ziegler](#) (CERN)

“Redefining the Axion window” Phys.Rev.Lett. 118 (2017) no.3, 031801 [arXiv:1610.07593]

“Window for preferred Axion models” Phys.Rev. D96 (2017) no.7, 075003 [arXiv:1705.05370]

In collaboration with [Luca Di Luzio](#) (IPPP, Durham) and [Federico Mescia](#) (Barcelona U.)

“The KLASH Conceptual Design Report: The physics case” (to appear)

In collaboration with [Maurizio Giannotti](#) (Barry U.) and [Luca Visinelli](#) (Uppsala U. & Nordita)

Enrico Nardi



15th PATRAS Workshop - Freiburg, June 3-7, 2019

Strong CP: a small value problem

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\cancel{P} \ \& \ \cancel{T}] \quad \leftarrow \quad d_n \lesssim 3 \cdot 10^{-26} e \text{ cm}$$

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\cancel{P} \ \& \ \cancel{T}] \quad \leftarrow \quad d_n \lesssim 3 \cdot 10^{-26} \text{ e cm}$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} (i \not{D} - m_q e^{i\theta_q}) q - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a - \theta \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a$$

$$\bar{\theta} = \theta - \theta_q$$

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\cancel{P} \ \& \ \cancel{T}] \quad \leftarrow \quad d_n \lesssim 3 \cdot 10^{-26} \text{ e cm}$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} (i \not{D} - m_q e^{i\theta_q}) q - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a - \theta \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a$$

$$\bar{\theta} = \theta - \theta_q$$

- This is qualitatively different from other small values problems:

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\cancel{P} \ \& \ \cancel{T}] \quad \leftarrow \quad d_n \lesssim 3 \cdot 10^{-26} \text{ e cm}$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} (i \not{D} - m_q e^{i\theta_q}) q - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a - \theta \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a$$

$$\bar{\theta} = \theta - \theta_q$$

- This is qualitatively different from other small values problems:

- In the SM $\bar{\theta}$ receives the first finite **Log** corrections at $O(\alpha^2)$ [Ellis, Gaillard (1979) Khriplovich, Vainshtein (1994)]

Unlike m_H^2 that is quadratically sensitive to Λ_{UV}^2

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\cancel{P} \ \& \ \cancel{T}] \quad \leftarrow \quad d_n \lesssim 3 \cdot 10^{-26} \text{ e cm}$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} (i \not{D} - m_q e^{i\theta_q}) q - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a - \theta \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a$$

$$\bar{\theta} = \theta - \theta_q$$

This is qualitatively different from other small values problems:

- In the SM $\bar{\theta}$ receives the first finite **Log** corrections at $O(\alpha^2)$ [Ellis, Gaillard (1979), Khriplovich, Vainshtein (1994)]
Unlike m_H^2 that is quadratically sensitive to Λ_{UV}^2
- Unlike Λ_{cc}** , $y_{e,u,d} \sim 10^{-6} \div 10^{-5}$ it evades explanations based on environmental selection [Ubbaldi, 0811.1599, Kaloper & Terning, 1710.01740, Dine, Stephenson Haskins, Ubbaldi, & Di Xu 1801.03466]

Strong CP: a small value problem

- QCD is defined in terms of two dimensionless parameters which are not predicted by the theory. Measurements yield:

$$\alpha_s \sim O(0.1-1) \quad \text{and} \quad \bar{\theta} < 10^{-10} \quad [\text{P} \ \& \ \text{T}] \quad \leftarrow d_n \lesssim 3 \cdot 10^{-26} \text{ e cm}$$

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{q} (i \not{D} - m_q e^{i\theta_q}) q - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a - \theta \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a$$

$$\bar{\theta} = \theta - \theta_q$$

- This is qualitatively different from other small values problems:

- In the SM $\bar{\theta}$ receives the first finite **Log** corrections at $O(\alpha^2)$ [Ellis, Gaillard (1979), Khriplovich, Vainshtein (1994)]

Unlike m_H^2 that is quadratically sensitive to Λ_{UV}^2

- Unlike Λ_{cc} , $y_{e,u,d} \sim 10^{-6} \div 10^{-5}$ it evades explanations based on environmental selection [Ubbaldi, 0811.1599, Kaloper & Terning, 1710.01740, Dine, Stephenson Haskins, Ubbaldi, & Di Xu 1801.03466]

If Multiverse, $\bar{\theta}$ becomes the most urgent s.v. problem in the SM

The Peccei-Quinn solution

The Peccei-Quinn solution

[Peccei, Quinn (1977), Weinberg (1978), Wilczek (1978)]

- Assume a global $U(1)_{PQ}$:

- (i) Spontaneously broken
- (ii) QCD anomalous

The Peccei-Quinn solution

[Peccei, Quinn (1977), Weinberg (1978), Wilczek (1978)]

- Assume a global $U(1)_{PQ}$:

- (i) Spontaneously broken
- (ii) QCD anomalous

- Implies a PNGB of $U(1)_{PQ}$: the Axion.

Axion field comes equipped with a shift symmetry: $a(x) \rightarrow a(x) + \delta\alpha f_a$

$$\mathcal{L}_{\text{eff}} = \left(\bar{\theta} + \frac{a}{f_a} \right) \frac{\alpha_s}{8\pi} G_a^{\mu\nu} \tilde{G}_{\mu\nu}^a - \frac{1}{2} \partial^\mu a \partial_\mu a + \mathcal{L}(\partial_\mu a, \psi)$$

$\theta_{\text{eff}}(x)$

and with a periodic potential $V(\theta_{\text{eff}}) \rightarrow V_{\text{min}}$ when $\theta_{\text{eff}} \rightarrow 0$

Axion models

- **PQWW axion:**

Axion identified with the phase of the Higgs in a 2HDM
($f_a \sim V_{EW}$ was quickly ruled out long ago)

[Peccei, Quinn (1977),
Weinberg (1978), Wilczek (1978)]

Viabale models require $f_a \gg V_{EW}$:

Most viable axion models fall in two classes:

- **DSFZ Axion:** SM quarks and Higgses, charged under PQ.

Requires 2HDM + 1 scalar singlet. SM leptons are also PQ charged.

[Dine, Fischler, Srednicki (1981), Zhitnitsky (1980)]

- **KSVZ Axion** (or hadronic axion):

All SM fields are neutral under PQ. QCD anomaly induced by new quarks, vectorlike under SM, chiral under PQ + 1 scalar singlet

[Kim (1979), Shifman, Vainshtein, Zakharov (1980)]

We need to know where to search

The "canonical" axion window

Central value:

$$g_{a\gamma\gamma} \sim \frac{\alpha}{2\pi} \frac{m_a}{f_\pi m_\pi} \sim \frac{10^{-10}}{\text{GeV}} \left(\frac{m_a}{\text{eV}} \right)$$

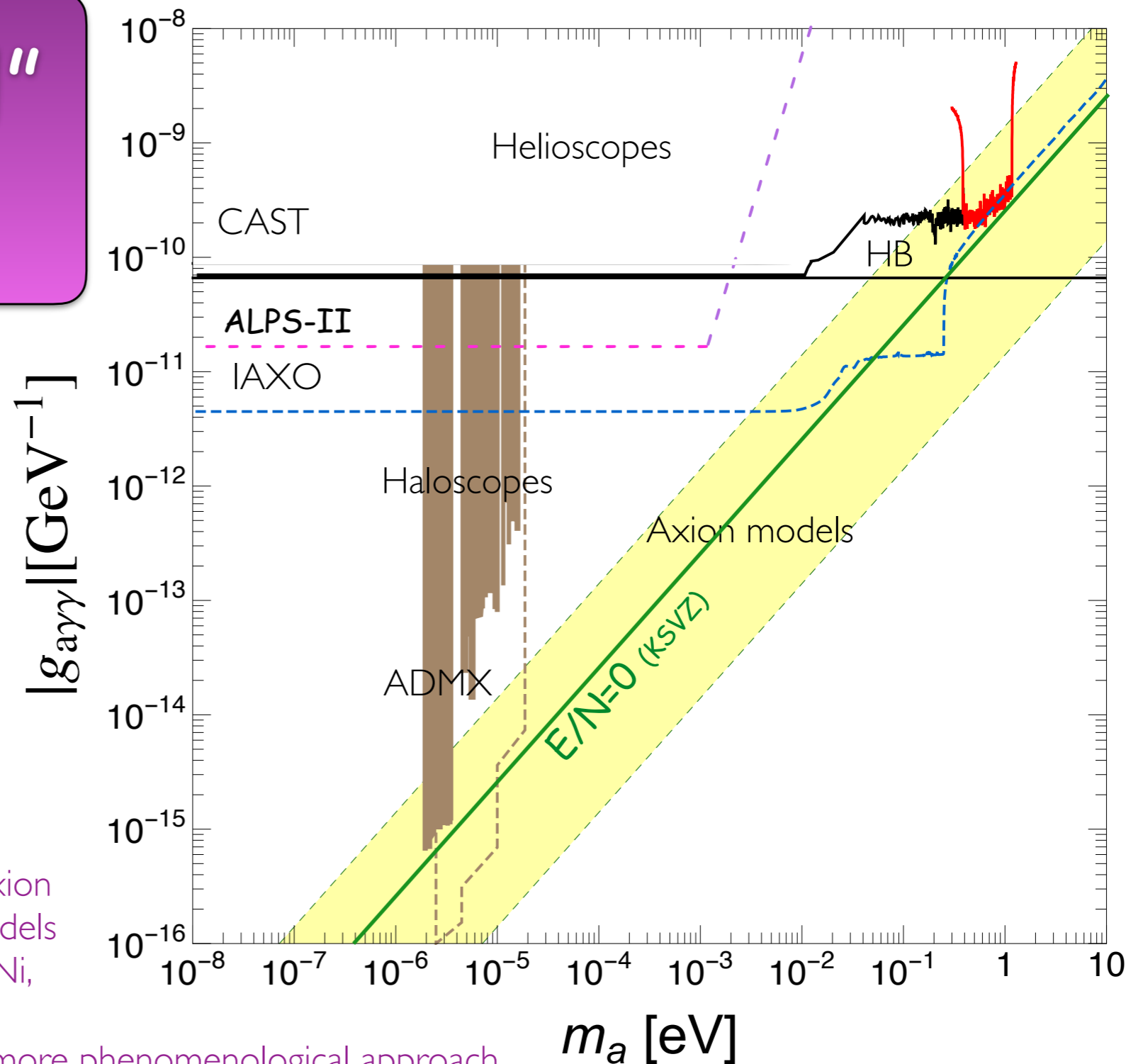
Model dependence:

$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92 \right)$$

$$|E/N - 1.92| \in [0.07, 7]$$

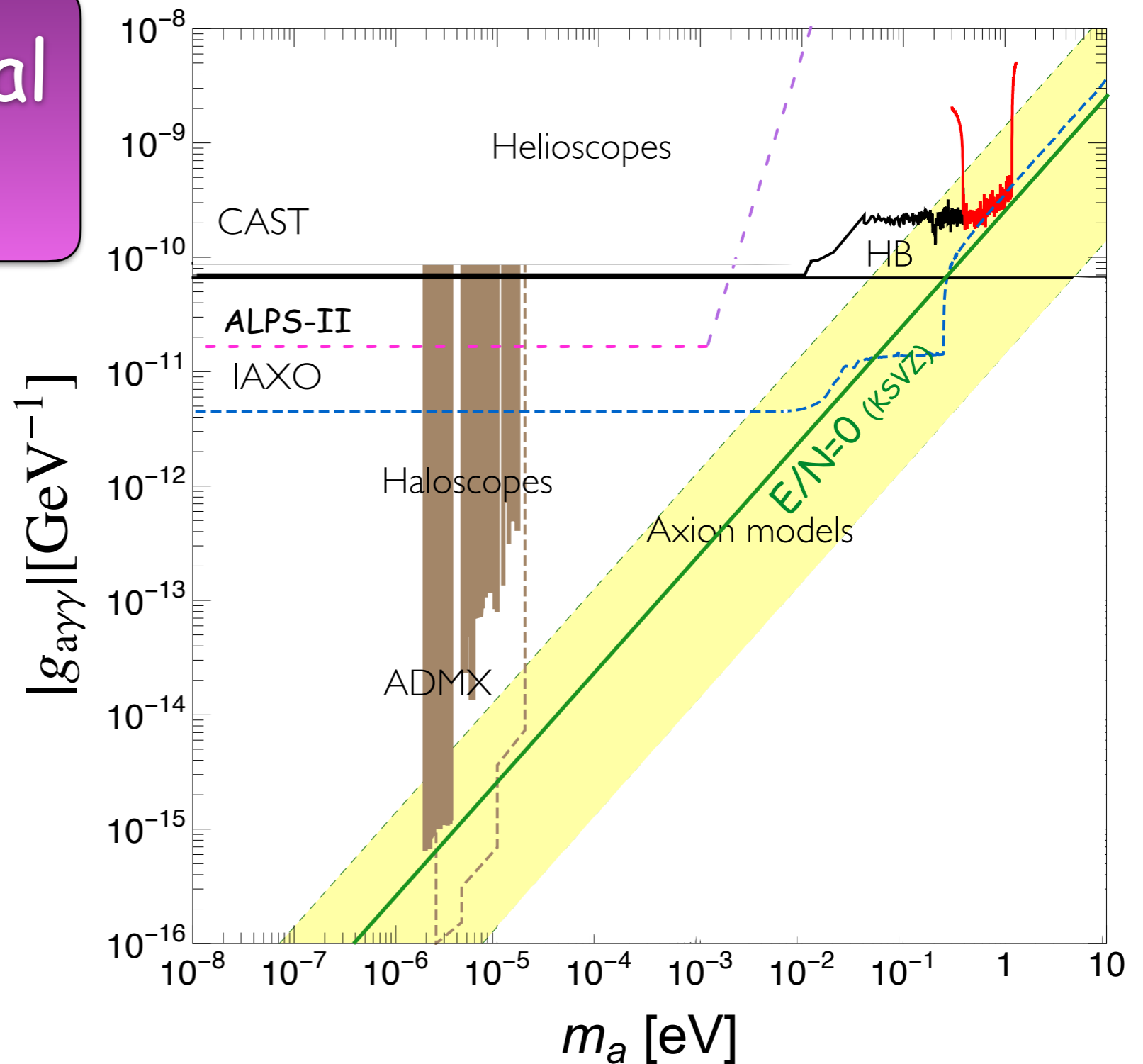
Particle Data Group: Since the end of 90's, Axion window chosen to include representative models from: Kaplan, NPB 260 (1985); Cheng, Geng, Ni, PRD 52 (1995); Kim, PRD 58 (1998).

Since 2017: $E/N \in (5/3, 44/3)$ based on a more phenomenological approach



We need to know where to search

$\Omega_a \approx \Omega_{\text{CDM}}$ canonical
mass windows:

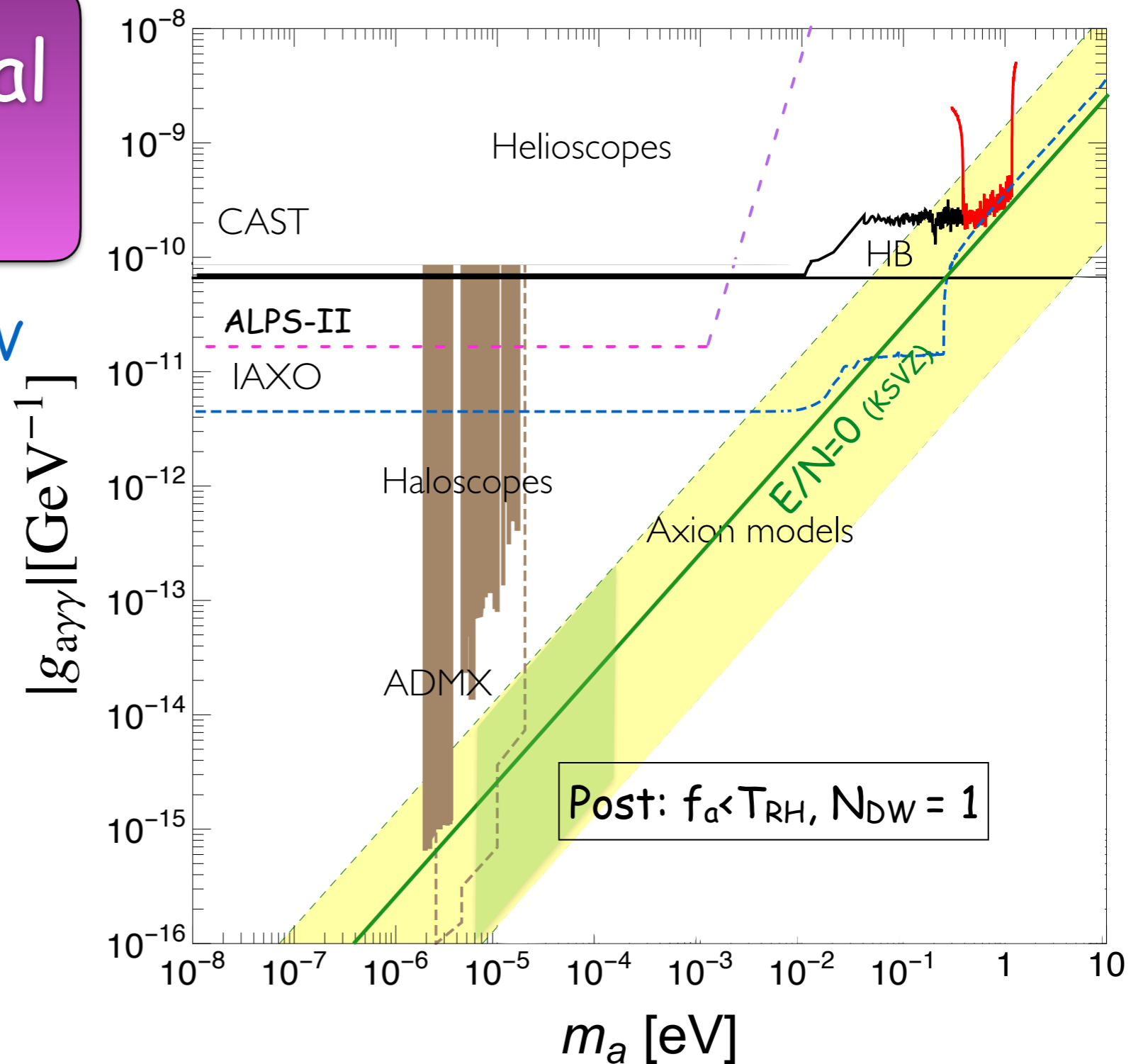


We need to know where to search

$\Omega_a \approx \Omega_{\text{CDM}}$ canonical
mass windows:

Preferred: $m_a \in (10, 100) \mu\text{eV}$

post-inflationary $\langle \theta_i \rangle \simeq 2$, $N_{\text{DW}} = 1$



We need to know where to search

$\Omega_a \approx \Omega_{\text{CDM}}$ canonical
mass windows:

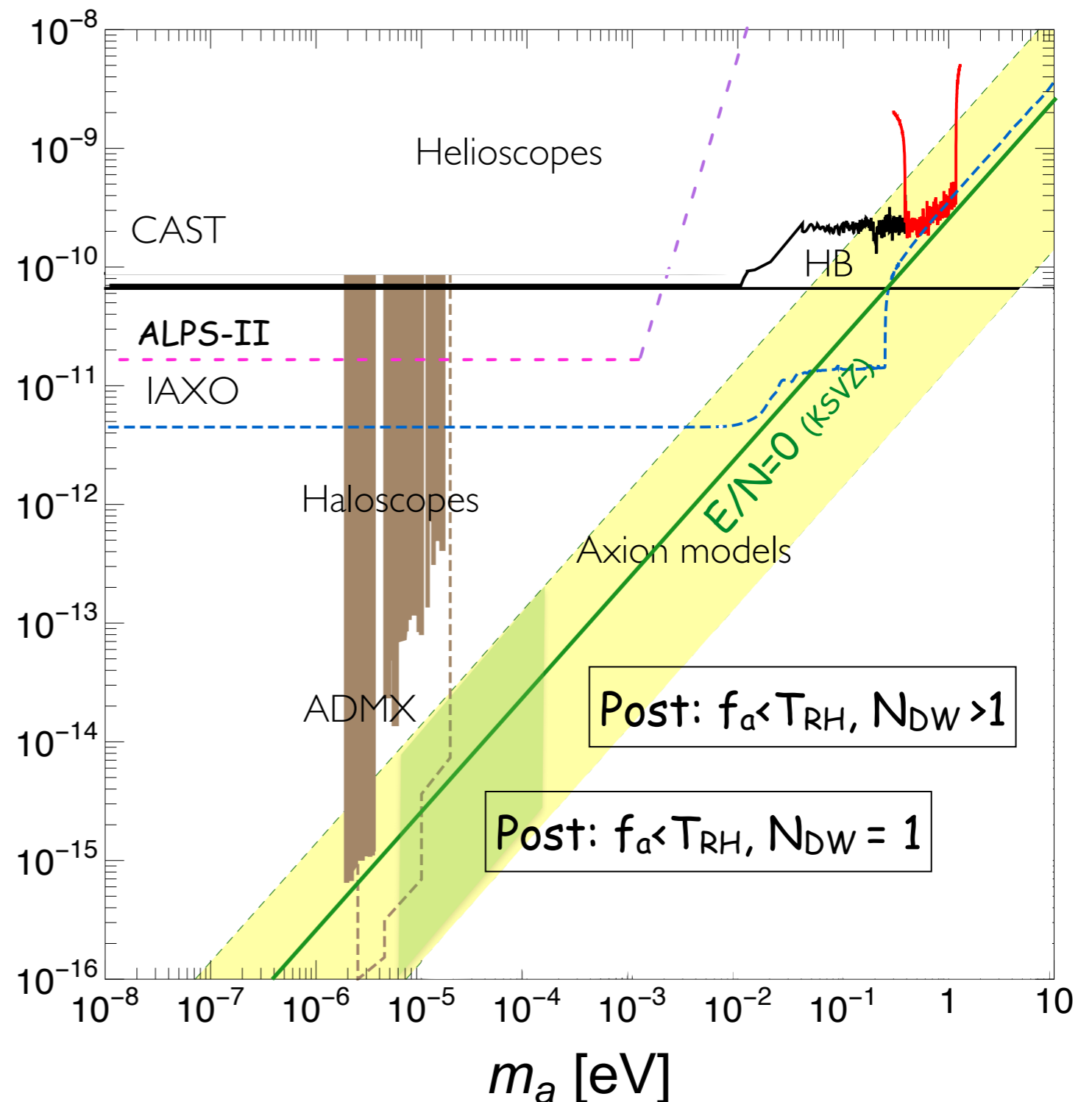
Preferred: $m_a \in (10, 100) \mu\text{eV}$

post-inflationary $\langle \theta_i \rangle \approx 2$, $N_{\text{DW}} = 1$

Larger: $m_a \gtrsim \text{few } 100 \mu\text{eV}$

$\Omega_a \approx \Omega_{\text{CDM}}$ needs $N_{\text{DW}} > 1$
contributions (cosmological issues)

$|g_{a\gamma\gamma}| [\text{GeV}^{-1}]$



We need to know where to search

$\Omega_a \approx \Omega_{\text{CDM}}$ canonical
mass windows:

Preferred: $m_a \in (10, 100) \mu\text{eV}$

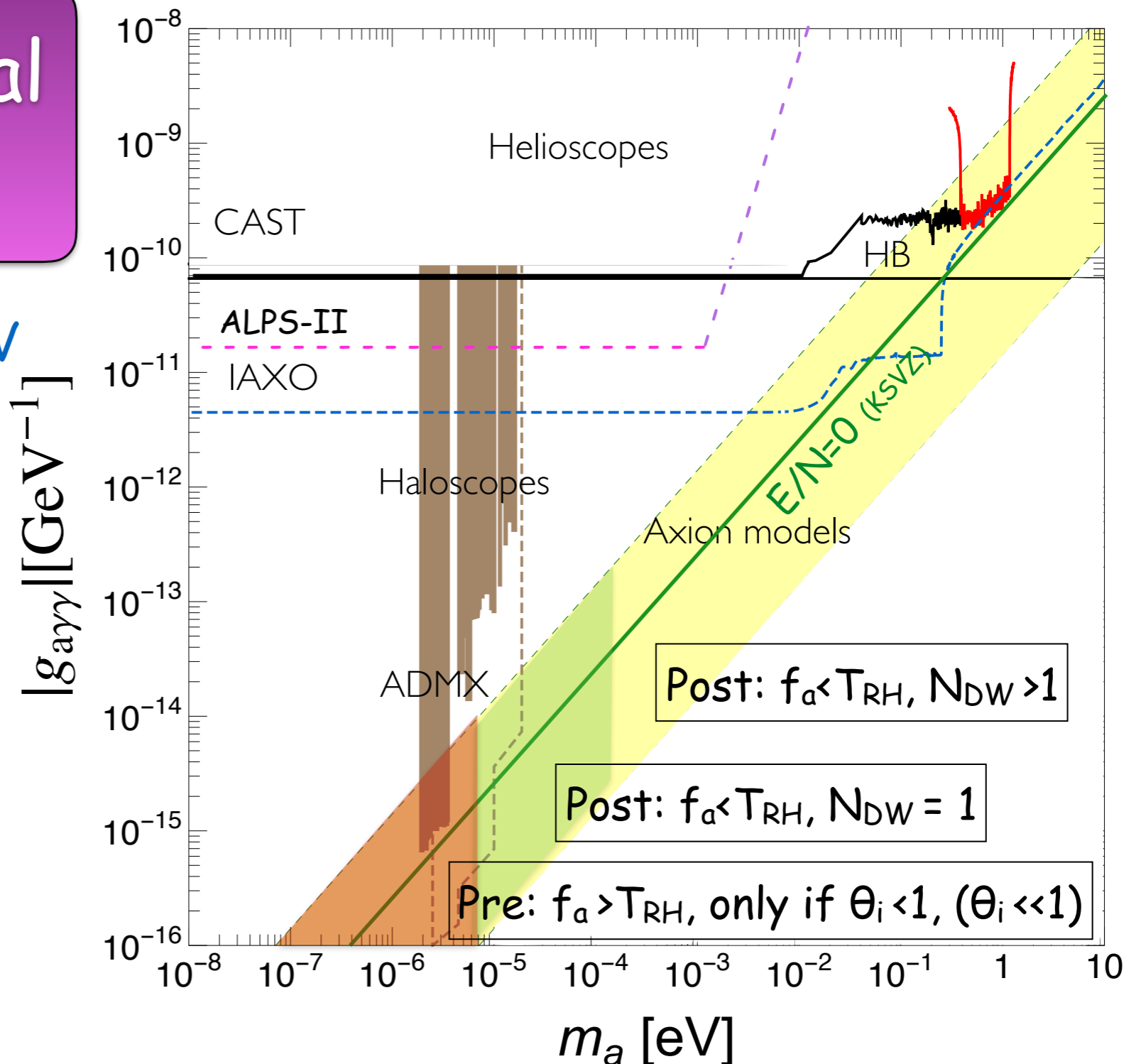
post-inflationary $\langle \theta_i \rangle \simeq 2, N_{\text{DW}} = 1$

Larger: $m_a \gtrsim \text{few } 100 \mu\text{eV}$

$\Omega_a \approx \Omega_{\text{CDM}}$ needs $N_{\text{DW}} > 1$
contributions (cosmological issues)

Smaller: $m_a \in (0.1, \text{few}) \mu\text{eV}$

pre-inflationary $|\theta_i| \simeq (0.2, 2)$



We need to know where to search

$\Omega_a \approx \Omega_{\text{CDM}}$ canonical
mass windows:

Preferred: $m_a \in (10, 100) \mu\text{eV}$

post-inflationary $\langle \theta_i \rangle \simeq 2$, $N_{\text{DW}} = 1$

Larger: $m_a \gtrsim \text{few } 100 \mu\text{eV}$

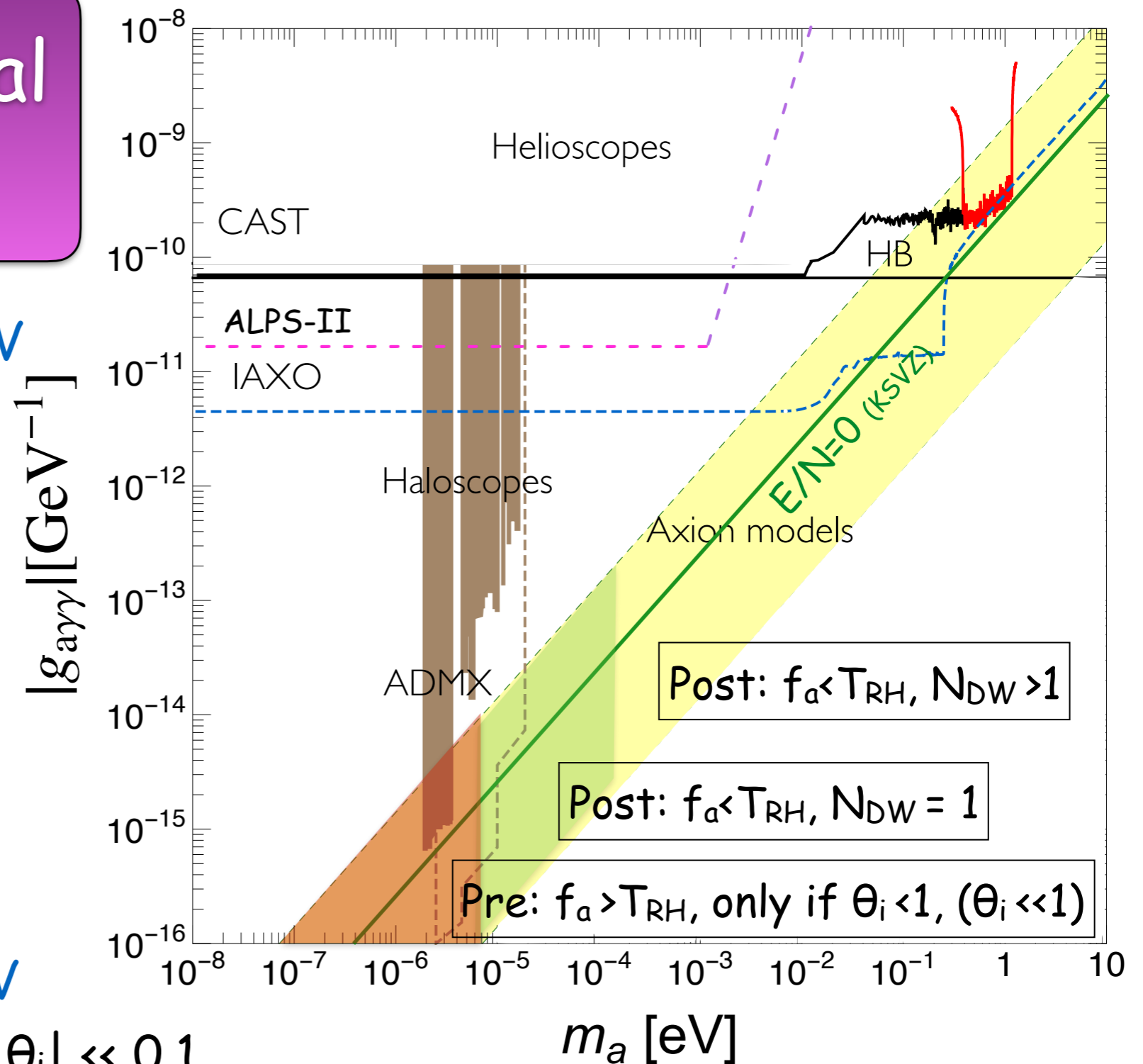
$\Omega_a \approx \Omega_{\text{CDM}}$ needs $N_{\text{DW}} > 1$
contributions (cosmological issues)

Smaller: $m_a \in (0.1, \text{few}) \mu\text{eV}$

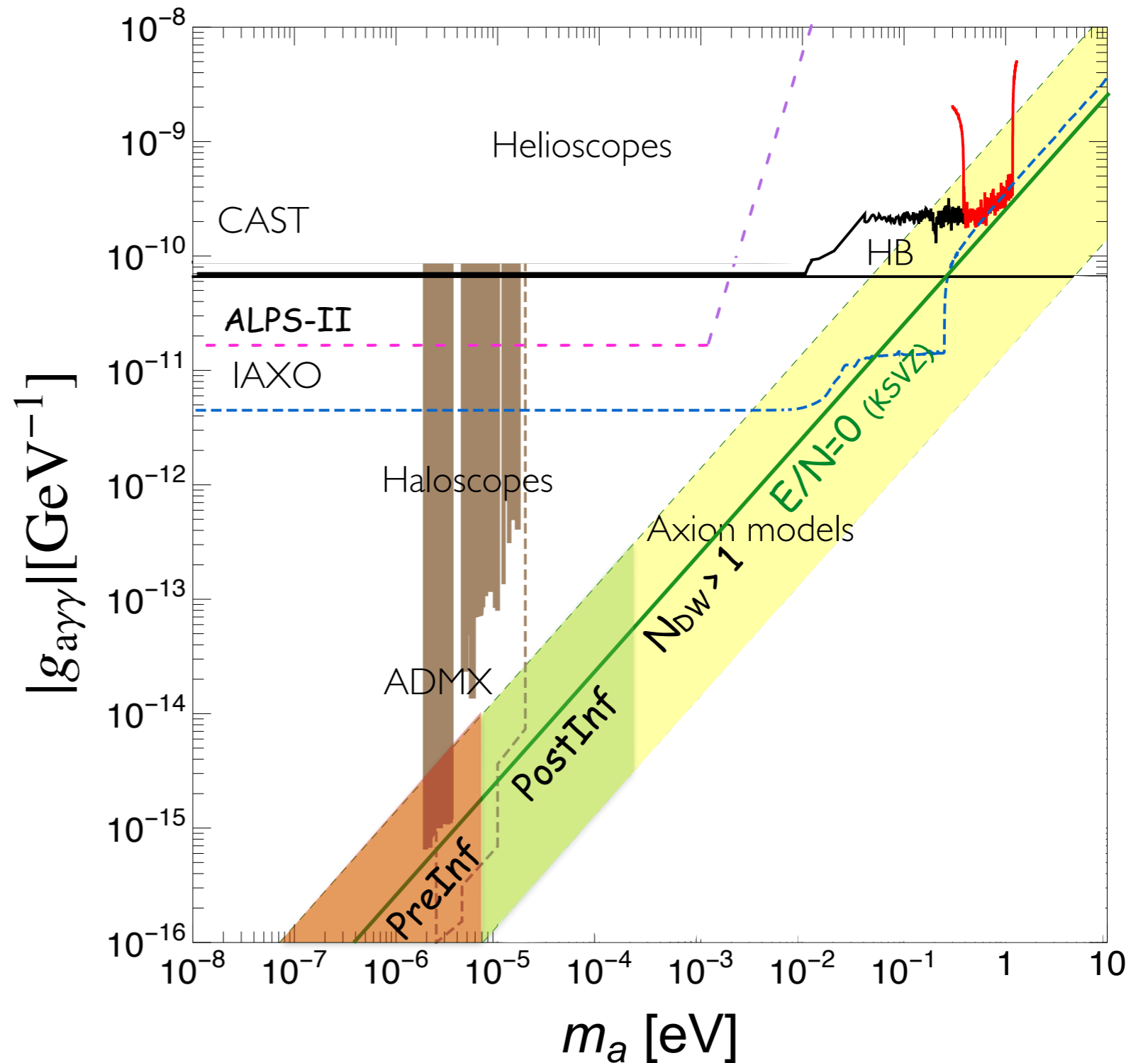
pre-inflationary $|\theta_i| \simeq (0.2, 2)$

Much smaller: $m_a \ll 0.1 \mu\text{eV}$

pre-inflationary, "anthropic" window: $|\theta_i| \ll 0.1$



Sometimes limits from astrophysics on $g_{ae}, g_{aN}, g_{a\pi}$ are fed into $g_{a\gamma}-m_a$ plots

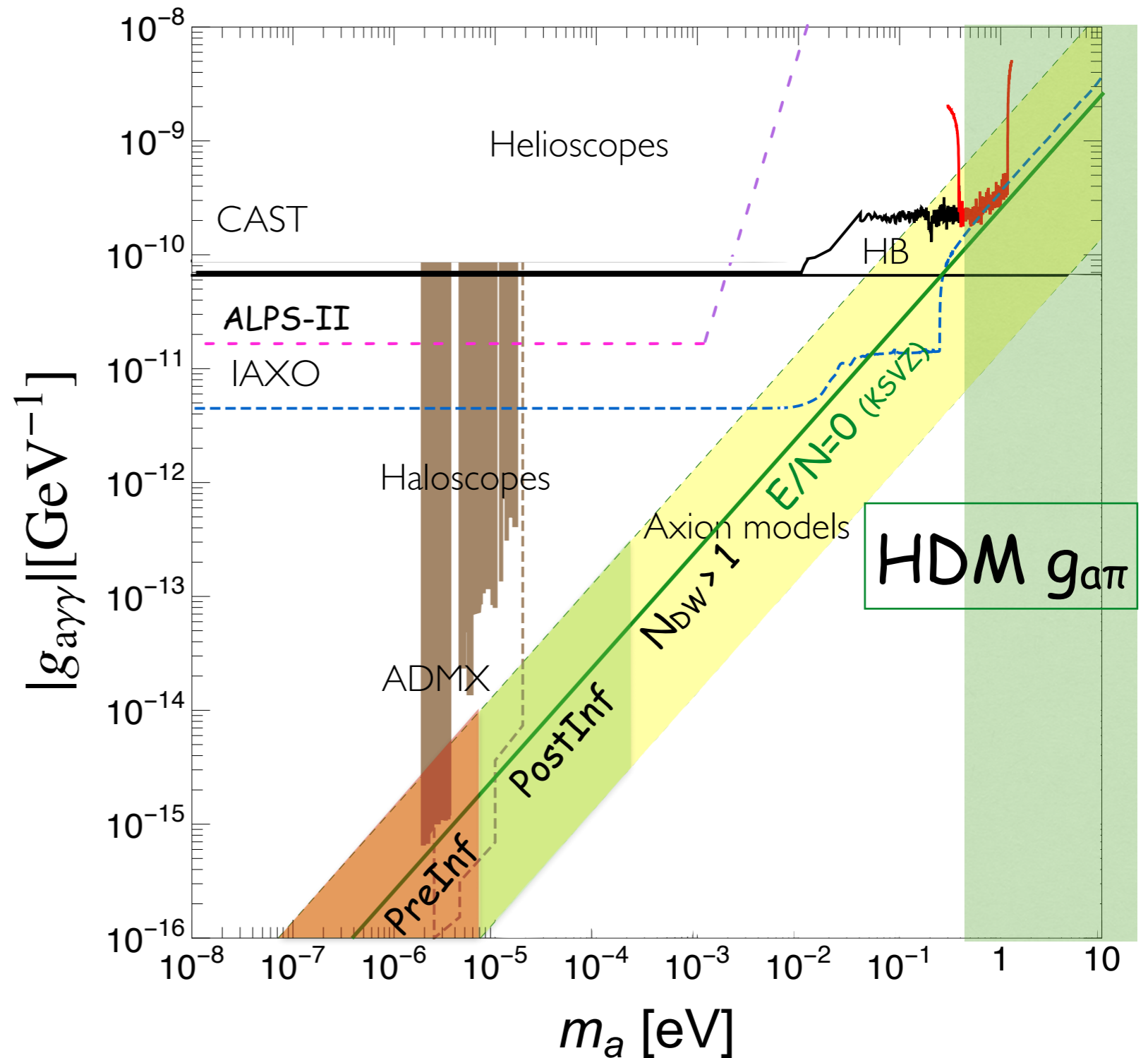


Sometimes limits from astrophysics on $g_{ae}, g_{aN}, g_{a\pi}$ are fed into $g_{a\gamma}-m_a$ plots

Hot DM ($\pi\pi \leftrightarrow \pi a$):

$$g_{a\pi} \sim (f_a f_\pi)^{-1} \lesssim 2 \times 10^{-7} \text{ GeV}^{-2}$$

$$g_{a\pi} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$



Sometimes limits from astrophysics on $g_{ae}, g_{aN}, g_{a\pi}$ are fed into $g_{a\gamma}-m_a$ plots

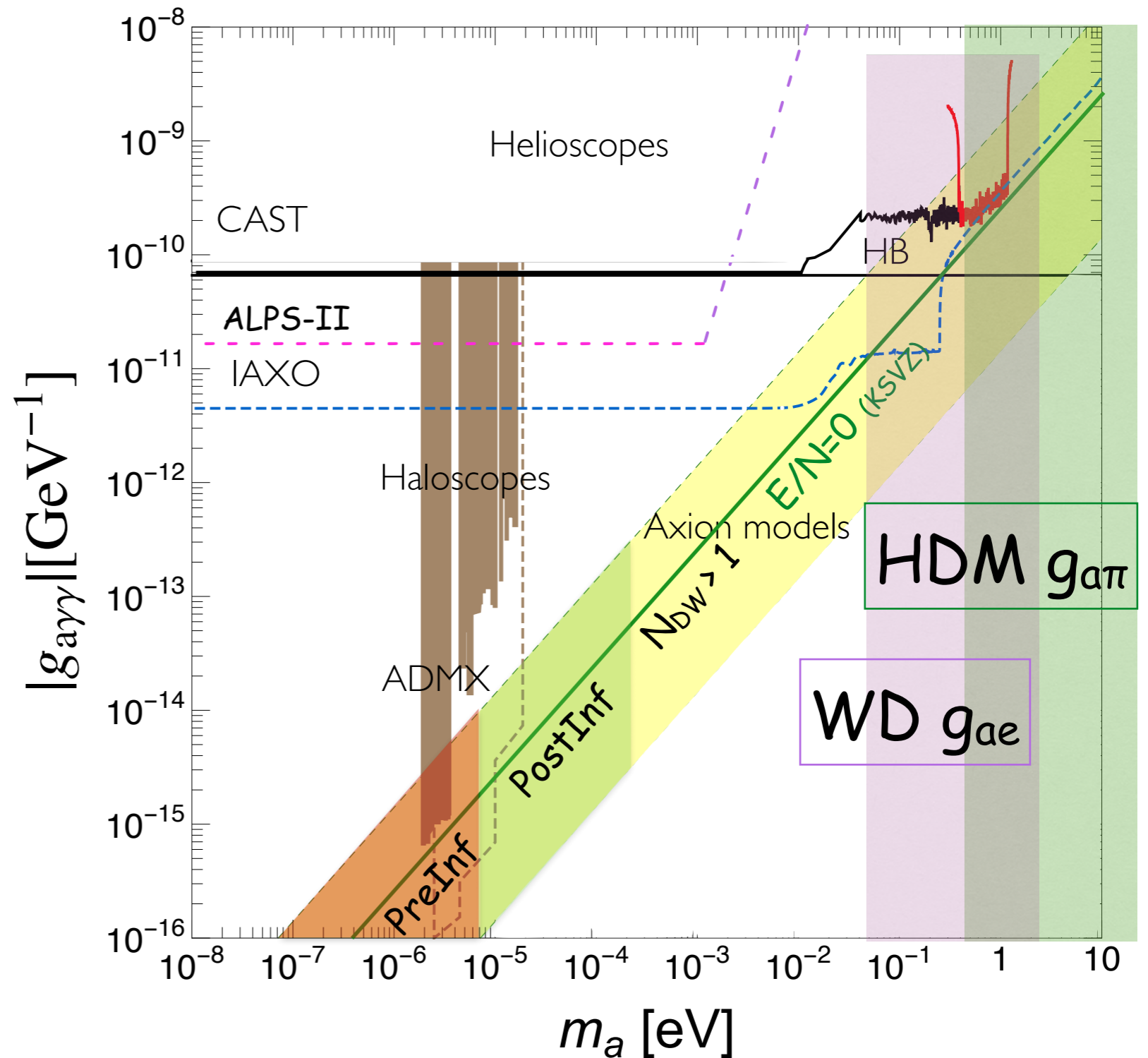
Hot DM ($\pi\pi \leftrightarrow \pi a$):

$$g_{a\pi} \sim (f_a f_\pi)^{-1} \lesssim 2 \times 10^{-7} \text{ GeV}^{-2}$$

$$g_{a\pi} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

WD/RG: $g_{ae} \lesssim 3 \times 10^{-13}$

$$g_{ae} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$



Sometimes limits from astrophysics on $g_{ae}, g_{aN}, g_{a\pi}$ are fed into $g_{a\gamma}-m_a$ plots

Hot DM ($\pi\pi \leftrightarrow \pi a$):

$$g_{a\pi} \sim (f_a f_\pi)^{-1} \lesssim 2 \times 10^{-7} \text{ GeV}^{-2}$$

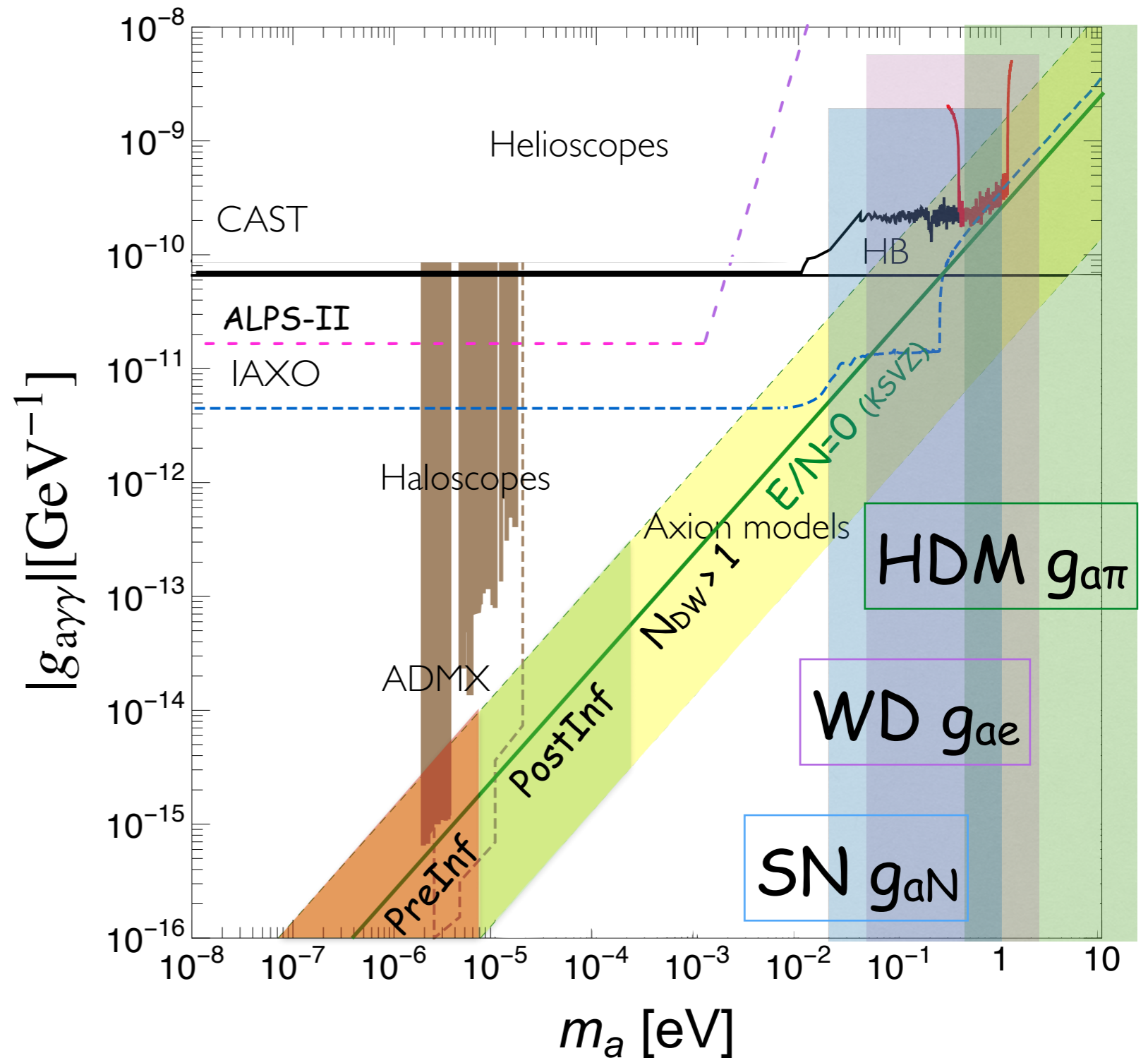
$$g_{a\pi} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

WD/RG: $g_{ae} \lesssim 3 \times 10^{-13}$

$$g_{ae} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

SN1987A: $g_{aN} \lesssim 0.9 \times 10^{-9}$
(see P.Carenza talk)

$$g_{aN} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$



Sometimes limits from astrophysics on $g_{ae}, g_{aN}, g_{a\pi}$ are fed into $g_{a\gamma}-m_a$ plots

Hot DM ($\pi\pi \leftrightarrow \pi a$):

$$g_{a\pi} \sim (f_a f_\pi)^{-1} \lesssim 2 \times 10^{-7} \text{ GeV}^{-2}$$

$$g_{a\pi} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

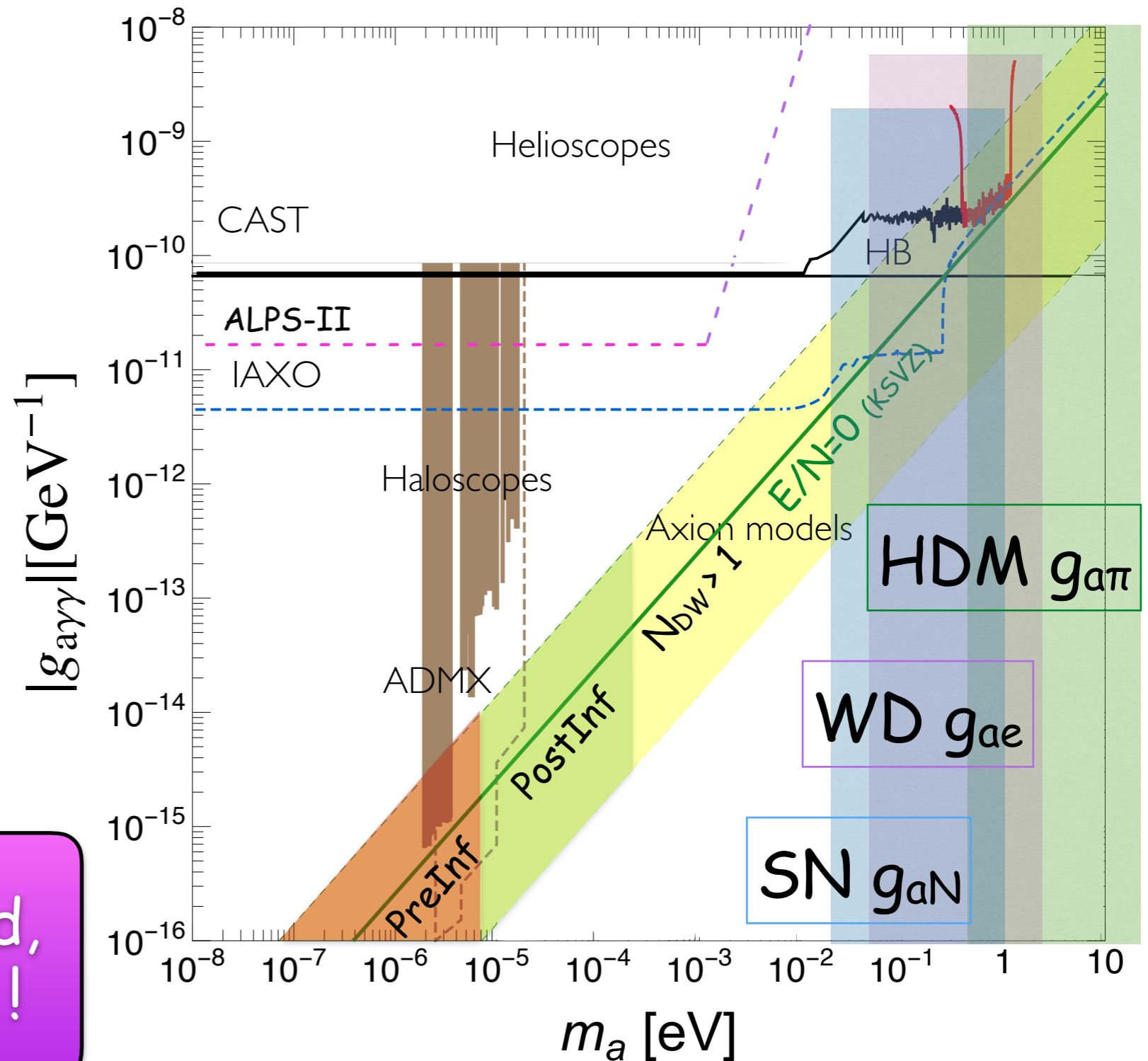
WD/RG: $g_{ae} \lesssim 3 \times 10^{-13}$

$$g_{ae} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

SN1987A: $g_{aN} \lesssim 0.9 \times 10^{-9}$
(see P.Carenza talk)

$$g_{aN} \xrightarrow{\text{(model)}} f_a \xrightarrow{\text{(...)}} m_a$$

This should be avoided,
it can be misleading!



Axion-photon coupling $g_{a\gamma}$: KSVZ-type

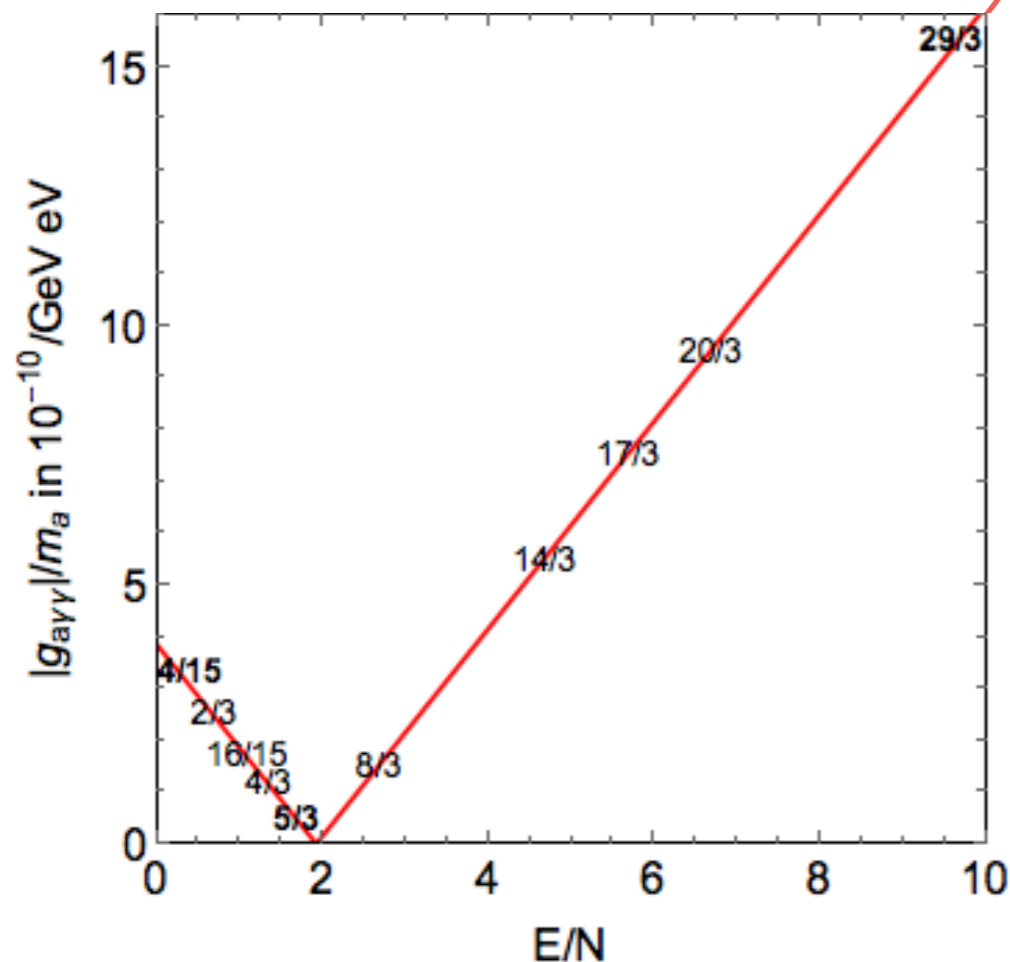
Original KSVZ: new Q's SM singlets: $E=0$
Simplest, but cosmo issues with Q stability

$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92(4) \right)$$

Axion-photon coupling $g_{a\gamma}$: KSVZ-type

Original KSVZ: new Q's SM singlets: $E=0$
Simplest, but cosmo issues with Q stability

Q in generic SM irreps.:
Only 15 are safe(τ_Q + no LP):



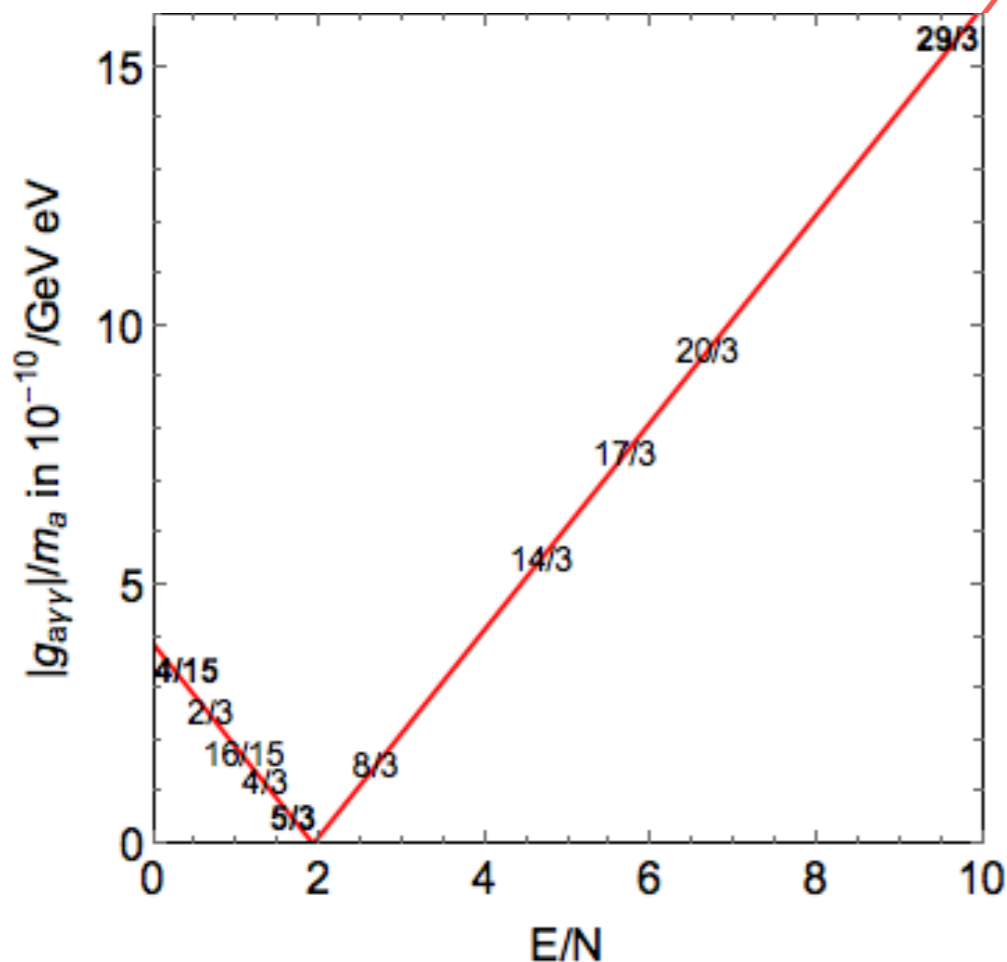
$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92(4) \right)$$

R_Q	\mathcal{O}_{Qq}	$\Lambda_{\text{Landau}}^{2\text{-loop}} [\text{GeV}]$	E/N
$(3, 1, -1/3)$	$\bar{Q}_L d_R$	$9.3 \cdot 10^{38} (g_1)$	$2/3$
$(3, 1, 2/3)$	$\bar{Q}_L u_R$	$5.4 \cdot 10^{34} (g_1)$	$8/3$
$(3, 2, 1/6)$	$\bar{Q}_R q_L$	$6.5 \cdot 10^{39} (g_1)$	$5/3$
$(3, 2, -5/6)$	$\bar{Q}_L d_R H^\dagger$	$4.3 \cdot 10^{27} (g_1)$	$17/3$
$(3, 2, 7/6)$	$\bar{Q}_L u_R H$	$5.6 \cdot 10^{22} (g_1)$	$29/3$
$(3, 3, -1/3)$	$\bar{Q}_R q_L H^\dagger$	$5.1 \cdot 10^{30} (g_2)$	$14/3$
$(3, 3, 2/3)$	$\bar{Q}_R q_L H$	$6.6 \cdot 10^{27} (g_2)$	$20/3$
$(3, 3, -4/3)$	$\bar{Q}_L d_R H^{\dagger 2}$	$3.5 \cdot 10^{18} (g_1)$	$44/3$
$(\bar{6}, 1, -1/3)$	$\bar{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$2.3 \cdot 10^{37} (g_1)$	$4/15$
$(\bar{6}, 1, 2/3)$	$\bar{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$5.1 \cdot 10^{30} (g_1)$	$16/15$
$(\bar{6}, 2, 1/6)$	$\bar{Q}_R \sigma_{\mu\nu} q_L G^{\mu\nu}$	$7.3 \cdot 10^{38} (g_1)$	$2/3$
$(8, 1, -1)$	$\bar{Q}_L \sigma_{\mu\nu} e_R G^{\mu\nu}$	$7.6 \cdot 10^{22} (g_1)$	$8/3$
$(8, 2, -1/2)$	$\bar{Q}_R \sigma_{\mu\nu} \ell_L G^{\mu\nu}$	$6.7 \cdot 10^{27} (g_1)$	$4/3$
$(15, 1, -1/3)$	$\bar{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$8.3 \cdot 10^{21} (g_3)$	$1/6$
$(15, 1, 2/3)$	$\bar{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$7.6 \cdot 10^{21} (g_3)$	$2/3$

Axion-photon coupling $g_{a\gamma}$: KSVZ-type

Original KSVZ: new Q's SM singlets: $E=0$
Simplest, but cosmo issues with Q stability

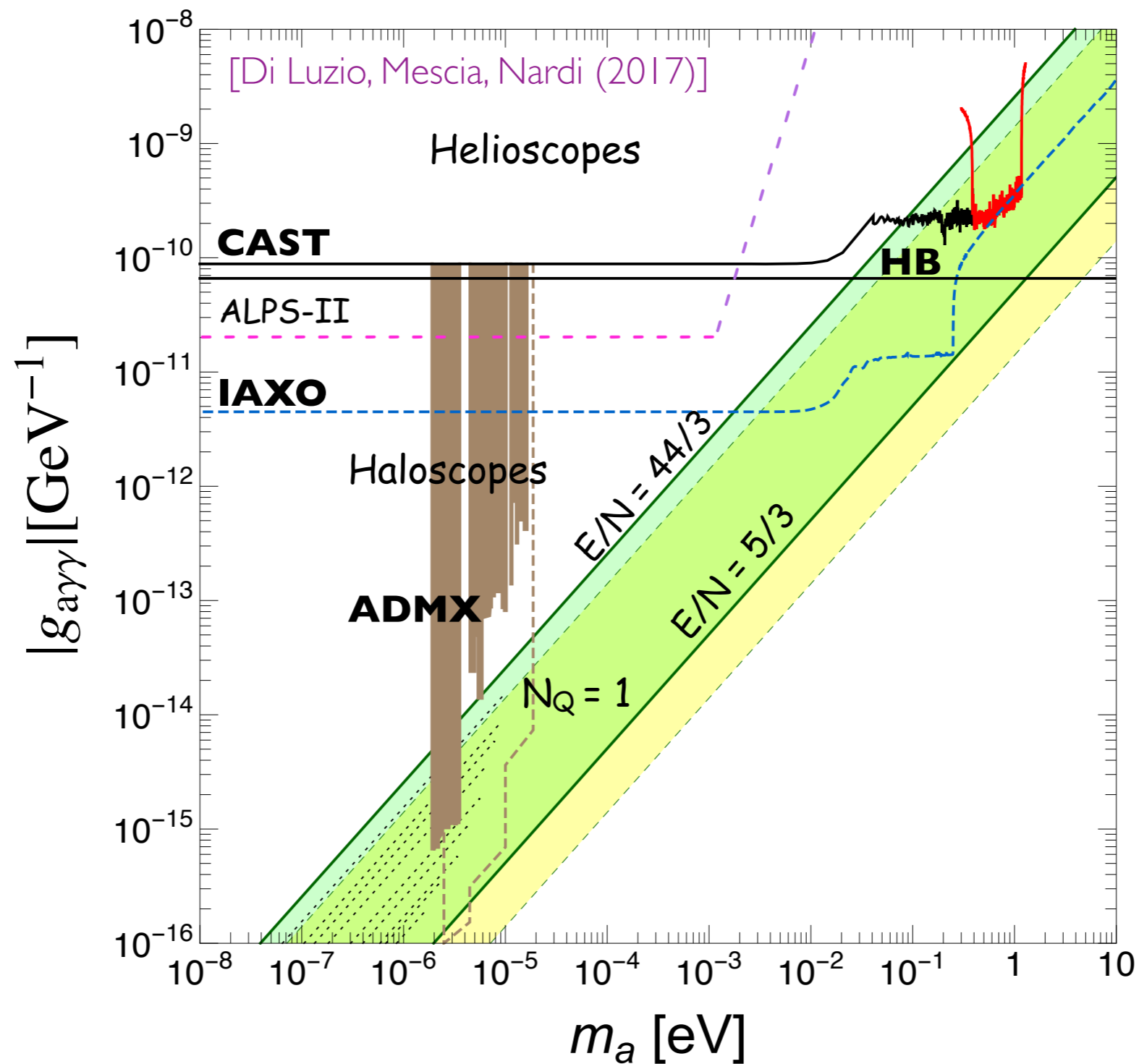
Q in generic SM irreps.:
Only 15 are safe(τ_Q + no LP):



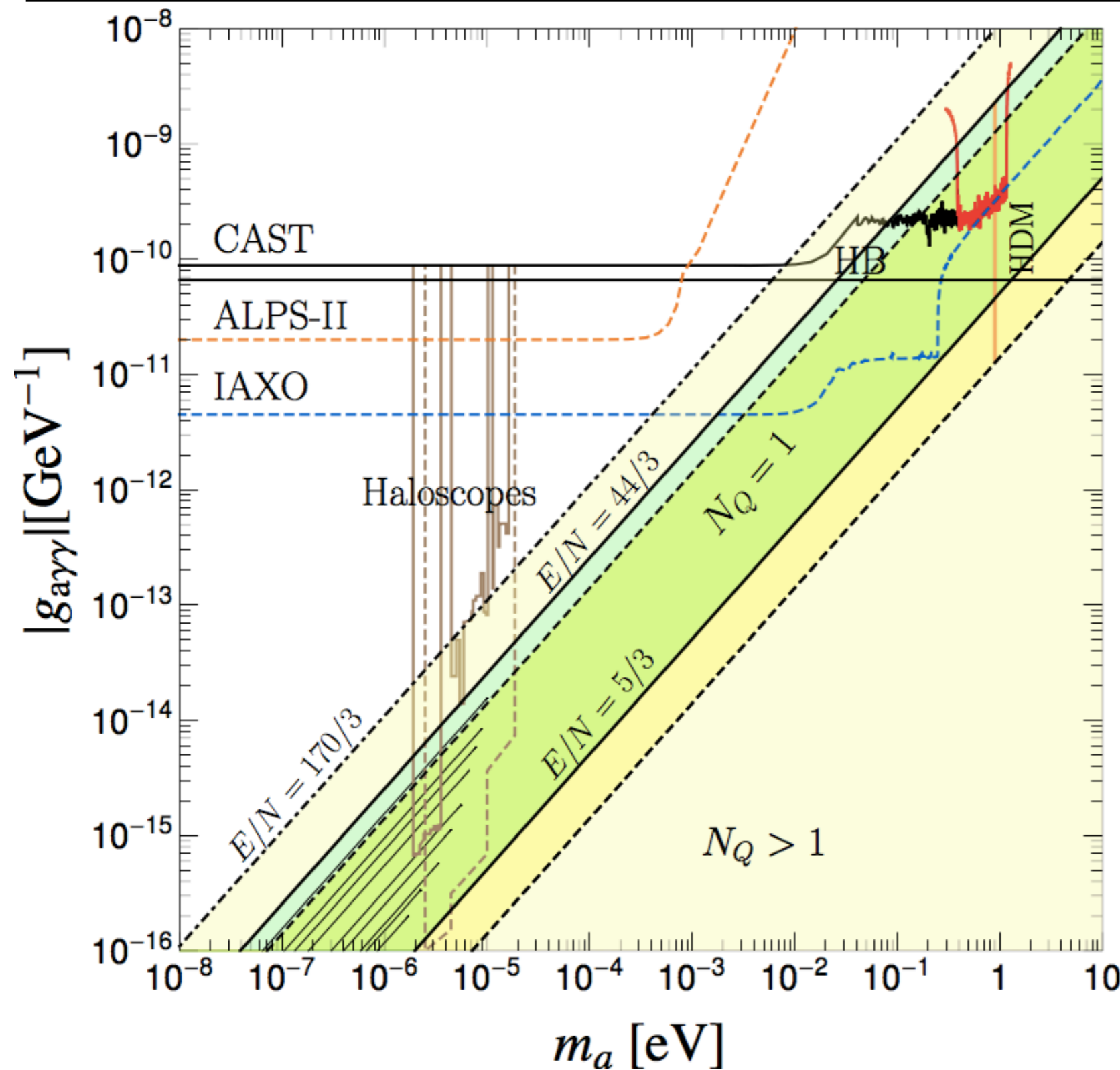
$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92(4) \right)$$

R_Q	\mathcal{O}_{Qq}	$\Lambda_{\text{Landau}}^{2-\text{loop}}[\text{GeV}]$	E/N	
R_Q^w	$(3, 1, -1/3)$	$\overline{Q}_L d_R$	$9.3 \cdot 10^{38}(g_1)$	2/3
	$(3, 1, 2/3)$	$\overline{Q}_L u_R$	$5.4 \cdot 10^{34}(g_1)$	8/3
	$(3, 2, 1/6)$	$\overline{Q}_R q_L$	$6.5 \cdot 10^{39}(g_1)$	5/3
	$(3, 2, -5/6)$	$\overline{Q}_L d_R H^\dagger$	$4.3 \cdot 10^{27}(g_1)$	17/3
	$(3, 2, 7/6)$	$\overline{Q}_L u_R H$	$5.6 \cdot 10^{22}(g_1)$	29/3
	$(3, 3, -1/3)$	$\overline{Q}_R q_L H^\dagger$	$5.1 \cdot 10^{30}(g_2)$	14/3
R_Q^s	$(3, 3, 2/3)$	$\overline{Q}_R q_L H$	$6.6 \cdot 10^{27}(g_2)$	20/3
	$(3, 3, -4/3)$	$\overline{Q}_L d_R H^{\dagger 2}$	$3.5 \cdot 10^{18}(g_1)$	44/3
	$(\overline{6}, 1, -1/3)$	$\overline{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$2.3 \cdot 10^{37}(g_1)$	4/15
	$(\overline{6}, 1, 2/3)$	$\overline{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$5.1 \cdot 10^{30}(g_1)$	16/15
	$(\overline{6}, 2, 1/6)$	$\overline{Q}_R \sigma_{\mu\nu} q_L G^{\mu\nu}$	$7.3 \cdot 10^{38}(g_1)$	2/3
	$(8, 1, -1)$	$\overline{Q}_L \sigma_{\mu\nu} e_R G^{\mu\nu}$	$7.6 \cdot 10^{22}(g_1)$	8/3
	$(8, 2, -1/2)$	$\overline{Q}_R \sigma_{\mu\nu} \ell_L G^{\mu\nu}$	$6.7 \cdot 10^{27}(g_1)$	4/3
	$(15, 1, -1/3)$	$\overline{Q}_L \sigma_{\mu\nu} d_R G^{\mu\nu}$	$8.3 \cdot 10^{21}(g_3)$	1/6
	$(15, 1, 2/3)$	$\overline{Q}_L \sigma_{\mu\nu} u_R G^{\mu\nu}$	$7.6 \cdot 10^{21}(g_3)$	2/3

Redefining the KSVZ axion window

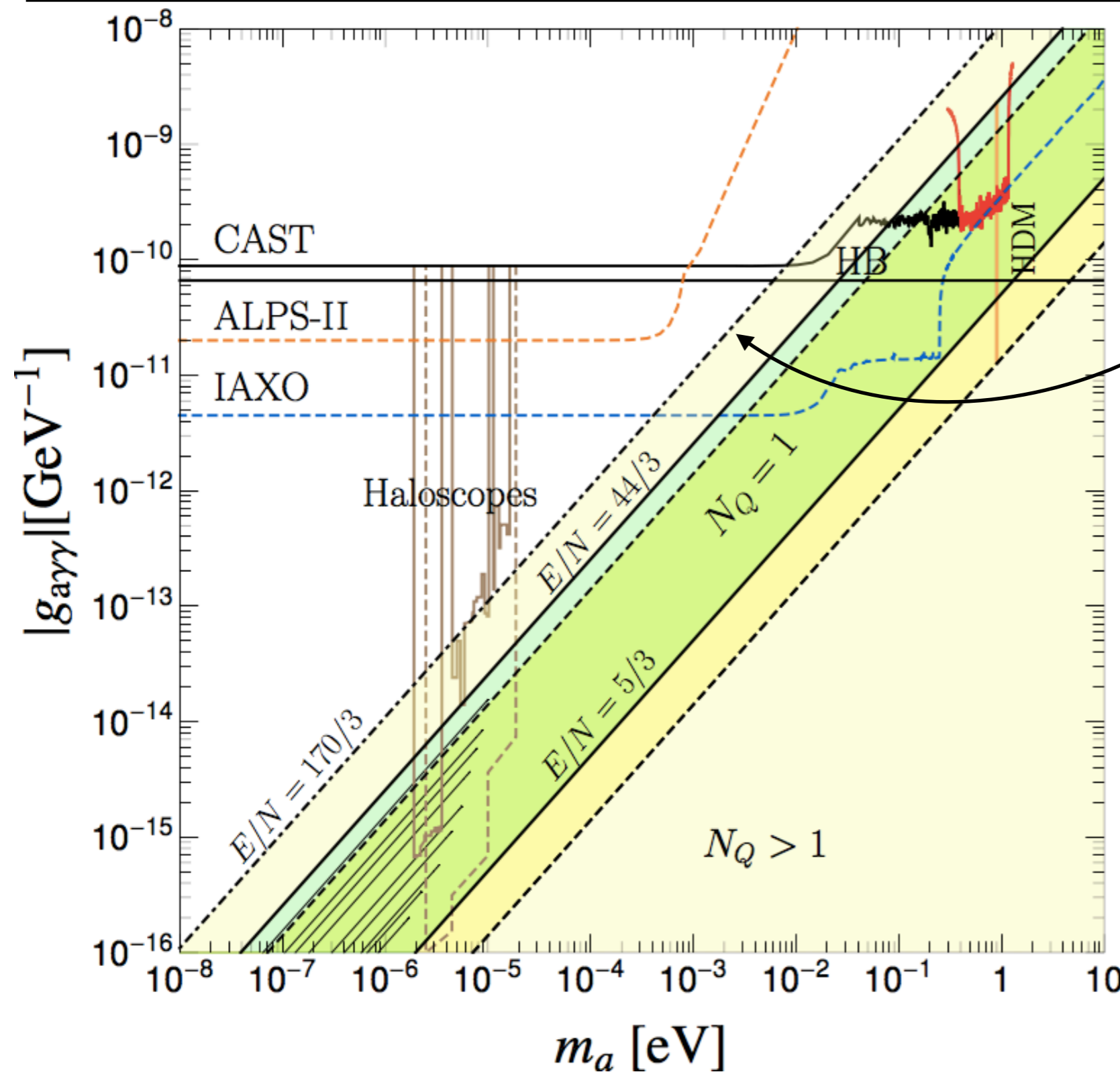


KSVZ: More Q representations



[Di Luzio, Mescia, Nardi (2017)]

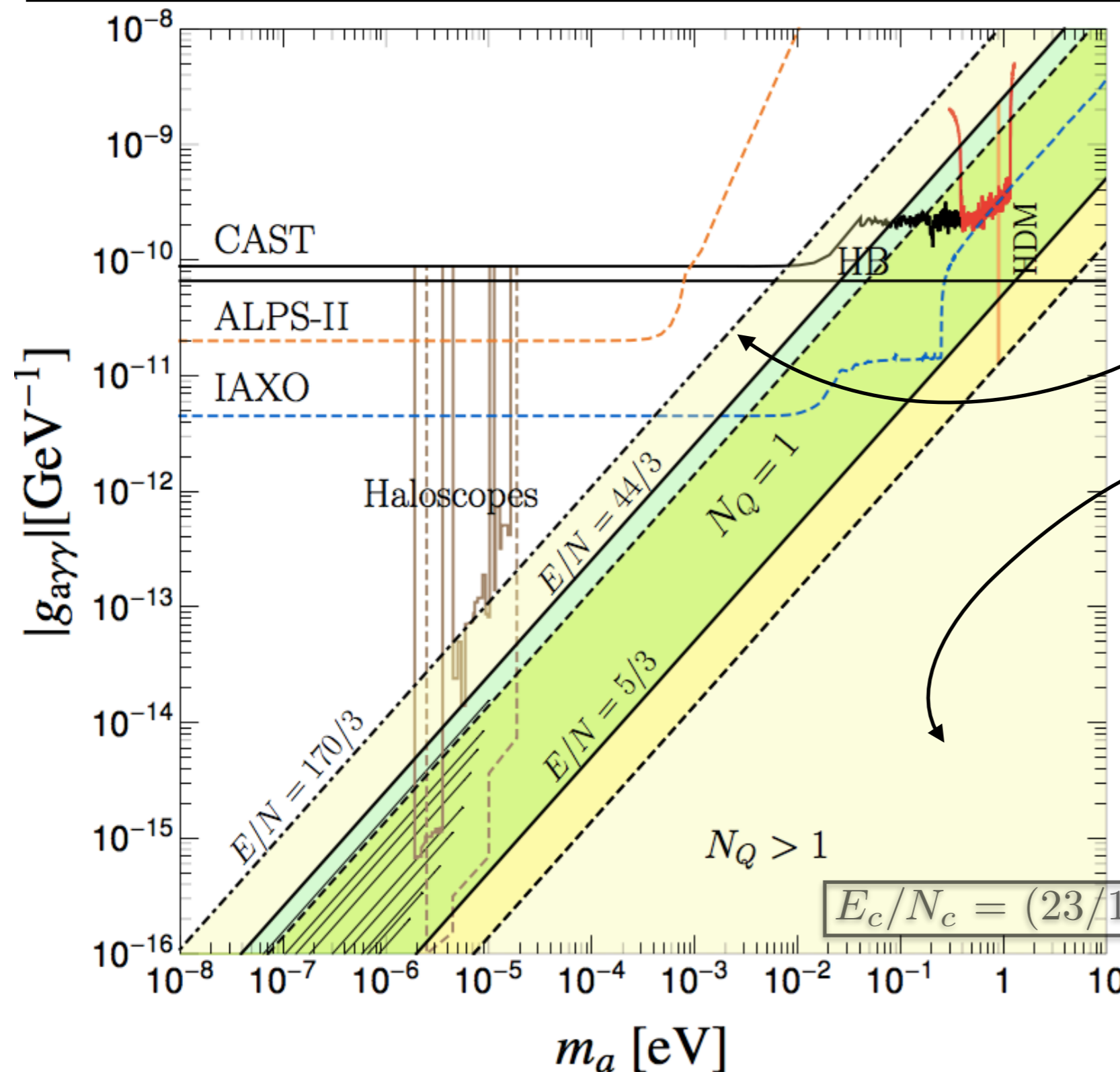
KSVZ: More Q representations



[Di Luzio, Mescia, Nardi (2017)]

The strongest possible coupling $E_c/N_c = 170/3$ ($3 R_Q$)

KSVZ: More Q representations



[Di Luzio, Mescia, Nardi (2017)]

The strongest possible coupling $E_c/N_c=170/3$ ($3 R_Q$)

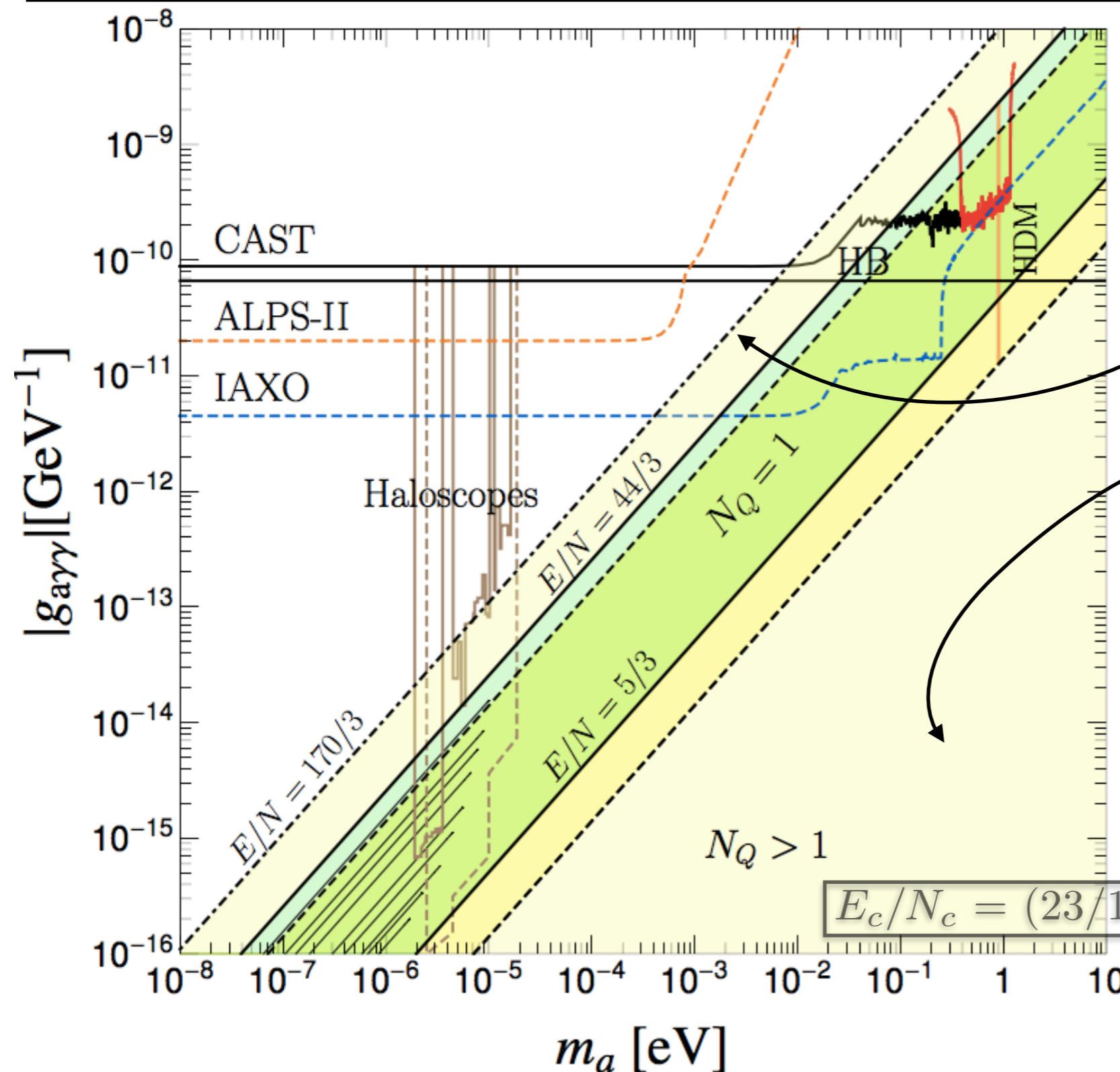
$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92(4) \right)$$

Axion-photon decoupling
(within theoretical errors)
can also occur !

$$\left. \begin{aligned} (3, 3, -1/3) \oplus (\bar{6}, 1, -1/3) \\ (\bar{6}, 1, 2/3) \oplus (8, 1, -1) \\ (3, 2, -5/6) \oplus (8, 2, -1/2) \end{aligned} \right\}$$

$$E_c/N_c = (23/12, 64/33, 41/21) \approx (1.92, 1.94, 1.95)$$

KSVZ: More Q representations



[Di Luzio, Mescia, Nardi (2017)]

The strongest possible coupling $E_c/N_c=170/3$ ($3 R_Q$)

$$g_{a\gamma\gamma} = \frac{m_a}{\text{eV}} \frac{2.0}{10^{10} \text{ GeV}} \left(\frac{E}{N} - 1.92(4) \right)$$

Axion-photon decoupling
(within theoretical errors)
can also occur !

$$\left. \begin{aligned} (3, 3, -1/3) \oplus (\bar{6}, 1, -1/3) \\ (\bar{6}, 1, 2/3) \oplus (8, 1, -1) \\ (3, 2, -5/6) \oplus (8, 2, -1/2) \end{aligned} \right\}$$

$$E_c/N_c = (23/12, 64/33, 41/21) \approx (1.92, 1.94, 1.95)$$

"Immoral but not fine tuned"

Axion-photon coupling $g_{a\gamma}$: DFSZ-type

DFSZ: Two (or more) Higgs doublet model plus one scalar singlet Φ

Axion-photon coupling $g_{a\gamma}$: DFSZ-type

DFSZ: Two (or more) Higgs doublet model plus one scalar singlet Φ

- In general each R-handed SM fermion can have a specific PQ charge X_{fj}

$$u_R^j \rightarrow \exp(iX_{uj}) u_R^j,$$

$$d_R^j \rightarrow \exp(iX_{dj}) d_R^j,$$

$$e_R^j \rightarrow \exp(iX_{ej}) e_R^j.$$

$$\frac{E}{N} = \frac{2}{3} + 2 \frac{\sum_j (X_{uj} + X_{ej})}{\sum_j (X_{uj} + X_{dj})}.$$

Axion-photon coupling $g_{a\gamma}$: DFSZ-type

DFSZ: Two (or more) Higgs doublet model plus one scalar singlet Φ

- In general each R-handed SM fermion can have a specific PQ charge X_{fj}

$$u_R^j \rightarrow \exp(iX_{uj}) u_R^j,$$

$$d_R^j \rightarrow \exp(iX_{dj}) d_R^j,$$

$$e_R^j \rightarrow \exp(iX_{ej}) e_R^j.$$

$$\frac{E}{N} = \frac{2}{3} + 2 \frac{\sum_j (X_{uj} + X_{ej})}{\sum_j (X_{uj} + X_{dj})}.$$

- For generation independent charges DFSZ remains within KSVZ window:

$$\text{DFSZ-I: } X_e = X_d, \quad E/N = 8/3$$

$$\text{DFSZ-II: } X_e = -X_u, \quad E/N = 2/3$$

$$\text{DFSZ-III: } X_e \neq X_{u,d}, \quad E/N_{(max)} = -4/3$$

$$\text{DFSZ-IV: } N_H=9 \quad (E/N)_{\max} = 524/3.$$

Axion-photon coupling $g_{a\gamma}$: DFSZ-type

DFSZ: Two (or more) Higgs doublet model plus one scalar singlet Φ

- In general each R-handed SM fermion can have a specific PQ charge X_{fj}

$$u_R^j \rightarrow \exp(iX_{uj}) u_R^j,$$

$$d_R^j \rightarrow \exp(iX_{dj}) d_R^j,$$

$$e_R^j \rightarrow \exp(iX_{ej}) e_R^j.$$

$$\frac{E}{N} = \frac{2}{3} + 2 \frac{\sum_j (X_{uj} + X_{ej})}{\sum_j (X_{uj} + X_{dj})}$$

- For generation independent charges DFSZ remains within KSVZ window:

DFSZ-I: $X_e = X_d, \quad E/N = 8/3$

DFSZ-II: $X_e = -X_u, \quad E/N = 2/3$

DFSZ-III: $X_e \neq X_{u,d}, \quad E/N_{(max)} = -4/3$

DFSZ-IV: $N_H=9 \quad (E/N)_{max} = 524/3.$

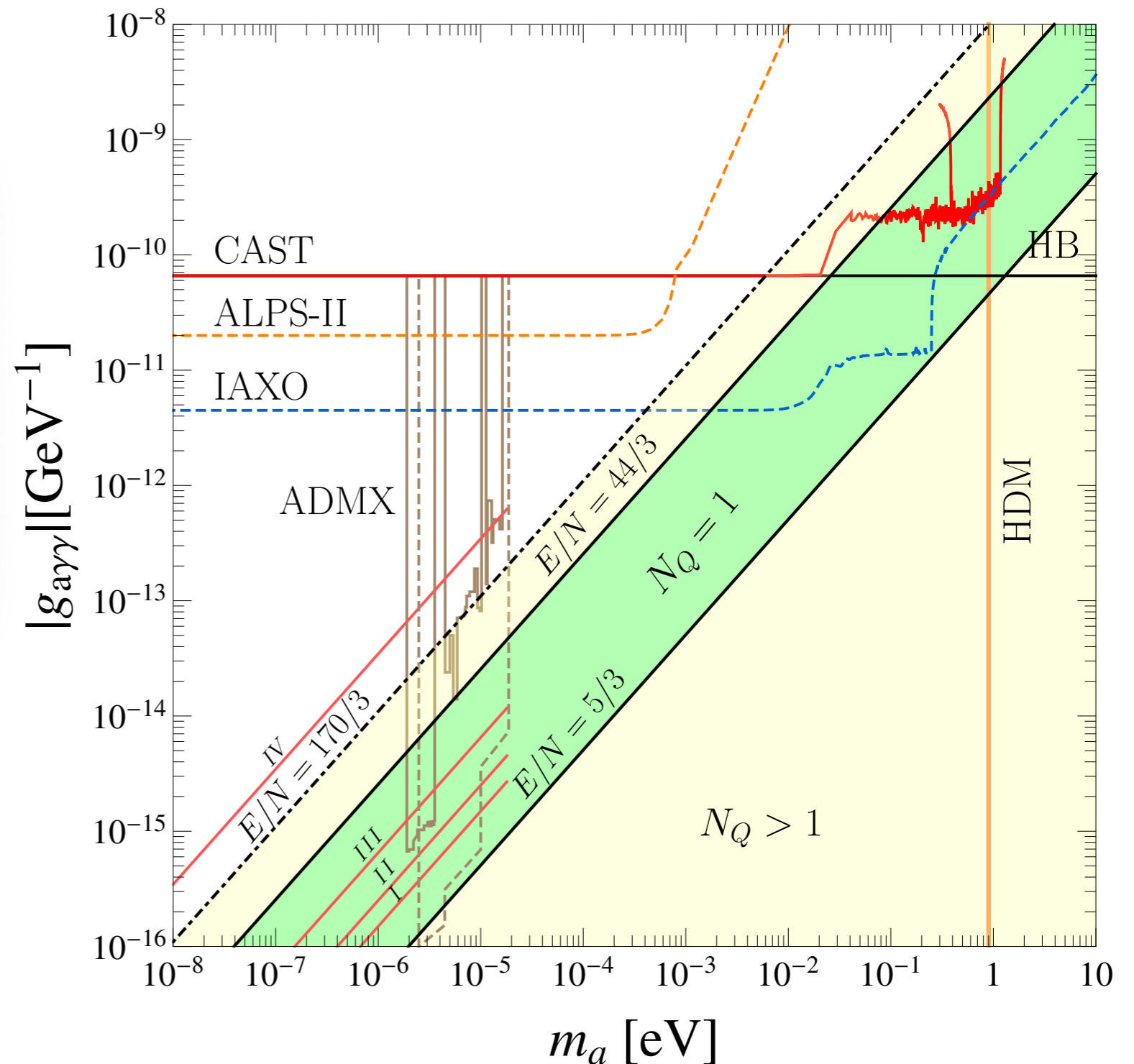
	$E/N(g_{a\gamma\gamma}^{max})$	$E/N(g_{a\gamma\gamma}^{min})$
KSVZ ($N_Q = 1$)	44/3	5/3
KSVZ ($N_Q > 1$)	170/3	23/12
DFSZ-I-II ($n_H = 2$)	2/3	8/3
DFSZ-III ($n_H = 3$)	-4/3	8/3
DFSZ-IV ($n_H = 9$)	524/3	23/12

- For generation dependent charges (max. of 9 Higgs doublets H_{fj}):

DFSZ-IV ($X_{ej} \ X_{dj}, X_{uj}$): $E/N_{(max)} = 524/3 = 3 \cdot E/N_{(max)}(\text{KSVZ})$

KSVZ/DFSZ: enlarged $g_{a\gamma}$ window

The region where
generic KSVZ
and DFSZ axion
models can live:



Suppressing axion-nucleon coupling: $g_{aN} \approx 0$ (DFSZ)

Suppressing axion-nucleon coupling: $g_{aN} \approx 0$ (DFSZ)

From the UV
theory we have:

$$\mathcal{L}_q = \frac{\partial_\mu a}{2f_a} c_q \bar{q} \gamma^\mu \gamma_5 q$$

Suppressing axion-nucleon coupling: $g_{aN} \approx 0$ (DFSZ)

From the UV
theory we have:

$$\mathcal{L}_q = \frac{\partial_\mu a}{2f_a} c_q \bar{q} \gamma^\mu \gamma_5 q$$

We want:

$$\mathcal{L}_N = \frac{\partial_\mu a}{2f_a} C_N \bar{N} \gamma^\mu \gamma_5 N$$

Suppressing axion-nucleon coupling: $g_{aN} \approx 0$ (DFSZ)

From the UV theory we have:

$$\mathcal{L}_q = \frac{\partial_\mu a}{2f_a} c_q \bar{q} \gamma^\mu \gamma_5 q$$

We want:

$$\mathcal{L}_N = \frac{\partial_\mu a}{2f_a} C_N \bar{N} \gamma^\mu \gamma_5 N$$

C_N in terms of c_q and of matrix elements $s^\mu \Delta_q = \langle N | \bar{q} \gamma^\mu \gamma_5 q | N \rangle$ by matching the matrix elements of \mathcal{L}_q and \mathcal{L}_N . One obtains:

$$\begin{aligned} (1) : \quad C_p + C_n &= (c_u + c_d) (\Delta_u + \Delta_d) - 2\delta_s & [\delta_s \approx O(10\%)] \\ (2) : \quad C_p - C_n &= (c_u - c_d) (\Delta_u - \Delta_d) \end{aligned}$$

Suppressing axion-nucleon coupling: $g_{aN} \approx 0$ (DFSZ)

From the UV theory we have:

$$\mathcal{L}_q = \frac{\partial_\mu a}{2f_a} c_q \bar{q} \gamma^\mu \gamma_5 q$$

We want:

$$\mathcal{L}_N = \frac{\partial_\mu a}{2f_a} C_N \bar{N} \gamma^\mu \gamma_5 N$$

C_N in terms of c_q and of matrix elements $s^\mu \Delta_q = \langle N | \bar{q} \gamma^\mu \gamma_5 q | N \rangle$ by matching the matrix elements of \mathcal{L}_q and \mathcal{L}_N . One obtains:

$$\begin{aligned} (1) : \quad C_p + C_n &= (c_u + c_d) (\Delta_u + \Delta_d) - 2\delta_s & [\delta_s \approx O(10\%)] \\ (2) : \quad C_p - C_n &= (c_u - c_d) (\Delta_u - \Delta_d) \end{aligned}$$

So that, independently of the matrix elements:

$$\begin{aligned} (1) : \quad C_p + C_n &\approx 0 & \text{if } c_u + c_d = 0 \\ (2) : \quad C_p - C_n &= 0 & \text{if } c_u - c_d = 0 \end{aligned}$$

First condition: $C_u + C_d \approx 0$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings:

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F \tilde{F}$$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings.

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F \tilde{F}$$

$$X_u = (X_{uR} - X_{uL})/2$$

$$\mathcal{L}_a \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \frac{a}{f_a} \frac{\alpha}{8\pi} \frac{E}{N} F\tilde{F} + \frac{\partial_\mu a}{v_a} [X_u \bar{u} \gamma^\mu \gamma_5 u + X_d \bar{d} \gamma^\mu \gamma_5 d]$$

$$\left(f_a = \frac{v_a}{2N} \right)$$

$$\frac{\partial_\mu a}{2f_a} \left[\frac{X_u}{N} \bar{u} \gamma^\mu \gamma_5 u + \frac{X_d}{N} \bar{d} \gamma^\mu \gamma_5 d \right]$$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings.

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F \tilde{F}$$

$$X_u = (X_{uR} - X_{uL})/2$$

$$\mathcal{L}_a \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \frac{a}{f_a} \frac{\alpha}{8\pi} \frac{E}{N} F\tilde{F} + \frac{\partial_\mu a}{v_a} [X_u \bar{u} \gamma^\mu \gamma_5 u + X_d \bar{d} \gamma^\mu \gamma_5 d]$$

$$\left(f_a = \frac{v_a}{2N} \right)$$

$$\frac{\partial_\mu a}{2f_a} \left[\frac{X_u}{N} \bar{u} \gamma^\mu \gamma_5 u + \frac{X_d}{N} \bar{d} \gamma^\mu \gamma_5 d \right]$$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings.

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F \tilde{F}$$

$$X_u = (X_{uR} - X_{uL})/2$$

$$\mathcal{L}_a \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \frac{a}{f_a} \frac{\alpha}{8\pi} \frac{E}{N} F\tilde{F} + \frac{\partial_\mu a}{v_a} [X_u \bar{u} \gamma^\mu \gamma_5 u + X_d \bar{d} \gamma^\mu \gamma_5 d]$$

$$\left(f_a = \frac{v_a}{2N}\right)$$

$$\frac{\partial_\mu a}{2f_a} \left[\frac{X_u}{N} \bar{u} \gamma^\mu \gamma_5 u + \frac{X_d}{N} \bar{d} \gamma^\mu \gamma_5 d \right]$$

model independent contributions

$$\frac{E}{N} \rightarrow \frac{E}{N} - 1.92(4); \quad \frac{X_u}{N} \rightarrow c_u = \frac{X_u}{N} - \frac{m_d}{m_d + m_u}; \quad \frac{X_d}{N} \rightarrow c_d = \frac{X_d}{N} - \frac{m_u}{m_d + m_u}$$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings.

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F \tilde{F}$$

$$X_u = (X_{uR} - X_{uL})/2$$

$$\mathcal{L}_a \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \frac{a}{f_a} \frac{\alpha}{8\pi} \frac{E}{N} F\tilde{F} + \frac{\partial_\mu a}{v_a} [X_u \bar{u} \gamma^\mu \gamma_5 u + X_d \bar{d} \gamma^\mu \gamma_5 d]$$

$$\left(f_a = \frac{v_a}{2N} \right)$$

$$\frac{\partial_\mu a}{2f_a} \left[\frac{X_u}{N} \bar{u} \gamma^\mu \gamma_5 u + \frac{X_d}{N} \bar{d} \gamma^\mu \gamma_5 d \right]$$

model independent contributions

$$\frac{E}{N} \rightarrow \frac{E}{N} - 1.92(4); \quad \frac{X_u}{N} \rightarrow c_u = \frac{X_u}{N} - \frac{m_d}{m_d + m_u}; \quad \frac{X_d}{N} \rightarrow c_d = \frac{X_d}{N} - \frac{m_u}{m_d + m_u}$$

Therefore:
$$c_u + c_d = \frac{X_u + X_d}{N} - 1$$

First condition: $C_u + C_d \approx 0$

Anomalous
axion couplings.

$$\frac{a}{v_a} \frac{g_s^2}{16\pi^2} N G\tilde{G}$$

$$\frac{a}{v_a} \frac{e^2}{16\pi^2} E F\tilde{F}$$

$$X_u = (X_{uR} - X_{uL})/2$$

$$\mathcal{L}_a \supset \frac{a}{f_a} \frac{\alpha_s}{8\pi} G\tilde{G} + \frac{a}{f_a} \frac{\alpha}{8\pi} \frac{E}{N} F\tilde{F} + \frac{\partial_\mu a}{v_a} [X_u \bar{u} \gamma^\mu \gamma_5 u + X_d \bar{d} \gamma^\mu \gamma_5 d]$$

$$\left(f_a = \frac{v_a}{2N}\right)$$

$$\frac{\partial_\mu a}{2f_a} \left[\frac{X_u}{N} \bar{u} \gamma^\mu \gamma_5 u + \frac{X_d}{N} \bar{d} \gamma^\mu \gamma_5 d \right]$$

model independent contributions

$$\frac{E}{N} \rightarrow \frac{E}{N} - 1.92(4); \quad \frac{X_u}{N} \rightarrow c_u = \frac{X_u}{N} - \frac{m_d}{m_d + m_u}; \quad \frac{X_d}{N} \rightarrow c_d = \frac{X_d}{N} - \frac{m_u}{m_d + m_u}$$

Therefore:

$$c_u + c_d = \frac{X_u + X_d}{N} - 1$$

$$\xrightarrow[N=n_g(X_u+X_d)]{\text{universality}} \quad \frac{1}{n_g} - 1 \neq 0$$

First conclusions from $C_u + C_d \approx 0$

First conclusions from $C_u + C_d \approx 0$

Nucleophobia unavoidably requires DFSZ-type of models with generation dependent PQ charges

such that the contribution to the anomaly from the two heavier generations vanishes: $N_{\text{tot}} = N_{(1^{\text{st}} \text{ gen})}$

First conclusions from $C_u + C_d \approx 0$

Nucleophobia unavoidably requires DFSZ-type of models with generation dependent PQ charges

such that the contribution to the anomaly from the two heavier generations vanishes: $N_{tot} = N_{(1^{st} \text{ gen})}$

E.g. in 1811.09637 (Bjorkeröth, Di Luzio, Mescia, EN) on U(1) flavour symmetries a number of such symmetries was serendipitously found

First conclusions from $C_u + C_d \approx 0$

Nucleophobia unavoidably requires DFSZ-type of models with generation dependent PQ charges

such that the contribution to the anomaly from the two heavier generations vanishes: $N_{\text{tot}} = N_{(1^{\text{st}} \text{ gen})}$

E.g. in 1811.09637 (Bjorkeröth, Di Luzio, Mescia, EN) on U(1) flavour symmetries a number of such symmetries was serendipitously found

Nucleophobia is not possible for KSVZ-type of models

Second Condition: $C_u - C_d \approx 0$

Second Condition: $C_u - C_d \approx 0$

Scalar content of DFSZ models: H_1, H_2, Φ_a with VEVs v_1, v_2, v_a ($v_1^2 + v_2^2 = v^2$) and PQ charges $X_1, X_2, X_a = (X_1 - X_2)(1/2)$

Second Condition: $c_u - c_d \approx 0$

Scalar content of DFSZ models: H_1, H_2, Φ_a with VEVs v_1, v_2, v_a ($v_1^2 + v_2^2 = v^2$) and PQ charges $X_1, X_2, X_a = (X_1 - X_2)(1/2)$

$$\begin{aligned} c_u - c_d &= \frac{(\mathcal{X}_{u_R} - \mathcal{X}_{u_L}) - (\mathcal{X}_{d_R} - \mathcal{X}_{d_L})}{2N_\ell} - \frac{m_d - m_u}{m_d + m_u} \\ &= -\frac{\mathcal{X}_1 + \mathcal{X}_2}{\mathcal{X}_2 - \mathcal{X}_1} - \left(\approx \frac{1}{3} \right) \end{aligned}$$

Second Condition: $c_u - c_d \approx 0$

Scalar content of DFSZ models: H_1, H_2, Φ_a with VEVs v_1, v_2, v_a ($v_1^2 + v_2^2 = v^2$) and PQ charges $X_1, X_2, X_a = (X_1 - X_2)(1/2)$

$$\begin{aligned} c_u - c_d &= \frac{(\mathcal{X}_{u_R} - \mathcal{X}_{u_L}) - (\mathcal{X}_{d_R} - \mathcal{X}_{d_L})}{2N_\ell} - \frac{m_d - m_u}{m_d + m_u} \\ &= -\frac{\mathcal{X}_1 + \mathcal{X}_2}{\mathcal{X}_2 - \mathcal{X}_1} - \left(\approx \frac{1}{3} \right) \end{aligned}$$

Goldstone of Hypercharge:

$$\varphi_Y = (v_2 \varphi_2 - v_1 \varphi_1)/v$$

Condition enforcing no a - φ_Y mixing:

$$\begin{aligned} \sum_i \mathcal{X}_i Y_i v_i^2 &= 0 \\ \Rightarrow \mathcal{X}_1 v_1^2 + \mathcal{X}_2 v_2^2 &= 0 \end{aligned}$$

Second Condition: $c_u - c_d \approx 0$

Scalar content of DFSZ models: H_1, H_2, Φ_a with VEVs v_1, v_2, v_a ($v_1^2 + v_2^2 = v^2$) and PQ charges $X_1, X_2, X_a = (X_1 - X_2)(1/2)$

$$\begin{aligned} c_u - c_d &= \frac{(\mathcal{X}_{u_R} - \mathcal{X}_{u_L}) - (\mathcal{X}_{d_R} - \mathcal{X}_{d_L})}{2N_\ell} - \frac{m_d - m_u}{m_d + m_u} \\ &= -\frac{\mathcal{X}_1 + \mathcal{X}_2}{\mathcal{X}_2 - \mathcal{X}_1} - \left(\approx \frac{1}{3} \right) \end{aligned}$$

Goldstone of Hypercharge:

$$\varphi_Y = (v_2 \varphi_2 - v_1 \varphi_1)/v$$

Condition enforcing no a - φ_Y mixing:

$$\begin{aligned} \sum_i \mathcal{X}_i Y_i v_i^2 &= 0 \\ \Rightarrow \mathcal{X}_1 v_1^2 + \mathcal{X}_2 v_2^2 &= 0 \end{aligned}$$

Couplings to the physical $a(x)$ are defined in terms of the charges

$$X_1 \propto -v_2^2/v^2 = -s_\beta^2$$

$$X_2 \propto v_1^2/v^2 = c_\beta^2$$

Second Condition: $c_u - c_d \approx 0$

Scalar content of DFSZ models: H_1, H_2, Φ_a with VEVs v_1, v_2, v_a ($v_1^2 + v_2^2 = v^2$) and PQ charges $X_1, X_2, X_a = (X_1 - X_2)(1/2)$

$$\begin{aligned} c_u - c_d &= \frac{(\mathcal{X}_{u_R} - \mathcal{X}_{u_L}) - (\mathcal{X}_{d_R} - \mathcal{X}_{d_L})}{2N_\ell} - \frac{m_d - m_u}{m_d + m_u} \\ &= -\frac{\mathcal{X}_1 + \mathcal{X}_2}{\mathcal{X}_2 - \mathcal{X}_1} - \left(\approx \frac{1}{3} \right) \end{aligned}$$

Goldstone of Hypercharge:

$$\varphi_Y = (v_2 \varphi_2 - v_1 \varphi_1)/v$$

Condition enforcing no a - φ_Y mixing:

$$\begin{aligned} \sum_i \mathcal{X}_i Y_i v_i^2 &= 0 \\ \Rightarrow \mathcal{X}_1 v_1^2 + \mathcal{X}_2 v_2^2 &= 0 \end{aligned}$$

Couplings to the physical $a(x)$ are defined in terms of the charges

$$X_1 \propto -v_2^2/v^2 = -s_\beta^2$$

$$X_2 \propto v_1^2/v^2 = c_\beta^2$$

With $v_2^2/v_1^2 \approx 2$ ($c_\beta^2 \approx 1/3$) $\rightarrow X_u - X_d = X_1 + X_2 \approx -1/3$
then $c_u - c_d \approx 0$ and approximate a - N, π decoupling $g_{aN} \approx 0$

DFSZ Electrophobia: $g_{ae} \approx 0$ ($H_{1,2} + H_3$)

Add a third H_3 coupled to leptons, relevant conditions:

DFSZ Electrophobia: $g_{ae} \approx 0$ ($H_{1,2} + H_3$)

Add a third H_3 coupled to leptons, relevant conditions:

$$\chi_1 v_1^2 + \chi_2 v_2^2 + \chi_3 v_3^2 = 0$$

α - $\psi\gamma$ decoupling condition:

DFSZ Electrophobia: $g_{ae} \approx 0$ ($H_{1,2} + H_3$)

Add a third H_3 coupled to leptons, relevant conditions:

$$\chi_1 v_1^2 + \chi_2 v_2^2 + \chi_3 v_3^2 = 0$$

α - ϕ γ decoupling condition:

$$H_3^\dagger H_2 \phi + H_3^\dagger H_1 \phi^2$$

(or $H_3^\dagger H_1 \phi + H_3^\dagger H_2 \phi^\dagger, H_3^\dagger H_1 \phi^2 + H_3^\dagger H_2 \phi^{\dagger 2}, \dots$)

Explicit breaking of $U(1)_{H_3}$ rephasing symmetry (no additional Goldstones)

DFSZ Electrophobia: $g_{ae} \approx 0$ ($H_{1,2} + H_3$)

Add a third H_3 coupled to leptons, relevant conditions:

$$\chi_1 v_1^2 + \chi_2 v_2^2 + \chi_3 v_3^2 = 0$$

α - ϕ γ decoupling condition:

$$H_3^\dagger H_2 \phi + H_3^\dagger H_1 \phi^2$$

(or $H_3^\dagger H_1 \phi + H_3^\dagger H_2 \phi^\dagger, H_3^\dagger H_1 \phi^2 + H_3^\dagger H_2 \phi^{\dagger 2}, \dots$)

Explicit breaking of $U(1)_{H_3}$ rephasing symmetry (no additional Goldstones)

$$C_u - C_d = \frac{\chi_1 + \chi_2}{\chi_1 - \chi_2} = \frac{m_d - m_u}{m_d + m_u}$$

2nd Nucleophobia condition

DFSZ Electrophobia: $g_{ae} \approx 0$ ($H_{1,2} + H_3$)

Add a third H_3 coupled to leptons, relevant conditions:

$$\chi_1 v_1^2 + \chi_2 v_2^2 + \chi_3 v_3^2 = 0$$

α - ϕ γ decoupling condition:

$$H_3^\dagger H_2 \phi + H_3^\dagger H_1 \phi^2$$

(or $H_3^\dagger H_1 \phi + H_3^\dagger H_2 \phi^\dagger, H_3^\dagger H_1 \phi^2 + H_3^\dagger H_2 \phi^{\dagger 2}, \dots$)

Explicit breaking of $U(1)_{H_3}$ rephasing symmetry (no additional Goldstones)

$$C_u - C_d = \frac{\chi_1 + \chi_2}{\chi_1 - \chi_2} = \frac{m_d - m_u}{m_d + m_u}$$

2nd Nucleophobia condition

Lepton-axion
decoupling:

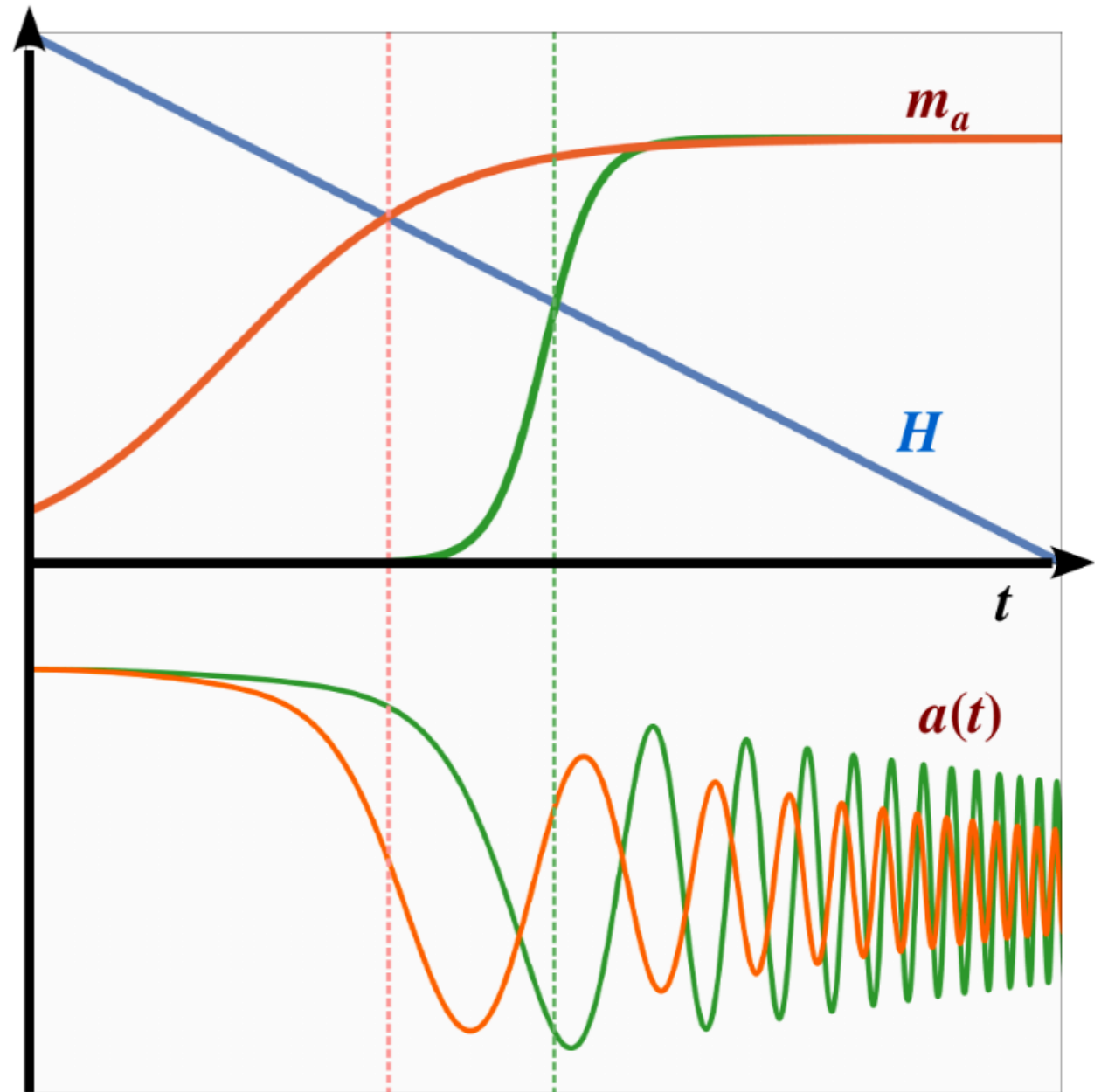
$$\chi_3 \approx 0$$

occurs for specific values
 $m_d/m_u \approx 2, 1, 1/2$ with no
additional tuning required

Axion relic density $\Omega_a = \rho_a / \rho_c$

$$\rho_a = m_a^2 a^2$$

Evaluated from integrating: $\ddot{a} + 3H\dot{a} + m_a^2(T)f_a \sin\left(\frac{a}{f_a}\right) = 0$



Credits to G. Villadoro

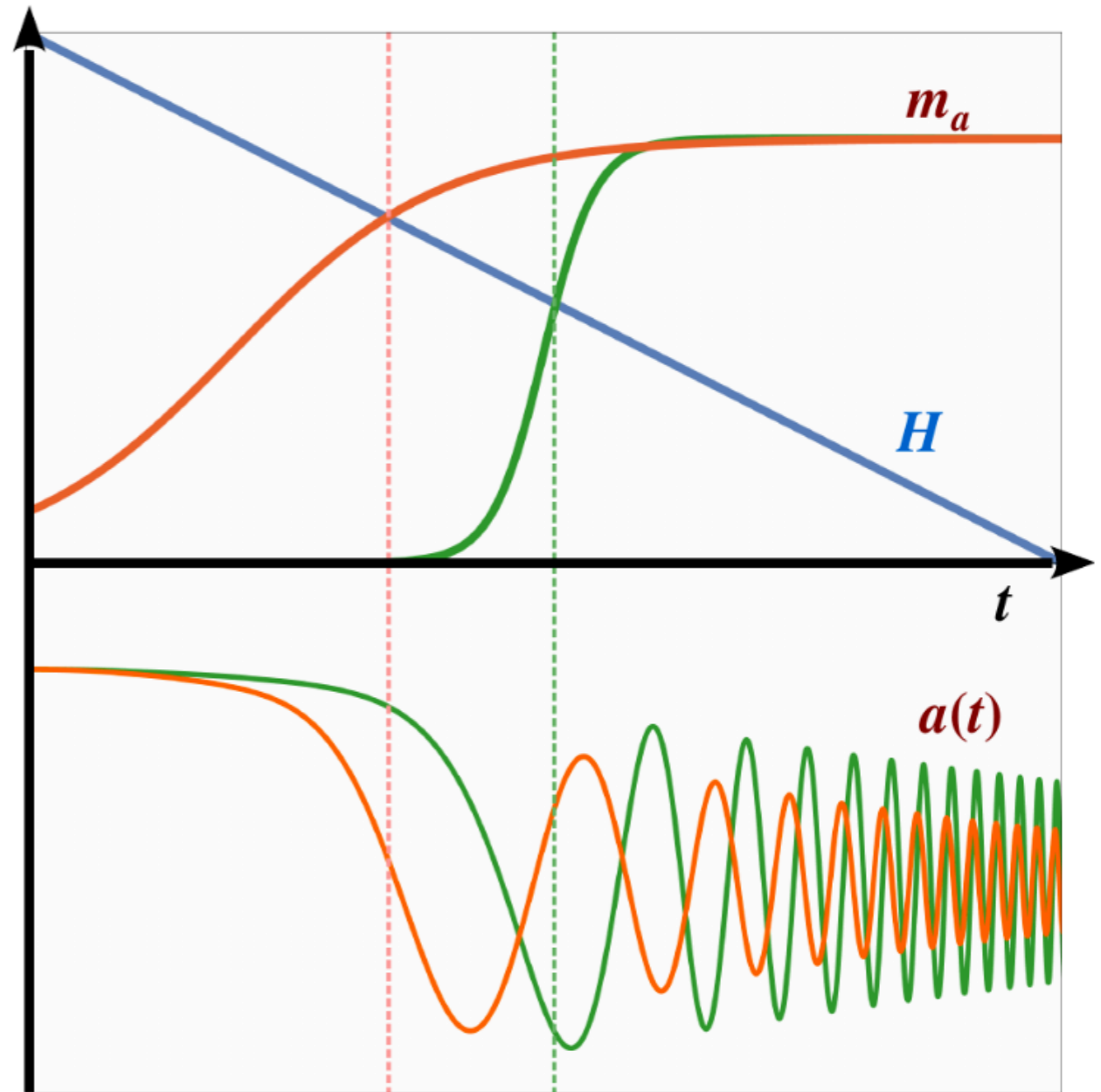
Axion relic density $\Omega_a = \rho_a / \rho_c$

$$\rho_a = m_a^2 a^2$$

Evaluated from integrating: $\ddot{a} + 3H\dot{a} + m_a^2(T)f_a \sin\left(\frac{a}{f_a}\right) = 0$

(1) boundary conditions:

➤ $a_i = \theta_i f_a; \quad (da/dt)_i = 0;$
 $f_a \gtrsim H_I, \quad H_I \gg \Lambda_{\text{QCD}}$



➤ Credits to G. Villadoro

Axion relic density $\Omega_a = \rho_a / \rho_c$

$$\rho_a = m_a^2 a^2$$

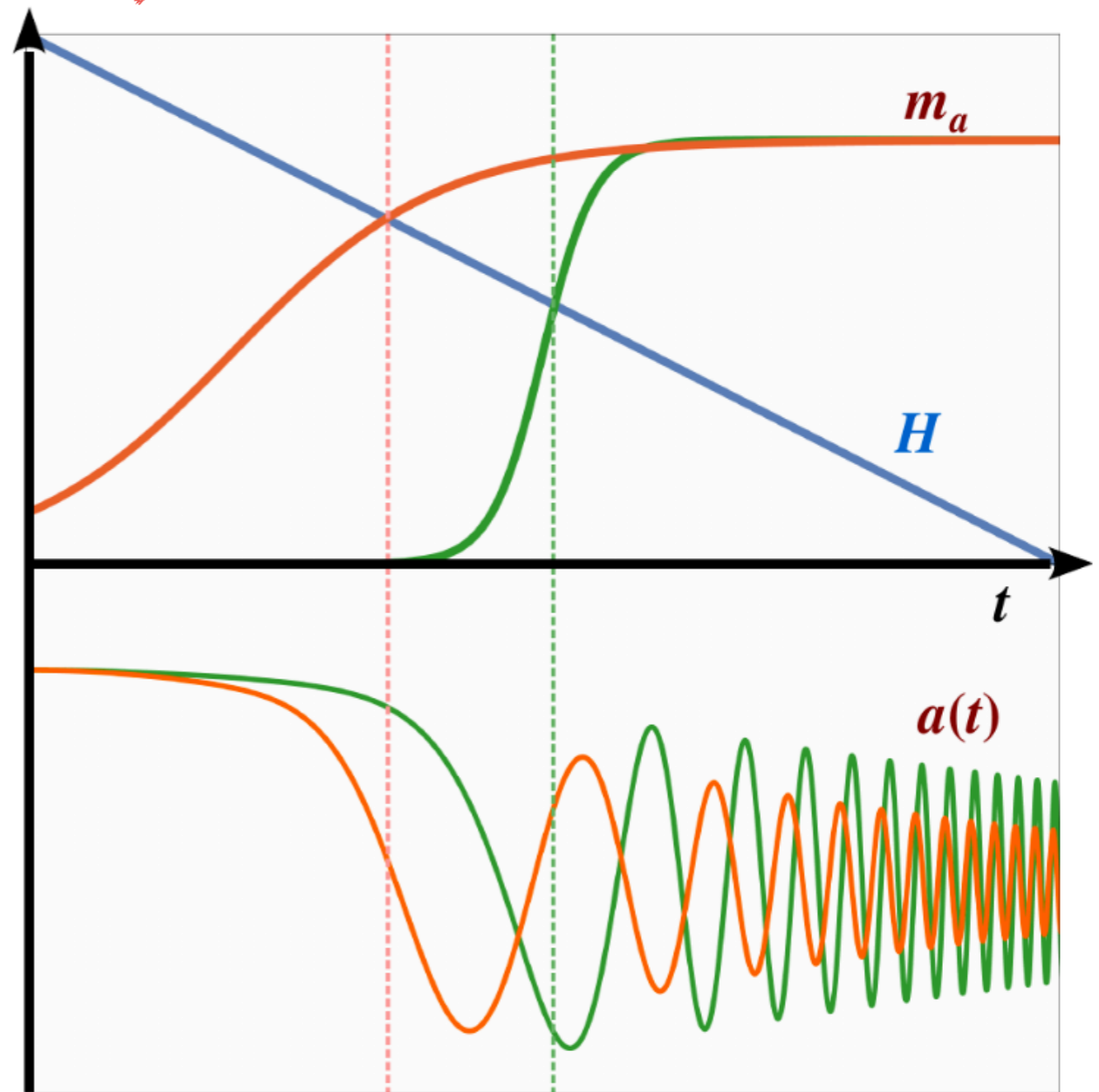
Evaluated from integrating: $\ddot{a} + 3H\dot{a} + m_a^2(T) f_a \sin\left(\frac{a}{f_a}\right) = 0$

(1) boundary conditions:

$\rightarrow a_i = \theta_i f_a; \quad (da/dt)_i = 0;$
 $f_a \gtrsim H_I, \quad H_I \gg \Lambda_{\text{QCD}}$

(2) $H(T) = 1.66 T^2/M_P$

(eq. st.: $w_{\text{eff}} = 1/3$, rad.dom.)



Credits to G. Villadoro

Axion relic density $\Omega_a = \rho_a / \rho_c$

$$\rho_a = m_a^2 a^2$$

Evaluated from integrating: $\ddot{a} + 3H\dot{a} + m_a^2(T) f_a \sin\left(\frac{a}{f_a}\right) = 0$

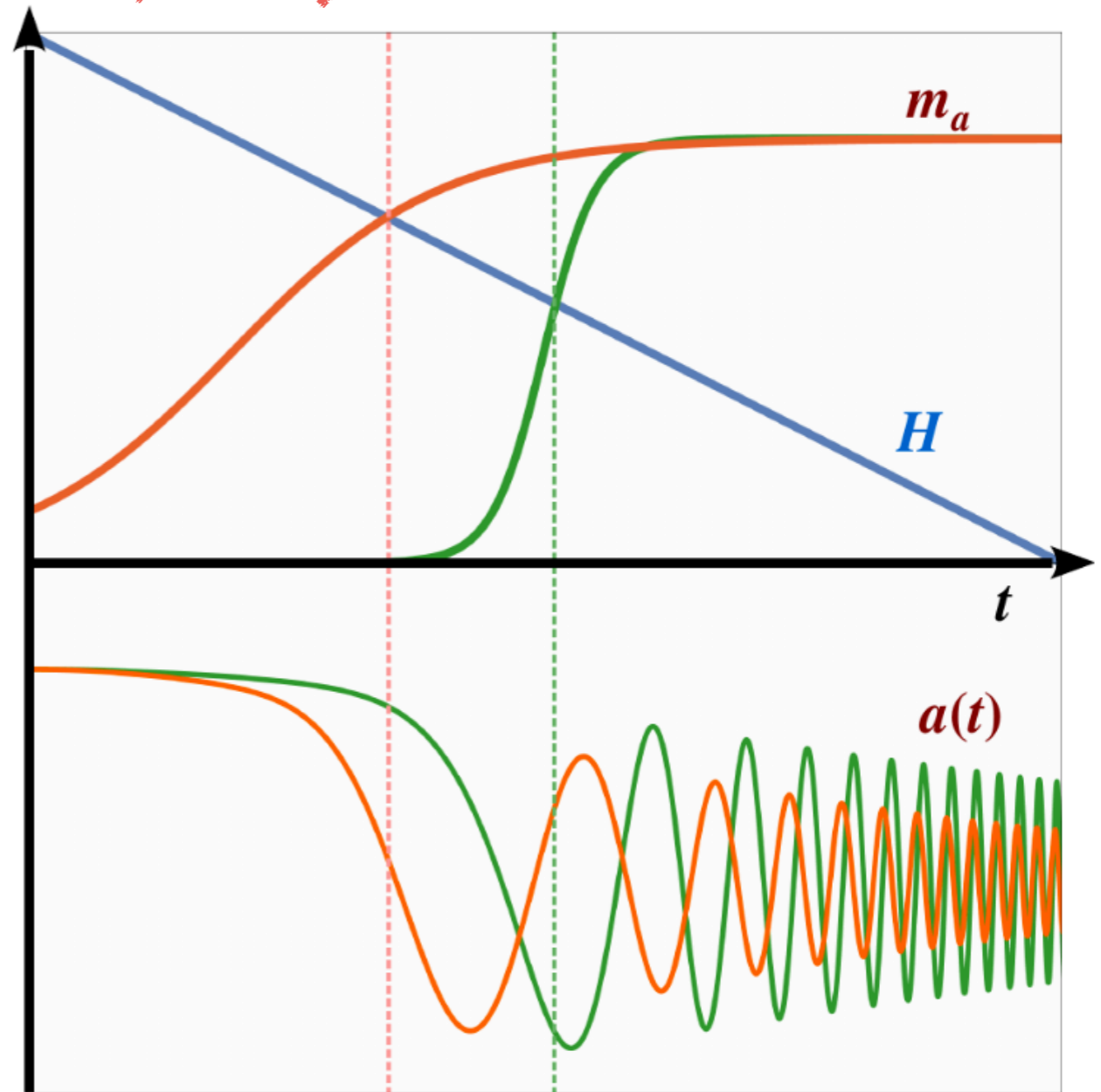
(1) boundary conditions:

$$\begin{aligned} \rightarrow a_i &= \theta_i f_a; \quad (da/dt)_i = 0; \\ f_a &\gtrsim H_I, \quad H_I \gg \Lambda_{\text{QCD}} \end{aligned}$$

(2) $H(T) = 1.66 T^2/M_P$
(eq. st.: $w_{\text{eff}}=1/3$, rad.dom.)

(3) $m_a^2(T) = m_a^2 \chi(T)/\chi$

(χ = QCD topological suscept.)
with $m_a \sim m_\pi f_\pi / f_a$



 Credits to G. Villadoro

Axion relic density $\Omega_a = \rho_a / \rho_c$

$$\rho_a = m_a^2 a^2$$

Evaluated from integrating: $\ddot{a} + 3H\dot{a} + m_a^2(T) f_a \sin\left(\frac{a}{f_a}\right) = 0$

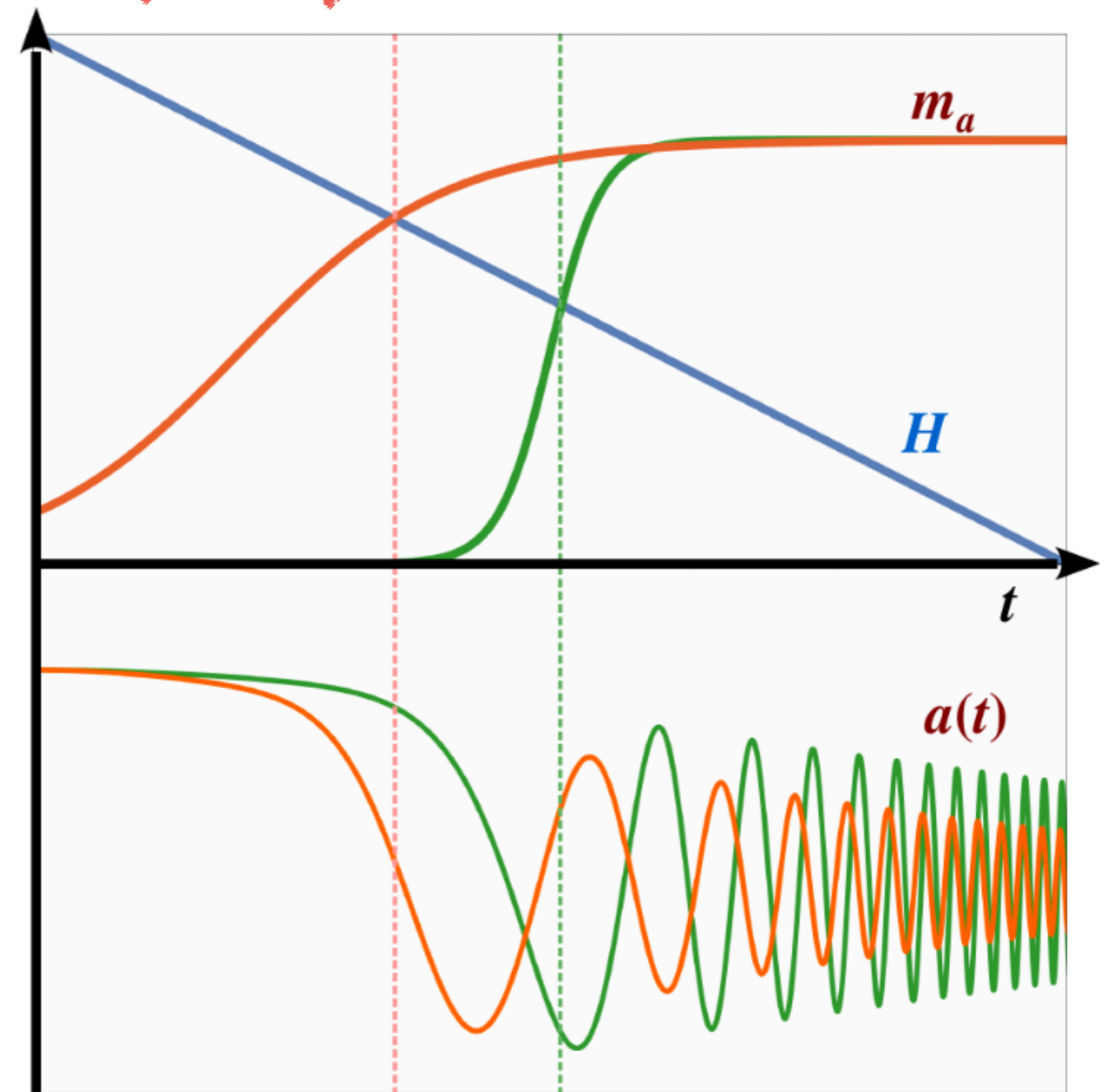
(1) boundary conditions:

$\rightarrow a_i = \theta_i f_a; \quad (da/dt)_i = 0;$
 $f_a \gtrsim H_I, \quad H_I \gg \Lambda_{\text{QCD}}$

(2) $H(T) = 1.66 T^2/M_P$
(eq. st.: $w_{\text{eff}}=1/3$, rad.dom.)

(3) $m_a^2(T) = m_a^2 \chi(T)/\chi$
(χ = QCD topological suscept.)
with $m_a \sim m_\pi f_\pi / f_a$

(4) assumed entropy conservation



Credits to G. Villadoro

$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

- θ_i : Long lasting, low scale inflation.

[P.W. Graham, A. Scherlis 1805.07362,
F.Takahashi, W.Yin, A. H.Guth 1805.08763]

$$H_I \lesssim \Lambda_{\text{QCD}}, \quad \theta_i \rightarrow \ll 1, \quad m_a \sim 10^{-12} \text{ eV}$$

$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

- θ_i : Long lasting, low scale inflation.

[P.W. Graham, A. Scherlis 1805.07362,
F.Takahashi,W.Yin, A. H.Guth 1805.08763]

$$H_I \lesssim \Lambda_{\text{QCD}}, \quad \theta_i \rightarrow \ll 1, \quad m_a \sim 10^{-12} \text{ eV}$$

- $H(T)$: GR + non standard thermal history

- Entropy generation $T < T_{\text{osc}}$: $m_a \approx (5-50) \mu\text{eV} / \Delta^{7/6}$

- ψ -domination, with $w_\psi \neq 1/3$: easy to arrange $m_a \ll 1 \mu\text{eV}$

[N. Ramberg, L. Visinelli, 1904.05707]

$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

- θ_i : Long lasting, low scale inflation.

[P.W. Graham, A. Scherlis 1805.07362,
F.Takahashi,W.Yin,A. H.Guth 1805.08763]

$$H_I \lesssim \Lambda_{\text{QCD}}, \quad \theta_i \rightarrow \ll 1, \quad m_a \sim 10^{-12} \text{ eV}$$

- $H(T)$: GR + non standard thermal history

- Entropy generation $T < T_{\text{osc}}$: $m_a \approx (5-50) \mu\text{eV} / \Delta^{7/6}$

- ψ -domination, with $w_\psi \neq 1/3$: easy to arrange $m_a \ll 1 \mu\text{eV}$

[N. Ramberg, L. Visinelli, 1904.05707]

Beyond GR: [work in progress]

- Scalar-tensor theories:

(conformal) boosted $H(T)$: $m_a \uparrow$

(disformal ?) quenched $H(T)$: $m_a \downarrow$

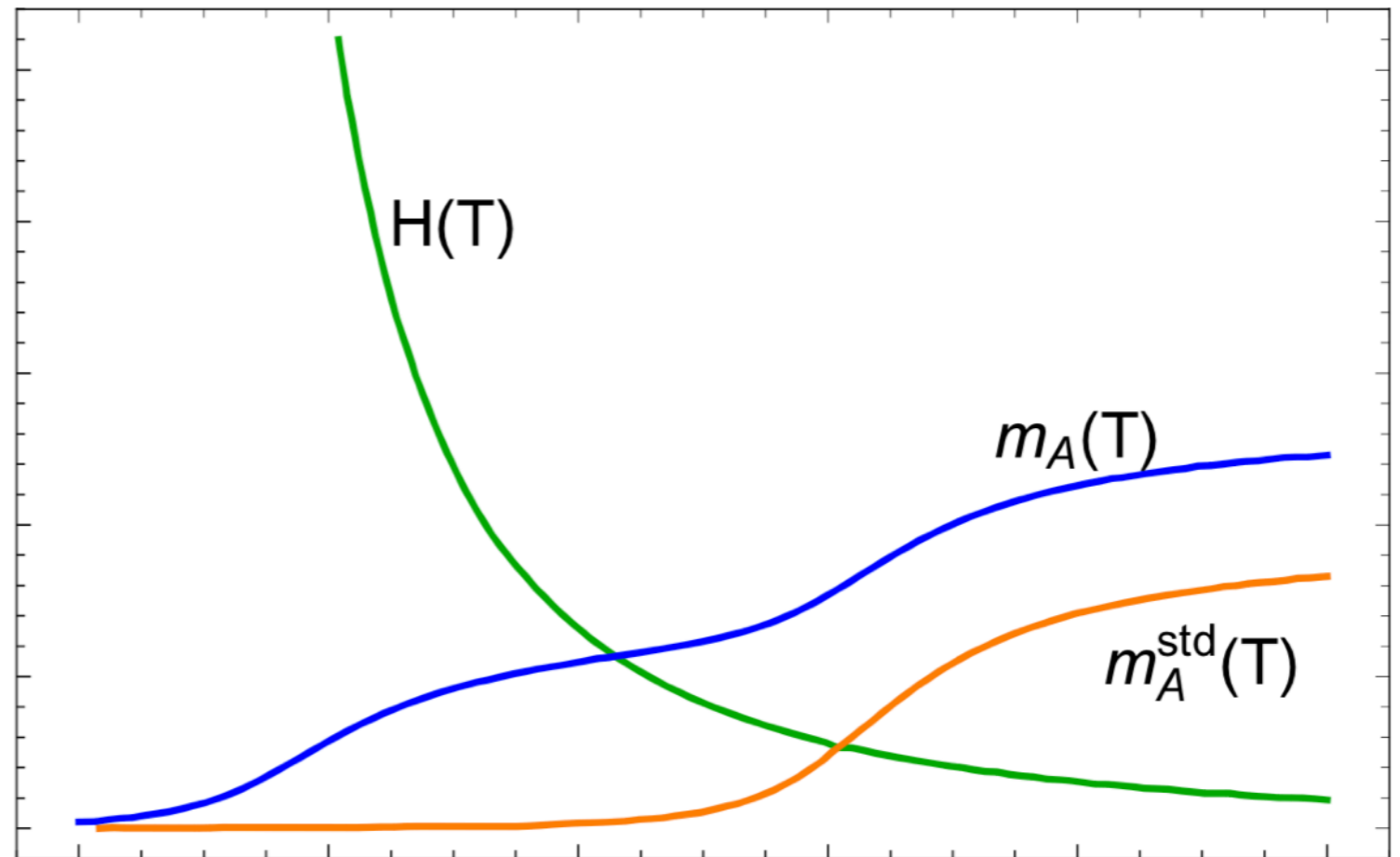
$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

- **Modify $m_a(T)$:**
Additional contributions
from mirror instantons:
(earlier onset of oscillations)

$m_a \downarrow$ [Giannotti, astro-ph/0504636]

[For $m_a \uparrow$ see P. Quilez talk]



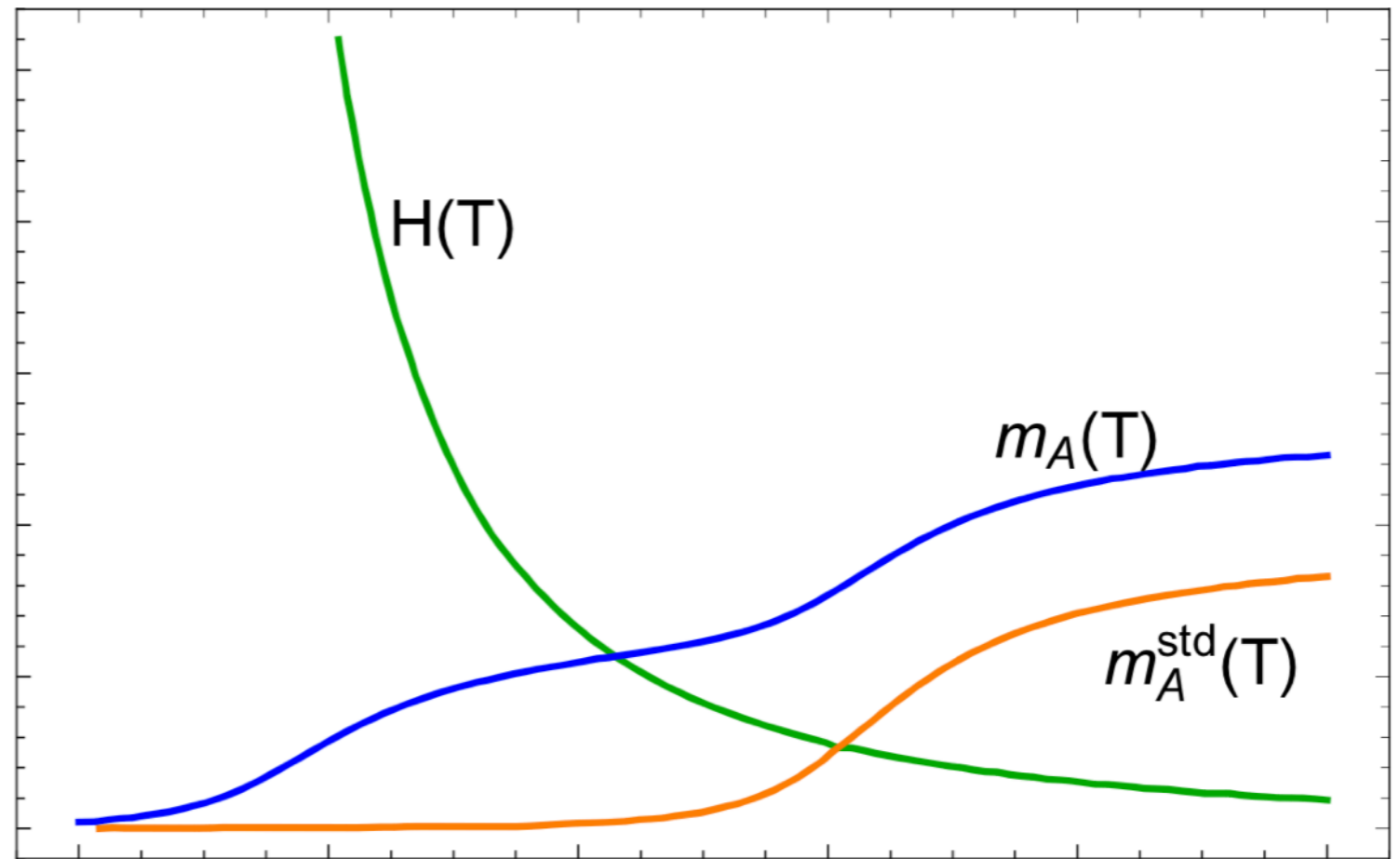
$\Omega_a \approx \Omega_{\text{CDM}}$ outside canonical m_a windows

- **Modify $m_a(T)$:**

Additional contributions
from mirror instantons:
(earlier onset of oscillations)

$m_a \downarrow$ [Giannotti, astro-ph/0504636]

[For $m_a \uparrow$ see P. Quilez talk]



- **Modify m_a :**

N copies of QCD related by a \mathbb{Z}_N symmetry

[A. Hook, 1802.10093]

$$m_a(N) = 2^{2-N/2} \times m_a, \quad \text{for } N > 4 \quad m_a \downarrow$$

Summary

Summary

- A certain number of quasi-model-independent features characterise axion models. Hard to circumvent, but possible.

Summary

- A certain number of quasi-model-independent features characterise axion models. Hard to circumvent, but possible.
- From any measurement we learn something: some scenarios can be ruled out, viable hypothesis can become no more viable, etc. Experimental results reaching into non-canonical regions are equally important for this learning process.

Summary

- A certain number of quasi-model-independent features characterise axion models. Hard to circumvent, but possible.
- From any measurement we learn something: some scenarios can be ruled out, viable hypothesis can become no more viable, etc. Experimental results reaching into non-canonical regions are equally important for this learning process.
- Working group to classify **axion models off the beaten tracks** (M. Giannotti, L. Di Luzio, EN, L. Visinelli). Full report to be expected after the summer.

Thanks for your attention