## VMB@CERN

ŠTĚPÁN KUNC ON BEHALF OF VMB@CERN COLLABORATION PATRAS WORKSHOP 2019

#### VMB@CERN collaboration

New approach to measure Vacuum Magnetic Birefringence Started in 2019

# PVLAS + Q&A + OSQAR-VMB + LIGO group (Cardiff)

about 15 members from 10 Institutes from Czech Republic, France, Italy, Poland, Republic of China and Wales

## VMB - Light propagating in vacuum in an external magnetic field

Linear Birefringence – polarization dependent index of refraction

- Linear Dichroism polarization dependent index of absorption
- ► In external perpendicular magnetic field  $\overline{B}$

 $\Delta n_{vac} = \Delta n_B + i \Delta k_B$ 

 $\Delta n_B \propto B^2$ 

Birefringence Predicted by QED, ALP processes, Millicharged particles (MCP)

Dichroism $\Delta k_B \propto B^2$ Predicted byALP processes,Millichargedparticles (MCP)

## **QED Vacuum Magnetic Birefringence**

- ► H. Euler, B. Kockel 1935
- They wrote an effective Lagrangian describing electromagnetic interactions in the presence of the virtual electron-positron sea

$$\mathcal{L}_{EK} = \frac{1}{2\mu_0} \left( \frac{E^2}{c^2} - B^2 \right) + \frac{A_e}{\mu_0} \left[ \left( \frac{E^2}{c^2} - B^2 \right)^2 + 7 \left( \frac{\vec{E}}{c} \cdot \vec{B} \right)^2 \right] + \cdots$$

$$A_e = \frac{2}{45\mu_0} \frac{\alpha^2 \lambda_e^3}{m_e c^2} = \mathbf{1.32} \times \mathbf{10^{-24} T^{-2}} \longrightarrow \Delta n_B = 3A_e B^2$$

No linear Dichroism is expected in the framework of QED

## Ellipticity

- We can not measure directly  $\Delta n_B$
- $\blacktriangleright$  What we measure is induced ellipticity  $\psi$
- ► We determine optical path difference  $D_n = \Delta n_B L$

To obtain maximum signal  $\psi$ we need high magnetic field Bin as long as possible region LPol. and B fields at 45° degree  $\delta$ 

ingle pass 
$$\psi_0 = rac{\pi 3A_e B^2 L}{\lambda} \sin 2d$$

#### **Ellipsometer parameters**

Long optical path – Fabry-Perot cavity

High amplification factor N, PVLAS F = 770000

Fabry Perot cavity 
$$\psi = \frac{\pi 3A_e B^2 L}{\lambda} \frac{2F}{\pi} \sin 2\delta$$

#### ► High magnetic field $B^2$

 $\mathbf{N}$  – amplification factor

Permanent magnets – PVLAS – B = 2.5 T

superconducting magnets - OSQAR - B = 9 T, L = 14.3 m

#### Time dependent effect - to achieve high sensitivity

Time dependent magnetic filed - modulation - BRFT, pulsed – BMV,OVAL Time dependent angle  $\delta$  – PVLAS I-II – rotating the permanent magnets

## Intrinsic noise

Sensitivity in optical path difference  $D_n$  between two orthogonal polarizations



#### **PVLAS** -

The ultimate level of shot noise was reached without the Fabry – Perot cavity.

But with the cavity there is an extra noise dominated by ellipticity due to the cavity mirrors even in the heterodyne modulation

#### Sensitivity in optical path difference $D_n$ does not depend on finesse

## Intrinsic noise

- Measured ellipticity noise and Cotton-Mouton signal as a function of the finesse
- ▶ Finesse range: 250'000 690'000



Noise and Cotton-Mouton ΔD signals are independent of finesse

► Above finesse of  $\approx$  20 000 the intrinsic noise is dominant at  $\approx$ 10 Hz

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Use LHC dipole magnet and separate magnet from modulation

#### 1. Increase signal

Use LHC magnet -  $B^2L \approx 1200 \text{ T}^2 \text{m}$  - almost impossible to modulate or rotate

#### 2. Time dependent angle $\delta$

Rotate the polarization instead of magnetic field Use rotating half wave plates to produce modulation in VMB ellipticity Profit from heterodyne measurement principle

## Rotation of half wave-plates

PVLAS: G. Zavattini at al. Eur. Phys. J. C(2016) 76:294

Rotate polarization inside the magnetic field Fix polarization on mirrors to avoid mirror birefringence signal Insert two co-rotating half wave plates at frequency f with a fixed relative angle Total losses  $\approx 0.4\%$  (commercial) finesse becomes F  $\approx 800$ Maybe 10 times lower loss is possible F  $\approx 1000$ 





#### Output intensity solution

 $I(t) = I_0 \{ \eta(t)^2 + 2\eta(t) N[\psi_0 \sin 4\theta(t) + \alpha_1 \sin 2\theta(t) + \alpha_2 \sin(2\theta(t) + 2\Delta\theta)] \}$ VMB signal appears at 4<sup>th</sup> harmonics of HWP rotation

#### Wave-plate defects $\alpha_{1,2}$

$$\alpha_{1,2}(t) = \alpha_{1,2}^{(0)} + \alpha_{1,2}^{(1)} \cos \theta(t) + \alpha_{1,2}^{(2)} \cos 2\theta(t) + \cdots$$

## LHC magnet

Signal

$$D_n = 4 \times 10^{-24} B^2 L$$

Intrinsic noise

$$S_{D_n}^{(intristic)} = 2.6 \times 10^{-18} \nu^{-0.77} \text{ m}/\sqrt{\text{Hz}}$$

Shot noise

$$S_{D_n}^{(shot)} = \sqrt{\frac{e}{I_0 q}} \frac{\lambda}{\pi N}$$

Measurement time

$$r = \left(\frac{S_{D_n}}{D_n}\right)^2$$

LHC magnet

 $B^2 L = 1200 \text{ T}^2 \text{m};$   $S_{D_n} = 10^{-18} \text{ m}/\sqrt{\text{Hz}}$  at 3Hz T = 12 h for S/N = 1

## $S_{D_n} \approx 10^{-18} \text{ m}/\sqrt{\text{Hz}}$ as goal sensitivity



Updated graph from G. Zavattini etal. Eur. Phys.J.C(2016)76:294

#### Conclusion

Rotating half wave-plates inside a Fabry-Perot cavity could be a viable solution to use high magnetic field of LHC magnet

Maximum finesse F ≈ 10'000

Defects may also be a limit but only at second order:  $\alpha_{1,2}^{(2)}$ With a sensitivity of  $D_n \approx 10^{-18} \text{ m}/\sqrt{\text{Hz}}$  and 1 LHC magnet, vacuum magnetic birefringence could be measured with S/N = 1 in about 1day. Lol has been submitted to CERN: CERN-SPSC-2018-036/SPSC-I-249 Joint effort between past Vacuum Magnetic Birefringence experiments + LIGO: 16 authors, 7 countries.