

Introduction to Accelerator Physics

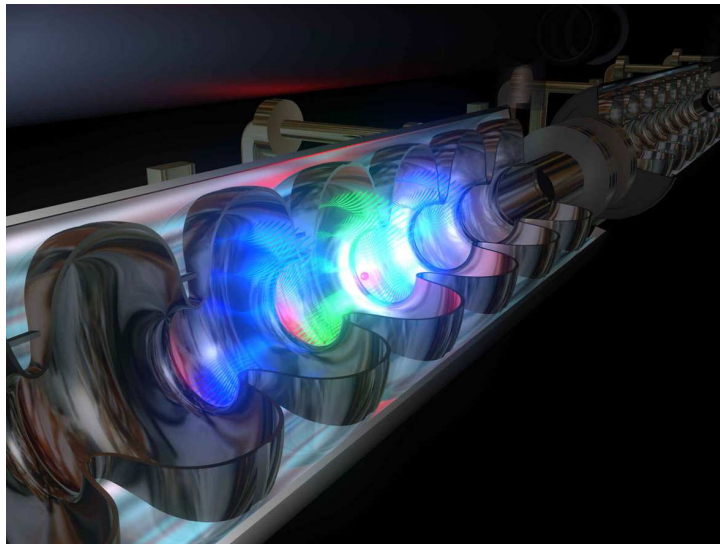
Part 3

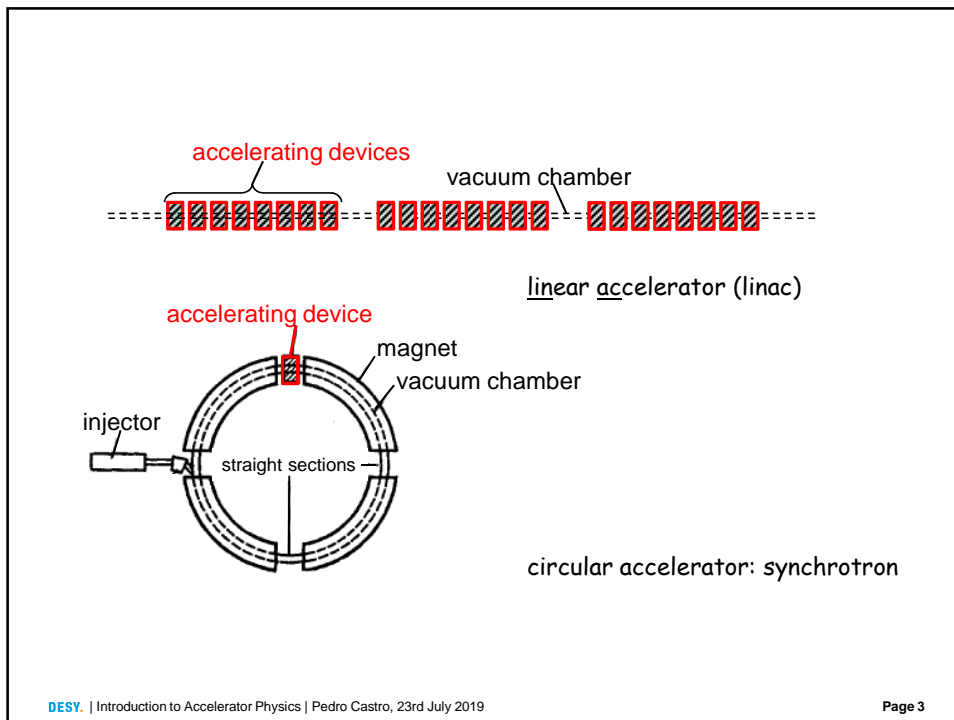
Pedro Castro / Accelerator Physics Group (MPY)
Hamburg, 23rd July 2019

HELMHOLTZ RESEARCH FOR GRAND CHALLENGES

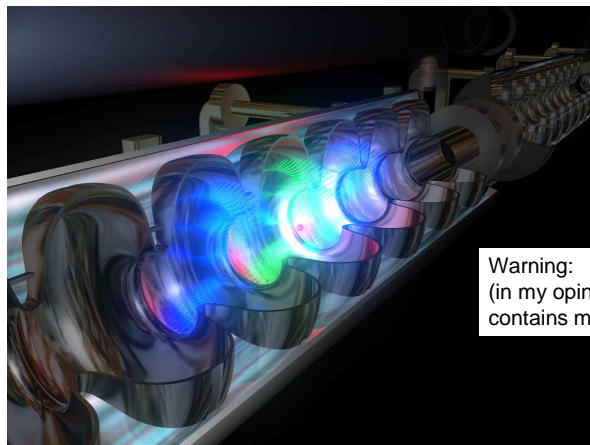


How electromagnetic fields accelerate particles





MEDIA DATABASE. "Electron acceleration – a virtual simulation"



Warning:
(in my opinion) this movie
contains many physical inaccuracies

DESY → Press → Media database → European XFEL (with filter: media type=movies)

<https://media.desy.de/DESYmediabank/?l=en#l=en&cid=3980&cname=European%20XFEL&f=2165&s=&p=&r=>

YouTube: https://www.youtube.com/watch?v=FJO_DmM4q7M
search text: electron acceleration

Motion in electric and magnetic fields

Equation of motion under Lorentz Force

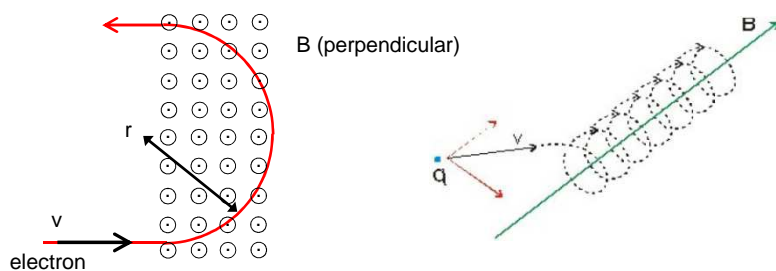
$$\frac{d\vec{p}}{dt} = \vec{F} = q (\vec{E} + \vec{v} \times \vec{B})$$

momentum charge velocity electric field magnetic field
 of the particle

Motion in magnetic fields

if the electric field is zero ($\vec{E} = 0$), then

$$\vec{F} = \frac{d\vec{p}}{dt} = q \cdot \vec{v} \times \vec{B} \rightarrow \vec{F} \perp \vec{v}$$



Magnetic fields do not change the particles energy

Motion in magnetic fields

if the electric field is zero ($E=0$), then

$$\vec{F} = \frac{d\vec{p}}{dt} = q \cdot \vec{v} \times \vec{B}$$

$$E^2 = \vec{p}^2 c^2 + E_0^2 \quad \text{energy-momentum relation in special relativity}$$

↑ total energy ↑ momentum ↑ energy at rest

Motion in magnetic fields

if the electric field is zero ($E=0$), then

$$\vec{F} = \frac{d\vec{p}}{dt} = q \cdot \vec{v} \times \vec{B}$$

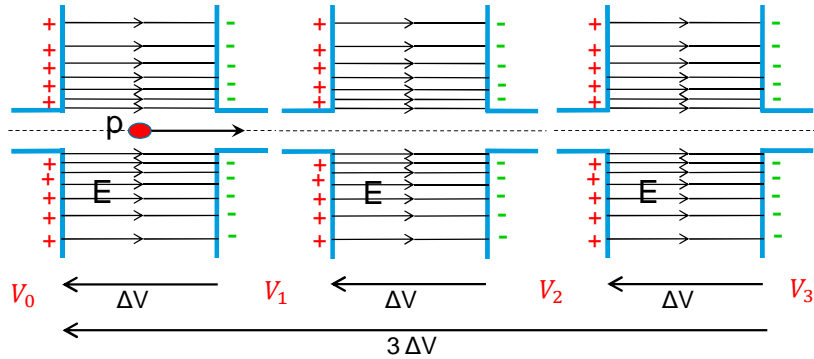
$$E^2 = \vec{p}^2 c^2 + E_0^2$$

$$E \frac{dE}{dt} = c^2 \vec{p} \frac{d\vec{p}}{dt} = c^2 q \vec{p} (\vec{v} \times \vec{B}) = c^2 q |\vec{p}| |\vec{v} \times \vec{B}| \cos \phi = 0$$

since $\vec{v} \times \vec{B} \perp \vec{v} \rightarrow \phi = 90^\circ$

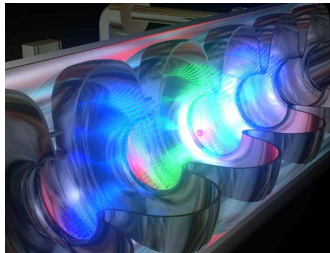
Magnetic fields do not change the particles energy, only electric fields do !

acceleration with DC electric fields



In general:

- Static magnetic fields \rightarrow to guide (bend + focus) particle beams
- Static electric fields \rightarrow accelerate particle beams (low energy)
- Radio-frequency EM fields \rightarrow accelerate particle beams (high E)



RF cavity basics: a cylindrical cavity

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RF cavity basics: a cylindrical cavity

LC circuit (or resonant circuit) analogy:

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RF cavity basics: a cylindrical cavity

a quarter of a period later: half a period later:

LC circuit (or resonant circuit) analogy:

angular frequency: $\omega = \frac{1}{\sqrt{LC}}$

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RF cavity basics: the pill box cavity

pill boxes

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Maxwell's equations
(differential formulation in SI units)

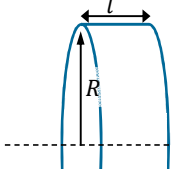
$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

+ boundary conditions



R : cavity radius
 l : cavity length

TM modes
(transverse magnetic modes)

set of solutions with $B_z = 0$ (that is, \vec{B} is transverse)

~~set of solutions with $E_z = 0$ (that is, \vec{E} is transverse)~~

~~**TE modes**
(transverse electric modes)~~

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Maxwell's equations
(differential formulation in SI units)

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

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+ boundary conditions

R : cavity radius
 l : cavity length

set of solutions with $B_z = 0$ (that is, \vec{B} is transverse)

$$E_z = E_0 J_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t}$$

$$E_r = -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t}$$

$$E_\theta = -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t}$$

$$B_z = 0$$

$$B_r = -j\omega \frac{mR^2}{x_{mn}^2 r c^2} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t}$$

$$B_\theta = -j\omega \frac{R}{x_{mn} c^2} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t}$$

indices:

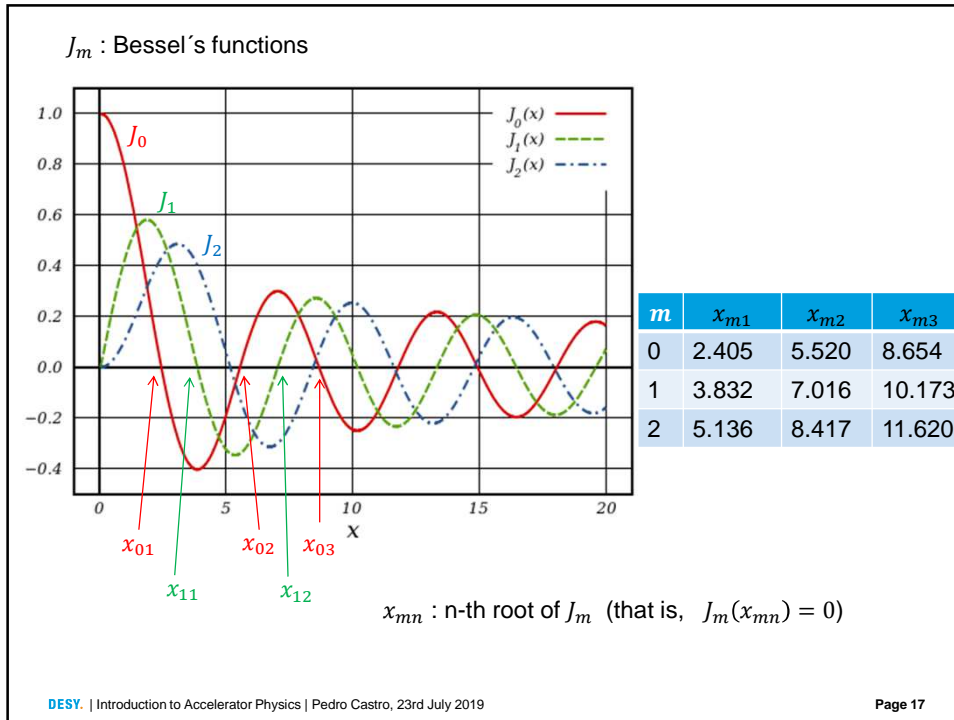
$m = 0, 1, 2, \dots$: number of full period variations in θ of the fields
 $n = 1, 2, \dots$: number of zeros of the axial field component in \tilde{r}
 $p = 0, 1, 2, \dots$: number of half period variations in z of the fields

J_m : Bessel's functions
 J'_m : derivative of the Bessel's functions

x_{mn} : n -th root of J_m (that is, $J_m(x_{mn}) = 0$)

angular frequency : $\omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$

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Maxwell's equations
(differential formulation in SI units)

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+ boundary conditions
 R : cavity radius
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set of solutions with $B_z = 0$ (that is, \vec{B} is transverse)

$$\begin{cases} E_z = E_0 J_m \left(\frac{x_{mn} r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_r = -\frac{p\pi R}{l} \frac{E_0}{x_{mn}} J'_m \left(\frac{x_{mn} r}{R} \right) \cos m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_\theta = -\frac{p\pi m R^2}{l} \frac{E_0}{x_{mn}^2 r} J'_m \left(\frac{x_{mn} r}{R} \right) \sin m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_z = 0 \\ B_r = -j\omega \frac{m R^2}{x_{mn}^2 r c^2} E_0 J_m \left(\frac{x_{mn} r}{R} \right) \sin m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_\theta = -j\omega \frac{R}{x_{mn} c^2} E_0 J'_m \left(\frac{x_{mn} r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \end{cases}$$

indices:
 $m = 0, 1, 2, \dots$: number of full period variations in θ of the fields
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Maxwell's equations
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$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

+ boundary conditions
R : cavity radius
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set of solutions with $B_z = 0$ (that is, \vec{B} is transverse)

$$\left\{ \begin{aligned} E_z &= E_0 J_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_r &= -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_\theta &= -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_z &= 0 \\ B_r &= -j\omega \frac{mR^2}{x_{mn}^2 r c^2} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_\theta &= -j\omega \frac{R}{x_{mn} c^2} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \end{aligned} \right.$$

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Maxwell's equations
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set of solutions with $B_z = 0$ (that is, \vec{B} is transverse)

$$\left\{ \begin{aligned} E_z &= E_0 J_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_r &= -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ E_\theta &= -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_z &= 0 \\ B_r &= -j\omega \frac{mR^2}{x_{mn}^2 r c^2} E_0 J_m \left(x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \\ B_\theta &= -j\omega \frac{R}{x_{mn} c^2} E_0 J'_m \left(x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left(\frac{p\pi}{l} z \right) e^{j\omega t} \end{aligned} \right.$$

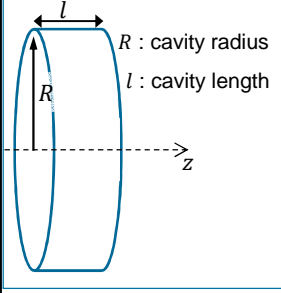
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angular frequency : $\omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$

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boundary conditions



R : cavity radius
 l : cavity length

fundamental solution with $B_z = 0$ (that is, \vec{B} is transverse)

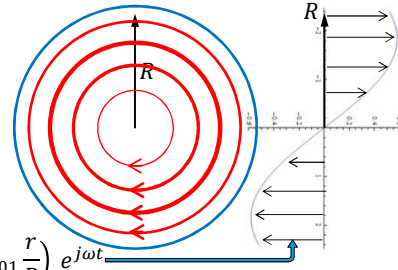
$$E_z = E_0 J_0 \left(x_{01} \frac{r}{R} \right) e^{j\omega t}$$

$$E_r = 0$$

$$E_\theta = 0$$

$$B_z = 0$$

$$B_r = 0$$

$$B_\theta = j\omega \frac{R}{x_{01} c^2} E_0 J_1 \left(x_{01} \frac{r}{R} \right) e^{j\omega t}$$


$m = 0$: rotation symmetry of the fields
 $n = 1$: no zeros of the axial field component in \vec{r}
 $p = 0$: no variation in z of the fields

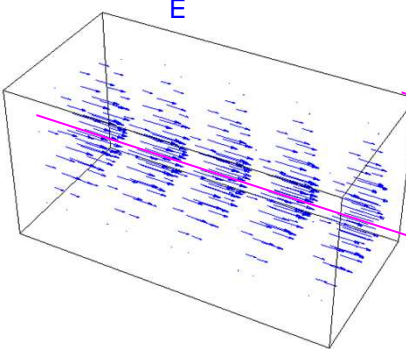
J_m : Bessel's functions
 J'_m : derivative of the Bessel's functions

angular frequency : $\omega = c \frac{x_{01}}{R}$ $x_{01} = 2.405$

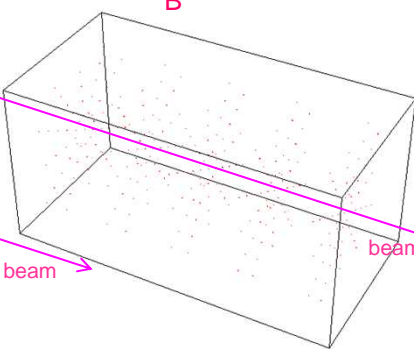
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Pill box cavity: 3D visualisation of E and B


E



B



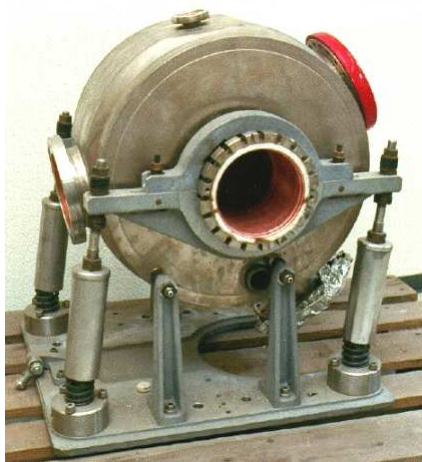
beam



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Examples of pill box cavities

DESY cavity (pill box)



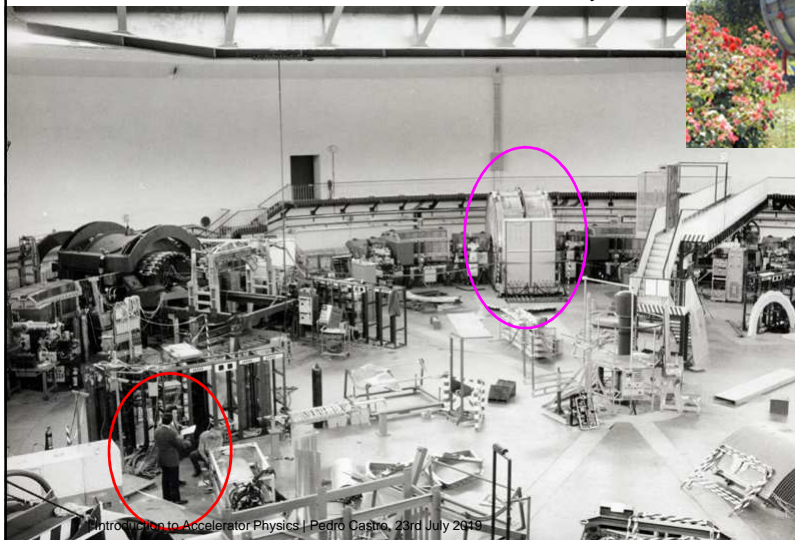
ADONE cavity 51 MHz (pill box)
Frascati lab, Italy



Examples of pill box cavities

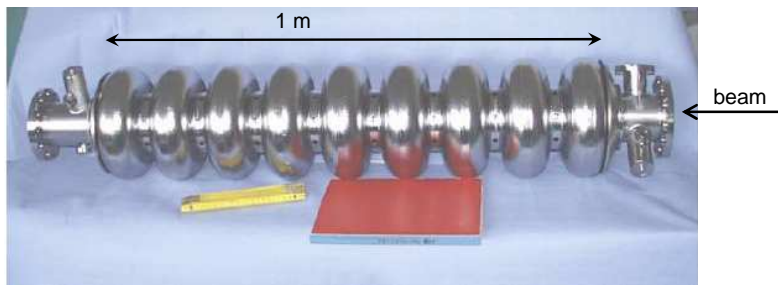
ADONE cavity 51 MHz (pill box)
Frascati lab, Italy

ADONE in 1963, Laboratori Nazionali di Frascati, Italy



Superconducting cavity used at DESY

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)



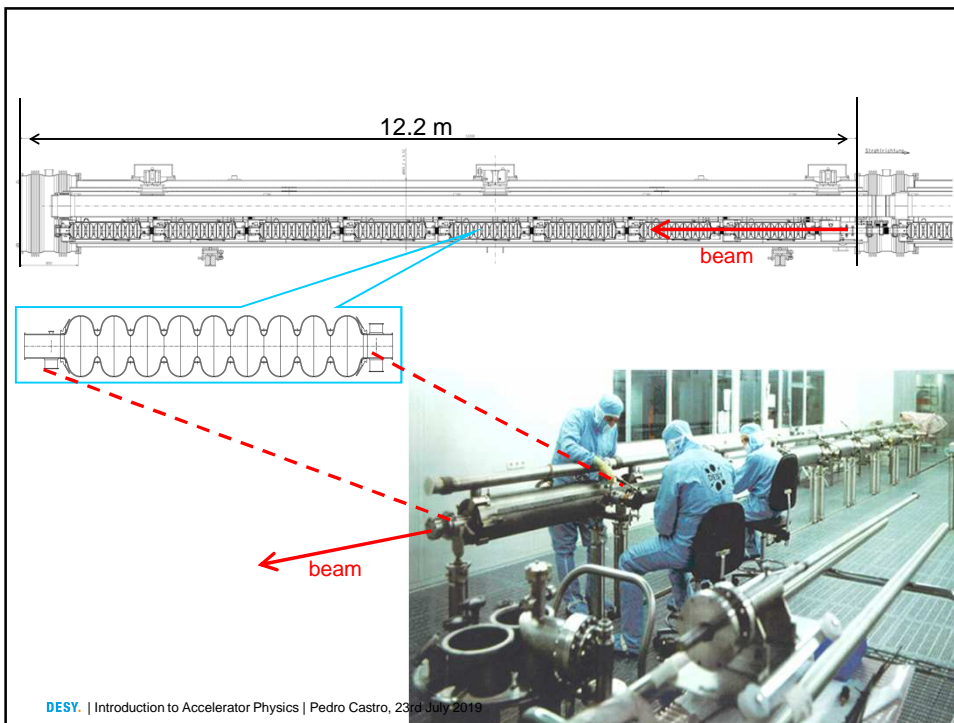
Free-electron <u>L</u> ASer in <u>H</u> amburg	0.3 km	DESY	2004-	?	e-	1.2 GeV
European <u>X</u> -ray Free-Electron <u>L</u> aser	3 km	DESY	2016-	?	e-	17.5 GeV
International <u>L</u> inear <u>C</u> ollider	30 km	?	?	?	e-/e+	2x250 GeV

Cavities inside a cryostat



Number of cavities 8
 Cavity length 1.038 m
 Operating frequency 1.3 GHz
 Operating temperature 2 K
 Accelerating Gradient 23.35 MV/m





Cavities inside a cryostat

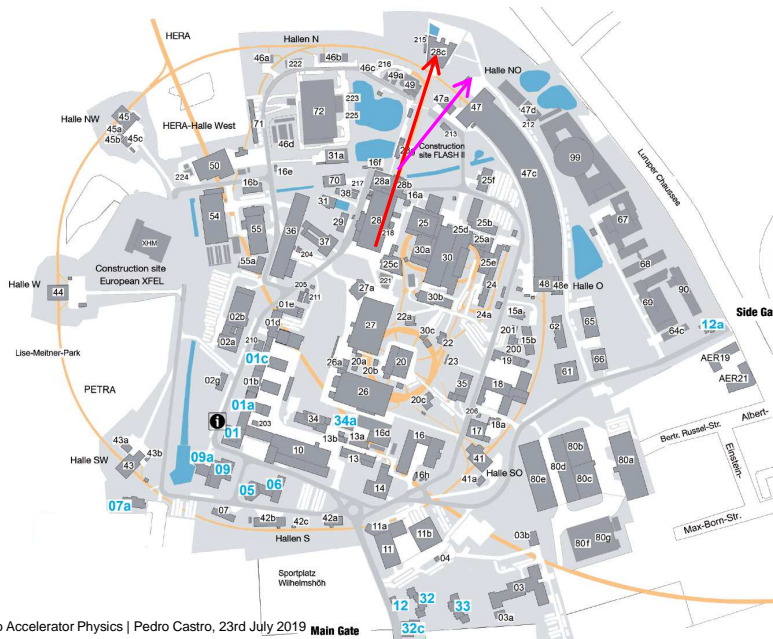


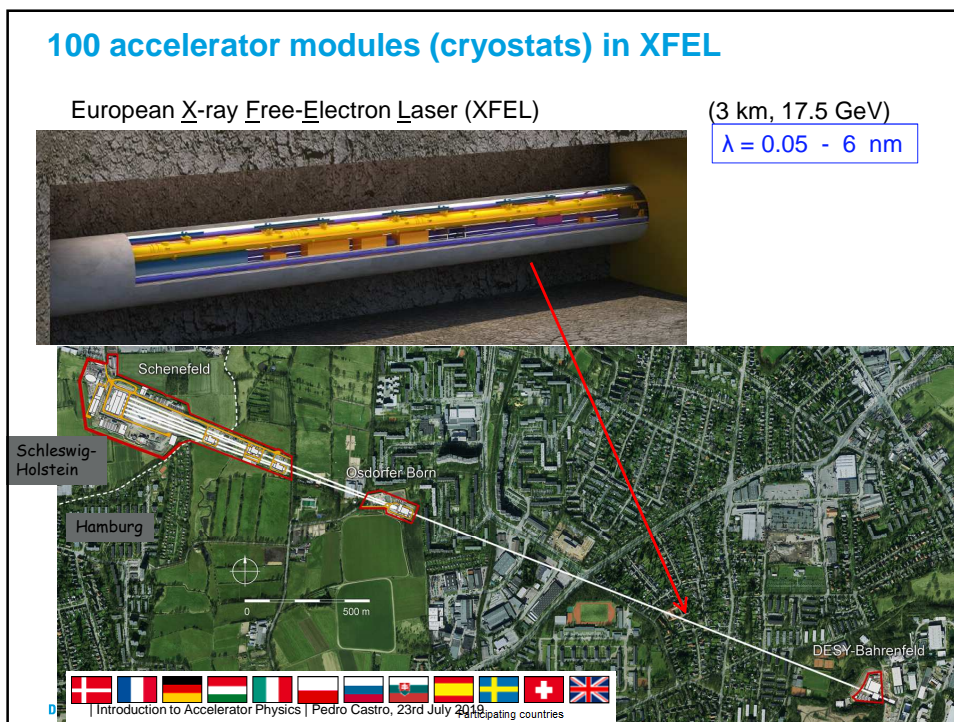
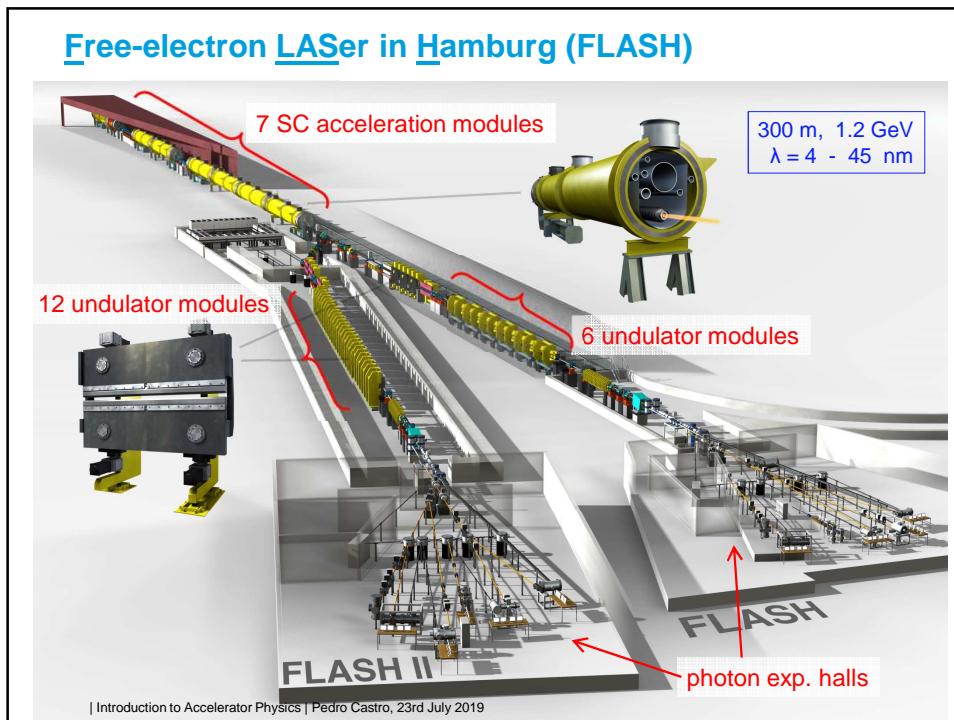
Cavities inside an accelerator module (cryostat)



module installation in FLASH (2004)

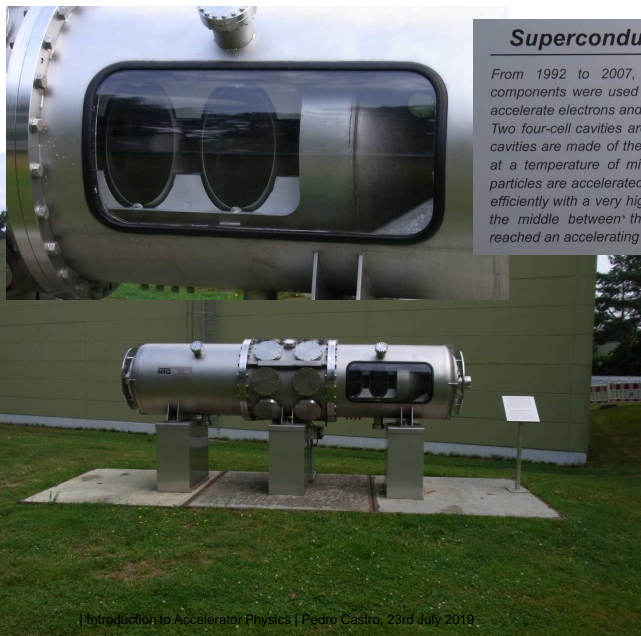
Free-electron LASer in Hamburg (FLASH)





Superconducting cavities at HERA

16 cavities
500 MHz



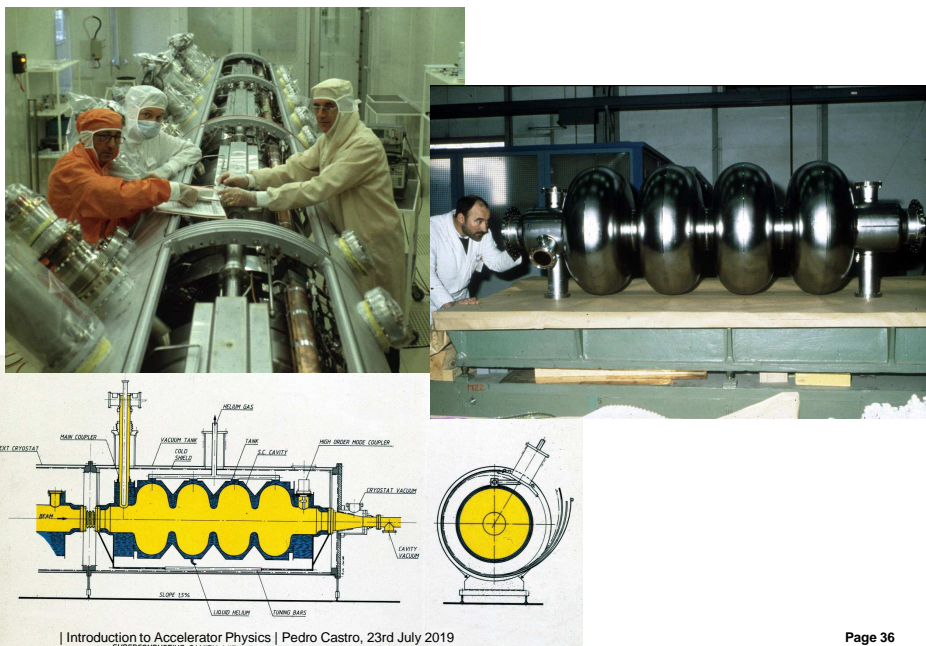
Superconducting Particle Accelerator

From 1992 to 2007, eight of these superconducting accelerator components were used in the 6.3-kilometre long storage ring HERA to accelerate electrons and their antiparticles, positrons. Two four-cell cavities are arranged in one thermal vessel (cryostat). The cavities are made of the metal niobium which becomes superconducting at a temperature of minus 269 degrees Celsius. At this temperature, particles are accelerated almost without electric resistance and thus very efficiently with a very high electric alternating voltage which is injected in the middle between the cavities. During HERA operation, this cavity reached an accelerating gradient of 5 million volts per metre.

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Superconducting cavities at LEP

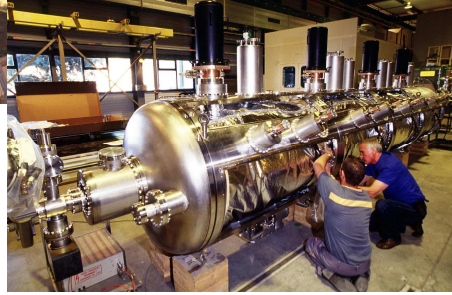
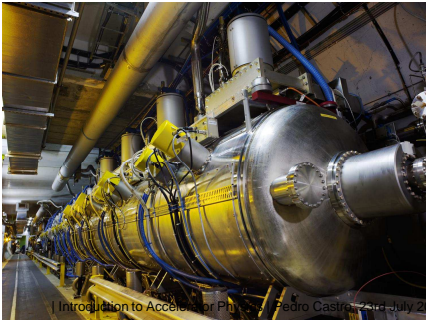
272 cavities
352 MHz



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SUPERCONDUCTING CAVITY WITH ITS CRYOSTAT

Superconducting cavities at LHC

16 cavities
400 MHz



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Other accelerators using superconducting cavities

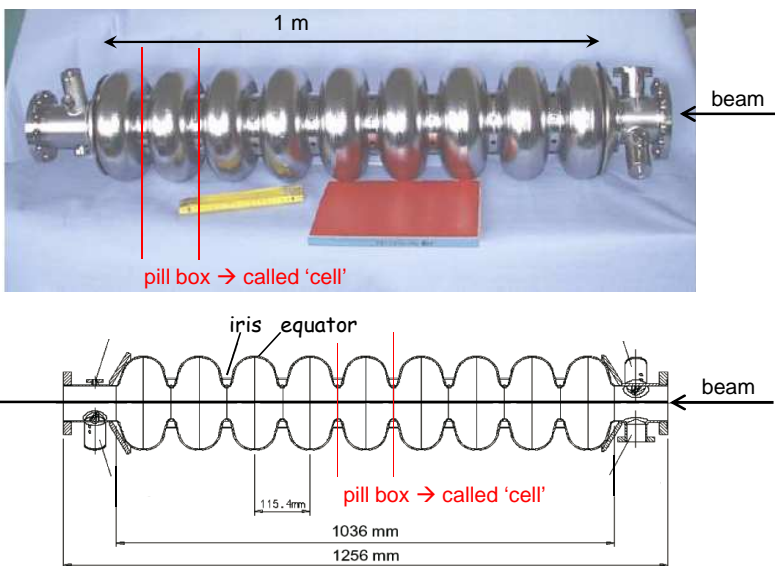
- 5 de-commissioned
- 11 in operation
- 4 in construction
- 9 in design phase

Total = 29

full list: http://tesla-new.desy.de/srf_accelerators

Superconducting cavity used in FLASH and in XFEL

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)

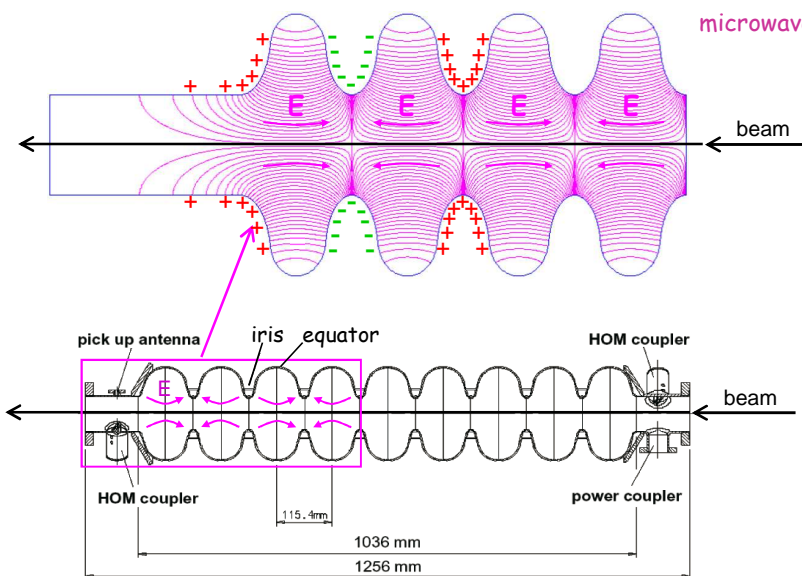


Accelerating field map

Simulation of the fundamental mode: electric field lines

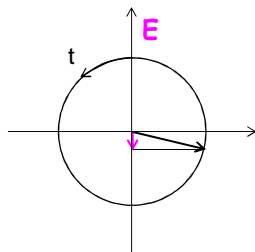
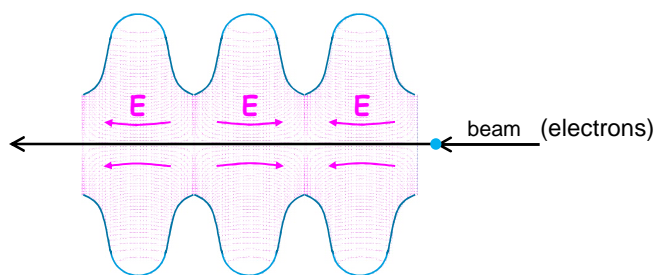
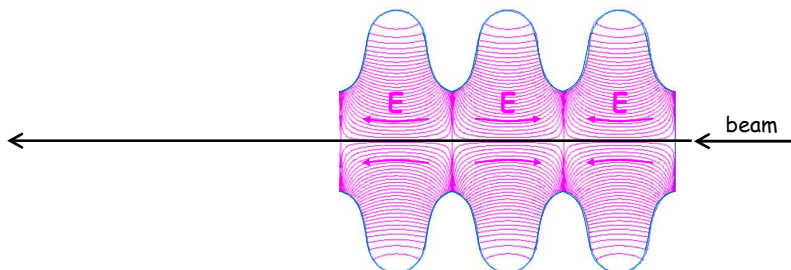
$f_{RF} = 1.3 \text{ GHz}$

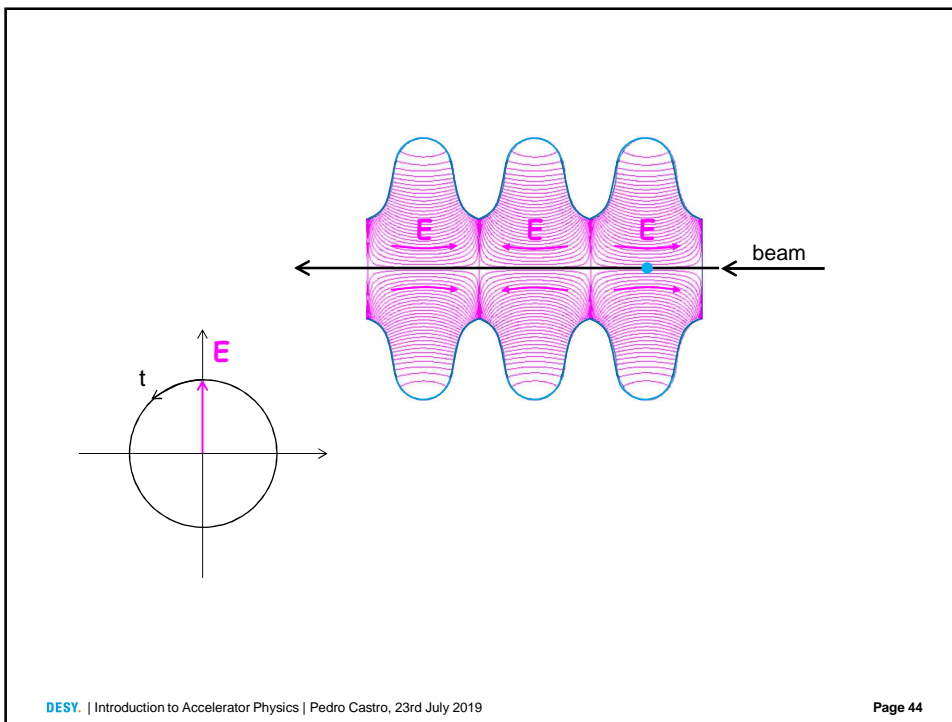
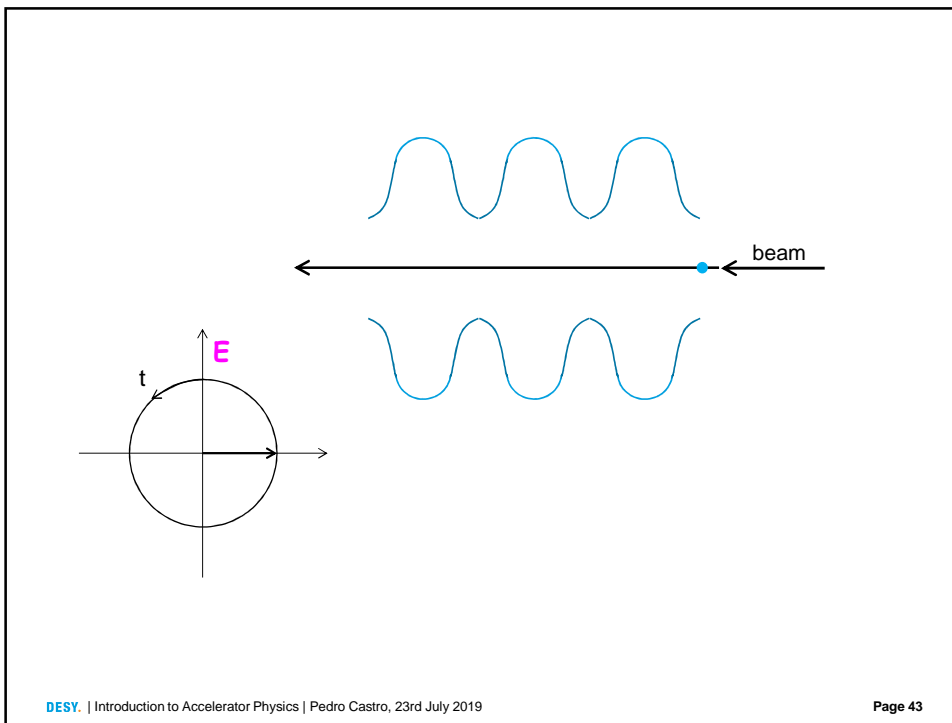
microwaves: (L-band)

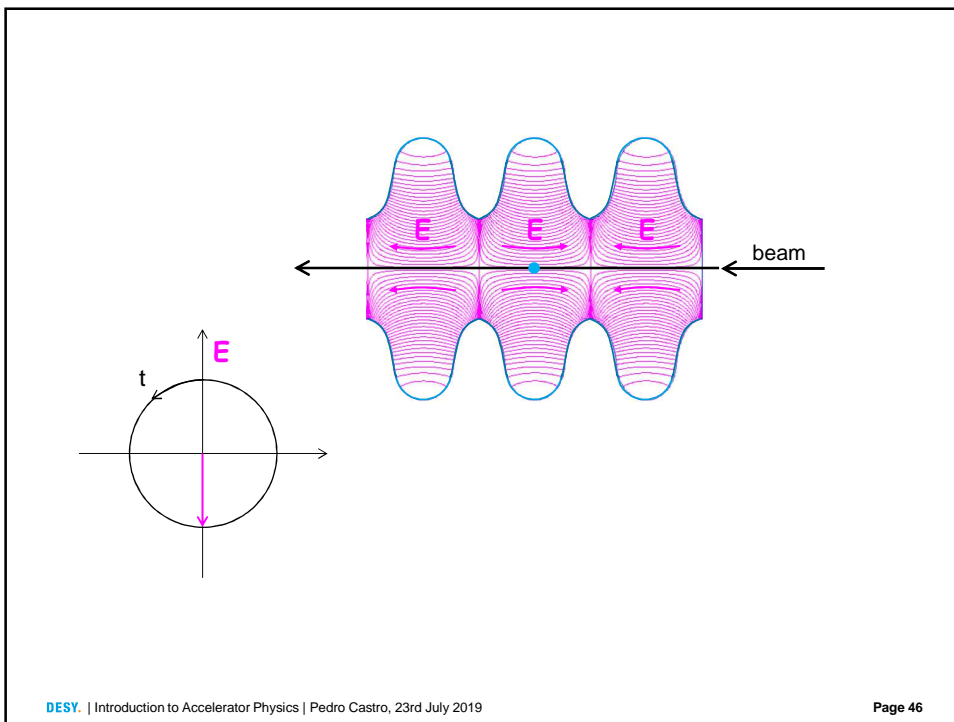
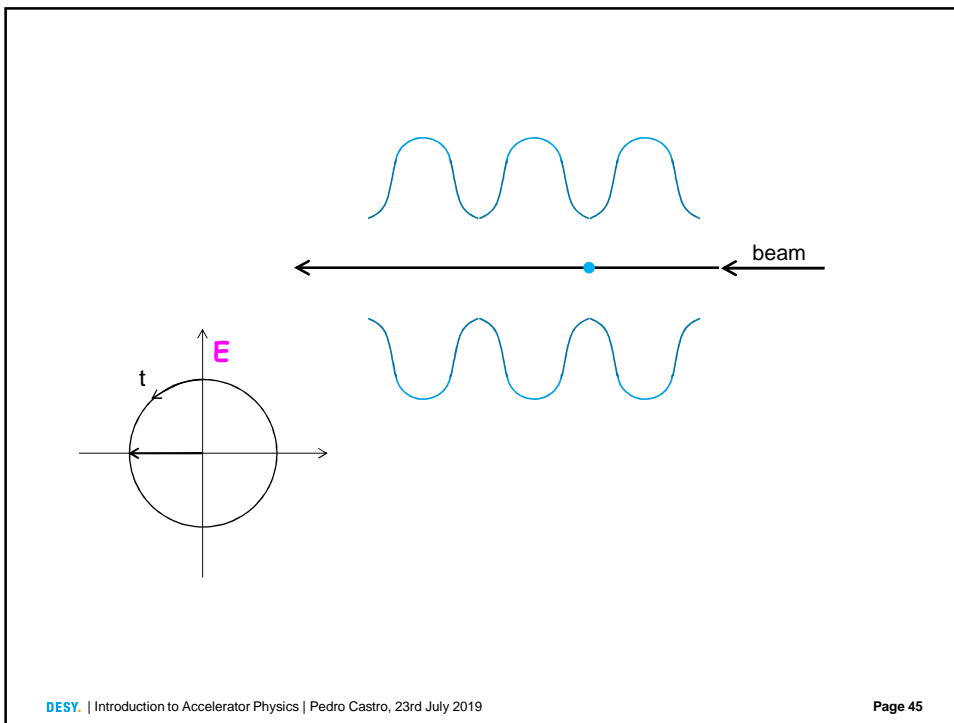


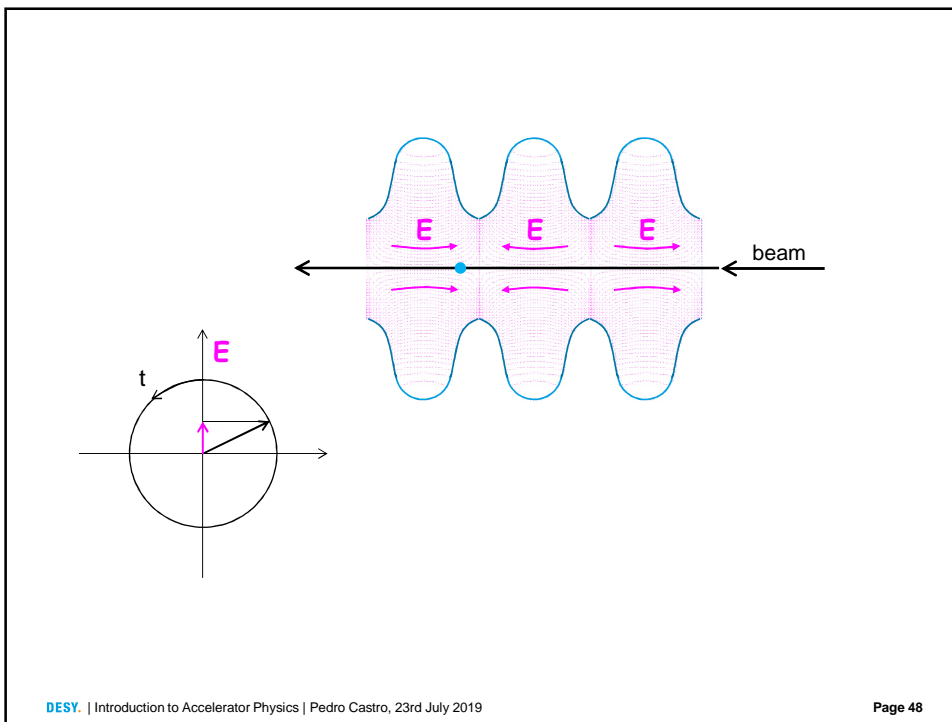
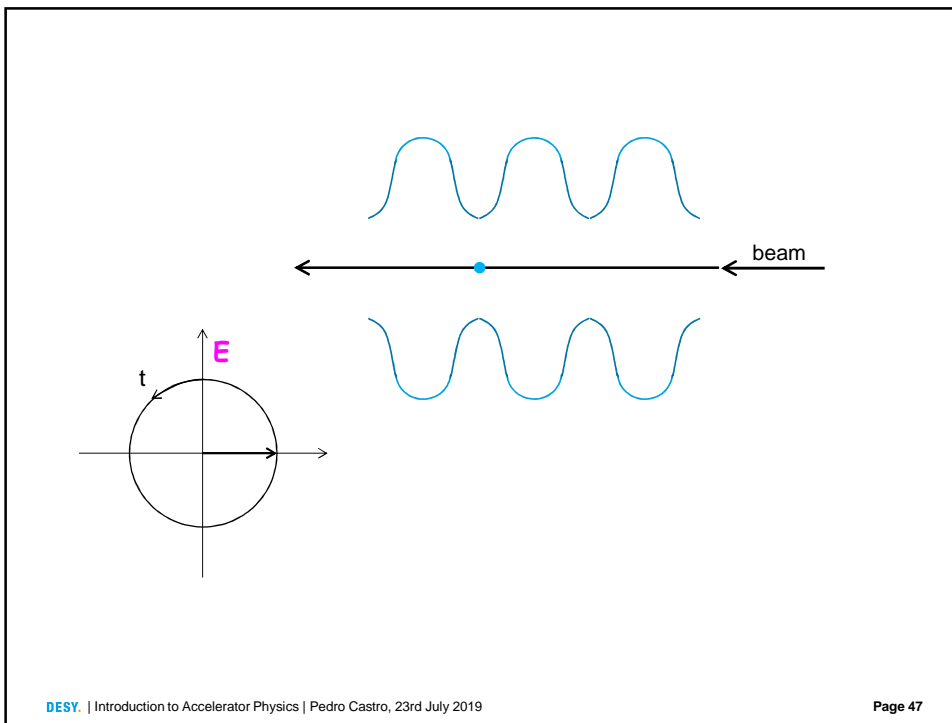
Is there a net acceleration?

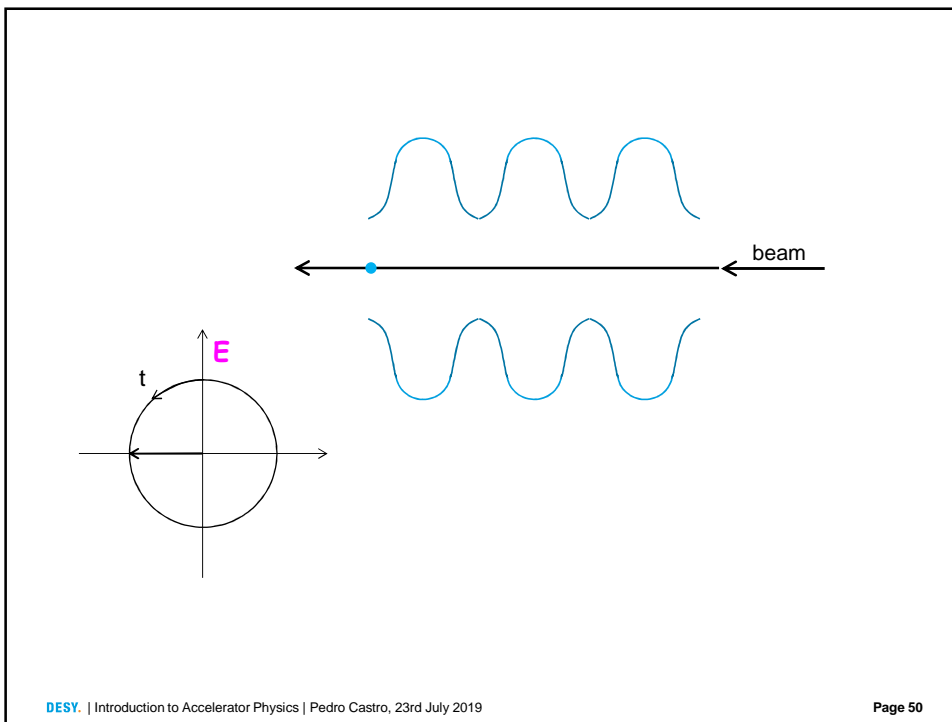
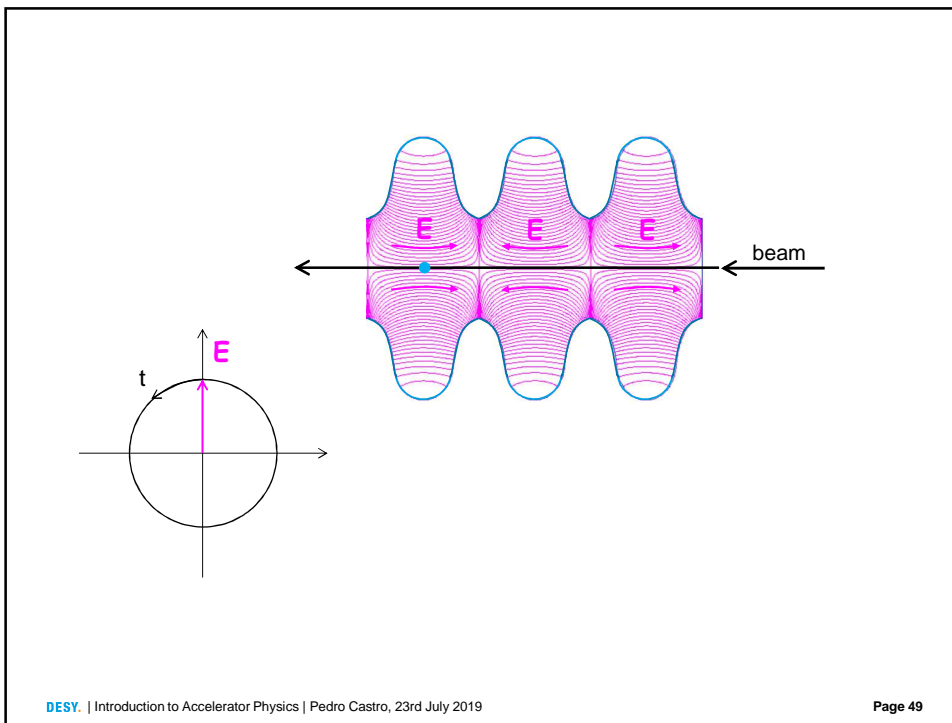
Simulation of the fundamental mode: electric field lines











Is there a net acceleration? timing is the key

for electrons, $\beta \cong 1$

$l = cT = \frac{c}{f} = \frac{3 \cdot 10^8}{1.3 \cdot 10^9} = 0.23 \text{ m} \quad (2 \text{ cells})$

1 cavity (1.038 m) / 9 cells = 0.115 m

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Superconducting cavity used at DESY

1 m

beam

material: pure Niobium

operating temperature: 2 K

accelerating field gradient: up to 35 MV/m

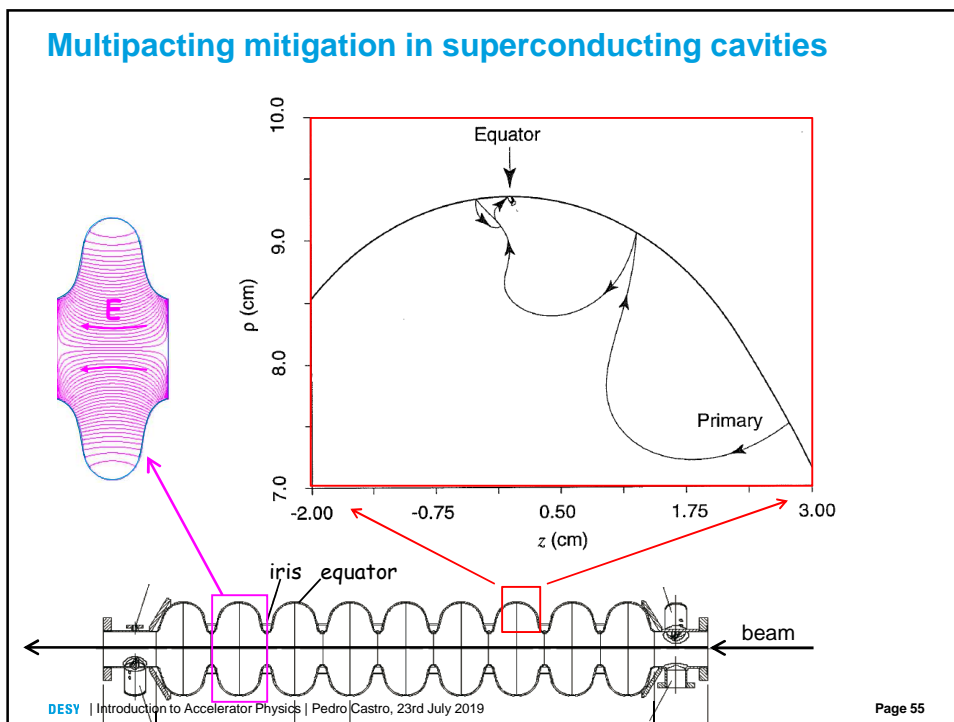
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Thermal conductivity

	$[W m^{-1} K^{-1}]$
water	0.56 – 0.61
copper (at 20 °C)	385 – 401
helium II	$> 10^5$

Frequently Asked Questions

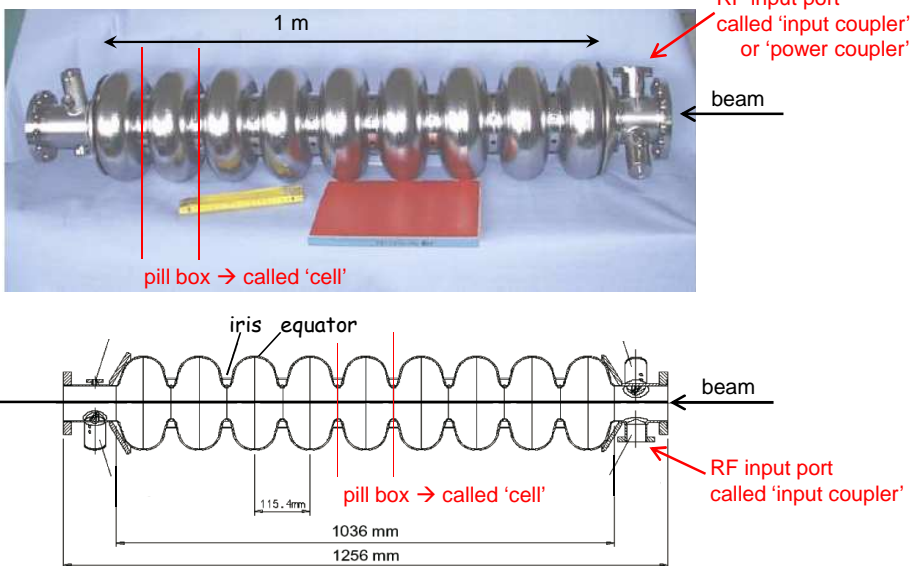
- 1) Why this shape?
- 2)
- 3)
- 4)



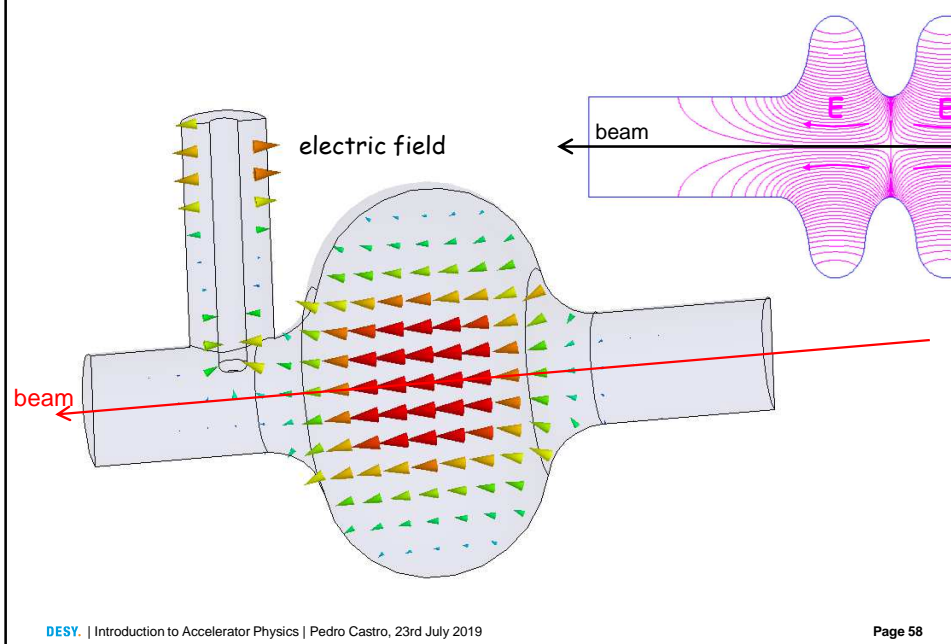
- 1) Why this shape? to reduce/avoid multipacting
- 2) How to feed \vec{E} in?
- 3)
- 4)

Superconducting cavity used in FLASH and in XFEL

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)



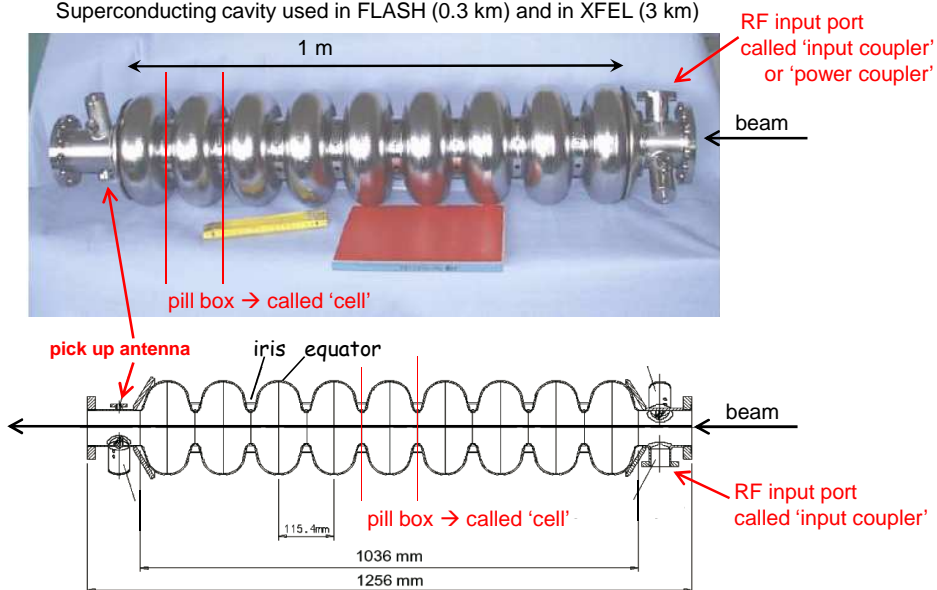
Fundamental mode coupler (input coupler)



- 1) Why this shape? to reduce/avoid multipacting
- 2) How to feed \vec{E} in? with input couplers
- 3) How to measure \vec{E} ?
- 4)

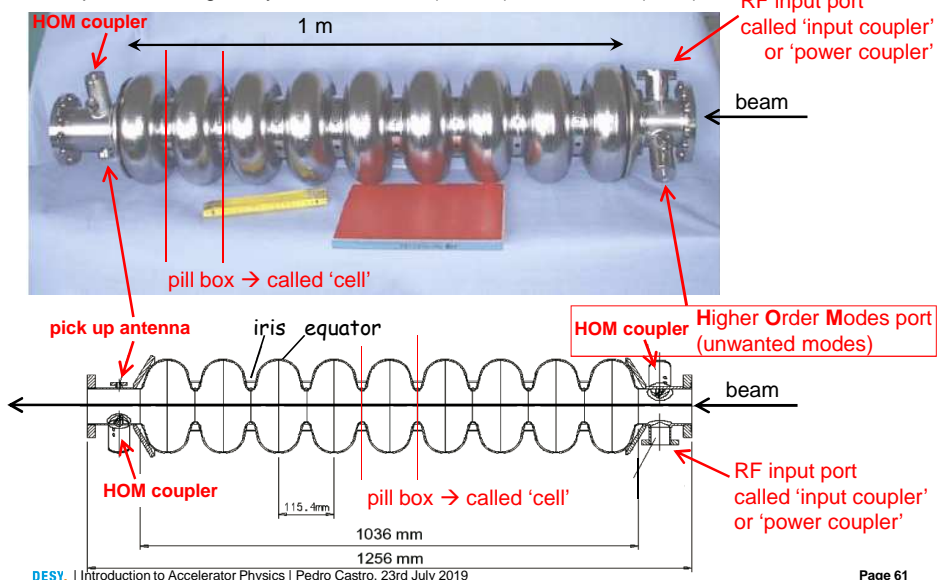
Superconducting cavity used in FLASH and in XFEL

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)

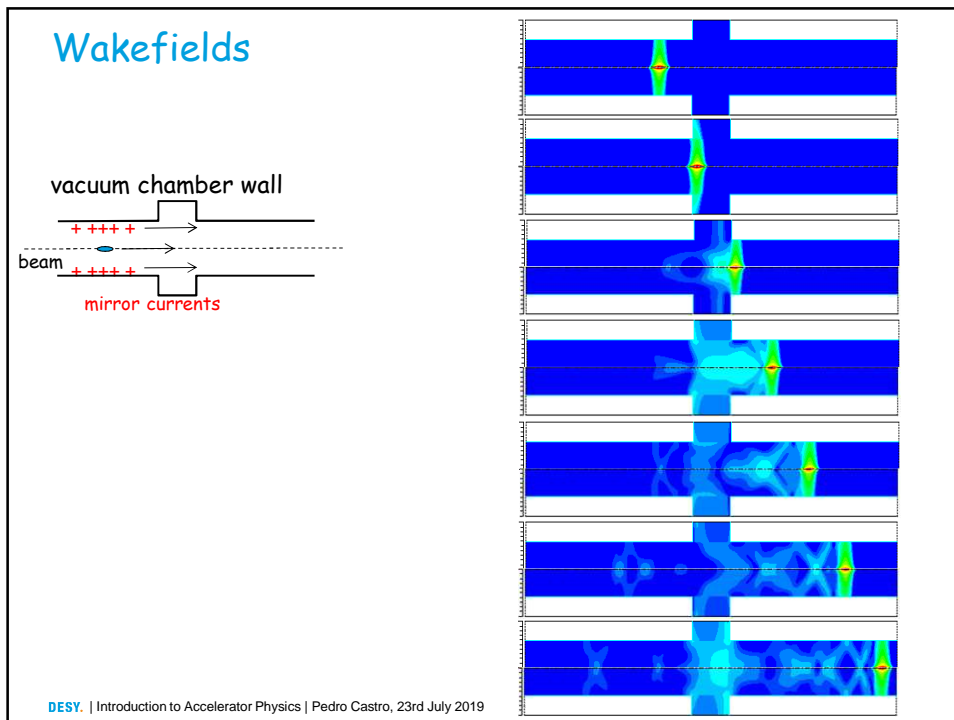


Superconducting cavity used in FLASH and in XFEL

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)



- 1) Why this shape? to reduce/avoid multipacting
- 2) How to feed \vec{E} in? with input couplers
- 3) How to measure \vec{E} ? with pick up antennas
- 4) What are HOM couplers for?



- 1) Why this shape? to reduce/avoid multipacting
- 2) How to feed \vec{E} in? with input couplers
- 3) How to measure \vec{E} ? with pick up antennas
- 4) What are HOM couplers for? to reduce HOM (wakefields)

Summing-up of this part

Particle acceleration using radio-frequency fields:

basic cavity: pill box	{	analogy to an LC circuit infinite number of solutions for \vec{E} and \vec{B} eq. for the fundamental solution for \vec{E} and \vec{B}
superconducting cavity	{	multipacting mitigation RF couplers and antennas wakefields and HOMs FLASH and XFEL

Contact

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