

Numerical Support for LUXE

Nina Elkina

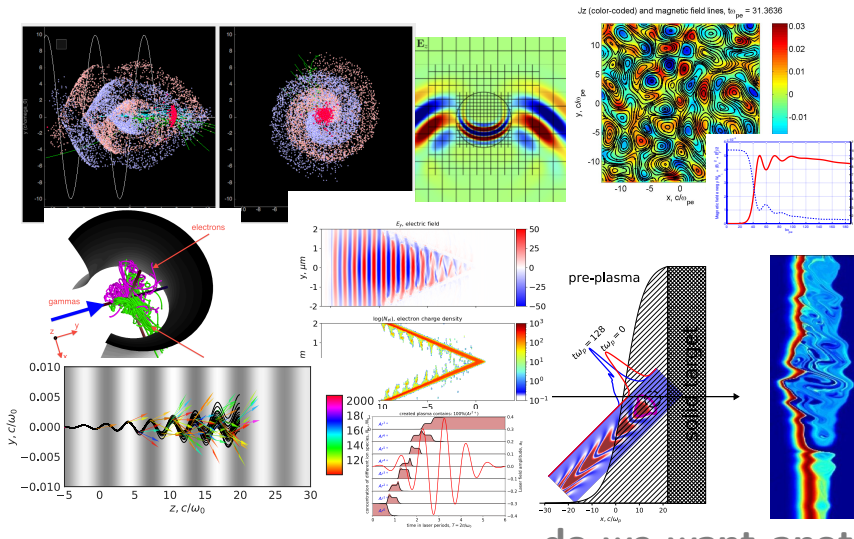
in collaboration with
Tom Teter and Matt Zepf

thanks for valuable discussions
Ben King and Felix Karbstein

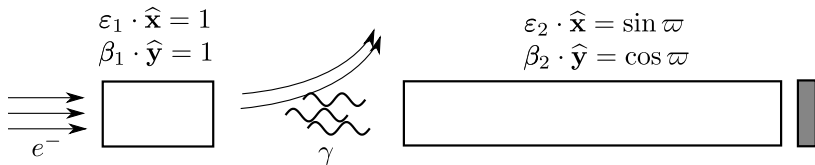
seed γ HI-Jena

April 16, Jena

Plasma simulations at HI-Jena



scattering simulation (vacuum birefringence) ¹



STAGE I

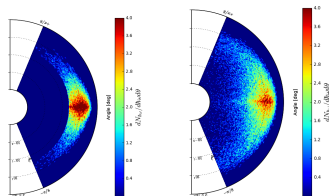
stage I: make γ via nonlinear Compton scattering

STAGE II

stage II: measure refractive indexes experienced by γ

cut off $\epsilon_{\min} = 1$

no cut off



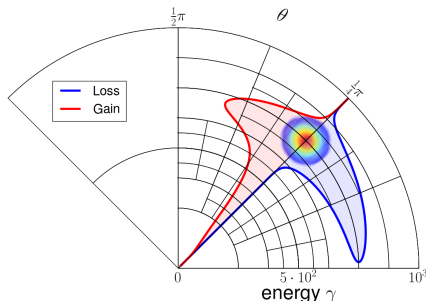
¹B. King, N. Elkina, Vacuum birefringence in high-energy laser-electron collisions, Phys. Rev. A 94, 062102, 2016

features of PIC-ANTARES code

$$\left(\frac{\partial}{\partial t} + \frac{\vec{p}}{\gamma m} \cdot \frac{\partial}{\partial \vec{r}} + \vec{F} \cdot \frac{\partial}{\partial \vec{p}} \right) f(\vec{r}, \vec{p}, t)$$

$$= \underbrace{\int d^3k w_{rad}^{\vec{E}, \vec{B}}(\vec{k}, \vec{p} + \vec{k}) f(\vec{r}, \vec{p} + \vec{k}, t)}_{\text{GAIN}} - \overbrace{f(\vec{r}, \vec{p}, t) \int d^3k w_{rad}^{\vec{E}, \vec{B}}(\vec{k}, \vec{p})}^{\text{LOSS}}$$

- Specific features of PIC-ANTARES code
- Descent event generators
- polarization effects
- norm-conserving integrators
- Adaptive mesh and particles

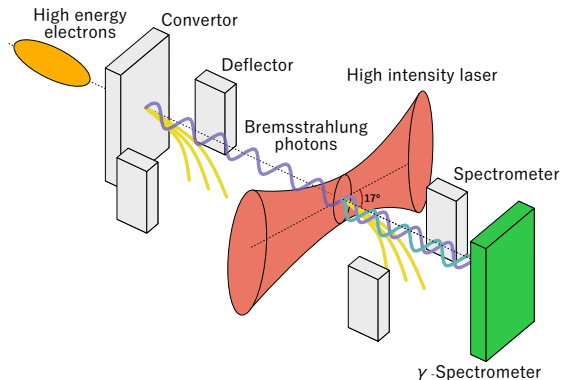


Plan

- overview of LUXE
- update on bremsstrahlung spectrum
- lessons from SLAC-144
- Measure Laser Pulse with Compton scattering

LUXE and SLAC

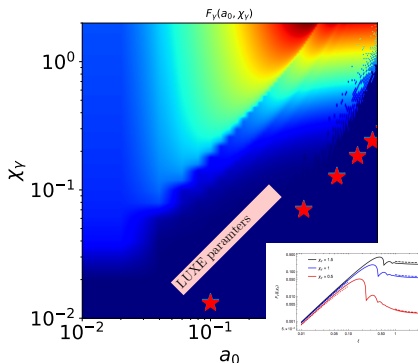
	λ	I, a_0	ε_e	ε_γ	χ_e/χ_γ	a_0
SLAC	527 nm	10^{18} W/cm^2	47 GeV	29.2 GeV	0.3/0.2	0.3
LUXE	800 nm	10^{20} W/cm^2	17.5 GeV	17.5 GeV (?)	0.80109	4-6



matching perturbative/non-perturbative QED

$$\Gamma = \frac{\alpha m_e^2}{4\omega_i} F_\gamma(a_0, \chi_\gamma),$$

$$F_\gamma(\xi_0, \chi_\gamma) = \sum_{n>n_0}^{\infty} \int_1^{v_n} \frac{dv}{v\sqrt{v(v-1)}} \left[2J_n^2(z_v) - \xi^2(2v-1)(J_{n+1}^2(z_v) + J_{n-1}^2(z_v) - 2J_n^2(z_v)) \right]$$

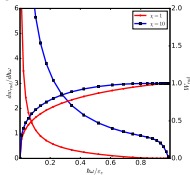


First order processes (with polarization effects)²

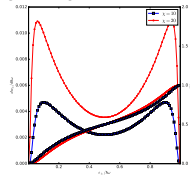
$$\frac{dW_\gamma}{d\varepsilon_\gamma} = \frac{-\alpha m^2}{\varepsilon_e^2} \left\{ Ai_1(x) + \left[\frac{g(\phi)}{x} + \chi_\gamma \sqrt{x} \right] Ai'_1(x) \right\}, \quad (1)$$

$$\frac{dW_e}{d\varepsilon_e} = \frac{+\alpha m^2}{\varepsilon_\gamma^2} \left\{ Ai_1(x) + \left[\frac{g(\phi)}{x} - \chi_\gamma \sqrt{x} \right] Ai'_1(x) \right\}, \quad (2)$$

where polarization effects are in $g(\phi) = 2 \cos^2 \phi + 1$
 photon emission

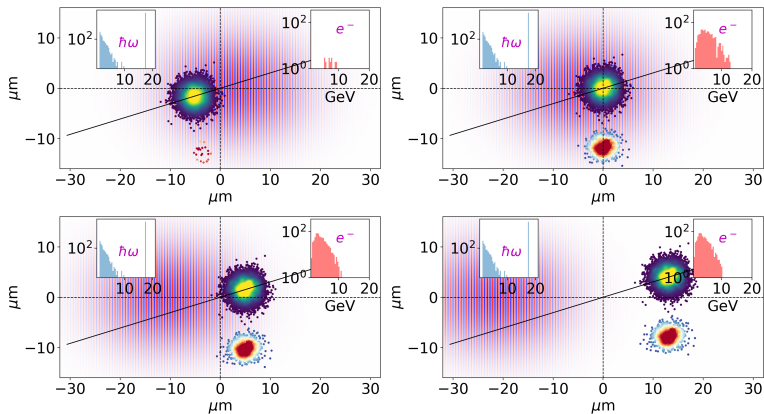


pair production



²B. King, N. Elkina, H. Ruhl, Photon polarization in electron-seeded pair-creation cascades, PhysRev. B, 2013 Elkina&Fedotov, 2011

Setup $\theta = 17^\circ$, $a_0 = 6$, $\hbar\omega = 17.5\text{GeV}$, ³

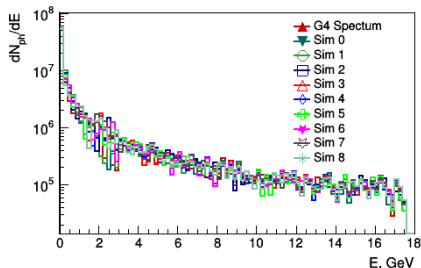


³partially simulated on Hemera cluster

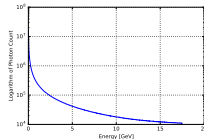
Bremstrahlung spectra, e^- : 17.5 GeV (GEANT4)

4

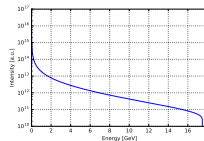
Sasha's Simulation



Tom&Evgenii's simulation

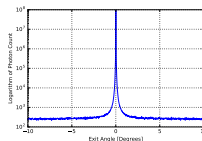


Geant4



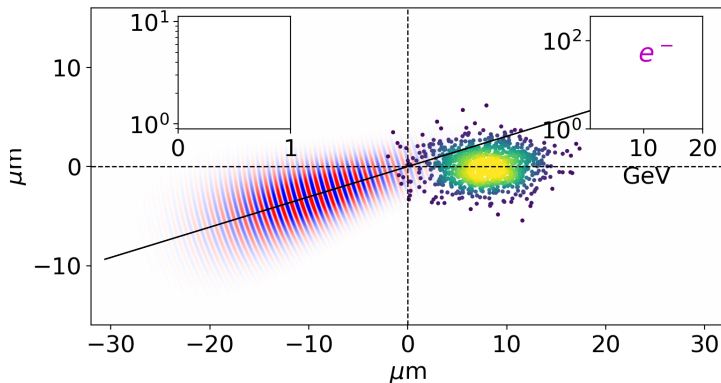
Tsai76:

Angular distribution



⁴O. Borysov, slides Dec.13, 2018

Pair yield depends strongly on laser pulse structure



Support to

Abstract

Reconstruction of Laser Pulse at SLAC

1. Laser diagnostics [Shmakov's PhD thesis]
 - laser energy: measured by a leakage monitor behind mirror
 - pulse duration: single shot autocorrelation , streak camera
 - focal spot: attenuated refocused pulse equivalent target method after interaction
2. Particle side [Horton's PhD thesis]
 - Cerenkov monitors to measure forward going photons
 -

Nonlinear Compton scattering

$$\frac{dW_n(\omega')}{d\omega'} = \frac{e^2 m^2 n_e}{16\pi\epsilon^2} \left\{ -4J_n^2(z) + a_0^2 \left(2 + \frac{u^2}{1+u} \right) \left[J_{n-1}^2(z) + J_{n+1}^2(z) - 2J_n^2(z) \right] \right\}$$

where $u = \frac{k_\mu k'^\mu}{k_\mu p'^\mu} \simeq \frac{\omega'}{\epsilon - \omega'}, \quad u_1 = \frac{2k_\mu p^\mu}{m^2} \simeq 2 \frac{\epsilon \omega_0 (1 - \cos \alpha)}{m^2}$

energy of scattered photon

$$\omega'(\theta) \simeq \epsilon \cdot \frac{u_1}{(1 + u_1) + \gamma^2 \theta^2}$$

maximum at $\theta = 0$

$$\omega'_{\max} \simeq \epsilon \cdot \frac{u_1}{(1 + u_1)}, \quad \epsilon = \gamma m_e c^2$$

1st harmonics+mass shift

first harmonic with mass shift $\overline{m}^2 = m^2(1 + a_0^2)$

Klein-Nishina term

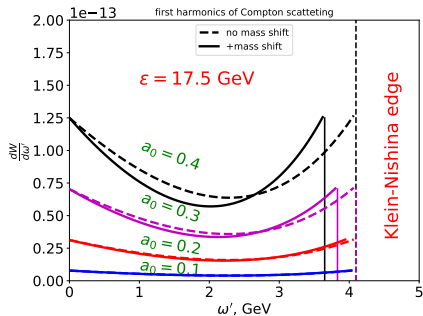
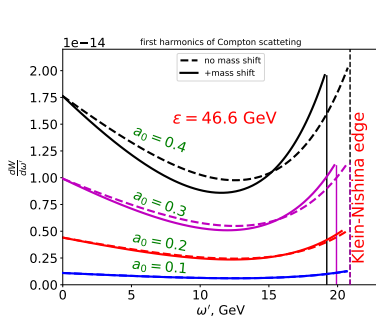
$$\frac{dW_1(\omega')}{d\omega'} = \frac{e^2 m^2 n_e}{16\pi\epsilon^2} = \left\{ \overbrace{a_0^2 \left[2 + \frac{u^2}{1+u} - 4\frac{u}{u_1} \left(1 - \frac{u}{u_1} \right) \right]}^{\text{Klein-Nishina term}} - a_0^4 \frac{u}{u_1} \left(1 - \frac{u}{u_1} \right) \left[1 + \frac{u^2}{1+u} - \frac{u}{u_1} \left(1 - \frac{u}{u_1} \right) \right] \right\}$$

true Klein-Nishina (no mass shift)

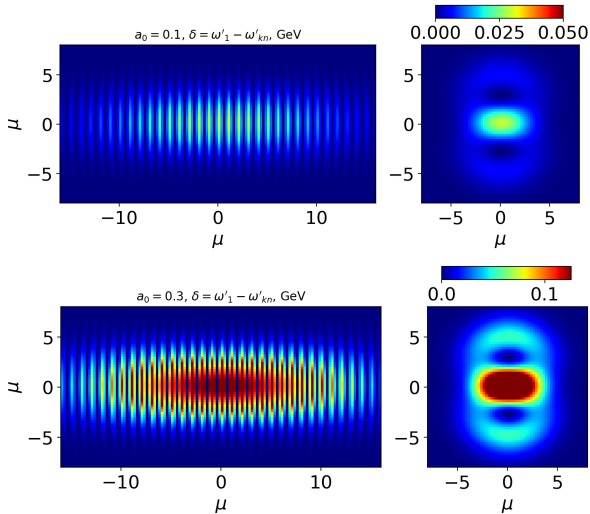
$$\frac{dW(\omega')}{d\omega'} = \frac{e^2 m_e^2 n_e}{16\pi\epsilon^2} a_0^2 \left\{ 2 + \frac{u^2}{1+u} - \frac{4u}{u_1} \left(1 - \frac{u}{u_1} \right) \right\}$$

$$\text{where } u \simeq \frac{\omega'}{\epsilon - \omega'}, \quad u_1 \simeq 2 \frac{\epsilon\omega_0(1 - \cos\alpha)}{m^2}$$

Laser diagnostic in a linear regime



Relative difference for Gaussina-like pulse ⁵



⁵Narozhny Fofanov, 2000

Summary&Conclusion

- Luxe requires to improve our numerical tools to deal with
 - low statistics because
 - local crossed field approximation
 - predict experimental observables
- Two major sources of uncertainty
 - Initial bremsstrahlung spectrum
 - laser focus parameter
- Mapping laser field strength with near-linear regime of Compton scattering