Boiling the Vacuum in Photon Laser Collisions

Laser-assisted Bremsstrahlung Photon Pair Production

Andreas Ringwald LUXE Workshop DESY Hamburg, Germany 16-17 April 2019

[Anthony Hartin, AR, Natalia Tapia, Phys.Rev. D99 (2019) no.3, 036008 [arXiv:1807.10670 [hep-ph]]





QED beyond ordinary perturbation theory

- Quantum Electrodynamics (QED) is the most accurately tested theory in physics
- Its predictions, for observables accessible by ordinary perturbation theory, have been verified up to very high accuracy
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Theory: $a_e = 0.001 \ 159 \ 652 \ 181 \ 643(764)$

Experiment: $a_e = 0.001\ 159\ 652\ 180\ 73(28)$

[Hanneke,Fogwell Hoogerheide,Gabrielse 11]

[[]Aoyama,Hayakawa,Kinoshita,Ńio 12,15]

QED beyond ordinary perturbation theory

- Quantum Electrodynamics (QED) is the most accurately tested theory in physics
- Its predictions, for observables accessible by ordinary perturbation theory, have been verified up to very high accuracy
 - Anomalous magnetic moment of electron: theory (five loop) agrees with experiment to more than 10 significant digits
- There are, however, also observables which are inaccessible by ordinary perturbation theory and whose prediction lacks an experimental verification
 - Vacuum electron-positron pair production (VPP) in a strong static electric field

[Sauter 31; Euler, Heisenberg 36; Schwinger 51]



QED beyond ordinary perturbation theory

• Rate per unit volume:

[Schwinger 51]

$$\frac{\Gamma_{\rm VPP}}{V} = \frac{m_e^4}{(2\pi)^3} \left(\frac{|\mathbf{E}|}{\mathbf{E}_{\rm c}}\right)^2 \exp\left(-\pi \frac{\mathbf{E}_{\rm c}}{|\mathbf{E}|}\right)$$

• Schwinger critical field

$$\mathcal{E}_c \equiv \frac{m_e^2}{e} \simeq 1.3 \times 10^{18} \text{ V/m}$$

• Rate for VPP non-perturbative in QED coupling e, as typical for process which can occur, for $|\mathbf{E}| \leq E_c$, only via quantum tunnelling:

$$\Gamma_{\rm VPP} \propto \exp\left(-\pi \frac{m_e^2}{e|\mathbf{E}|}\right)$$



Beyond QED

- Schwinger effect and its analogues have been suggested to play a role in many problems of phenomenological and cosmological interest:
 - Particle production in hadronic collisions

[Casher,Neuberger,Nussinov 79]

Particle Production

<u>QED Example</u>: Schwinger mechanism of particle pair production



Quantum tunneling from the negative energy continuum

<u>QCD Example</u>: Color electric flux tube (string model)



Beyond QED

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 - Particle production in hadronic collisions

[Casher,Neuberger,Nussinov 79]

- Black hole quantum evaporation
 [Hawking 75]
- Gibbons-Hawking radiation during inflation





Laser-Assisted Vacuum Pair Production

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Laser-Assisted Vacuum Pair Production

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- Consider vacuum pair production at the focus of an intense laser:

• Need optical laser in hundreds of exawatt range to have observable effect!

• Required power can be reduced if one superimposes the field at the focus with further weak HF field DESY. | Boiling the Vacuum in Photon Laser Collisions | Andreas Ringwald, LUXE Workshop, DESY, Hamburg, Germany, 16 April 2019 [Schützhold, Gies, Dunne 08] Page 9

• Consider one photon pair production at the focus of an intense laser:



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[Reiss 62, Narozhnyi, Nikishov, Ritus 65]





• Rate: $\Gamma_{\text{laOPPP}} = \frac{\alpha m_e^2}{4 \omega_{\text{i}}} F_{\gamma}(\xi, \chi_{\gamma})$

(circular polarisation, infinite plane wave approximation) [Narozhnyi,Nikishov,Ritus 65]

$$F_{\gamma}(\xi, \chi_{\gamma}) = \sum_{n>n_{o}}^{\infty} \int_{1}^{v_{n}} \frac{\mathrm{d}v}{v\sqrt{v(v-1)}} \left[2J_{n}^{2}(z_{v}) + \xi^{2}(2v-1)\left(J_{n+1}^{2}(z_{v}) + J_{n-1}^{2}(z_{v}) - 2J_{n}^{2}(z_{v})\right) \right]$$
$$n_{0} \equiv \frac{2\xi(1+\xi^{2})}{\chi_{\gamma}}, \quad z_{v} \equiv \frac{4\xi^{2}\sqrt{1+\xi^{2}}}{\chi_{\gamma}} \left[v(v_{n}-v)\right]^{1/2}, \quad v_{n} \equiv \frac{\chi_{\gamma}n}{2\xi(1+\xi^{2})}$$
Photon recoil parameter: $\chi_{\gamma} \equiv (1+\cos\theta) \frac{\omega_{i}}{m_{e}} \frac{|\mathbf{E}|}{\mathbf{E}_{c}} \simeq 0.5 (1+\cos\theta) \left(\frac{\omega_{i}}{17.5 \,\mathrm{GeV}}\right) \left(\frac{I}{10^{20} \mathrm{W/cm^{2}}}\right)^{1/2}$

• Consider one photon pair production at the focus of an intense laser:

•

•

[Reiss 62, Narozhnyi, Nikishov, Ritus 65]



$$\Gamma_{\text{IaOPPP}} = \frac{\alpha m_e^2}{4\omega_1} F_{\gamma}(\xi, \chi_{\gamma})$$

$$(\xi, \chi_{\gamma}) = \begin{cases} 2\xi^2 \left[\log\left(\frac{2\chi_{\gamma}}{\xi}\right) - 1 \right] + \mathcal{O}\left(\xi^3 \log \xi\right) & : \xi \ll 1, \\ \frac{3}{4}\sqrt{\frac{3}{2}}\chi_{\gamma} \exp\left[-\frac{8}{3\chi_{\gamma}}\left(1 - \frac{1}{15}\xi^{-2} + \mathcal{O}(\xi^{-4})\right) \right] & : \xi \gtrsim 1/\sqrt{\chi_{\gamma}} \gg 1 \end{cases}$$

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• At $\xi \gtrsim 1/\sqrt{\chi_{\gamma}} \gg 1$, rate asymptotes to

$$\Gamma_{\text{laOPPP}} \to \frac{3}{16} \sqrt{\frac{3}{2}} \alpha \, m_e \left(1 + \cos \theta\right) \, \frac{|\mathbf{E}|}{\mathbf{E}_c} \exp\left[-\frac{8}{3} \frac{1}{1 + \cos \theta} \frac{m_e}{\omega_i} \frac{\mathbf{E}_c}{|\mathbf{E}|}\right]$$

- Ressembles rate of VPP in constant electric field
- In fact, $\xi \gg 1$ corresponds to quasi-static electric field, $\omega \ll e \left| {f E} \right| / m_e$
- Rate enhanced tremendously in OPPP in comparison to VPP due to boost factor $\,\,\omega_i/m_e\gg 1$

• Experimental set up to measure laser assisted one photon pair production: [Reiss 71,Keitel et al., 10,Turcu et al., 16]



[Hartin, AR, Tapia 1807.10670]

• Rate:

$$\Gamma_{\text{laBPPP}} = \frac{\alpha m_e^2}{4} \int_0^{E_e} \frac{\mathrm{d}\omega_i}{\omega_i} \frac{\mathrm{d}N_\gamma}{\mathrm{d}\omega_i} F_\gamma(\xi, \chi_\gamma(\omega_i))$$

• For a thin target, $X \ll X_0$, and high laser intensity, $\xi \gtrsim 1/\sqrt{\chi_e} \gg 1$, with $\chi_e \equiv (1 + \cos \theta) \omega E_e \xi/m_e^2$:

$$\Gamma_{\text{laBPPP}} \to \frac{\alpha \, m_e^2}{E_e} \frac{9}{128} \sqrt{\frac{3}{2}} \, \chi_e^2 \, e^{-\frac{8}{3\chi_e} \left(1 - \frac{1}{15\xi^2}\right)} \frac{X}{X_0} \qquad \text{[Hartin, AR, Tapia 1807.10670]}$$

$$\rightarrow \frac{9}{128} \sqrt{\frac{3}{2}} \alpha E_e \left(1 + \cos\theta\right)^2 \left(\frac{|\mathbf{E}|}{\mathbf{E}_c}\right)^2 \exp\left[-\frac{8}{3} \frac{1}{1 + \cos\theta} \frac{m_e}{E_e} \frac{\mathbf{E}_c}{|\mathbf{E}|}\right] \frac{X}{X_0}$$





Summary and Outlook

- Measurement of laser-assisted bremsstrahlung photon pair production exploiting primary electron beam from European XFEL and optical laser with intensity above about $10^{19} \,\mathrm{W/cm^2}$ range allows test of QED in the non-perturbative tunneling regime (sub-Schwinger electric fields involved)
- Ongoing:
 - GEANT model of converter target (cf. Sasha's talk)
 - Strong field QED particle-in-cell simulation (cf. Tony's talk)
 - Study effects of finite pulse length

BackUp: Infinite plane wave approximation

- The infinite plane wave approximation (IPWA) which we have exploited in order to determine the OPPP transition rate neglects the finite length and shape of the laser pulse.
- If the pulse length is not too short in comparison to the laser wavelength, the slowly varying envelope approximation (SVEA) can be employed. This requires the assumption that the pulse envelope does not change much over the course of one wavelength of the laser field. For LUXE, an 800 nm laser with a 35 fs pulse length is envisaged. This gives around N=12 periods within the pulse, meaning that the SVEA may be a reasonable assumption.
- In [Nousch:2012xe,Titov:2012rd,Titov:2013kya,Titov:2015pre] it was shown that the IPWA is a very good assumption if the pulse length is above N=10