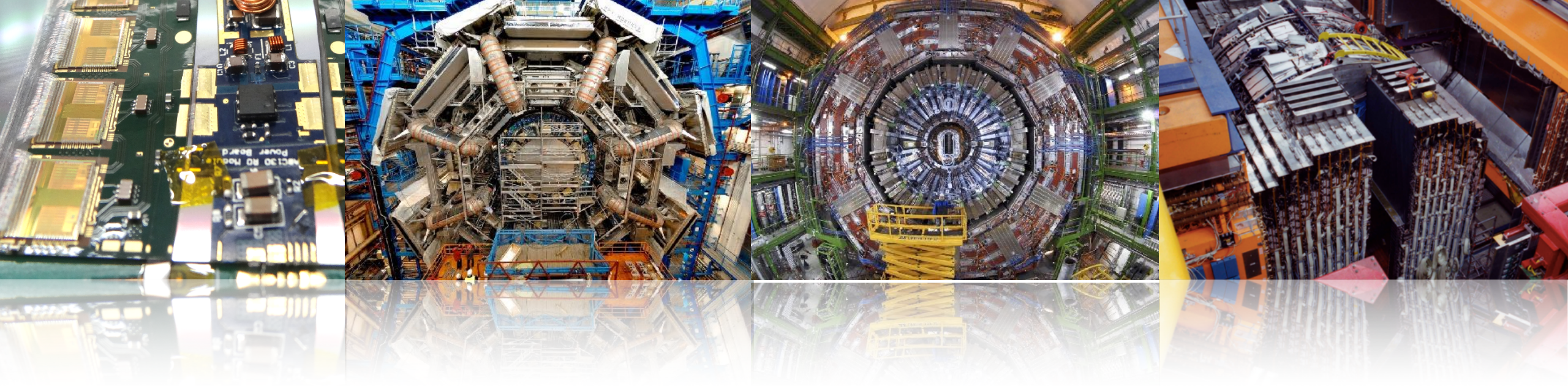


# DETECTORS FOR HIGH ENERGY PHYSICS



Ingrid-Maria Gregor  
DESY/Universität Bonn  
Summerstudents 2019  
30.07.2019





# DETECTORS FOR HIGH ENERGY PHYSICS

Part 2



Ingrid-Maria Gregor  
DESY/Universität Bonn  
Summerstudents 2019  
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# OVERVIEW

I. Detectors for Particle Physics

II. Interaction with Matter

III. Calorimeters

IV. Tracking Detectors

- Gas detectors
- Semiconductor trackers

V. Examples from the real life

}

*Tuesday*

}

*Wednesday*

### III. CALORIMETERS



# CALORIMETRY



# CALORIMETRY



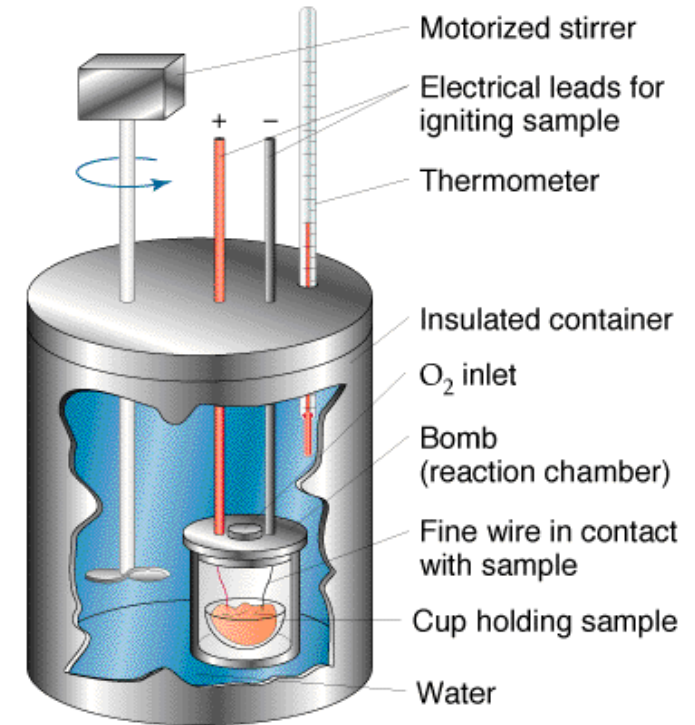


# CALORIMETRY: THE IDEA BEHIND IT ....



- Calorimetry originated in thermo-dynamics
  - The total energy released within a chemical reaction can be measured by measuring the temperature difference

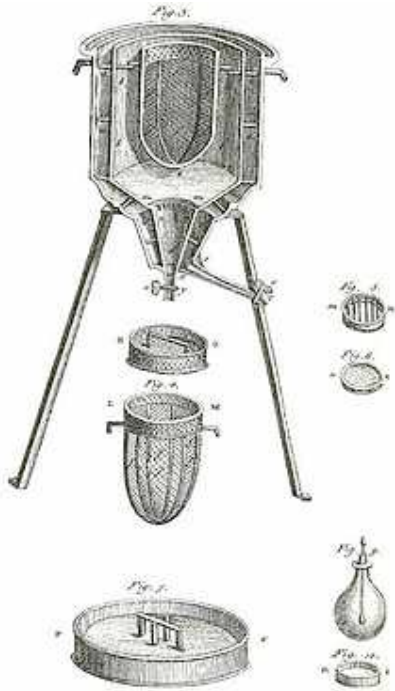
Ice-calorimeter from Antoine Lavoisier's 1789 *Elements of Chemistry*.



- What is the effect of a 1 GeV particle in 1 litre water (at 20°C)?

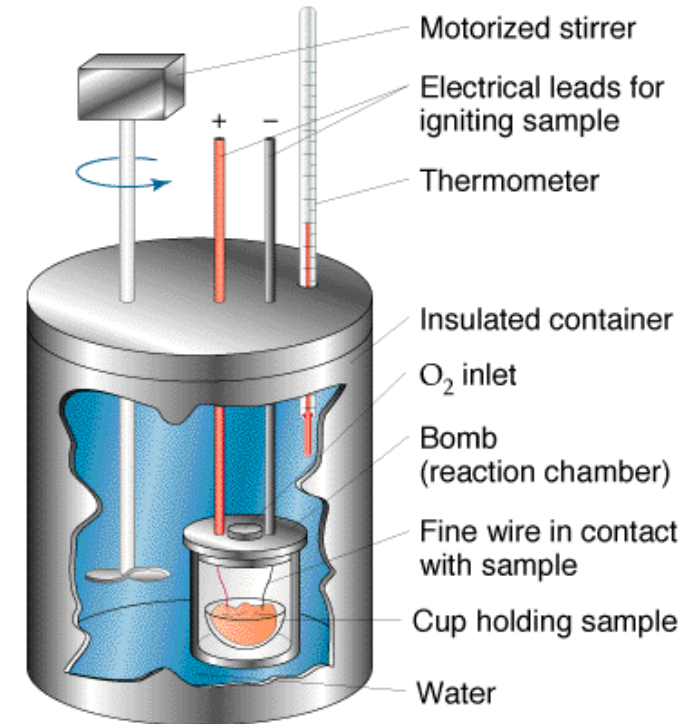
Picture: Francois G. Amar

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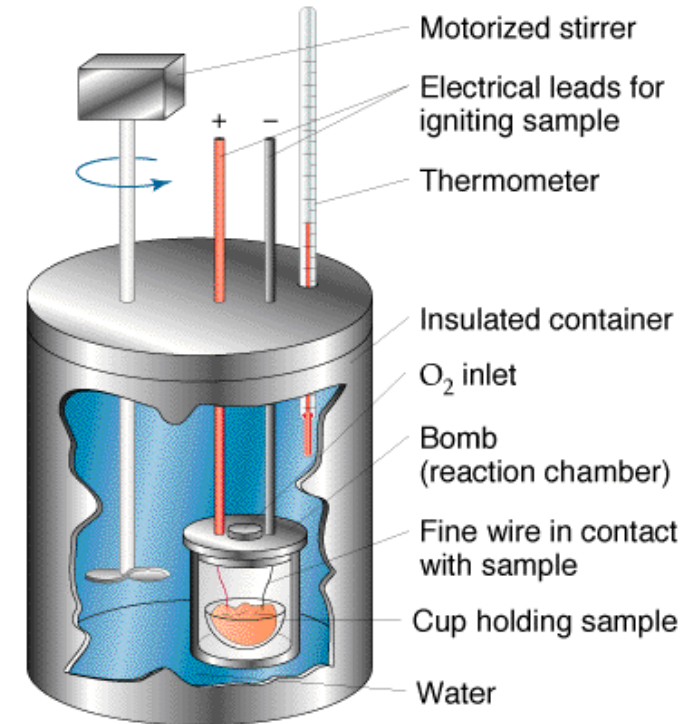


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$$\Delta T = E / (c \cdot M_{\text{water}}) = 3.8 \cdot 10^{-14} \text{K} !$$



- In particle physics:
  - Measurement of the energy of a particle by measuring the total absorption

# CALORIMETRY: OVERVIEW

- Basic mechanism for calorimetry in particle physics:
  - formation of electromagnetic
  - or hadronic showers.
- The energy is converted into ionisation or excitation of the matter.



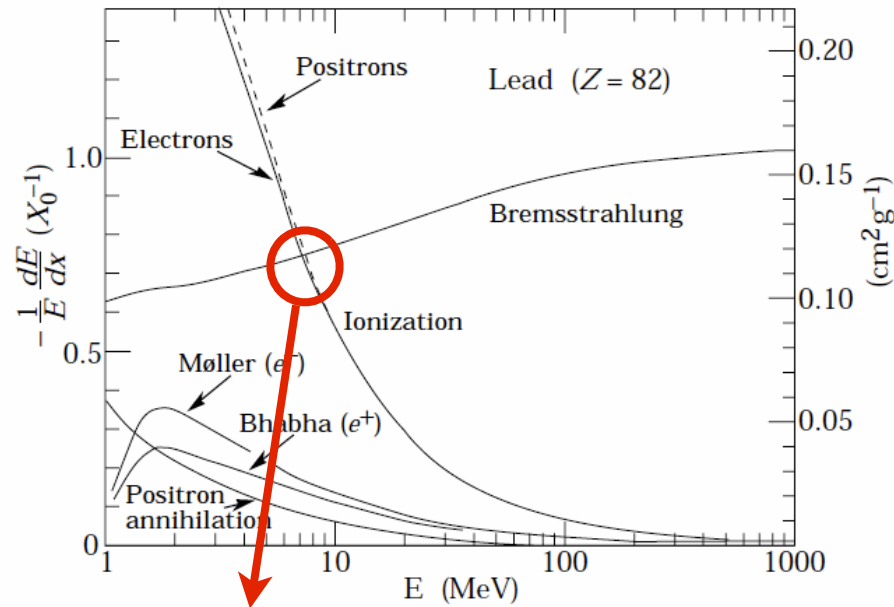
- Calorimetry is a “destructive” method. The energy and the particle get absorbed!
- Detector response  $\propto E$
- Calorimetry works both for charged ( $e^\pm$  and hadrons) and neutral particles ( $n, \gamma$ ) !



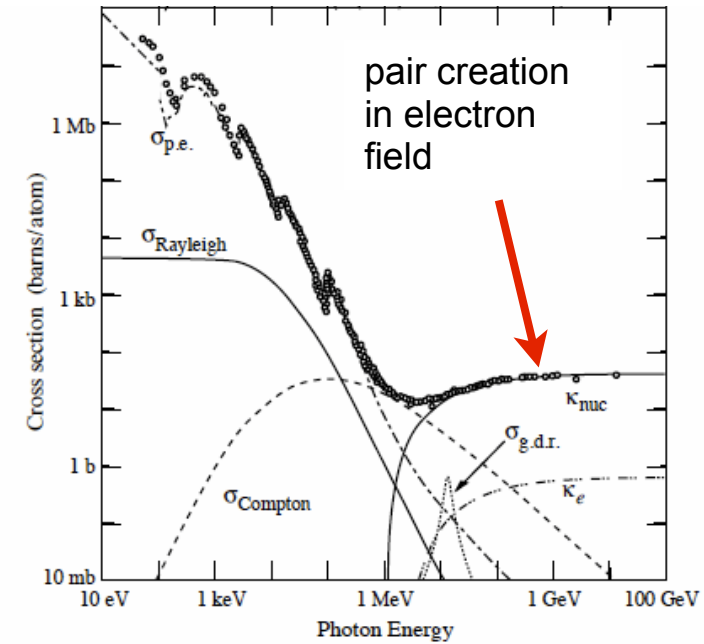


# REMINDER

## Electrons



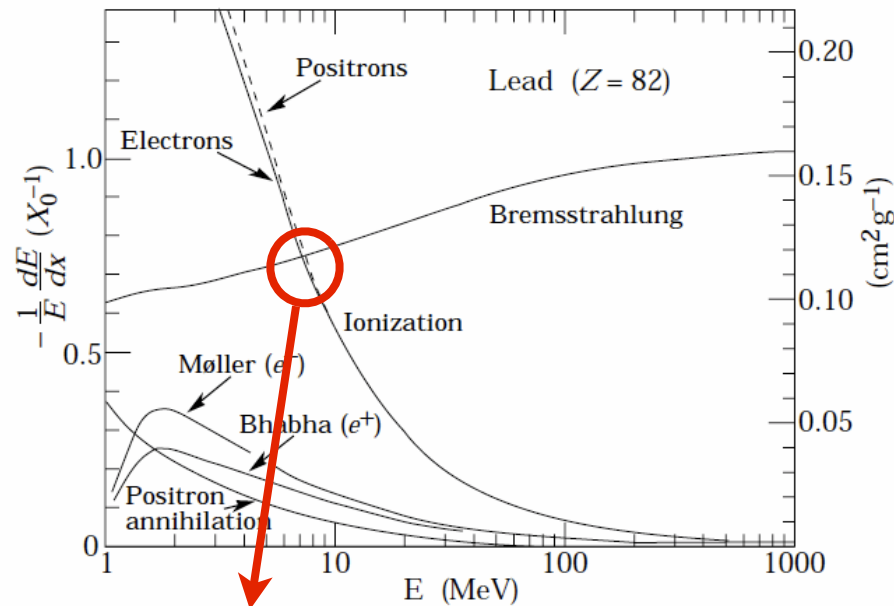
## Photons



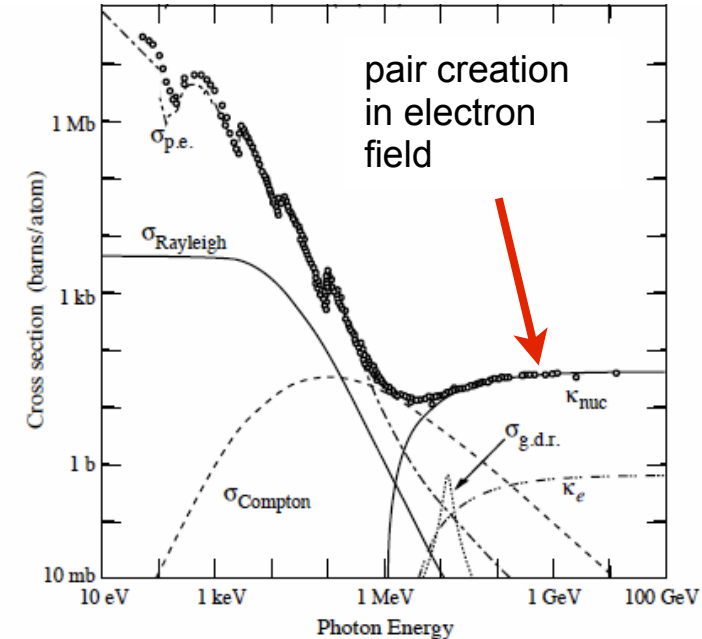
- **Critical energy:** the energy at which the losses due to ionisation and Bremsstrahlung are equal
- **Radiation length** defines the amount of material a particle has to travel through until the energy of an electron is reduced by Bremsstrahlung to 1/e of its original energy  $\langle E_e(x) \rangle \propto e^{-\frac{x}{X_0}}$

# REMINDER

## Electrons



## Photons



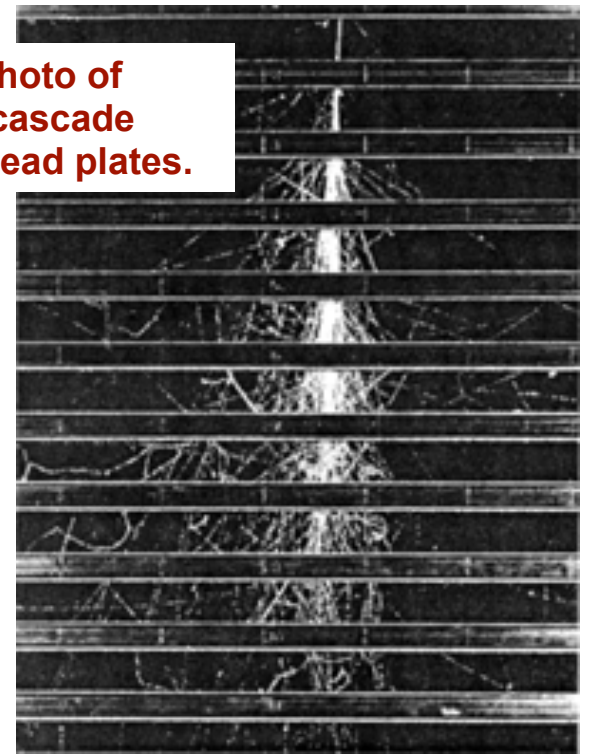
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empirical: 
$$X_0 = \frac{716.4 A}{Z(1+Z) \ln(287/\sqrt{Z})} \frac{g}{cm^2} \propto \frac{A}{Z^2}$$

# ELECTROMAGNETIC SHOWERS

- High energetic particles: forming a shower if passing through (enough) matter.
- An alternating sequence of interactions leads to a cascade:
  - Primary  $\gamma$  with  $E_0$  energy produces  $e^+e^-$  pair with 54% probability in layer  $X_0$  thick
  - On average, each has  $E_0/2$  energy
  - If  $E_0/2 > E_c$ , they lose energy by Bremsstrahlung

Cloud chamber photo of  
electromagnetic cascade  
between spaced lead plates.



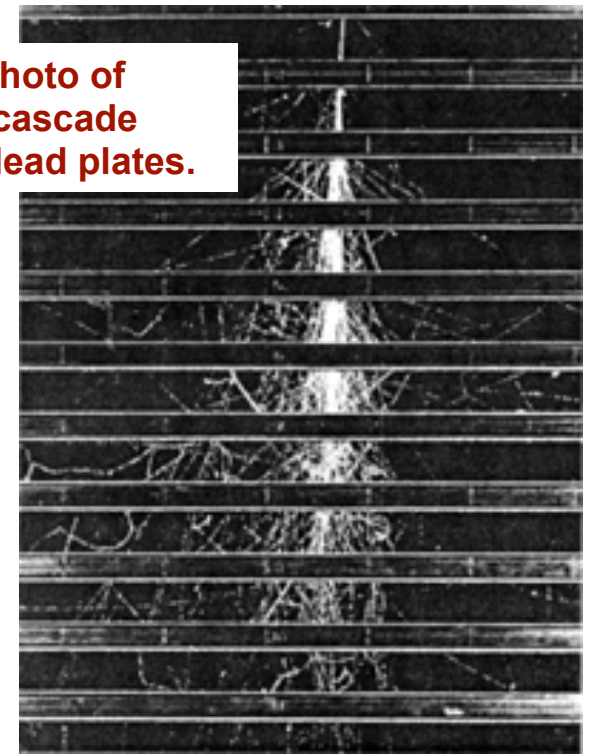
Pic: MIT cosmic ray group



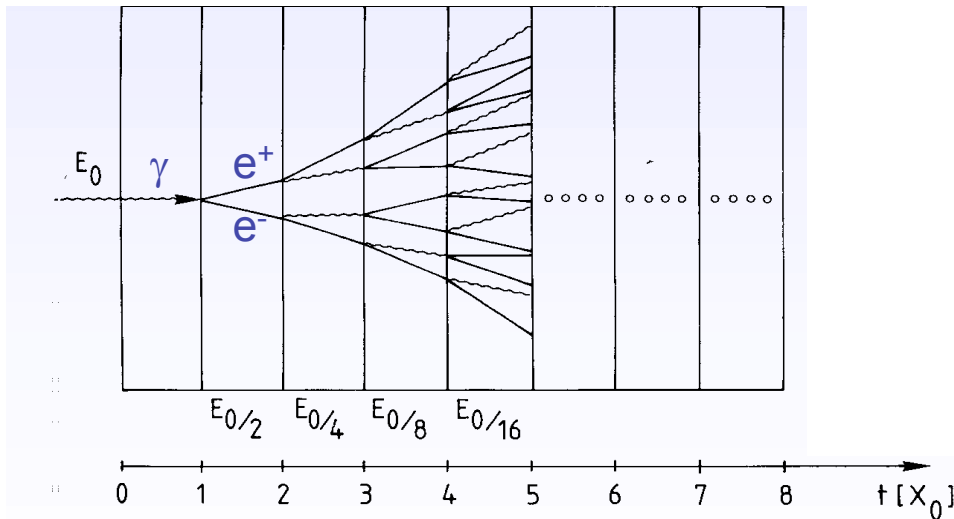
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Cloud chamber photo of electromagnetic cascade between spaced lead plates.



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- Next layer  $X_0$ , charged particle energy decreases to  $E_0/(2e)$
- Bremsstrahlung with an average energy between  $E_0/(2e)$  and  $E_0/2$  is radiated
- Radiated  $\gamma$ s produce again pairs
- After  $t$  radiation lengths
  - number of particles
  - each with average energy

$$N \simeq 2^t$$

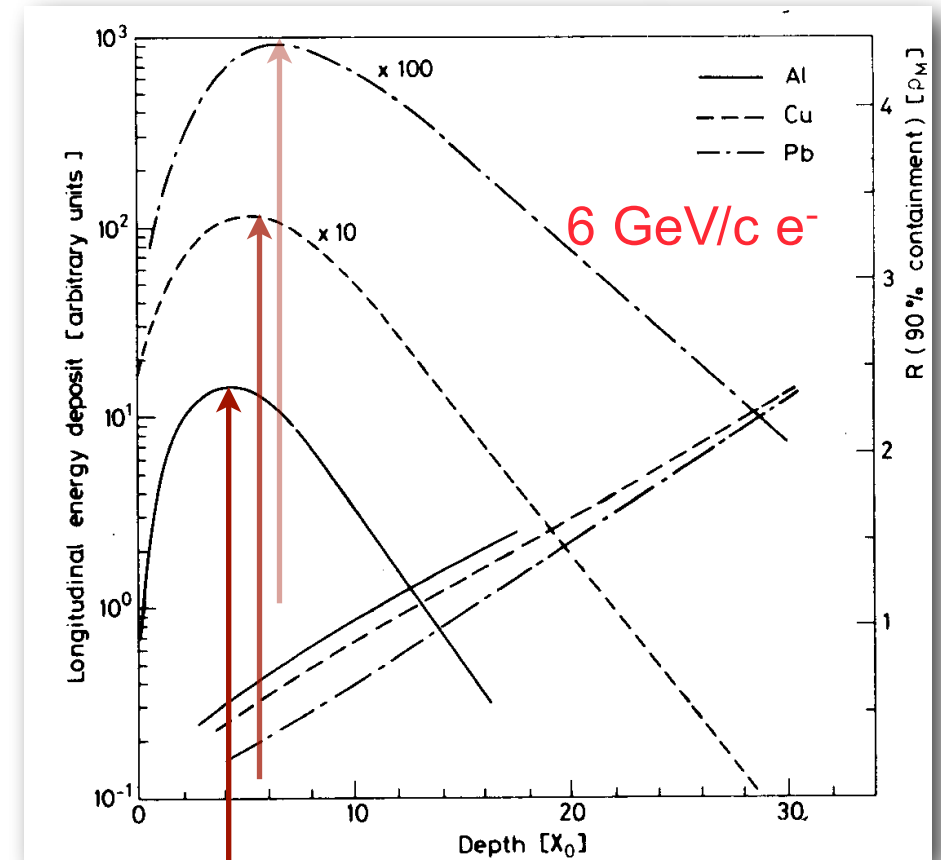
$$E_N \simeq \frac{E_0}{2^t}$$

# EM SHOWER PROPERTIES

- Shower continues until energy of particles below critical energy.

$$E(t_{max}) = \frac{E_0}{2^{t_{max}}} = E_c$$

$$t_{max} = \frac{\ln \frac{E_0}{E_c}}{\ln 2} \quad N_{max} \simeq \frac{E_0}{E_c}$$



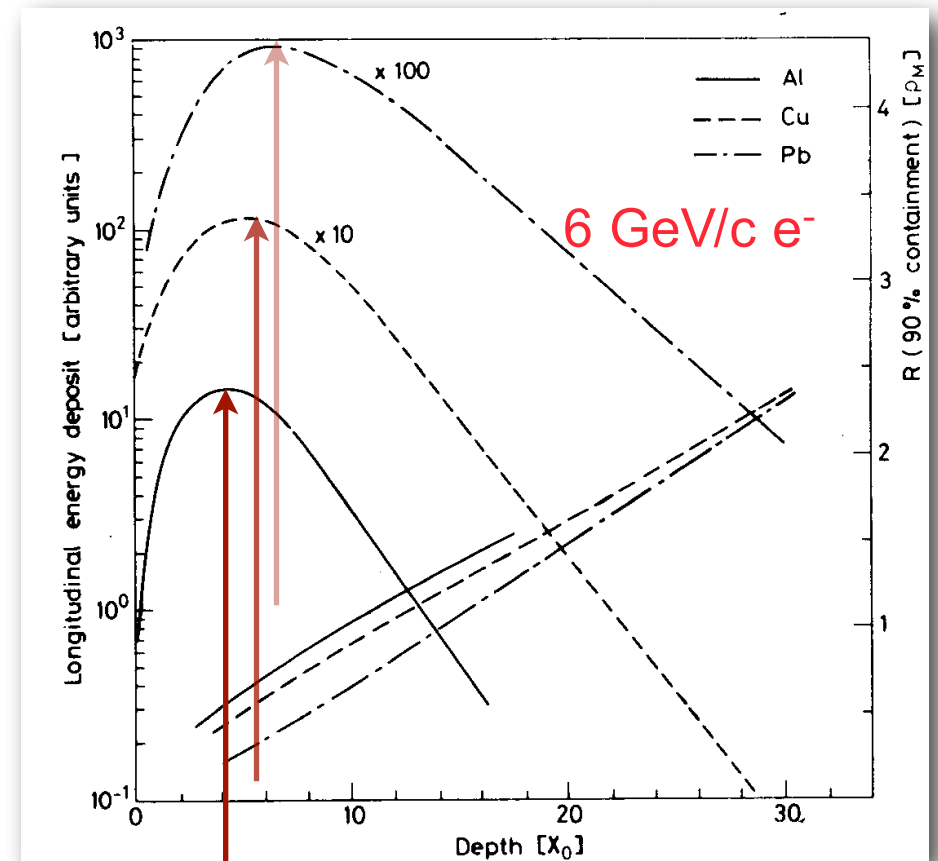
Shower maximum at  $t_{max}$

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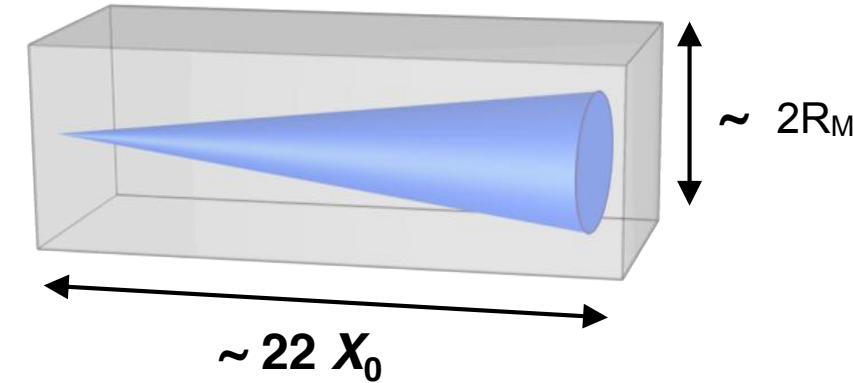


- Simple model only, for more details MC simulation required.
- Shower curve should rise rapidly to a peak value and then fall to zero.
- The broad peak of the experimental curve can be interpreted in terms of an **energy spread** of the incoming particles.
- Long tail due to **muon interactions** producing knock-on electrons capable of making a contribution to the cascade process.



# EM SHOWER PROPERTIES

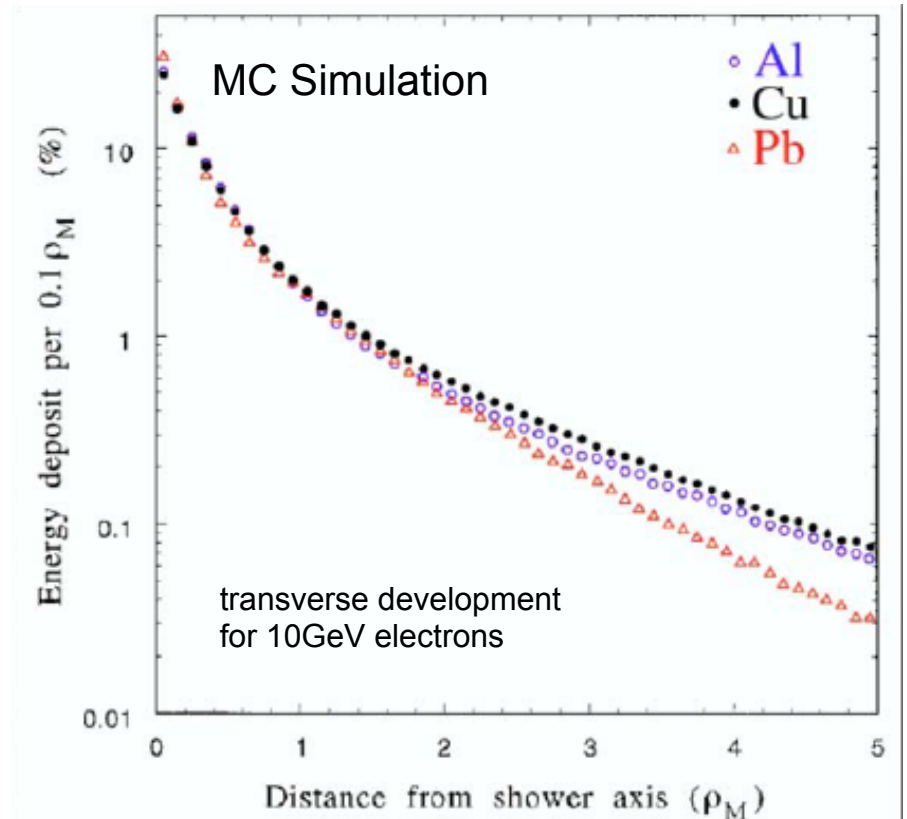
- Longitudinal development governed by the radiation length  $X_0$ .
- Lateral spread due to electron undergoing multiple Coulomb scattering:
  - 95% of the shower cone is located in a cylinder with radius  $2 R_M$
  - Beyond this point, electrons are increasingly affected by multiple scattering
- Lateral width scales with the **Molière radius  $R_M$** 
  - Important parameter for shower separation



$$R_M = X_0 \frac{E_s}{E_c} = 21.2 \text{ MeV} * \frac{X_0}{E_c}$$

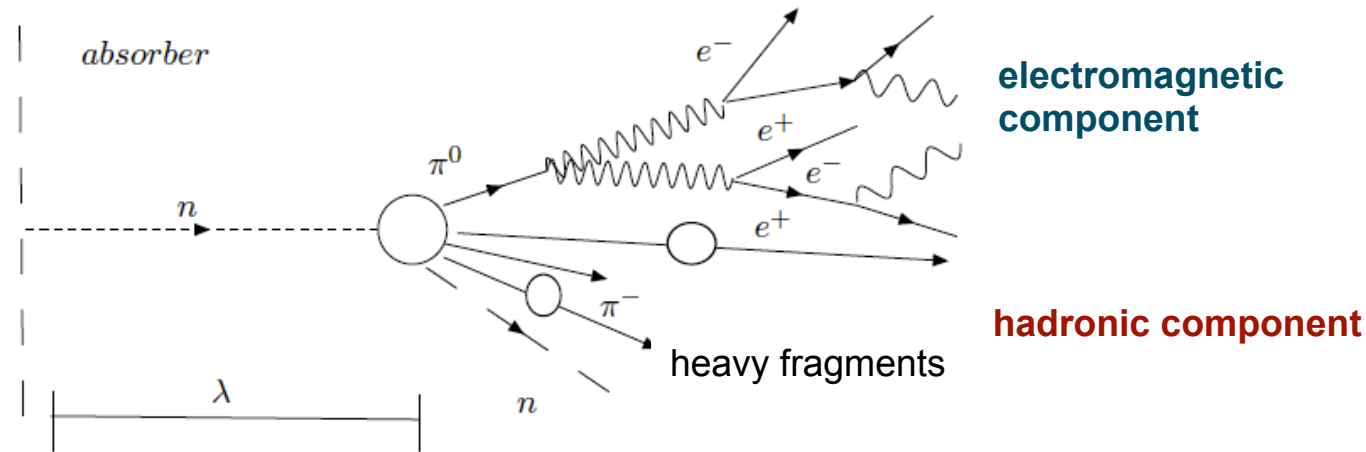
$$E_s = m_e c^2 \sqrt{4\pi/\alpha} = 21.2 \text{ MeV}$$

Example:  
 $E_0 = 100 \text{ GeV}$   
 in lead glass  $E_c = 11.8 \text{ MeV}$   
 $\rightarrow N_c \approx 13, t_{95\%} \approx 23$   
 $X_0 \approx 2 \text{ cm}, R_M = 1.8 \cdot X_0 \approx 3.6 \text{ cm}$



# HADRONIC CASCADE

- Within the calorimeter material a hadronic cascade is build up: in inelastic nuclear processes more hadrons are created



The length scale of the shower is given in means of the nuclear reaction length  $\lambda_l$

$$\lambda_l = \frac{A}{N_A \sigma_{total}}$$

total cross section for  
nuclear processes

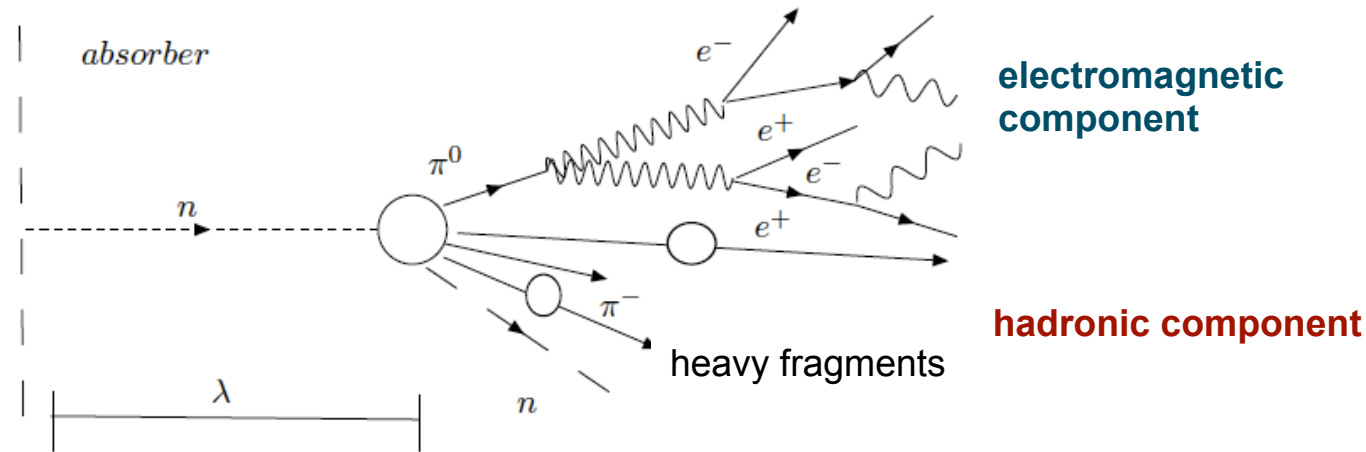
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Probability that no hadronic reaction happens on the path  $x$  happened:

$$P = e^{-\frac{x}{\lambda_I}}$$

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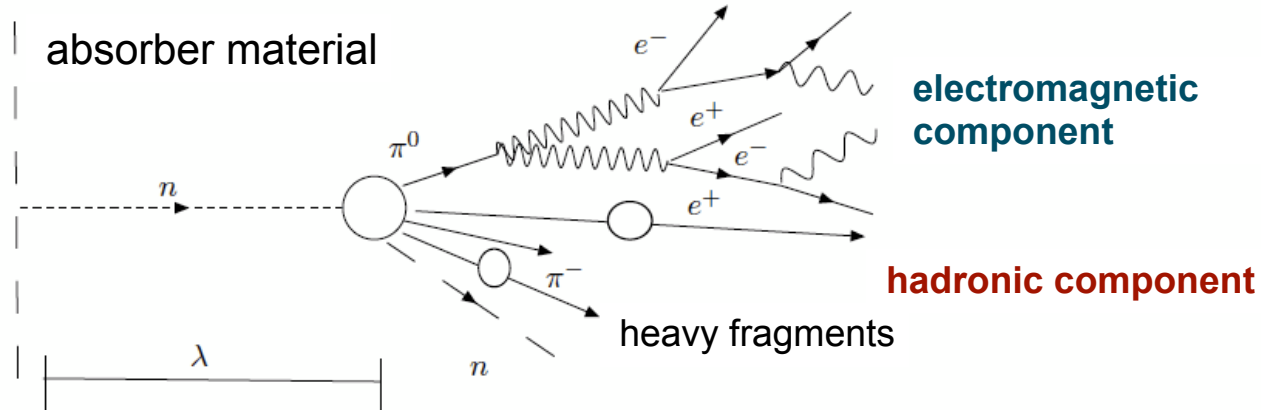
$$P = e^{-\frac{x}{\lambda_I}}$$

Compare  $X_0$  for high-Z materials, we see that the size needed for hadron calorimeters is large compared to EM calorimeters.

	$\lambda_I$	$X_0$
Polystyren	81.7 cm	43.8 cm
PbWO	20.2 cm	0.9 cm
Fe	16.7 cm	1.8 cm
W	9.9 cm	0.35 cm



# HADRONIC CASCADE: THE DETAILS



Hadronic showers are way more complicated than em showers.

- Different processes are created by the impinging hadron:
  - high energetic secondary hadrons taking a significant part of the momentum of the primary particle [e.g. O(GeV)]
  - a significant part of the total energy is transferred into nuclear processes: nuclear excitation, spallation, ...  $\Rightarrow$  Particles in the MeV range
  - neutral pions (1/3 of all pions), decay instantaneously into two photons  $\Rightarrow$  start of em showers
  - Breaking up of nuclei (binding energy) neutrons, neutrinos, soft  $\gamma$ 's, muons

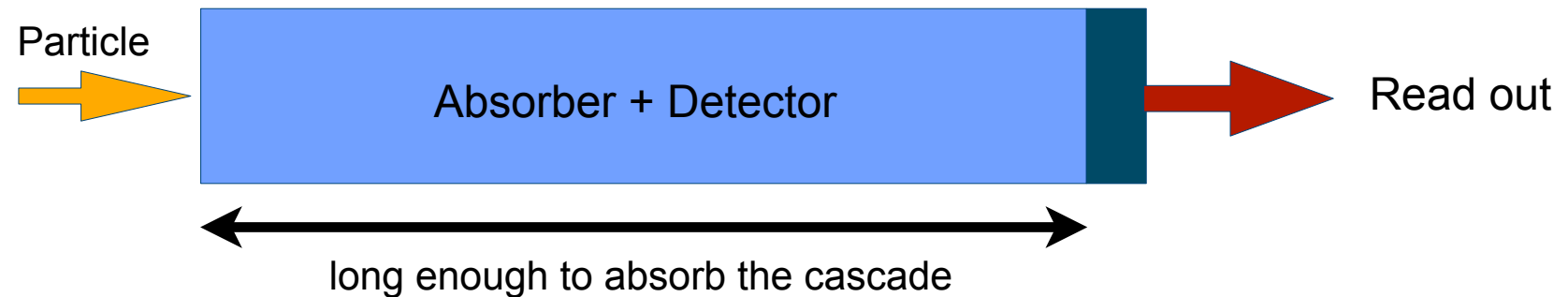
invisible energy  
 $\rightarrow$  large energy fluctuations  
 $\rightarrow$  limited energy resolution

# CALORIMETER TYPES

- Two different types of calorimeters are commonly used: Homogeneous and Sampling Calorimeter

## ● Homogeneous Calorimeter

- The absorber material is active; the overall deposited energy is converted into a detector signal
- Pro: very good energy resolution
- Contra: segmentation difficult, selection of material is limited, difficult to built compact calorimeters

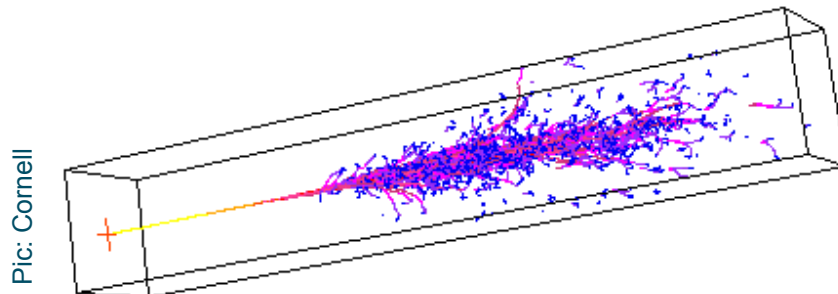
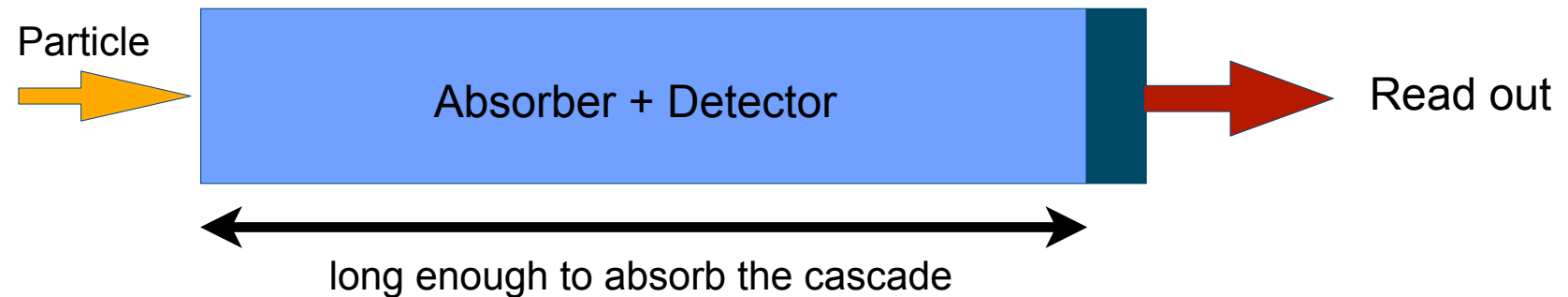


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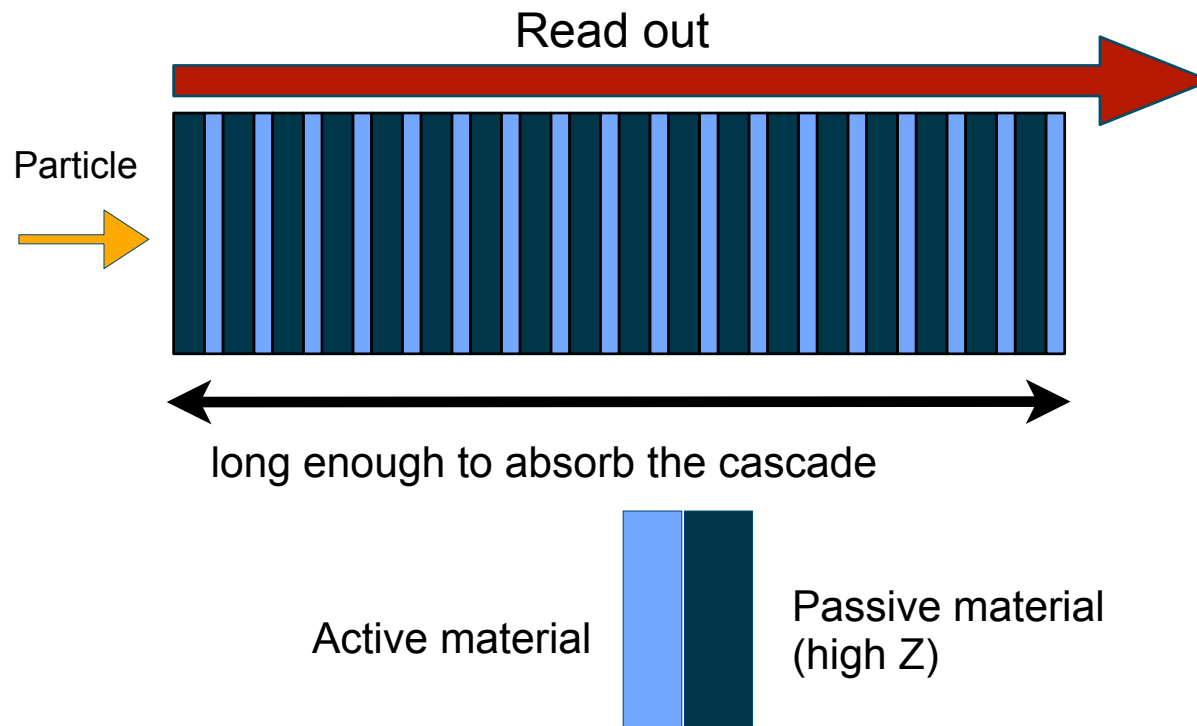
Example: Crystal calorimeter



# SAMPLING CALORIMETER

## Sampling Calorimeter

- A layer structure of passive material and an active detector material; only a fraction of the deposited energy is “registered”
- **Pro:** Segmentation (transversal and lateral), compact detectors by the usage of dense materials (tungsten, uranium,...)
- **Contra:** Energy resolution is limited by fluctuations



Important parameter:  
**Sampling Fraction**

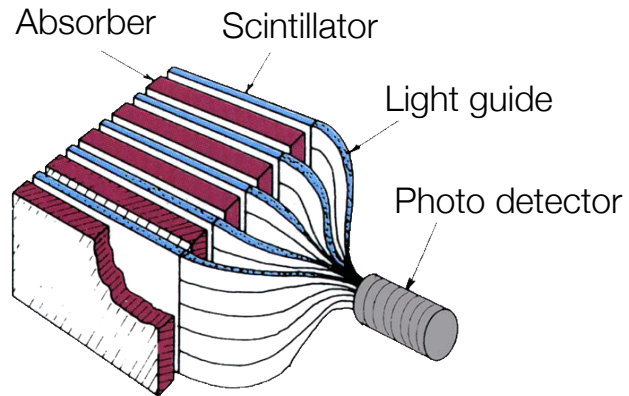
The fraction of the energy of a passing particle seen by the active material.

Typically in the percent range

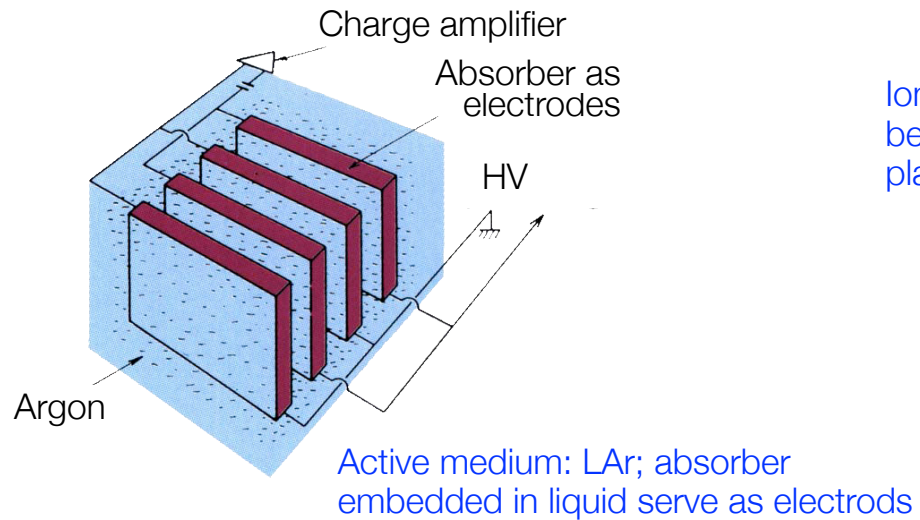
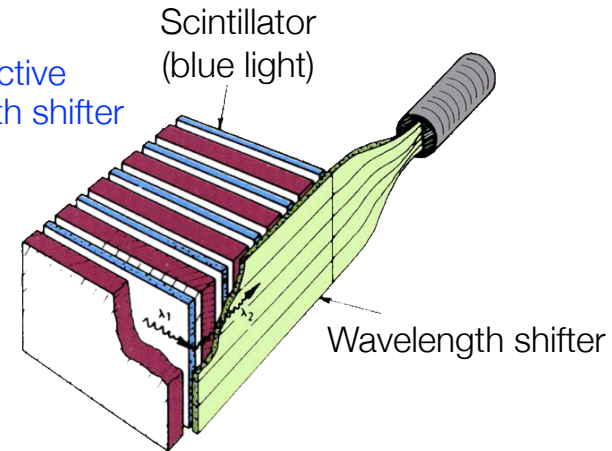
Example: ZEUS Uranium Calorimeter

# SAMPLING CALOS: POSSIBLE SETUPS

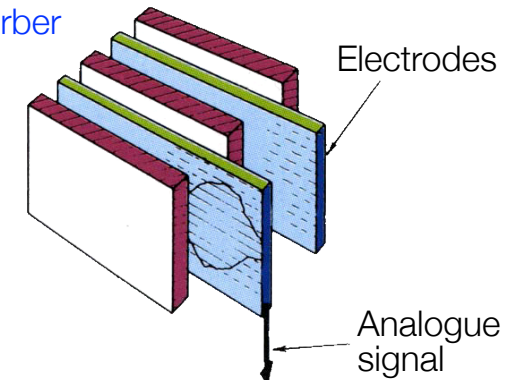
Scintillators as active layer;  
signal readout via photo multipliers



Scintillators as active layer; wave length shifter to convert light



Ionization chambers between absorber plates



# CALORIMETER: IMPORTANT PARAMETER (1)

- The relative **energy resolution** of a calorimeter is parametrised:

$$\left(\frac{\Delta E}{E}\right)^2 = \left(\frac{c_s}{\sqrt{E}}\right)^2 + \left(\frac{c_n}{E}\right)^2 + (c_c)^2$$

- **Stochastic term  $c_s$**

- the resolution depends on intrinsic shower fluctuations, photoelectron statistics, dead material in front of calo, and sampling fluctuations

- **Noise term  $c_n$**

- Electronic noise, radioactivity, i.e. dependent of the energy

- **Constant term  $c_c$**

- Energy independent term contributing to the resolution: due to inhomogeneities with in the detector sensitivity, calibration uncertainties and radiation damage

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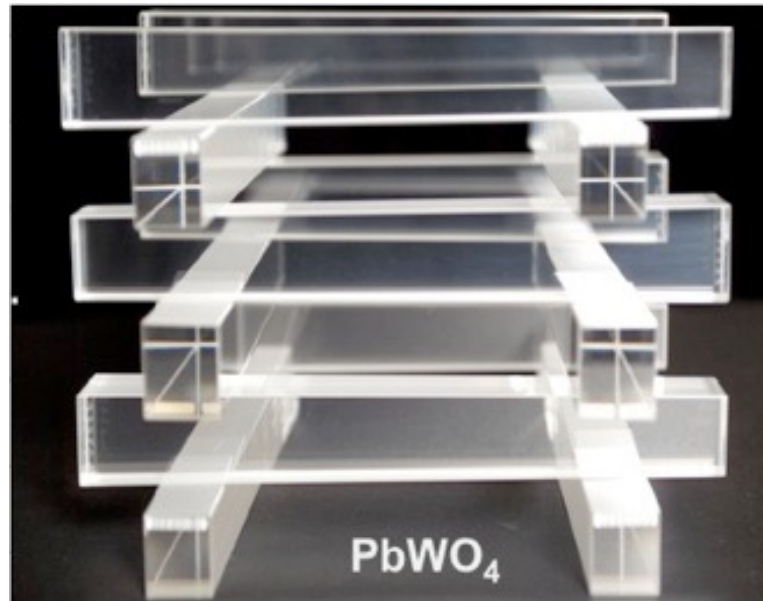
## Losses of Resolution:

- **Shower not contained** in detector → fluctuation of leakage energy; longitudinal losses are worse than transverse leakage.
- **Statistical fluctuations** in number of photoelectrons observed in detector.
- **Sampling fluctuations** if the counter is layered with inactive absorber.
- ....

# CALOS: ACTIVE MATERIAL

Active  
material

- Detectors based on registration of excited atoms
- Emission of photons by excited atoms, typically UV to visible light.
  - Observed in noble gases (even liquid !)
  - Polyzyclic Hydrocarbons (Naphtalen, Anthrazen, organic scintillators) -> Most important category.
  - Inorganic Crystals -> Substances with largest light yield. Used for precision measurement of energetic Photons.



- PbWO<sub>4</sub>: Fast, dense scintillator,
  - Density  $\sim 8.3 \text{ g/cm}^3$  (!)
  - $\rho_M$  2.2 cm,  $X_0$  0.89 cm
  - low light yield:  $\sim 100 \text{ photons / MeV}$

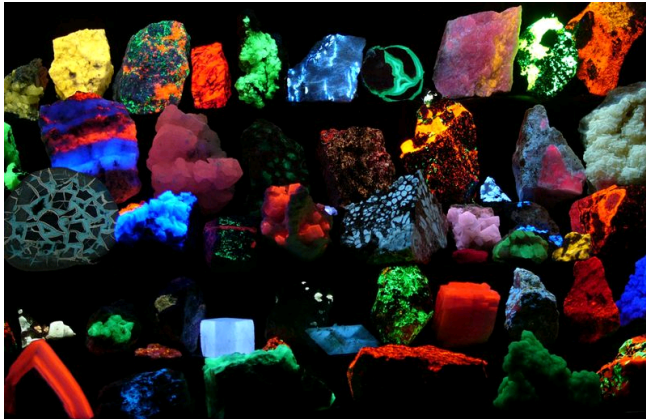


Picture: CDF@Fermilab



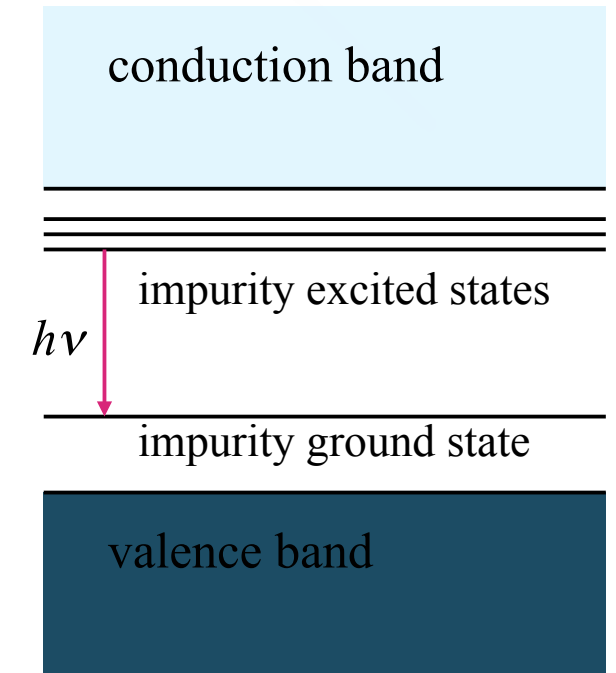
# SCINTILLATORS TO MEASURE THE ENERGY

- An incident photon or particle ionises the medium (on band structure level).
- Ionised electrons slow down causing excitation.
- Excited states immediately emit light.



## Inorganic scintillators

- Fluorescence is known in many natural crystals.
  - UV light absorbed
  - Visible light emitted
- Artificial scintillators can be made from many crystals.
  - Doping impurities added
  - Improve visible light emission



### Advantages:

- Good efficiency
- Good linearity
- Radiation tolerance

### Disadvantage:

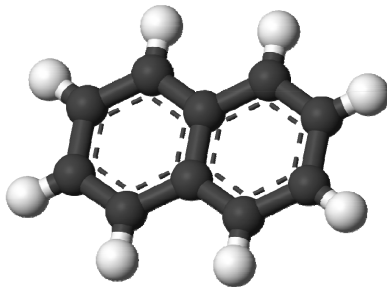
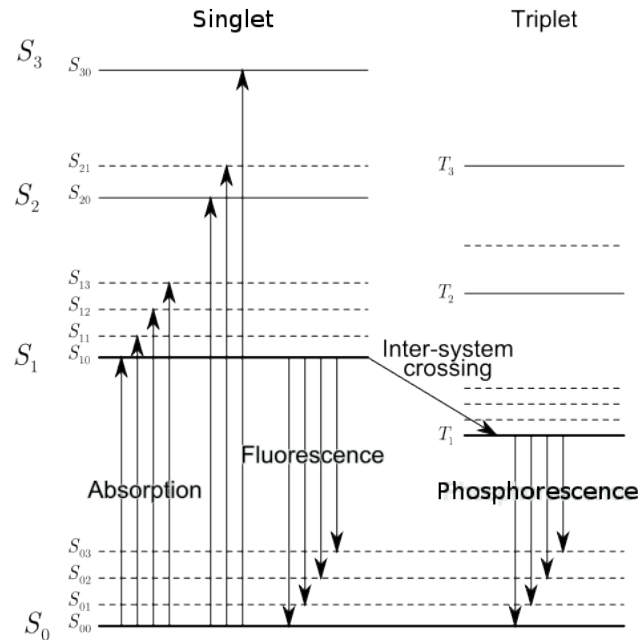
- Relatively slow
- Crystal structure needed (small and expensive)

# SCINTILLATORS TO MEASURE THE ENERGY

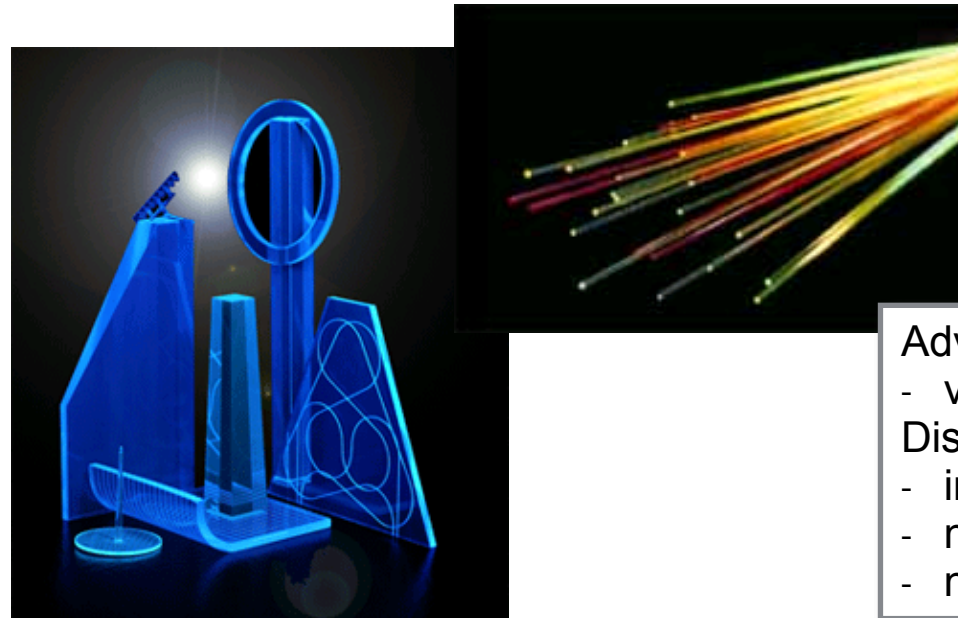
Active  
material

- Very common: Measurement of the deposited energy using scintillation

## Organic scintillators



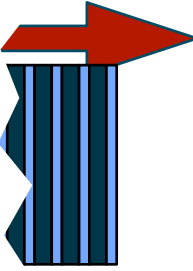
- Organic scintillators are aromatic hydrocarbon compounds (containing benzene ring compounds)
- The scintillation mechanism is due to the transition of electrons between molecular orbitals
  - organic scintillators are fast  $\sim$  few ns.
- Excited states radiate photons in the visible and UV spectra.
  - Fluorescence is the fast component
  - Phosphorescence is the slow component



Advantages:  
- very fast  
Disadvantage:  
- inefficient  
- non-linear  
- not good for photons

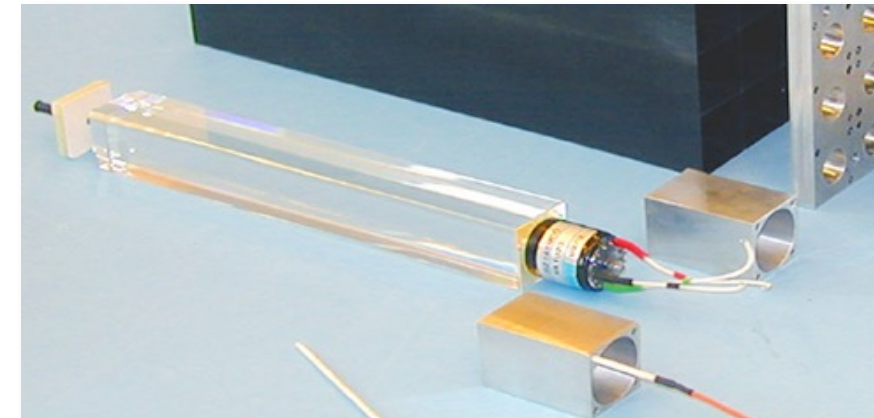
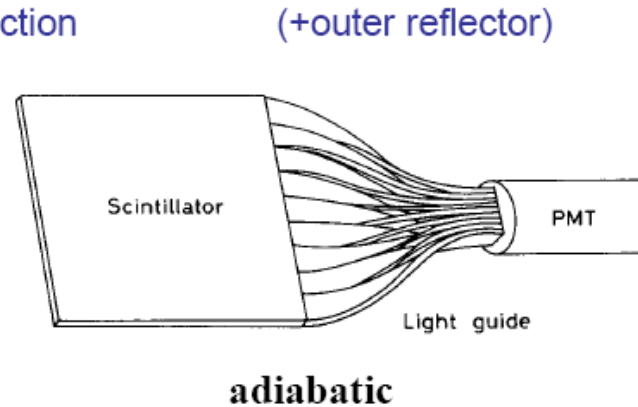
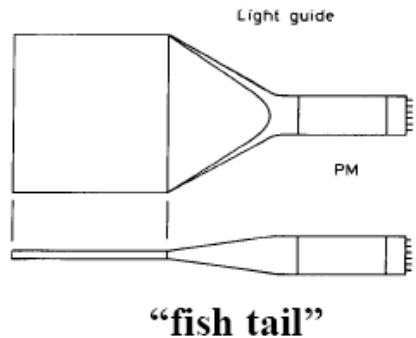
source: Wikipedia

# LIGHT TRANSPORT



- The photons are being reflected towards the end of the scintillator
- A light guide brings the light to a Photomultiplier

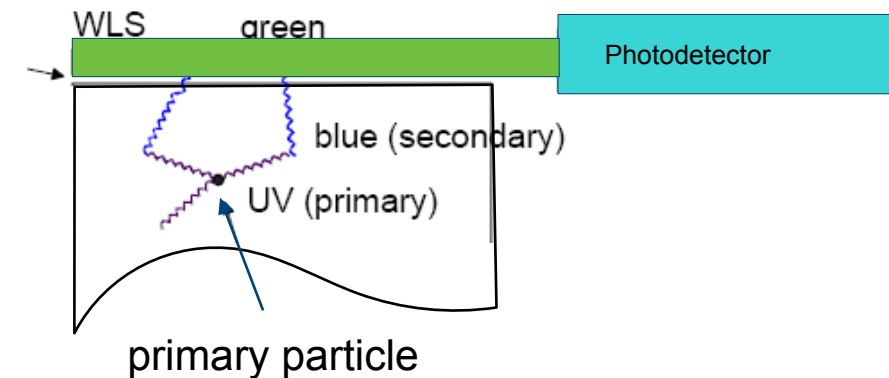
- Light guides: transfer by total internal reflection (+outer reflector)



- UV light enters the light guide material
- Light is transformed into longer wavelength (wavelength shifter)
- -> Total internal reflection inside the WLS material
- -> 'transport' of the light to the photo detector

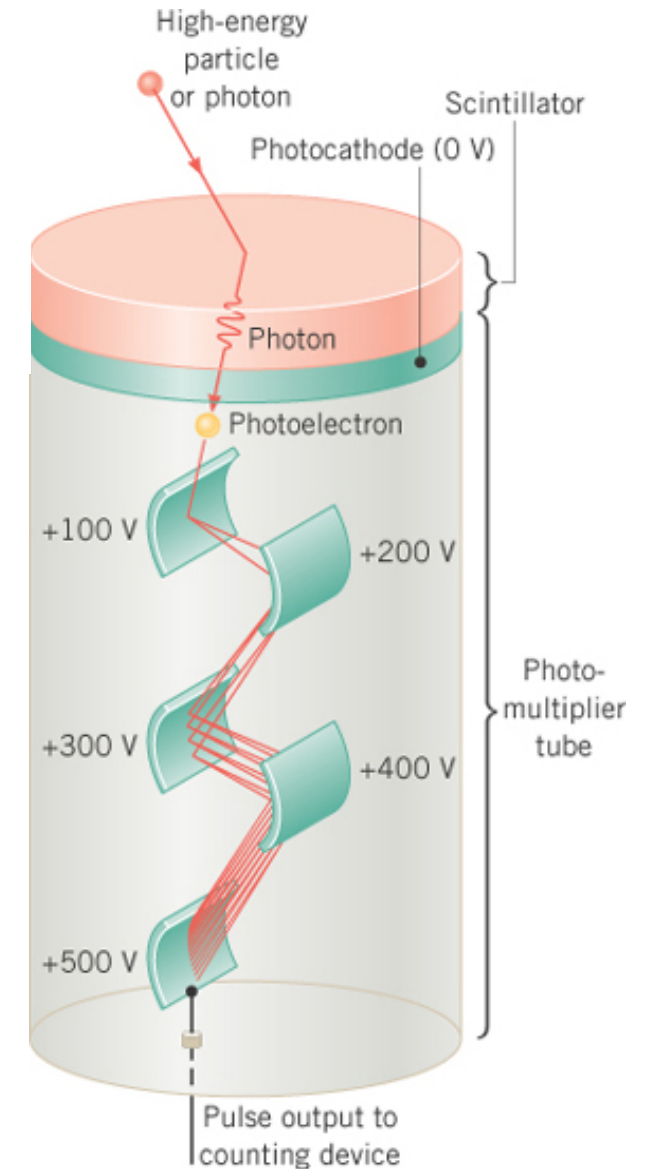
small air gap

scintillator



# DETECTING THE LIGHT

- The classic method to detect photons are photomultipliers
  - Conversion of a photon into electrons via photo-electric effect when the photon impinges on the photo cathode
  - The following dynode system is used to amplify the electron signal
  - Usable for a large range of wave lengths (UV to IR)
  - good efficiencies, single photon detection possible
  - large active area possible (SuperKamiokande O 46cm)

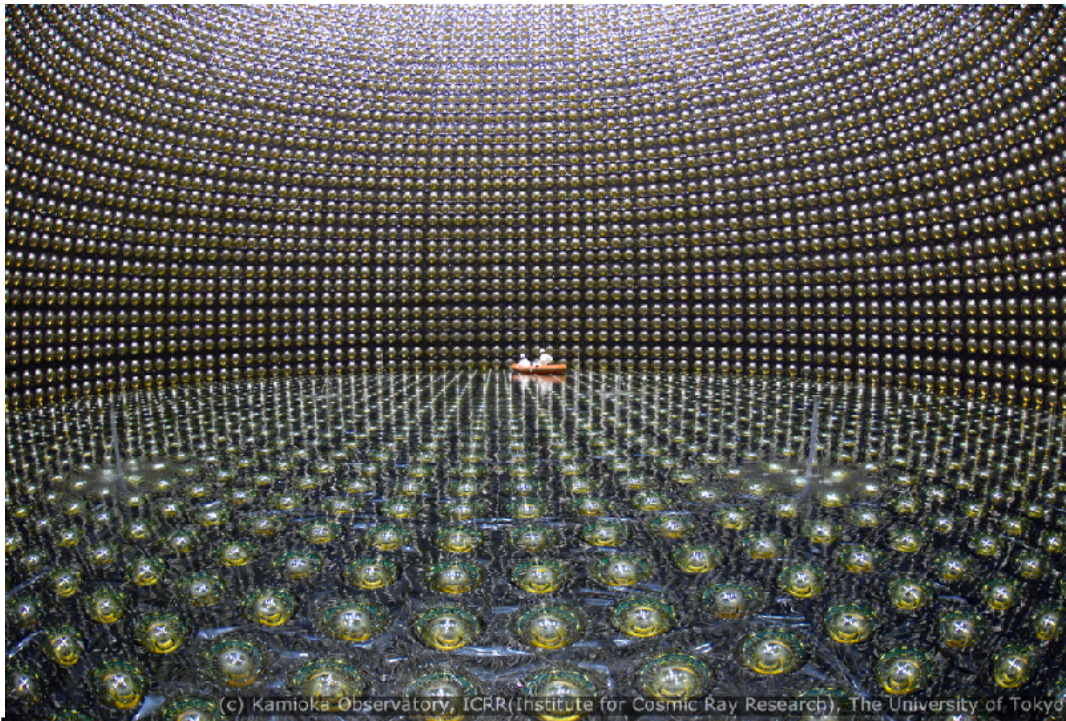
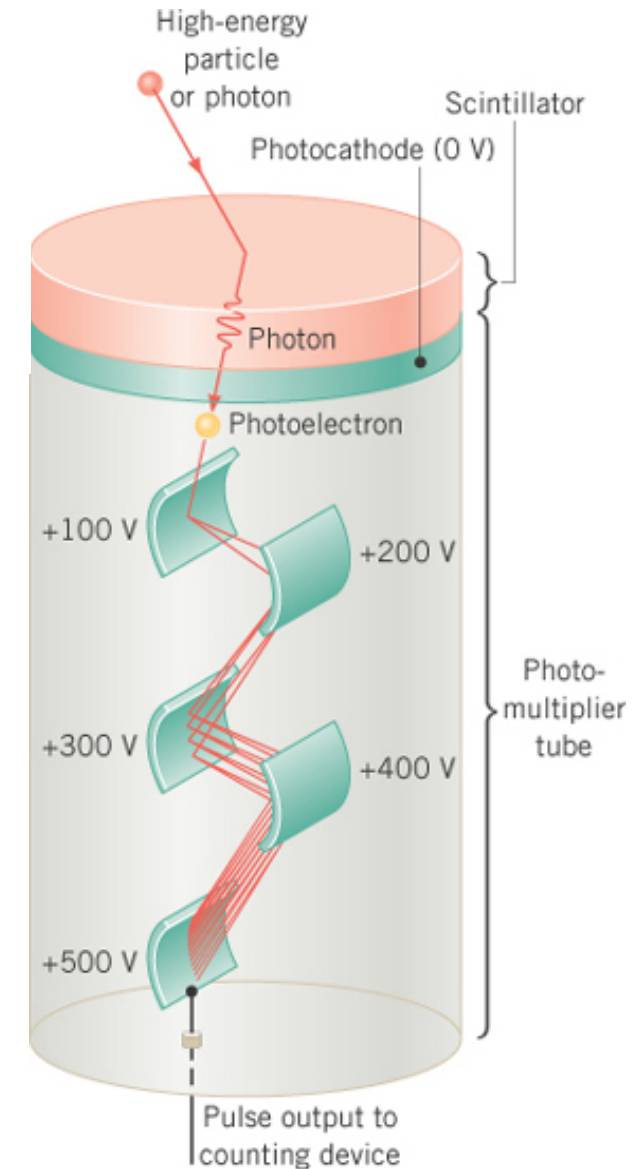


Pic: ICRR/University of Tokyo



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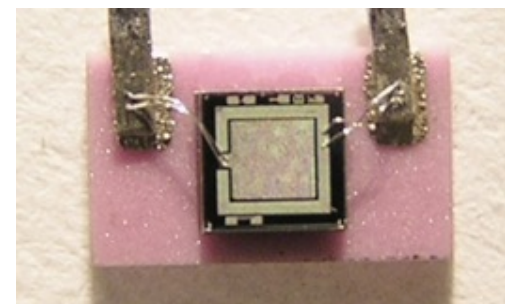
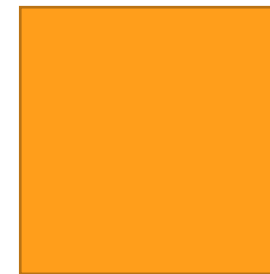
Pic: ICRR/University of Tokyo

(c) Kamioka Observatory, ICRR (Institute for Cosmic Ray Research), The University of Tokyo



# MODERN PHOTON DETECTION

- Silicon photomultipliers



Silicon  
photo-multiplier

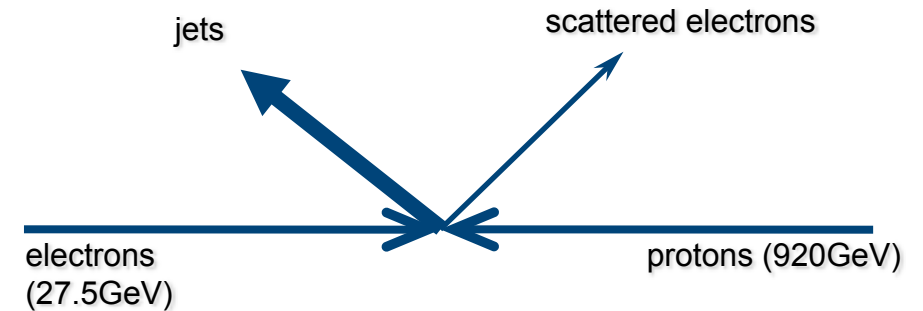
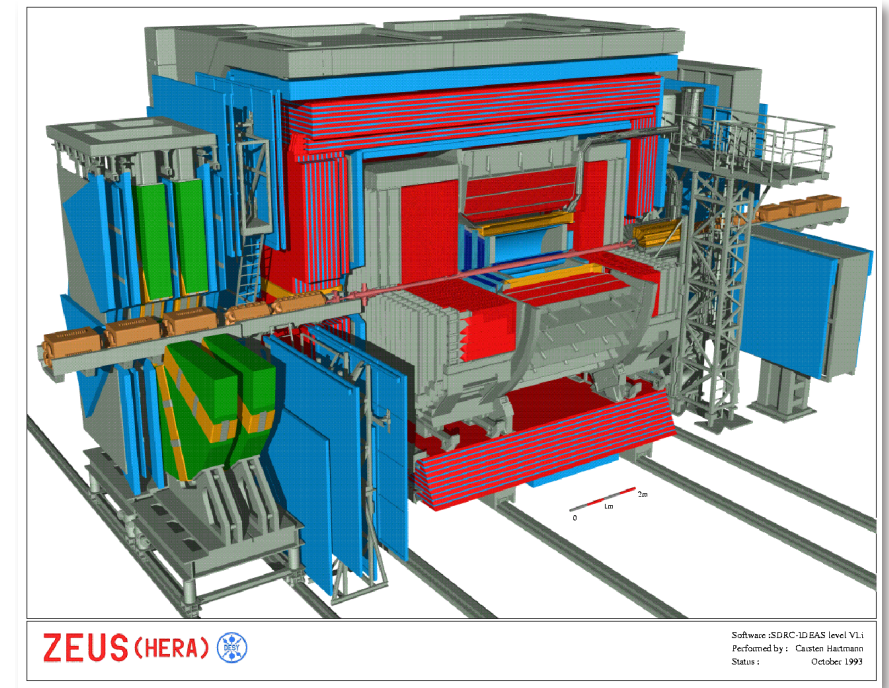
# EXAMPLE: ZEUS CALO

A rather hostile environment in ZEUS at HERA

- bunch crossing every 96ns
- high beam gas rate
- very energetic particles produced

## Requirements for the ZEUS calorimeter:

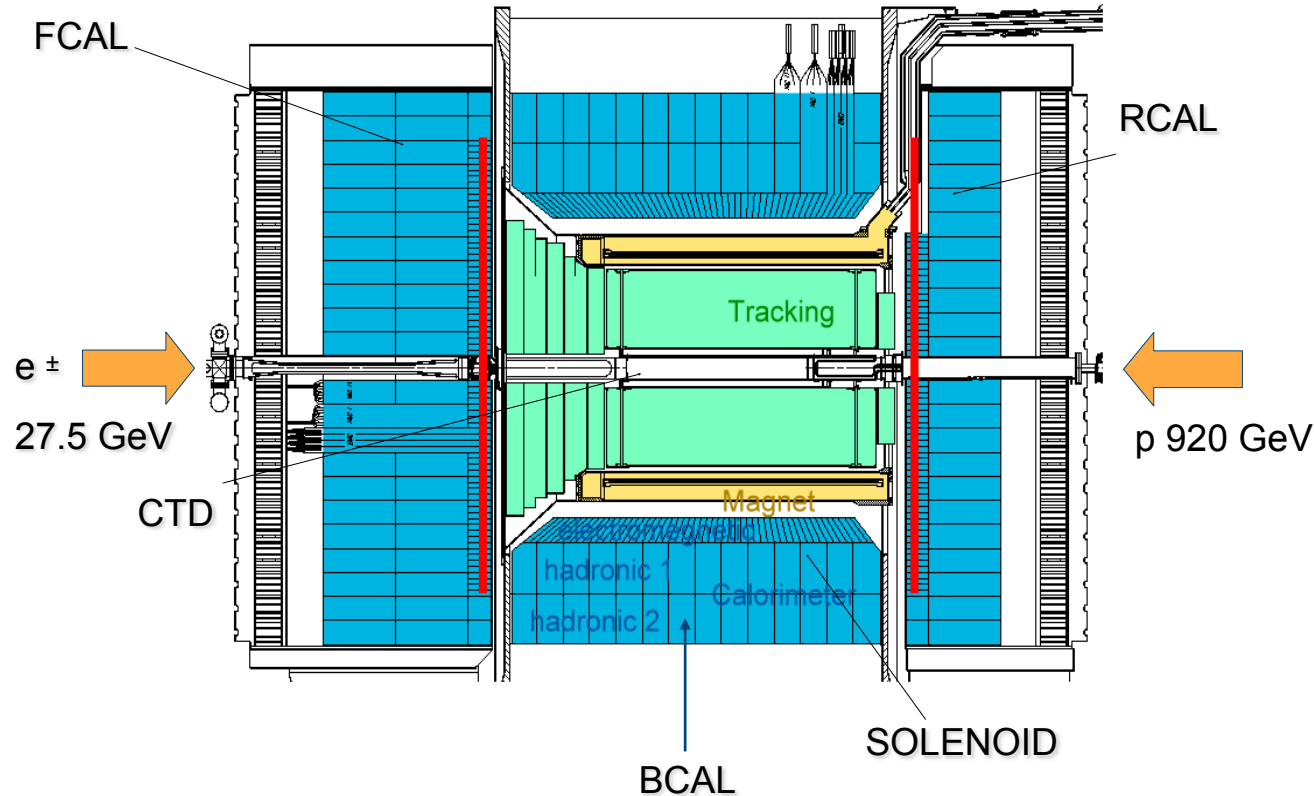
- hermeticity
- dead time free readout
- time resolution in nanosecond range
- uniform response
- radiation tolerance (15 years of running)
- electron-hadron separation
- good position resolution
- good electron and jet energy resolution



Keep in mind: this was developed in the middle of the 80s!

# THE ZEUS CALORIMETER - SOLUTION

- highly-segmented, uranium scintillator sandwich calorimeter read out with photomultiplier tubes (PMTs)



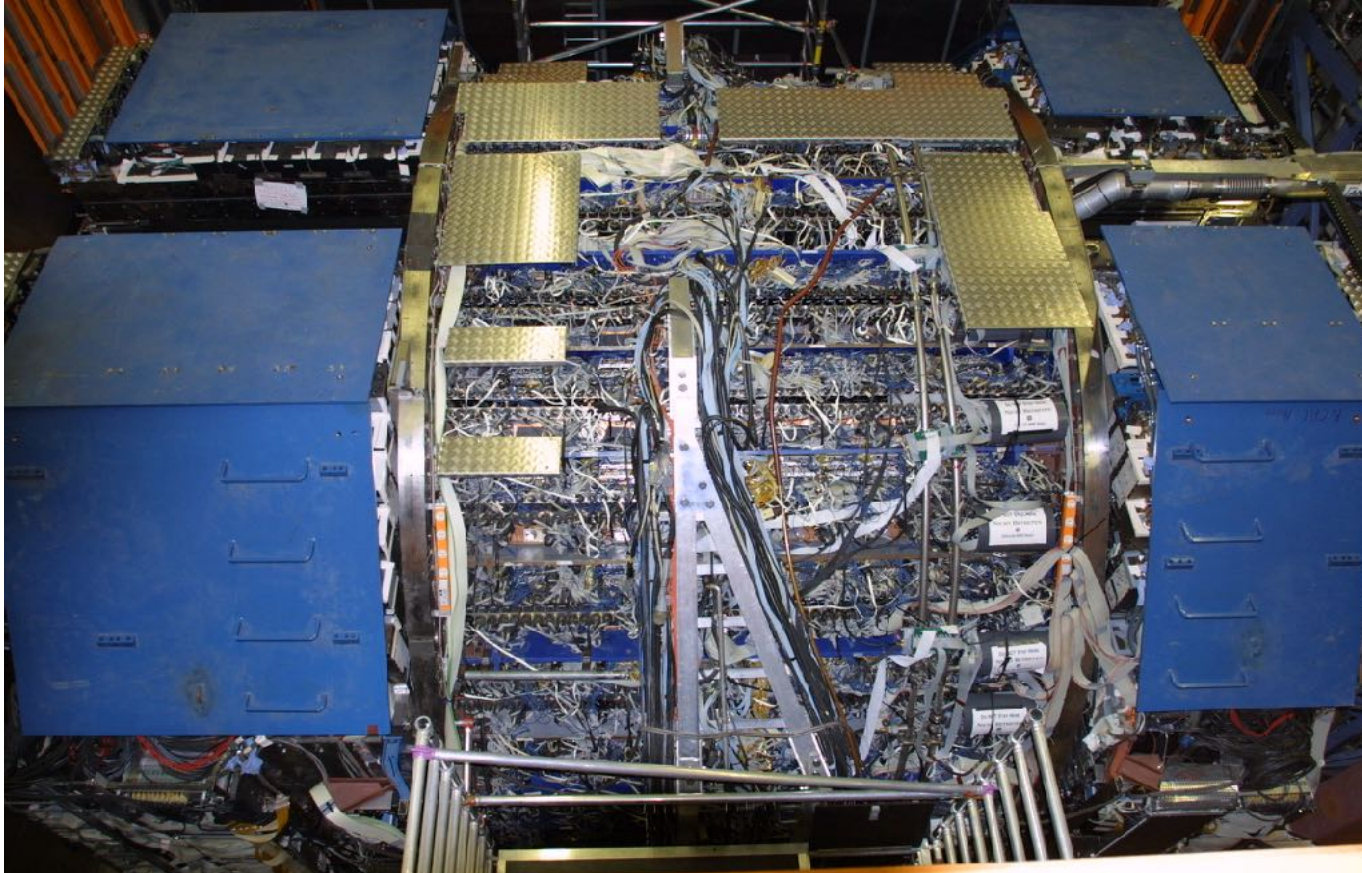
## Uranium + Scintillator:

- compensation
- high Z material  $\rightarrow$  more compact size of calorimeter
- natural radioactivity provides means of calibration

- Very hermetic: covering up to  $\eta < 4.2$  in the forward direction and  $\eta < -3.8$  in the rear direction.
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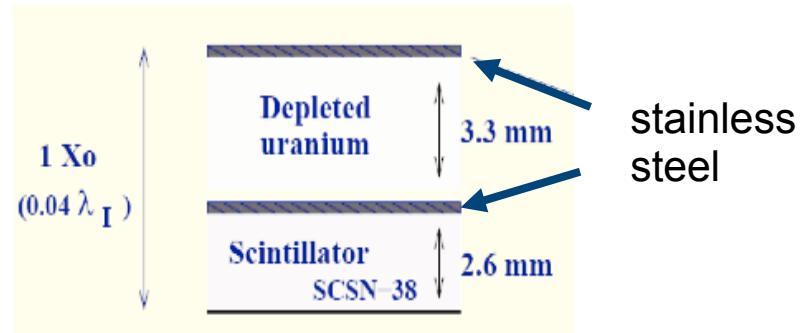
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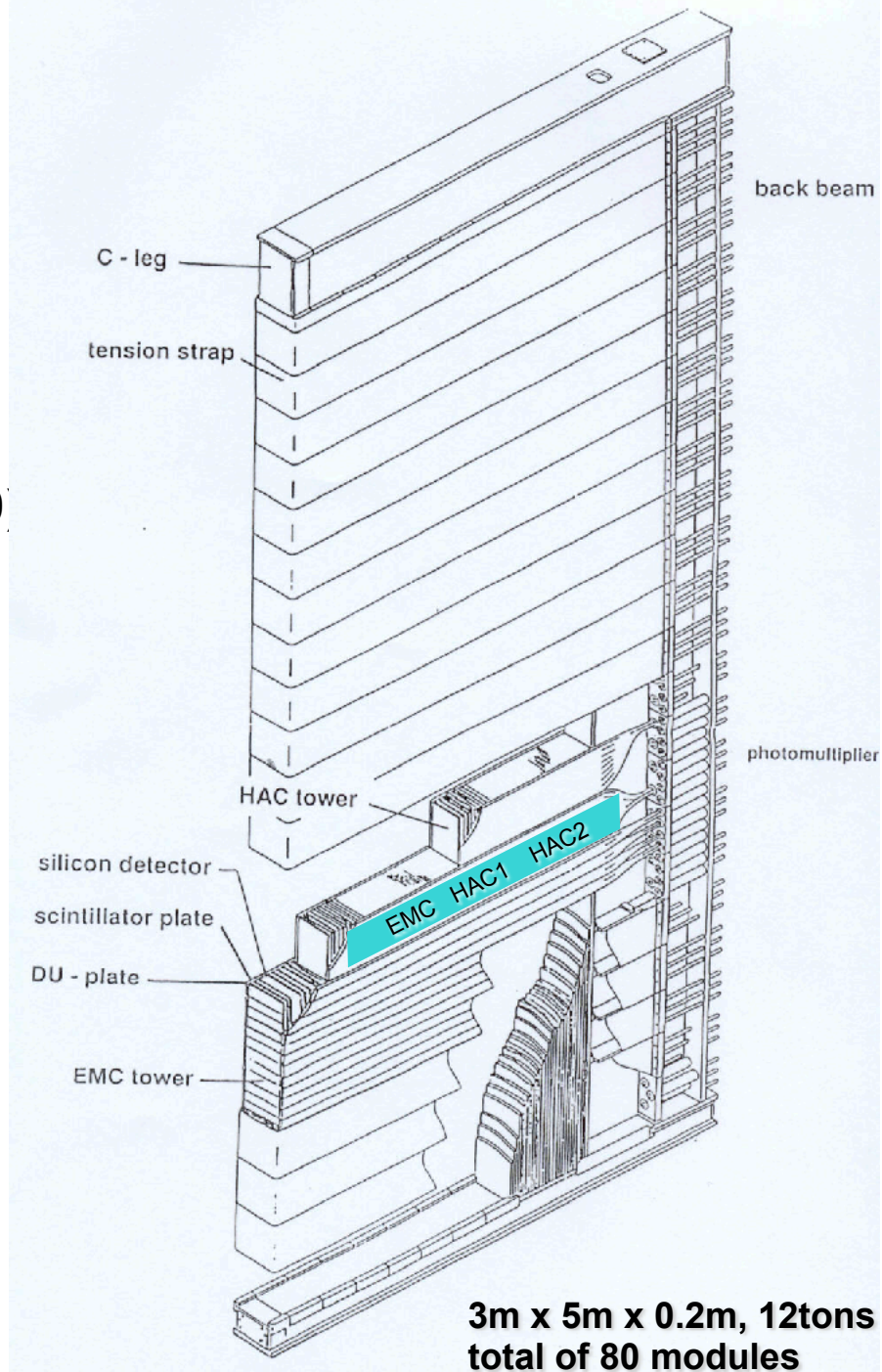
# DESIGN

- Layers:



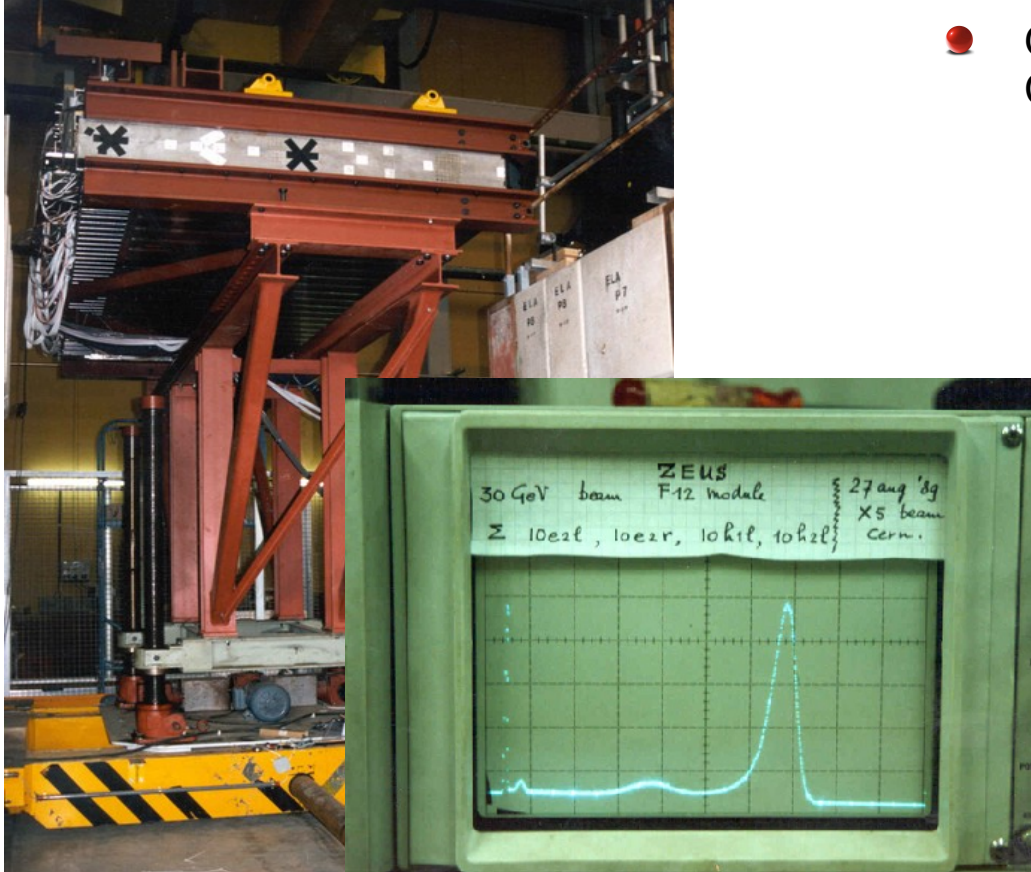
- Choice of active and passive thicknesses -> compensation ( $e/h = 1.0$ )
- Uniformity in structure + natural radioactivity -> good calibration
- F/B/RCAL with ~6000 cells
  - EM cell size: 5x20 (10x20) cm<sup>2</sup> in F/BCAL (RCAL)
  - HA cell size: 20x20 cm<sup>2</sup>

- Cell read out on both sides with wavelength shifters
  - redundancy
  - transverse position measurement within the cell

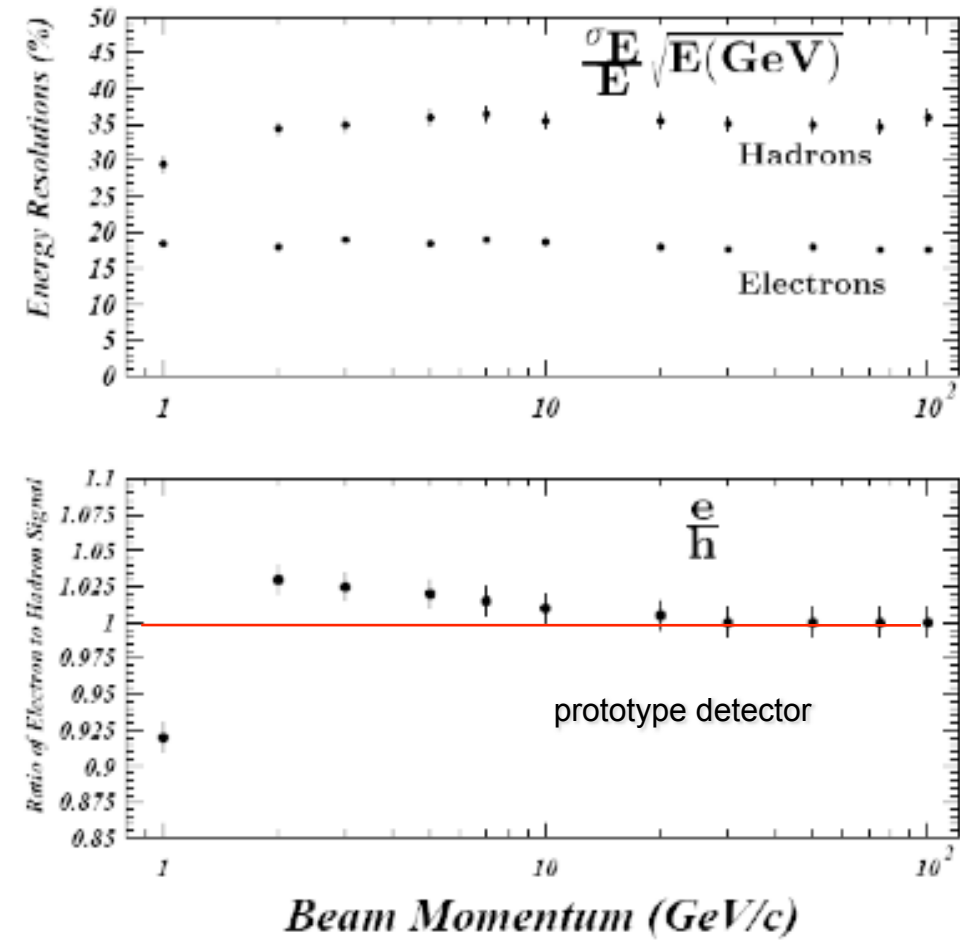




# TEST BEAM AT CERN



- Operation characteristics were determined in test beams at CERN (prototype detector)



Electrons:  $\frac{\sigma(E)}{E} = \frac{18\%}{\sqrt{E(\text{GeV})}}$

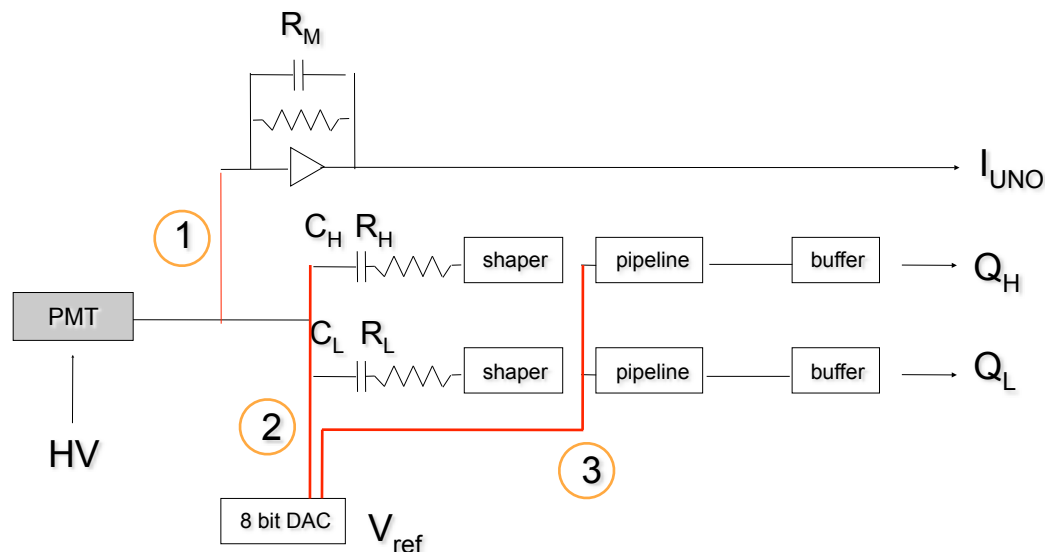
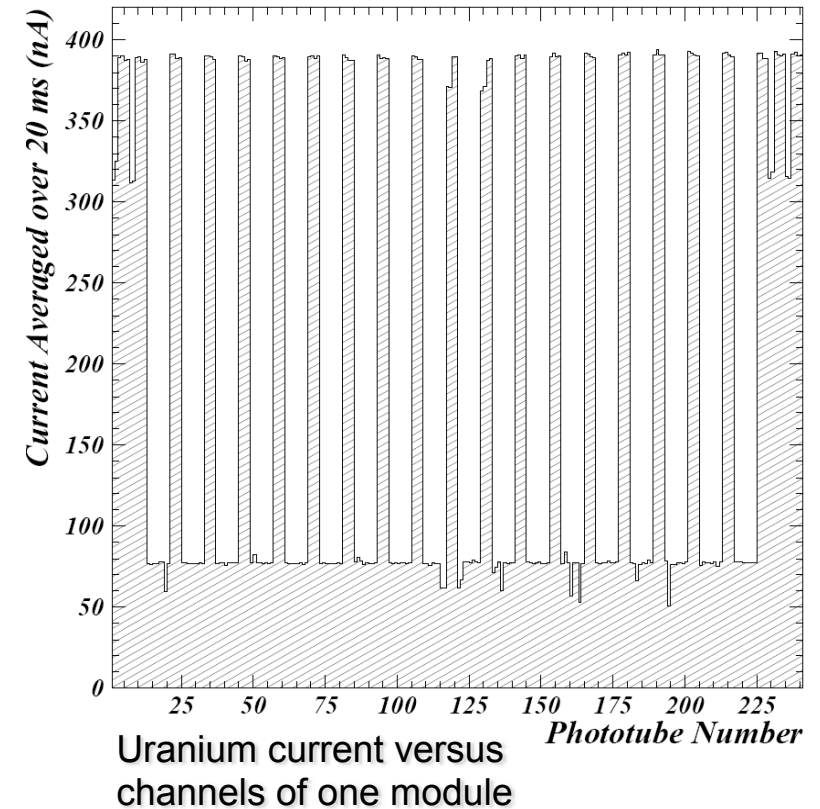
Hadrons:  $\frac{\sigma(E)}{E} = \frac{35\%}{\sqrt{E(\text{GeV})}}$

Production modules were all calibrated at CERN

# CALIBRATION METHODS

Stable radioactivity  
- good for calibration

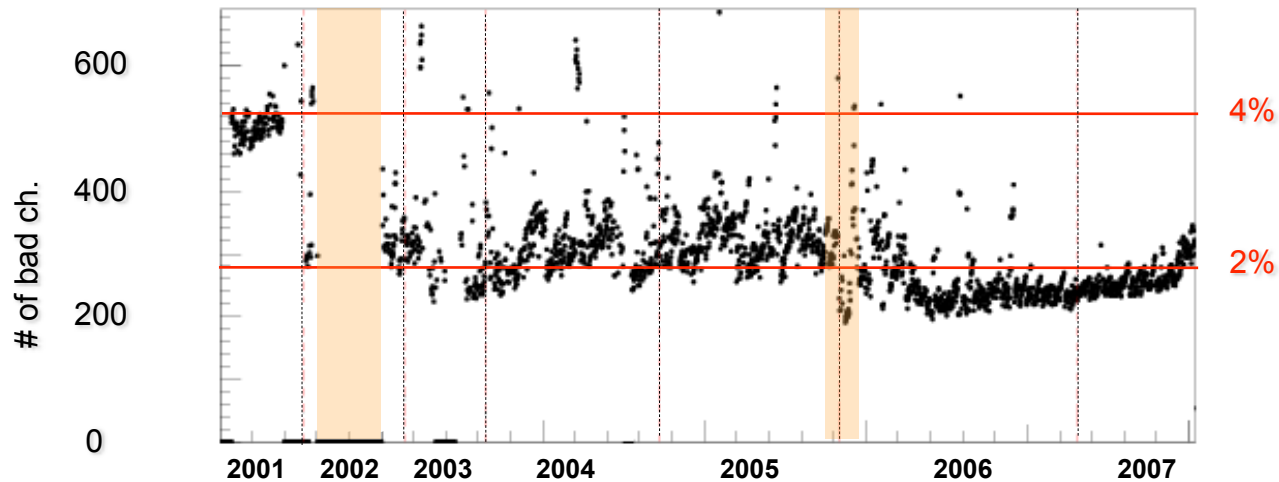
- Natural uranium activity provides absolute energy calibration in situ!
  - 98.1% U238 + 1.7% Nb + 0.2% U235
  - Half-Life of U238 is  $4.5 \cdot 10^9$  years
- Detectable uranium induced signal current
- Uranium noise signal
  - ~ 2MHz (EM Calo)
  - ~ 10MHz (Hadronic Calo)
  - with Uranium noise calibration can be tracked very easy



- 1 Uranium noise
- 2 Charge injection
- 3 Pedestals and Gains

Channels out of range  
-> declared as "bad" until  
readjusted

# HARDWARE PERFORMANCE



- Number of bad channels versus run number (over years)
- “Bad channels” are excluded from data taking -> reducing the calo performance in that area
- Read out from both sides -> bad channel is not complete loss of information
- Ups and downs visible in bad channel behaviour over the years

- At the time of the shutdown (30.06.2007):
  - only ~ 2% bad channels (one side) and only 2 holes (both sides failed) -> 0.3 per mille
- **In general very stable and robust system**

- Front End Cards:
  - About 1000 necessary for the running, ~10% spares
  - Main failure mode: buffer or pipeline chip (socketed)
  - Cards easy to debug and maintain
  - Failure rate: <1/month (12 channels – one side)
  - Very successful





# OVERVIEW OF CALORIMETERS

- In order to maximise the sensitivity for  $H \rightarrow \gamma\gamma$  decays, the experiments need to have an excellent  $e/\gamma$  identification and resolution

CMS

## CMS DETECTOR

Total weight : 14,000 tonnes  
Overall diameter : 15.0 m  
Overall length : 28.7 m  
Magnetic field : 3.8 T

STEEL RETURN YOKE  
12,500 tonnes

SILICON TRACKERS  
Pixel (100x150  $\mu\text{m}$ ) ~16m<sup>2</sup> ~66M channels  
Microstrips (80x180  $\mu\text{m}$ ) ~200m<sup>2</sup> ~9.6M channels

SUPERCONDUCTING SOLENOID  
Niobium titanium coil carrying ~18,000A

MUON CHAMBERS  
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers  
Endcaps: 468 Cathode Strip, 432 Resistive Plate Chambers

PRESHOWER  
Silicon strips ~16m<sup>2</sup> ~137,000 channels

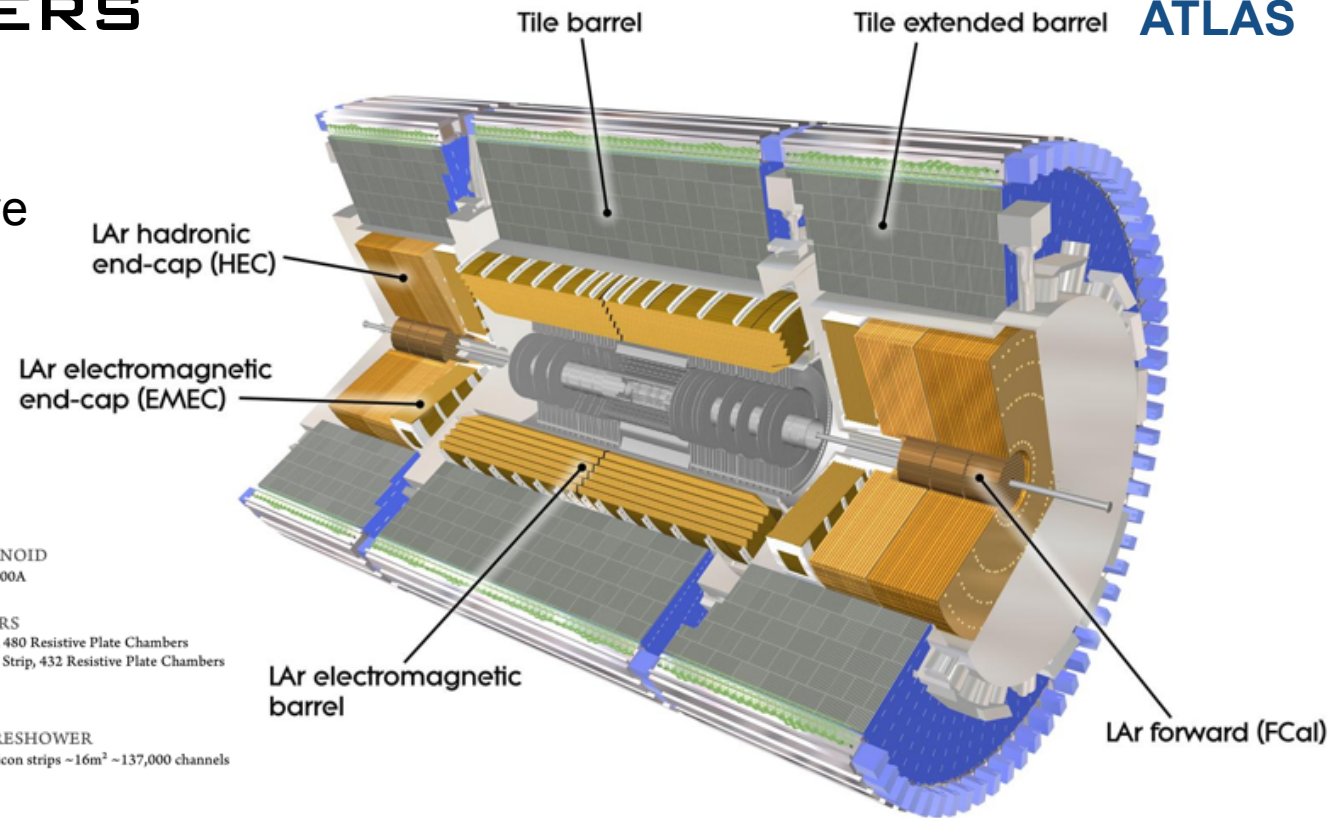
FORWARD CALORIMETER  
Steel + Quartz fibres ~2,000 Channels

CRYSTAL ELECTROMAGNETIC CALORIMETER (ECAL)  
~76,000 scintillating PbWO<sub>4</sub> crystals

HADRON CALORIMETER (HCAL)  
Brass + Plastic scintillator ~7,000 channels

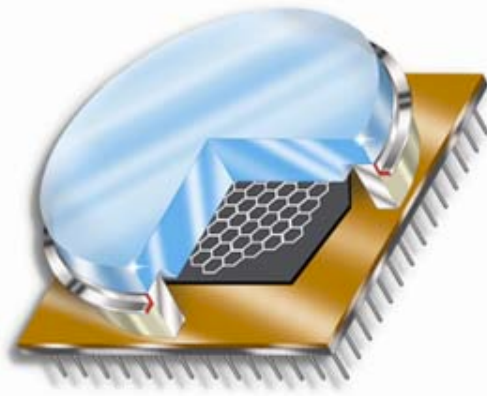


ATLAS

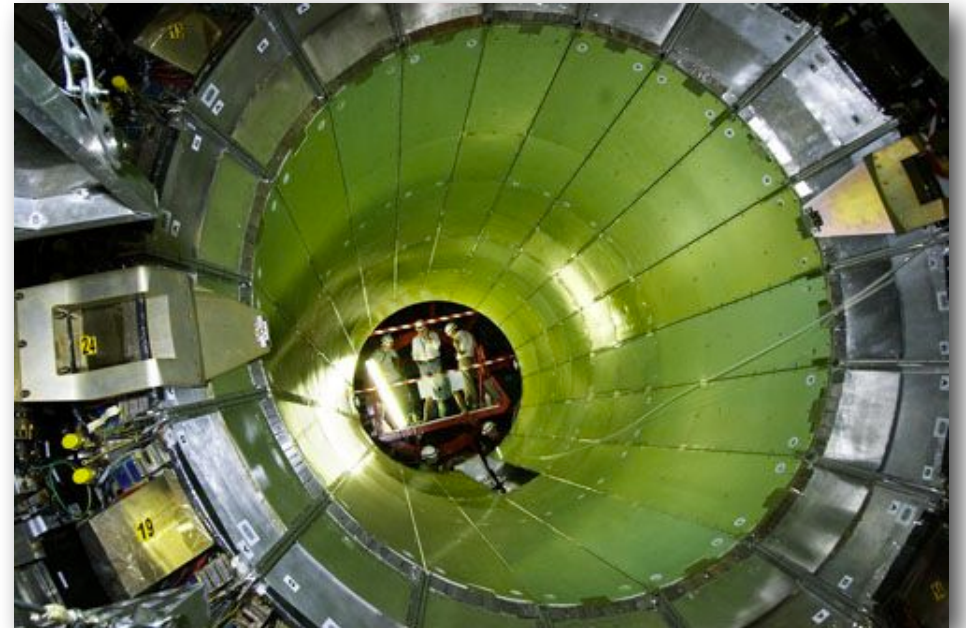


# CMS CALORIMETER

- **ECAL:** homogeneous calo
  - high resolution Lead Tungsten crystal calorimeter -> **higher intrinsic resolution**
  - 80000 crystals each read out by a photodetector
  - constraints of magnet -> HCAL absorption length not sufficient
  - tail catcher added outside of yoke
- **HCAL:** sampling calo
  - 36 barrel “wedges”, each weighing 26 tonnes
  - brass or steel absorber
  - plastic scintillators
  - read out by hybrid photodetectors



CMS Lead tungsten crystals, each 1.5kg (CERN)

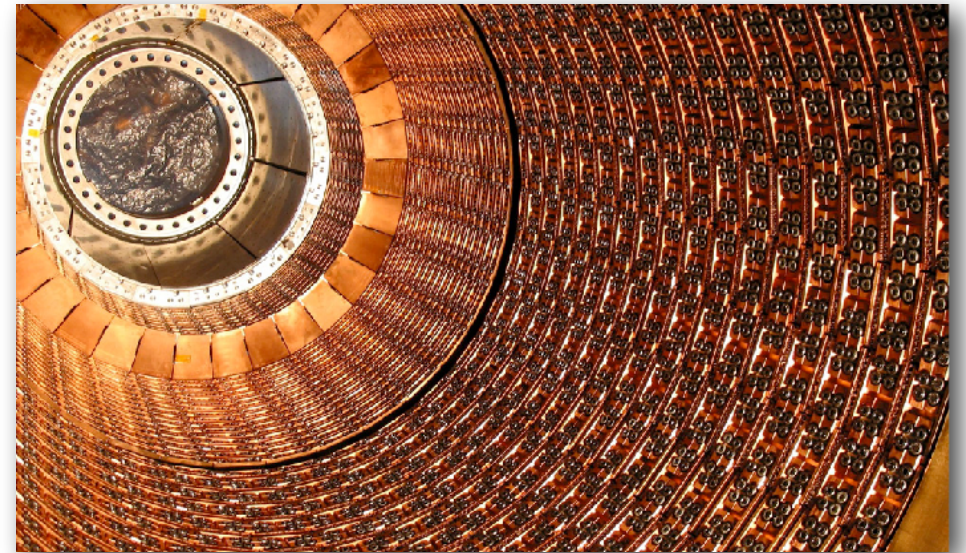
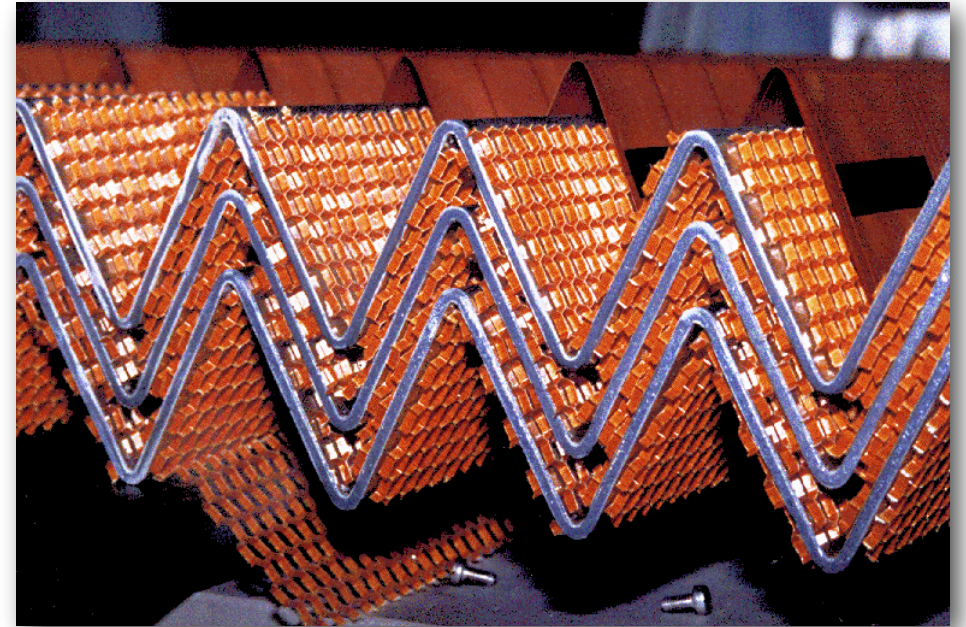


CMS ECAL during installation (CERN)



# ATLAS CALORIMETER

- **ECAL + HCAL:** sampling calo
  - Liquid argon LAr calorimeter > high granularity and longitudinally segmentation (better  $e/\gamma$  ID)
  - Electrical signals, high stability in calibration & radiation resistant (gas can be replaced)
  - Solenoid in front of ECAL -> a lot of material reducing energy resolution
  - Accordion structure chosen to ensure azimuthal uniformity (no cracks)
  - Liquid argon chosen for radiation hardness and speed
  - Tile calorimeter: covering outer region
  - “Conventional” steel absorber with plastic scintillators.



ATLAS Hadronic endcap Liquid Argon Calorimeter. (CERN)

# CALORIMETERS AT LHC

- All LHC experiments have a calorimetric system with at least an electromagnetic and a hadronic part

## Overview **EM** calorimeters at LHC

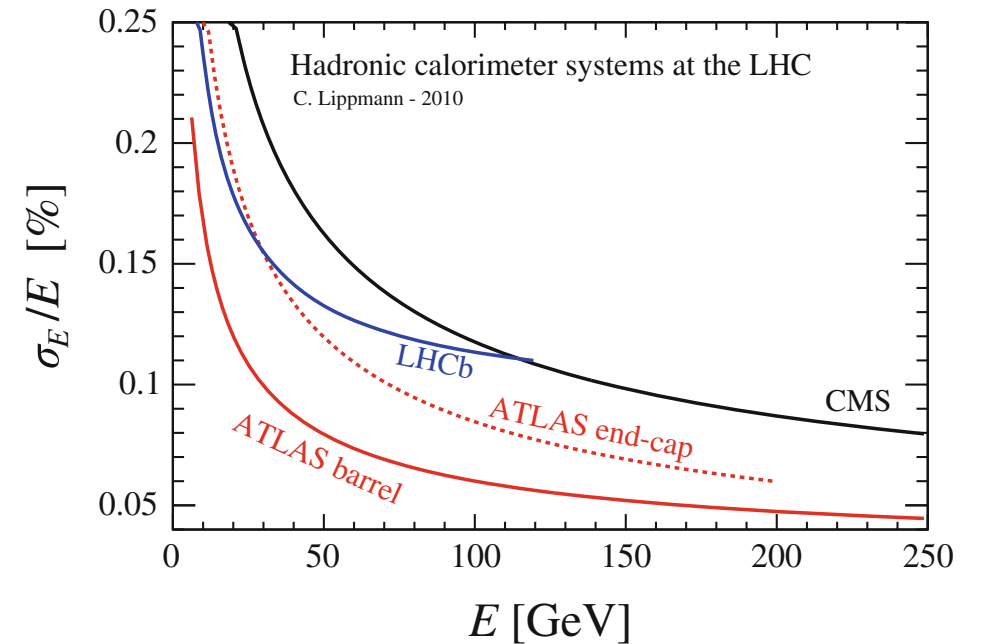
	Calorimeter	Material	Number of channels	Angular coverage	Energy resolution	
					$c_s$ (%)	$c_c$ (%)
ATLAS	EM barrel	$LAr + Pb$	109,568	$ \eta  < 1.475$	10	0.7
	EM end-cap	$LAr + Pb$	63,744	$1.375 <  \eta  < 3.2$	10	0.7
	FCal	$LAr + Cu$	2016	$3.1 <  \eta  < 4.9$	28.5	3.5
CMS	ECAL barrel	$PbWO_4$	61,200	$ \eta  < 1.479$	2.8	0.3
	ECAL end-cap	$PbWO_4$ homogeneous	14,648	$1.479 <  \eta  < 3.0$	2.8	0.3
LHCb	ECAL	Scint. + $Pb$	6016	$0.756 < \eta_x < 2.19$ $1.037 < \eta_y < 2.19$	9	0.8
ALICE	PHOS	$PbWO_4$	17,920	$ \eta  < 0.12$ , $220^\circ < \phi < 320^\circ$	3.3	1.1
	EMCal	Scint. + $Pb$	12,672	$ \eta  < 0.7$ , $80^\circ < \phi < 187^\circ$	10	2

- As expected, the sampling based on lead as absorber have a slightly worse resolution than the homogeneous crystal calorimeters.

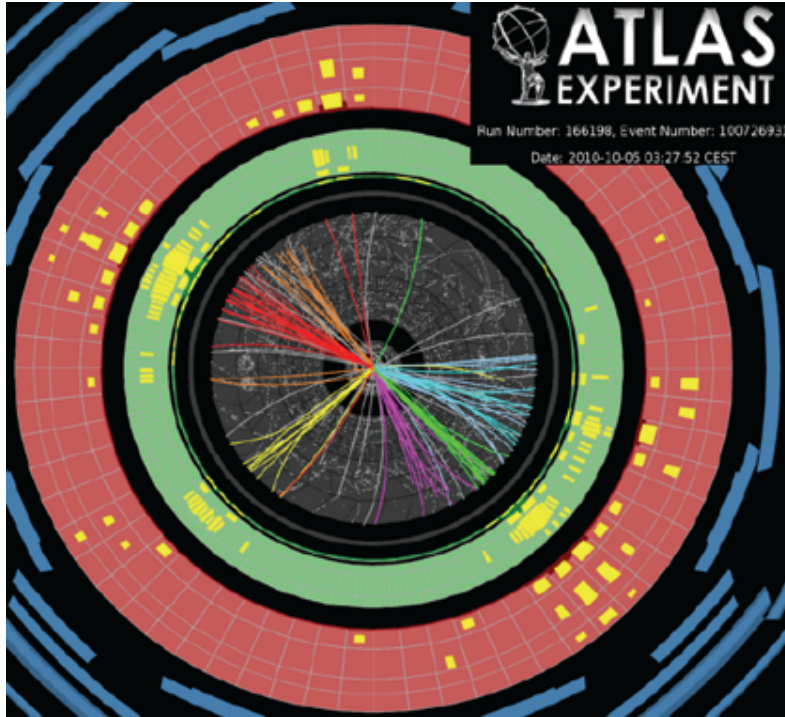
# HADRONIC CALOS AT LHC

	Calorimeter	Material	Number of channels	Angular coverage	Energy resolution	
					$c_s$ (%)	$c_c$ (%)
ATLAS	Tile	Scint. + $Pb$	9852	$ \eta  < 1.7$	52	3
	HEC	$LAr + Cu$	5632	$1.5 <  \eta  < 3.2$	84	–
	FCal	$LAr + W$	1508	$3.1 <  \eta  < 4.9$	94	7.5
CMS	HB	Scint. + steel/brass	2592	$ \eta  < 1.3$	90	9
	HE	Scint. + steel/brass	2592	$1.3 <  \eta  < 3$	90	9
	HO	Scint. + steel	2160	$ \eta  < 1.4$	–	–
	HF	Quartz fibre + steel	1728	$3 <  \eta  < 5.2$	120	–
LHCb	HCAL	Scint. + steel	1488	$ \eta_x  < 1.87$	69	9
				$ \eta_y  < 2.07$		

● All sampling calorimeter



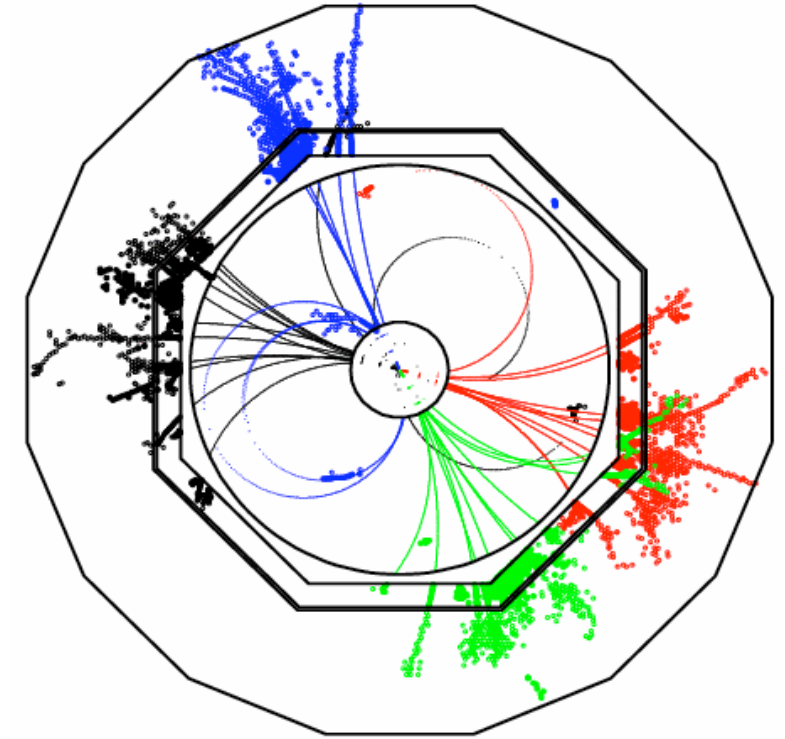
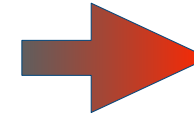
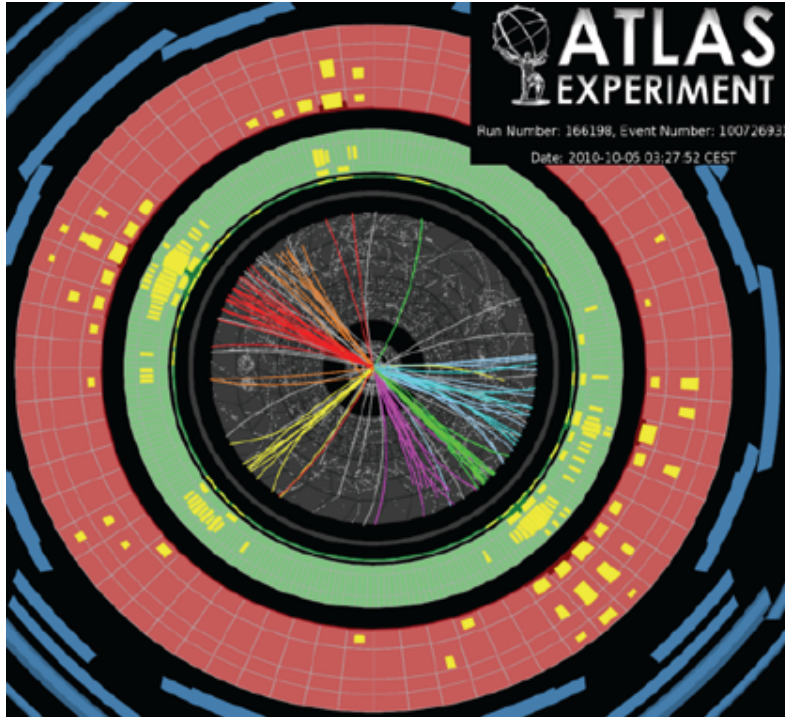
# CURRENT HADRON CALOS ... AND DREAMS



- Tower-wise readout: light from many layers of plastic scintillators is collected in one photon detector (typically PMT)  
O(10k) channels for full detectors



# CURRENT HADRON CALOS ... AND DREAMS

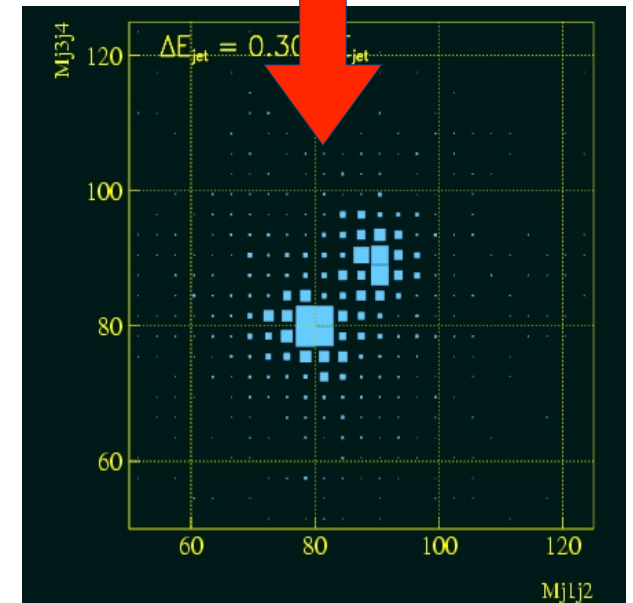
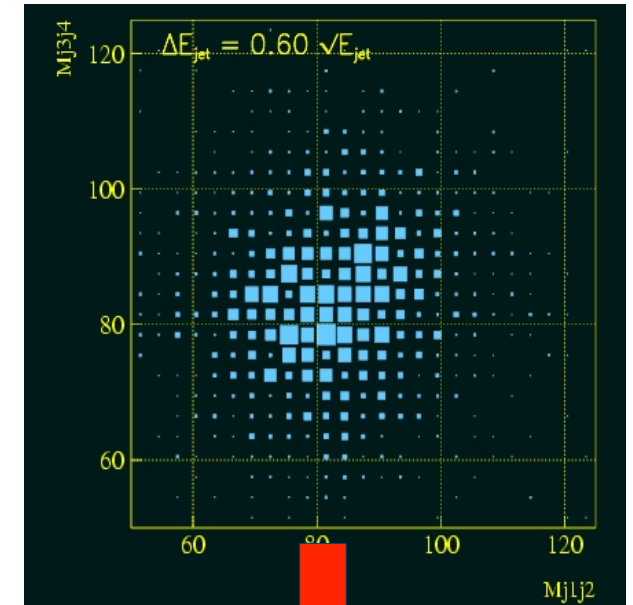


- Tower-wise readout: light from many layers of plastic scintillators is collected in one photon detector (typically PMT)  $O(10k)$  channels for full detectors

- Extreme granularity to see shower substructure: small detector cells with individual readout for Particle Flow  $O(10M)$  channels for full detectors

# PARTICLE FLOW CALORIMETER

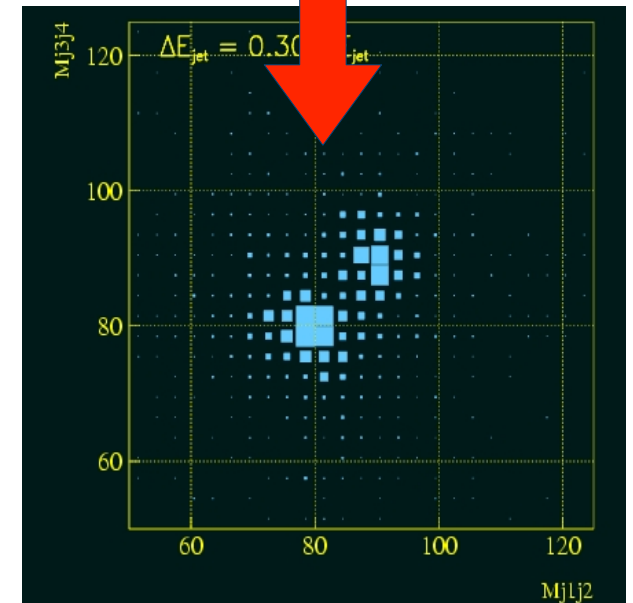
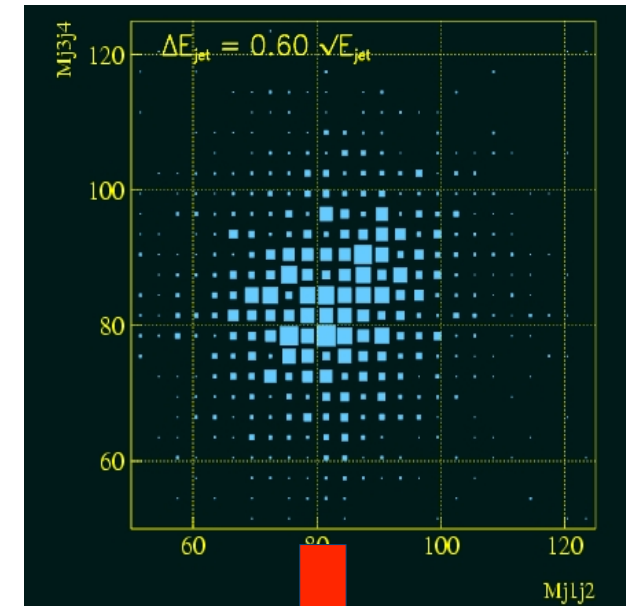
- Attempt to measure the energy/momentum of each particle with the detector subsystem providing the best resolution
- Used in three main contexts:
  - “Energy flow” -> Use tracks to correct jet energies
  - “Particle flow/Full event reconstruction” e.g. CMS
    - > Aim to reconstruct particles not just energy deposits
  - “High granularity particle flow” e.g. ILC
    - > Technique applied to detector concept optimised for particle flow



# PARTICLE FLOW CALORIMETER

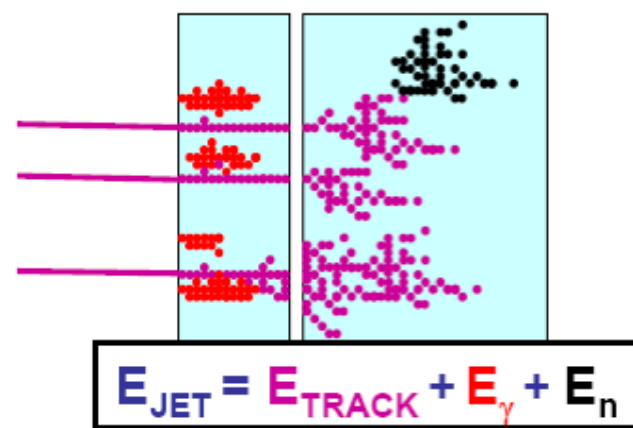
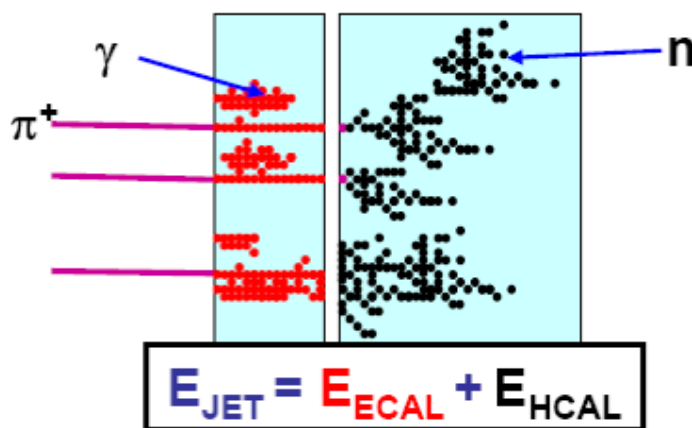
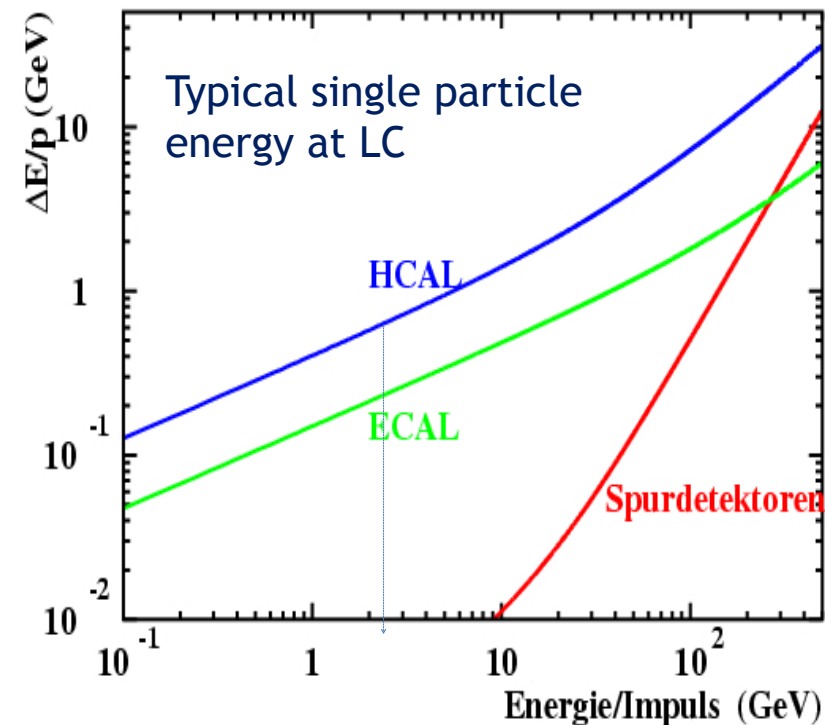
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- Need
  - a calorimeter optimised for photons: separation into ECAL + HCAL
  - to place the calorimeters inside the coil (to preserve resolution)
  - to minimise the lateral size of showers with dense structures
  - the highest possible segmentation of the readout
  - to minimise thickness of the active layer and the depth of the HCAL



# PARTICLE FLOW PARADIGM

- Reconstruct **every** particle in the event
- Up to ~100 GeV **Tracker** is superior to calorimeter
  - use tracker to reconstruct  $e^\pm, \mu^\pm, h^\pm$  ( $<65\%$ ) of  $E_{\text{jet}}$
  - use **ECAL** for  $\gamma$  reconstruction ( $<25\%$ )
  - use (**ECAL+**) **HCAL** for  $h^0$  reconstruction ( $<10\%$ )
- HCAL E resolution still dominates  $E_{\text{jet}}$  resolution
- But much improved resolution (only 10% of  $E_{\text{jet}}$  in HCAL)



PFLOW calorimetry = **Highly granular detectors**  
+ Sophisticated reconstruction software

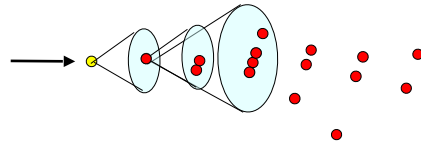


# STATE OF THE ART OF PARTICLE FLOW ALGORITHM

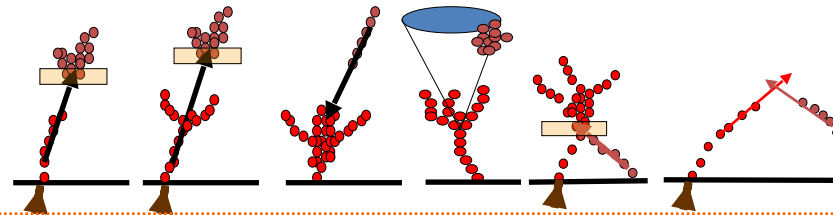
- High granularity Particle Flow reconstruction is highly non-trivial

many complex steps  
(not all shown)

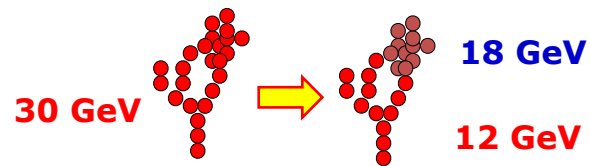
Clustering



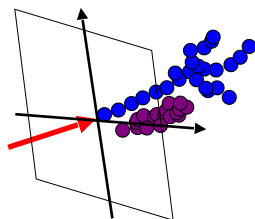
Topological Association



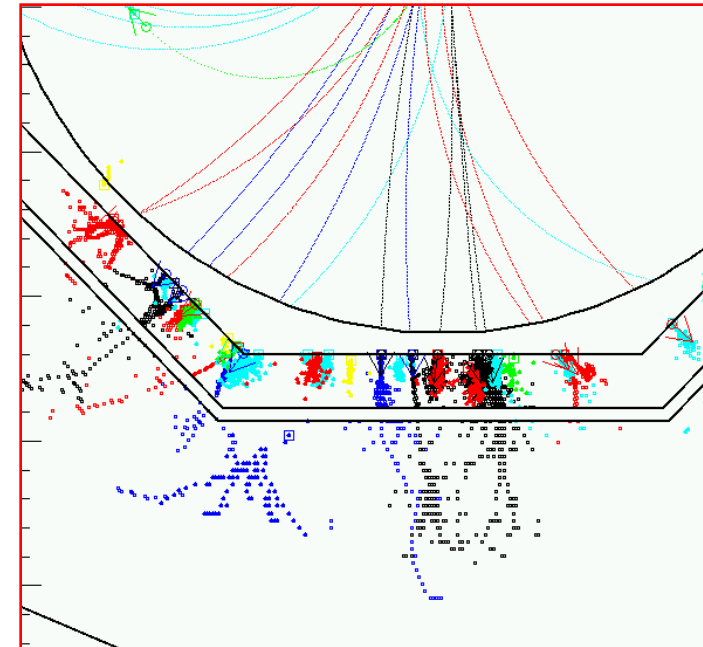
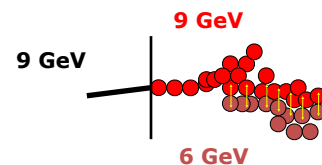
Iterative Reclustering



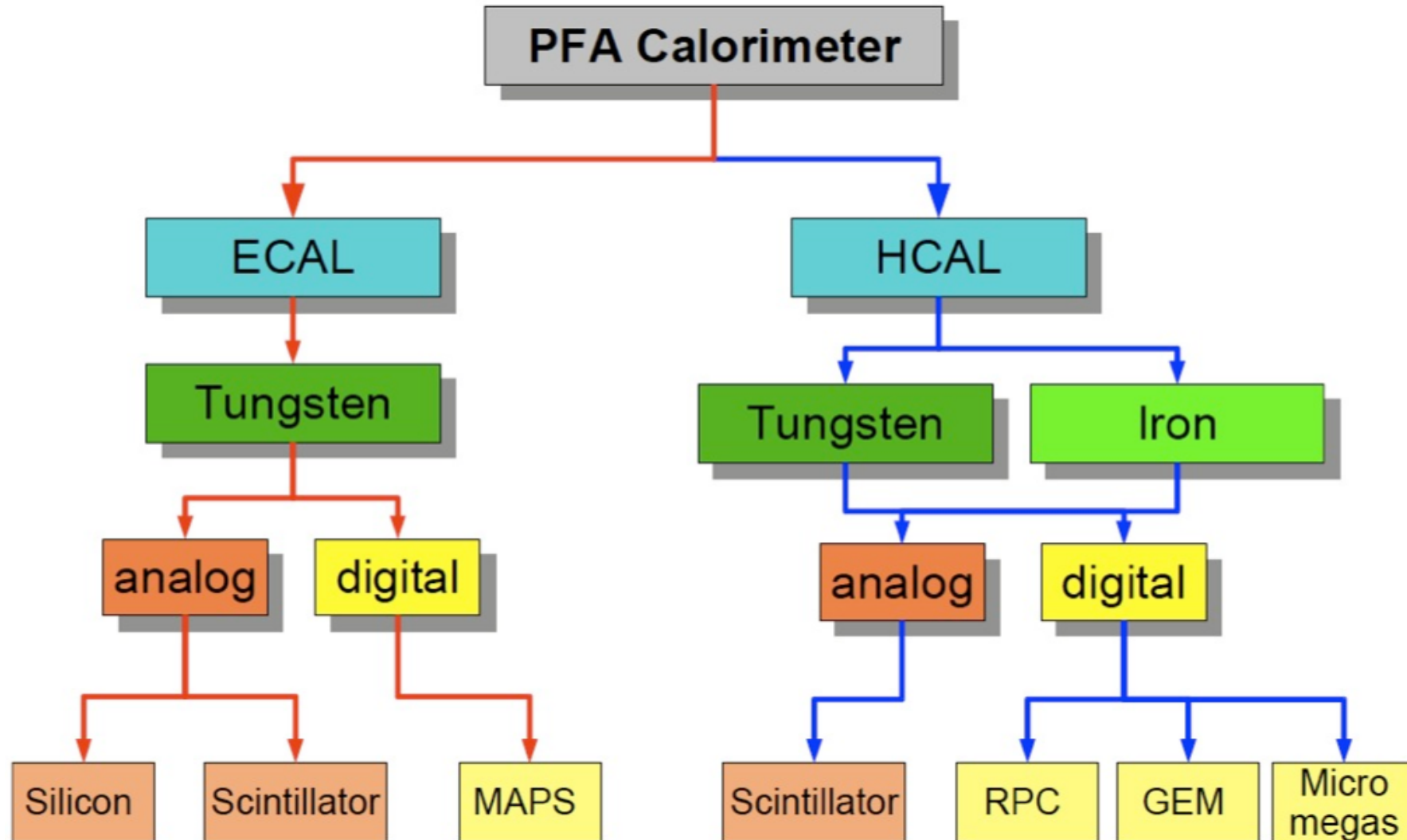
Photon ID



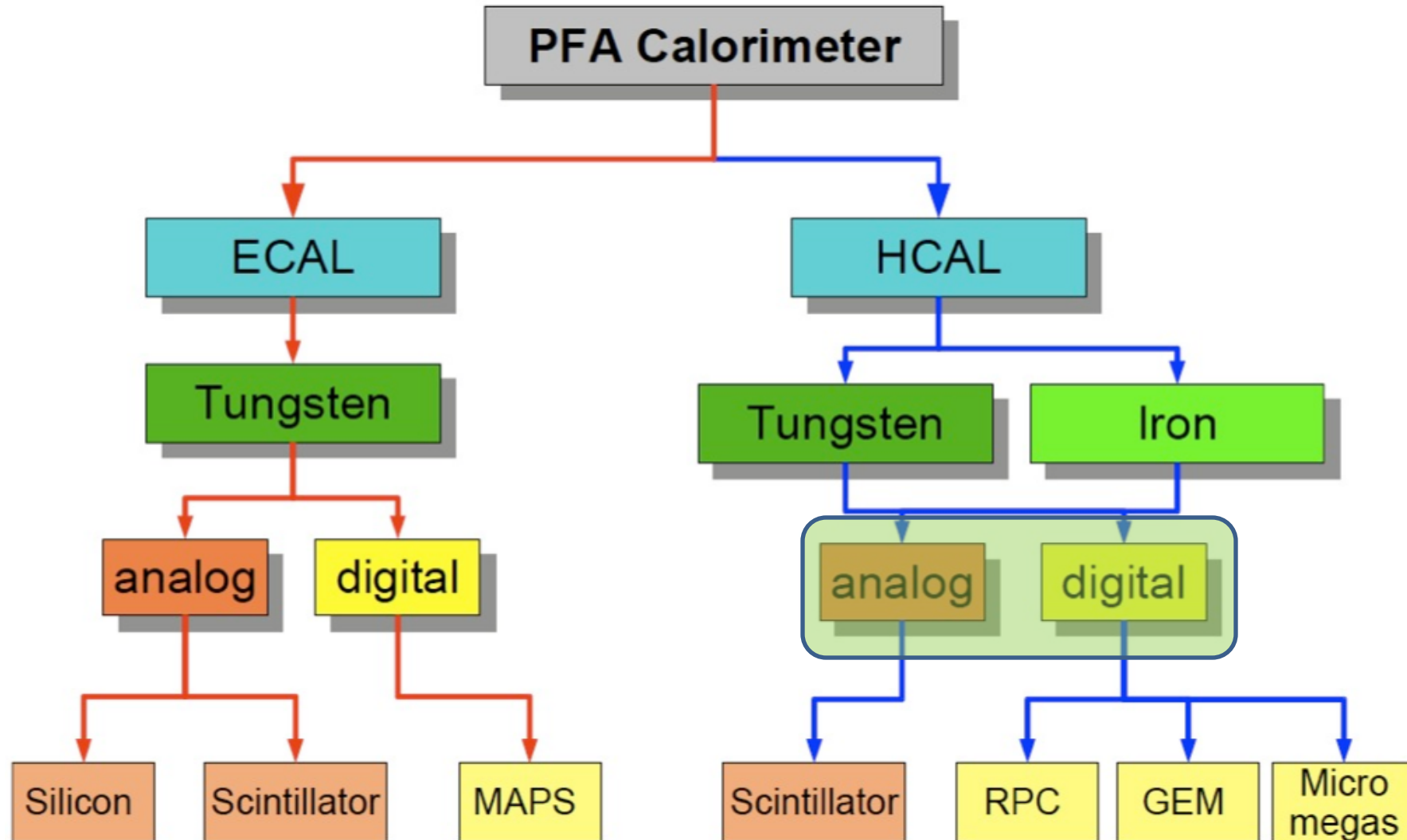
Fragment ID



# THE ZOO OF PFLOW CALORIMETERS



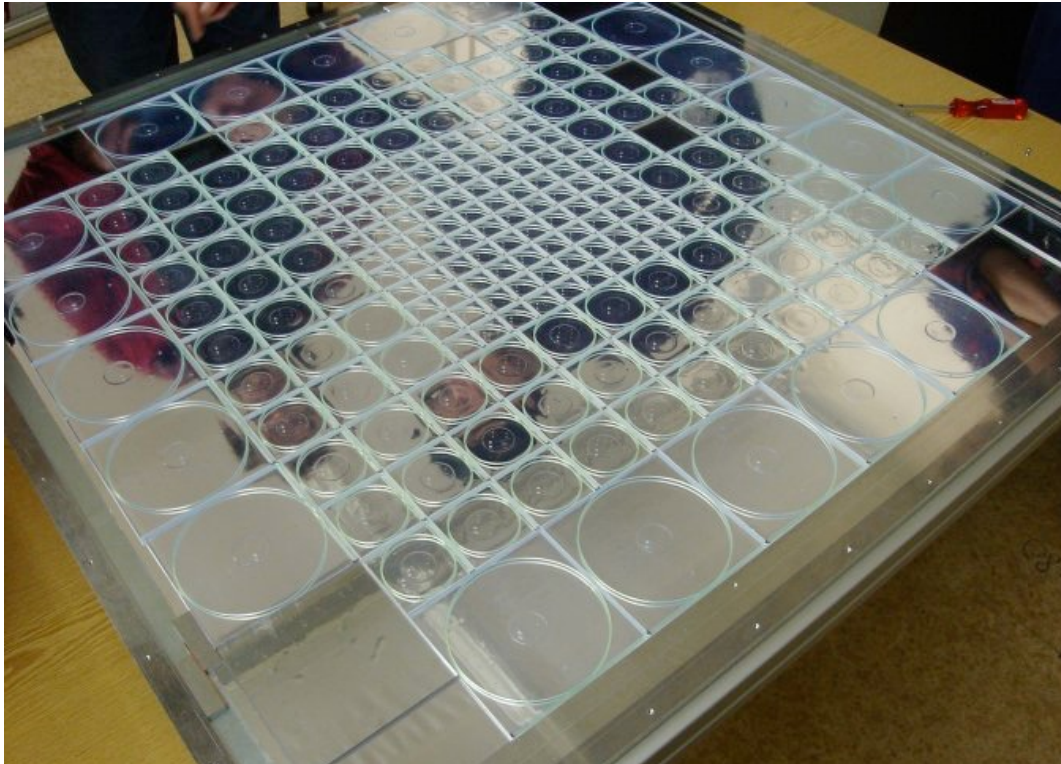
# THE ZOO OF PFLOW CALORIMETERS



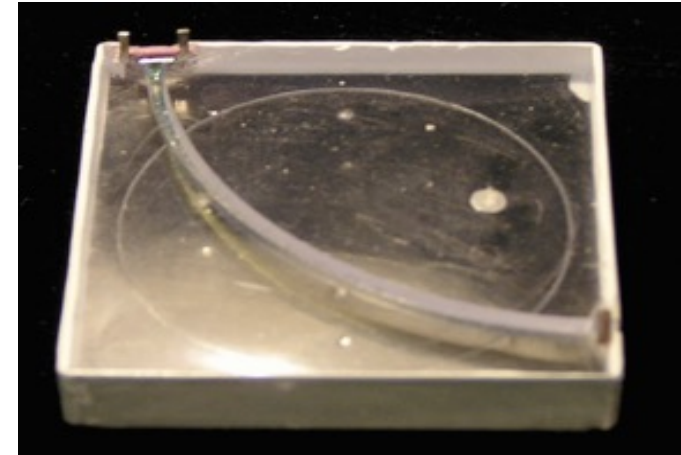
# NEW CONCEPTS: HIGHLY GRANULAR CALOS

- CALICE (CAlorimeter for a LInear Collider Experiment) HCAL prototype:
  - highly granular readout: 3 x 3 cm<sup>2</sup> scintillator tiles, 38 layers ( $\sim 4.7 \lambda_{\text{int}}$ ), each tile with individual SiPM readout

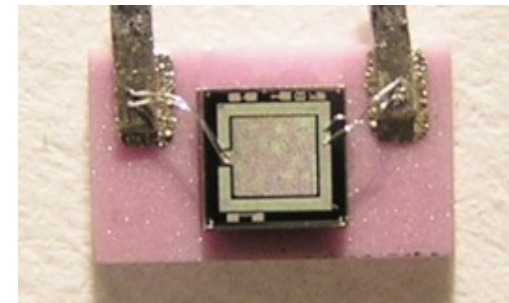
Pictures: CALICE collaboration



tiles in one layer



scintillator  
tile with  
WLS fiber

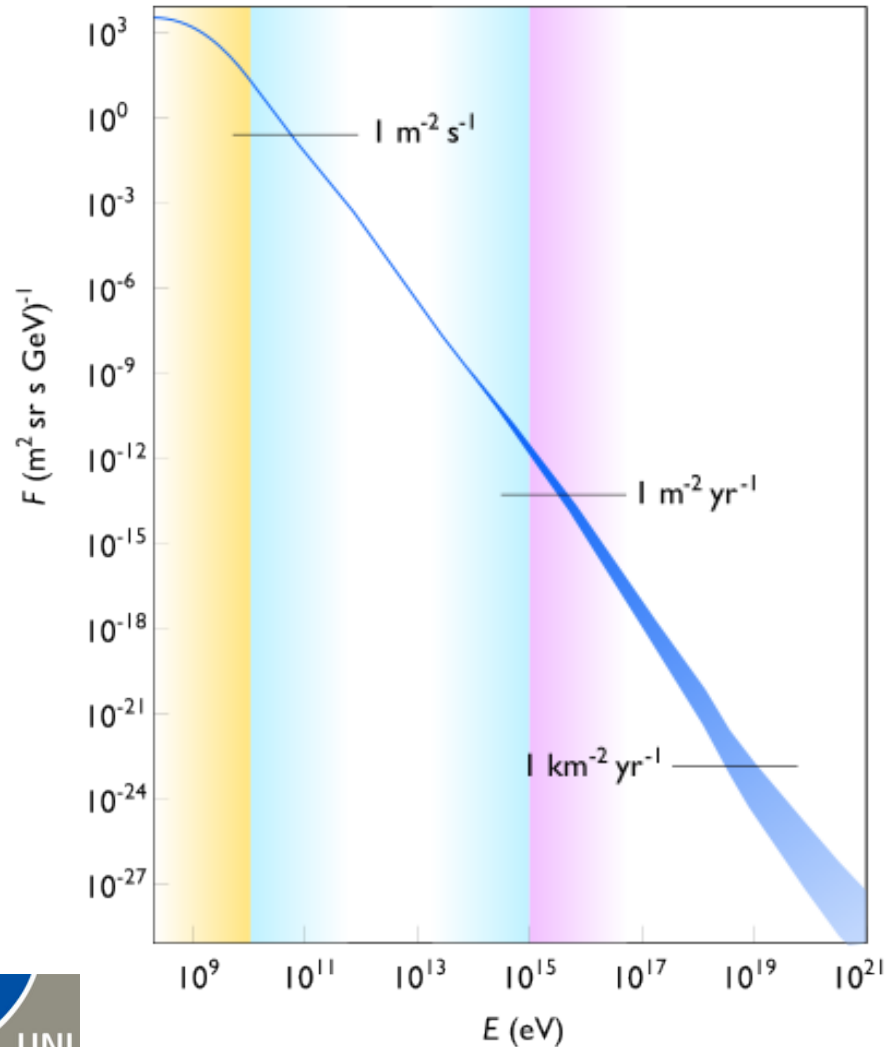


Silicon  
photo-multiplier



# CALOS: NOT ONLY AT ACCELERATORS!

- The methods used in particle physics are more and more used in astro particle physics.



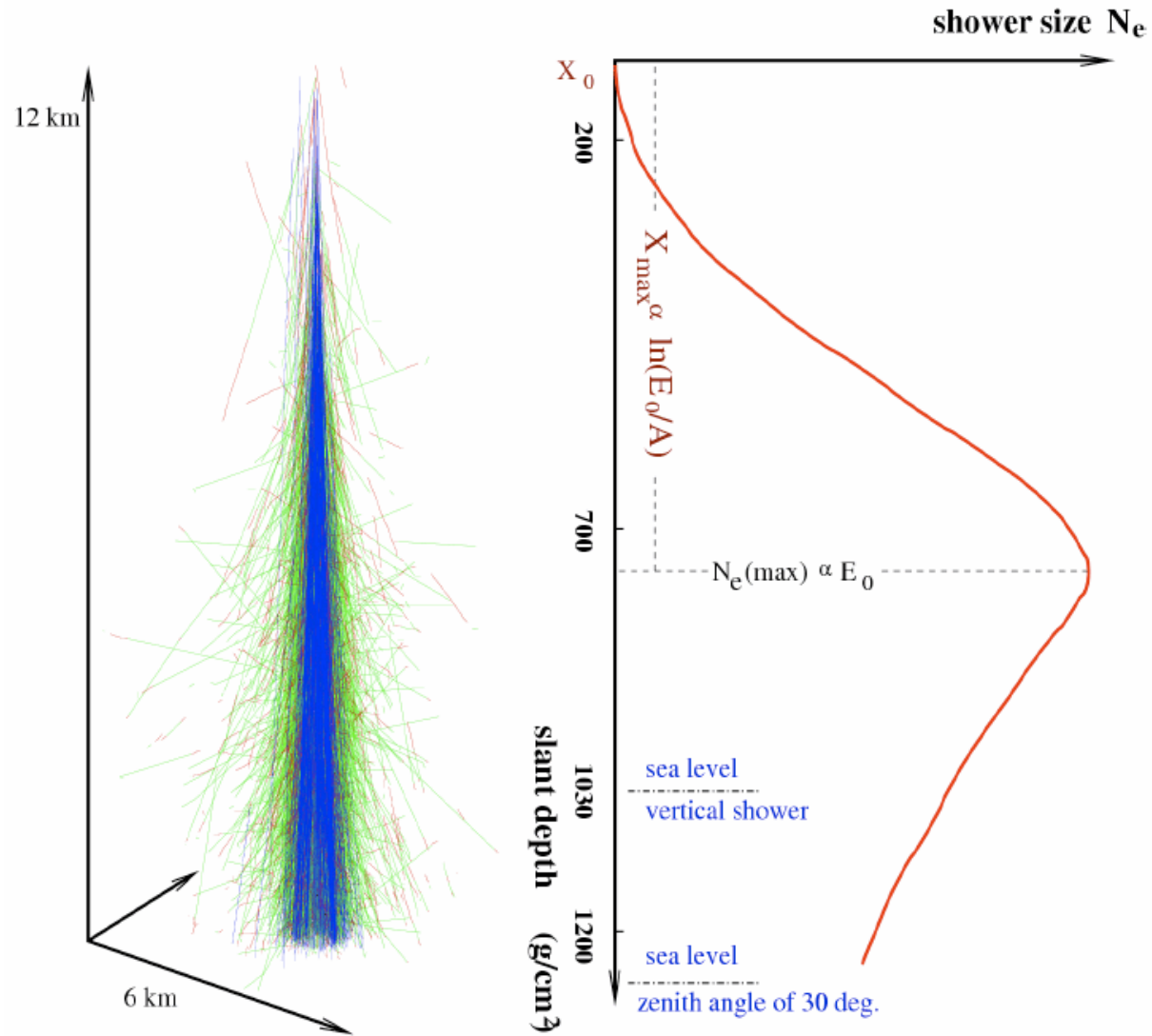
Pic: Wikipedia

## Requirements are different

- Search for extremely rare reactions
  - ▶ Large areas and volumina have to be covered
  - ▶ Background needs to be well suppressed
  - ▶ High efficiency: no event can be lost!
  - ▶ Data rate, radiation damage etc. are less of a problem

Flux of cosmic ray particles as a function of their energy.

# AIR SHOWER



- Mainly electromagnetic: photons, electrons
- Shower maximum:  
 $\sim \ln(E_0/A)$

## Use atmosphere as calorimeter

Nuclear reaction length  $\lambda_l \sim 90 \text{ g/cm}^2$

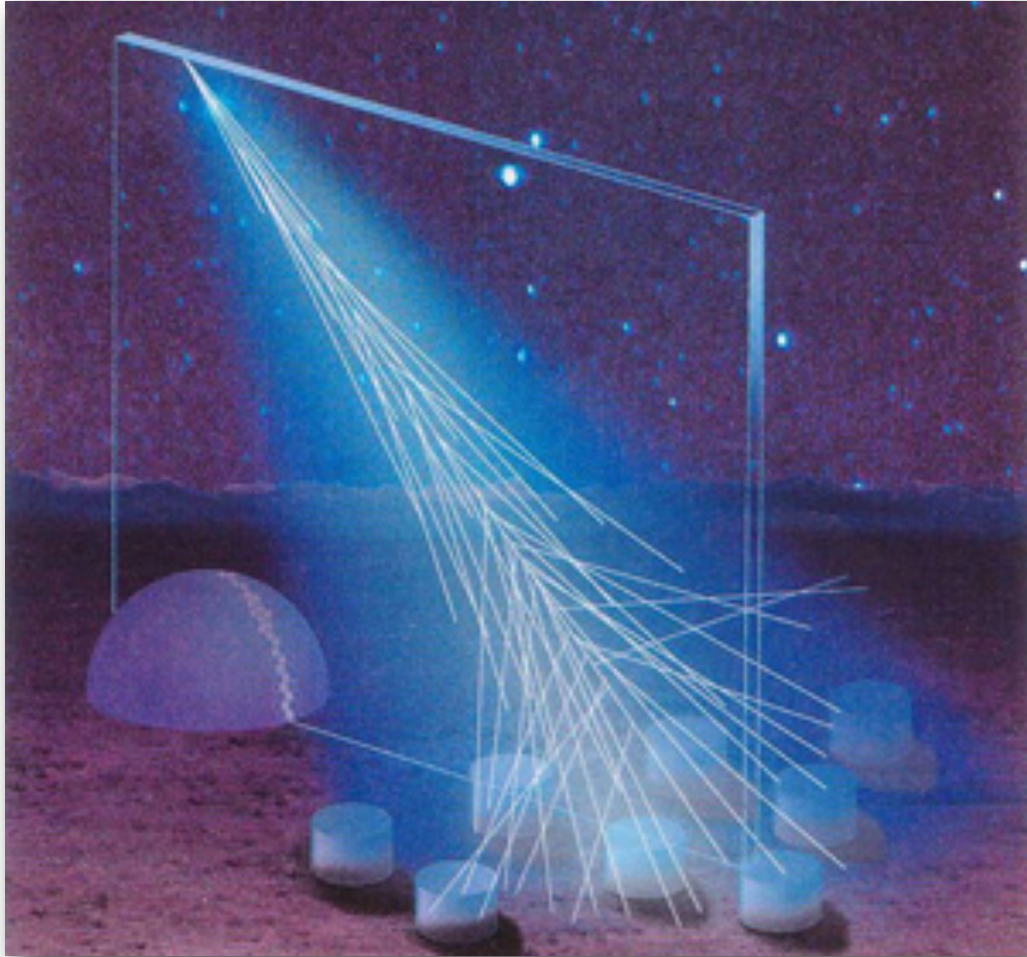
Radiation length  $X_0 \sim 36.6 \text{ g/cm}^2$

Density:  $\sim 1035 \text{ g/cm}^3$

$\sim 11 \lambda_l, \sim 28 X_0$

# TWO TECHNIQUES

Pic: Pierre Auger Observatory



- The atmosphere as homogeneous calorimeter:
  - Energy measurement by measuring the fluorescence light

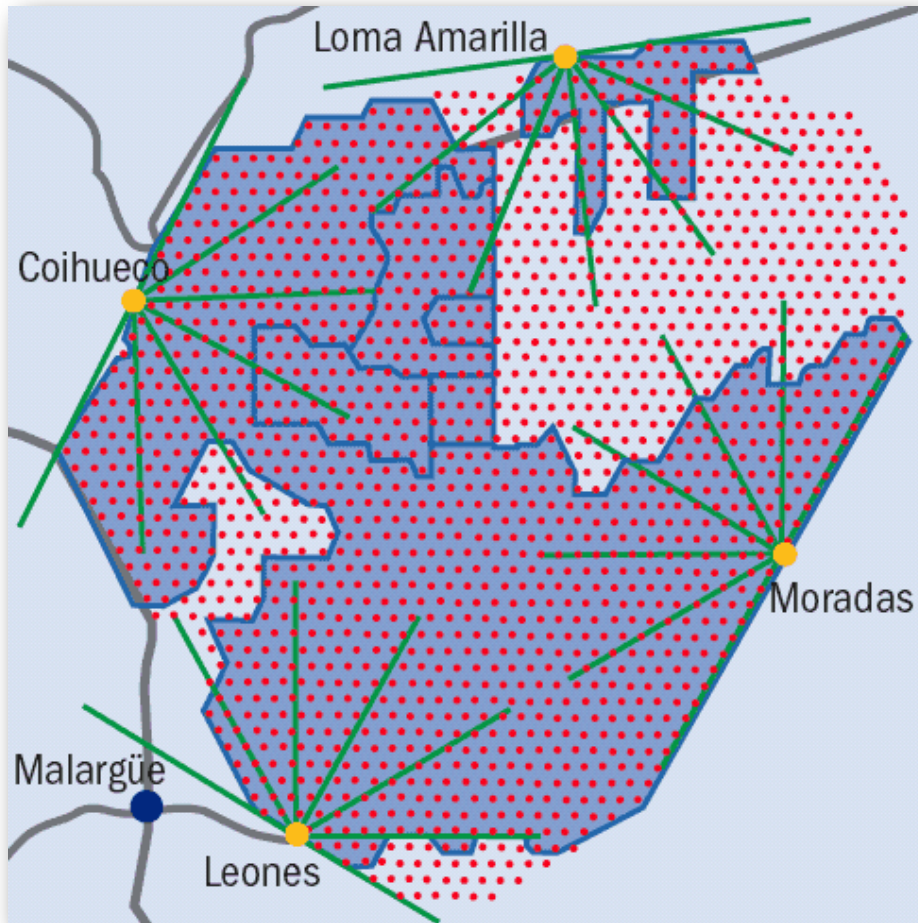
This is only possible with clear skies and darkness !

- A one-layer sampling calorimeter  $11 \lambda$  absorber
  - Energy measurement using particle multiplicity

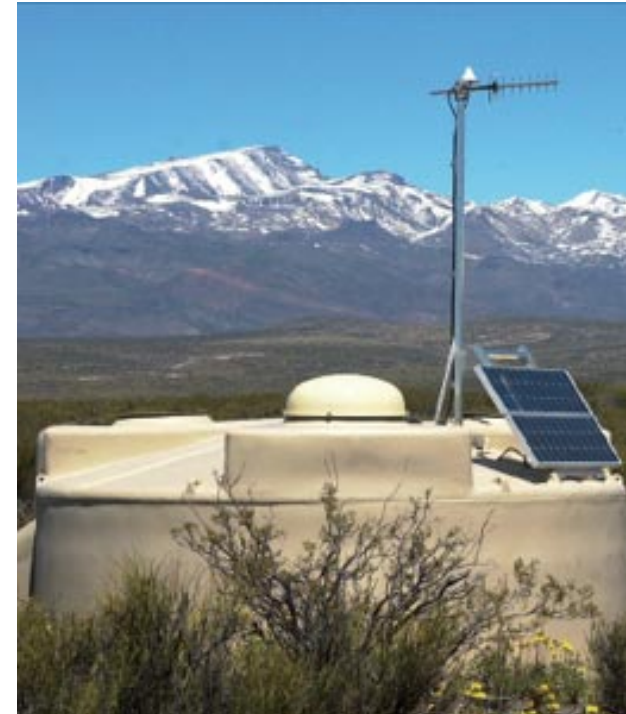
Always possible but has large uncertainties !

# AUGER-SOUTH: ARGENTINIAN PAMPA

Pics: Pierre Auger Observatory

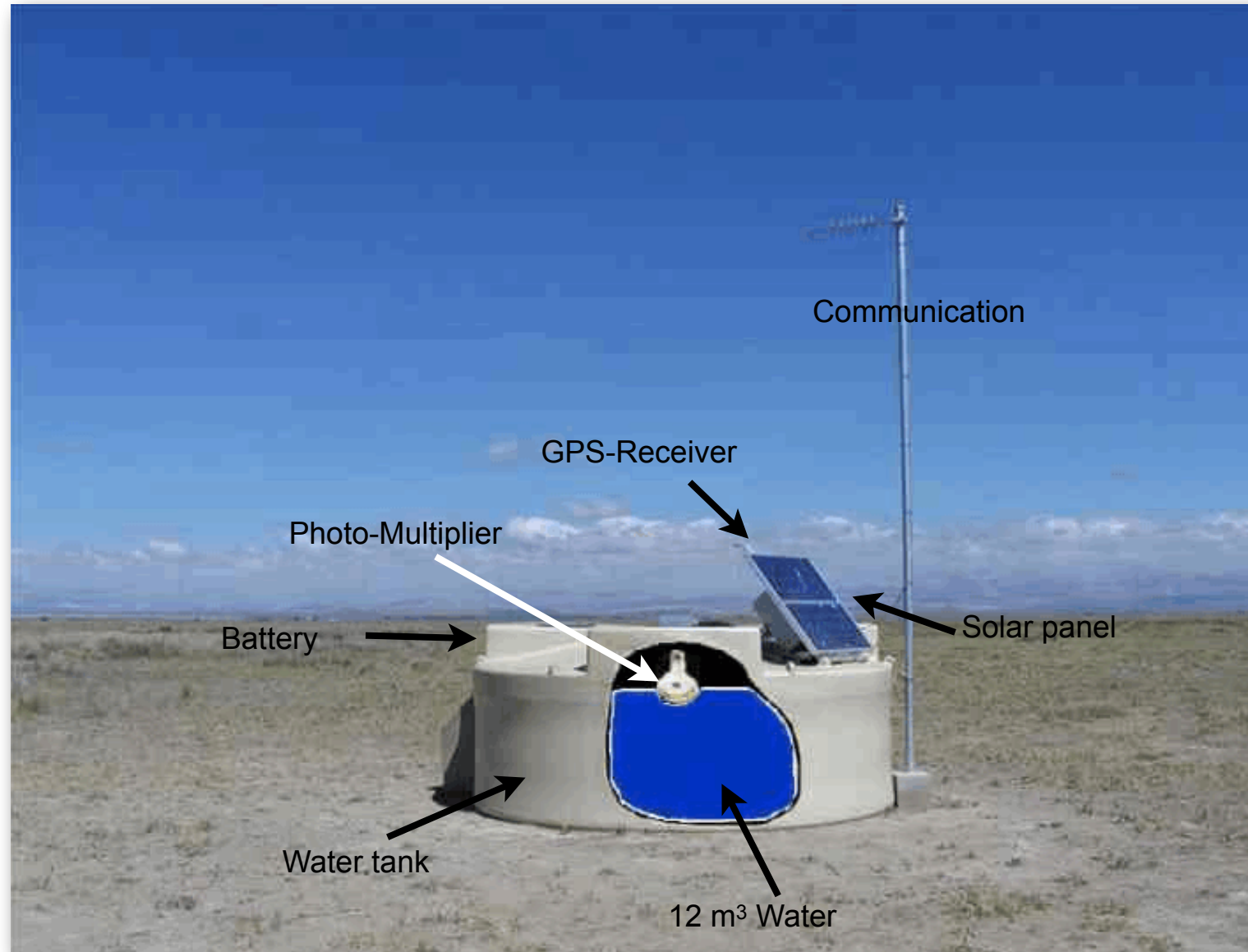


- 1600 water-Cherenkov detectors on ground
- 4 Fluorescence-stations with 6 telescopes
- Covered area:  
3000 km<sup>2</sup> (30 x Paris)
- Designed to measure energies above 10<sup>18</sup>eV

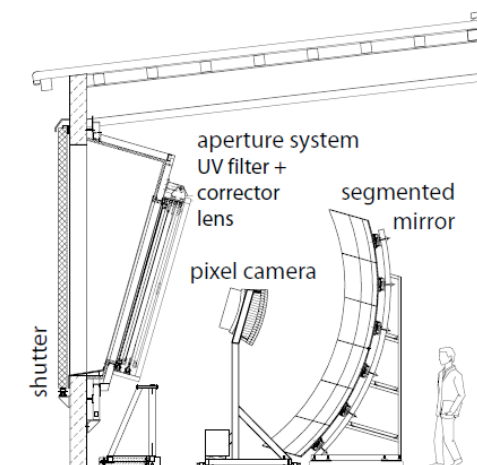
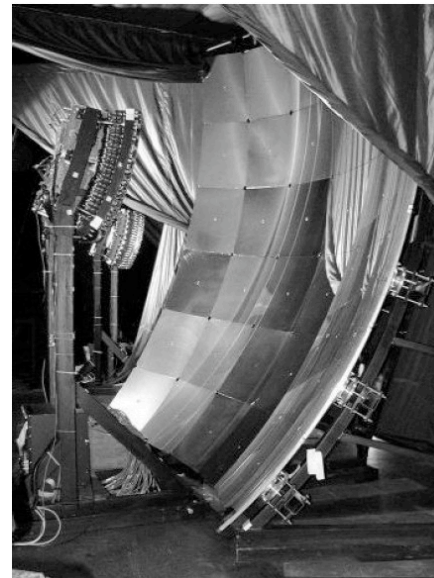
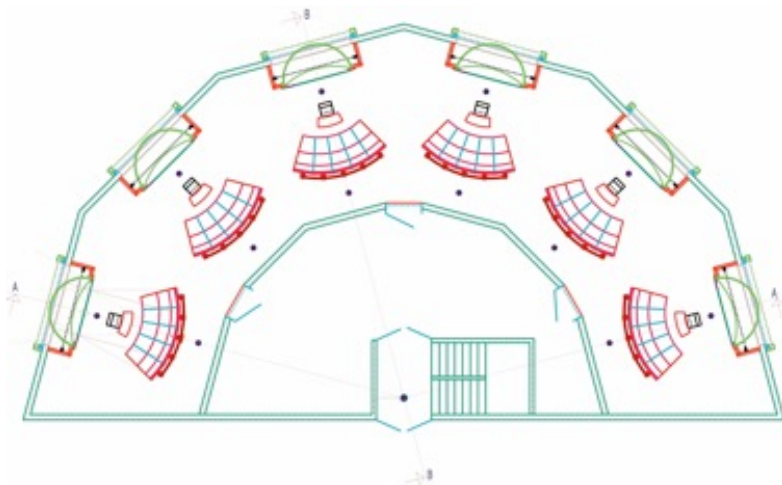
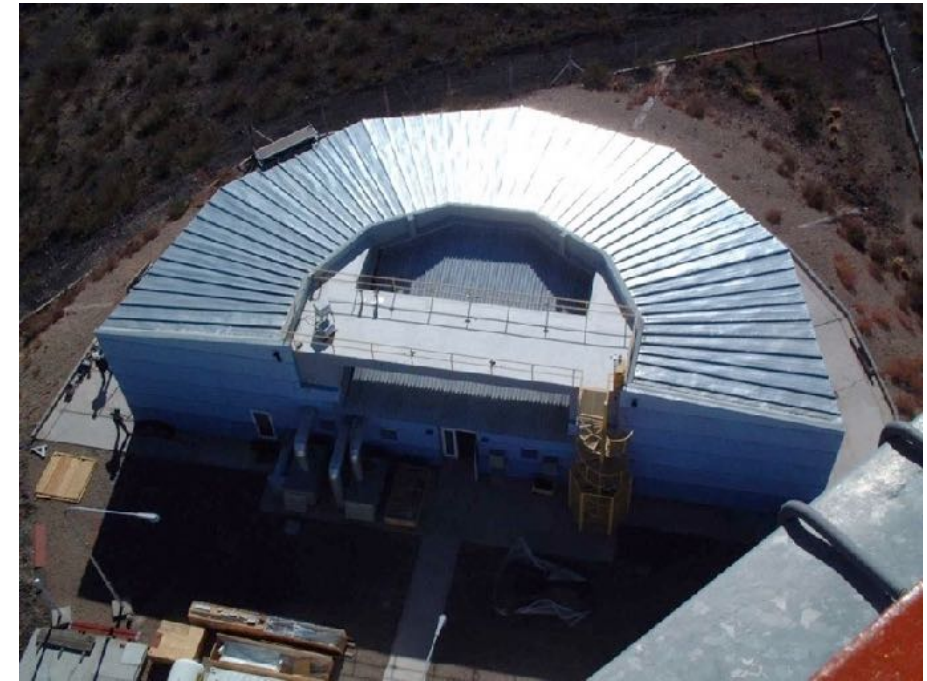




# AUGER-DETEKTOR: GROUND ARRAY



# AUGER HYBRID INSTALLATION



# SUMMARY CALORIMETERS

Calorimeters can be classified into:

## **Electromagnetic Calorimeters,**

- to measure electrons and photons through their EM interactions.

## **Hadron Calorimeters,**

- Used to measure hadrons through their strong and EM interactions.

The construction can be classified into:

## **Homogeneous Calorimeters,**

- that are built of only one type of material that performs both tasks, energy degradation and signal generation.

## **Sampling Calorimeters,**

- that consist of alternating layers of an absorber, a dense material used to degrade the energy of the incident particle, and an active medium that provides the detectable signal.

# OVERVIEW

I. Detectors for Particle Physics

II. Interaction with Matter

III. Calorimeters

V. Tracking Detectors

- Gas detectors
- Semiconductor trackers

VI. Examples from the real life

}

*Tuesday*

}

*Wednesday*