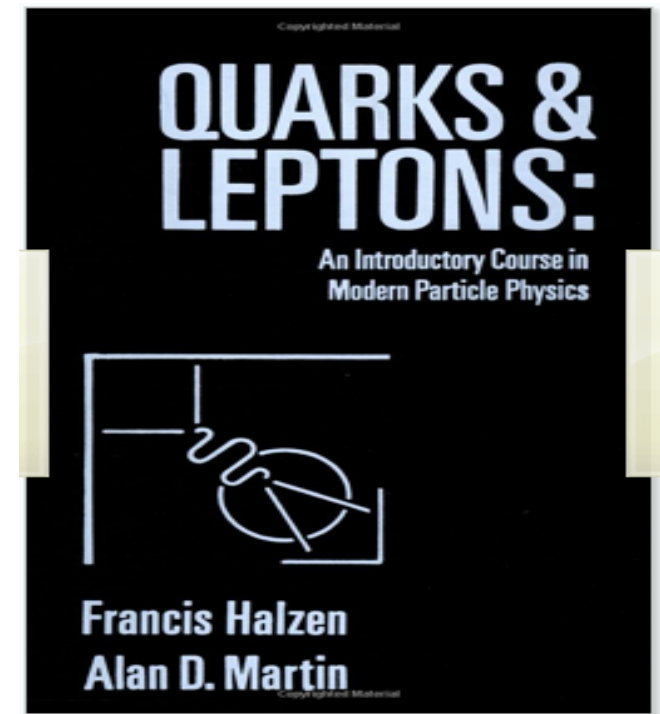
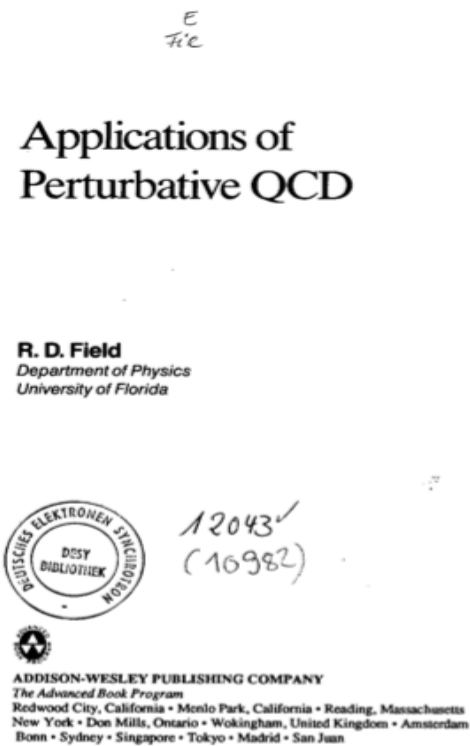
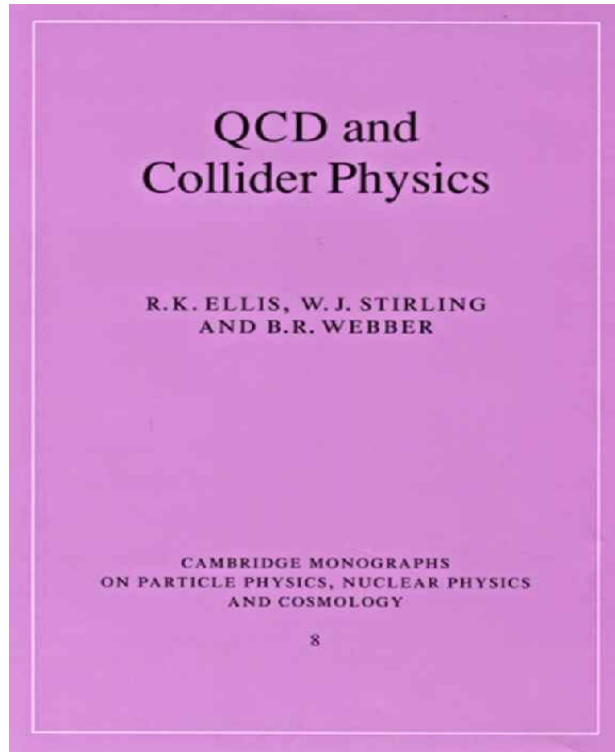


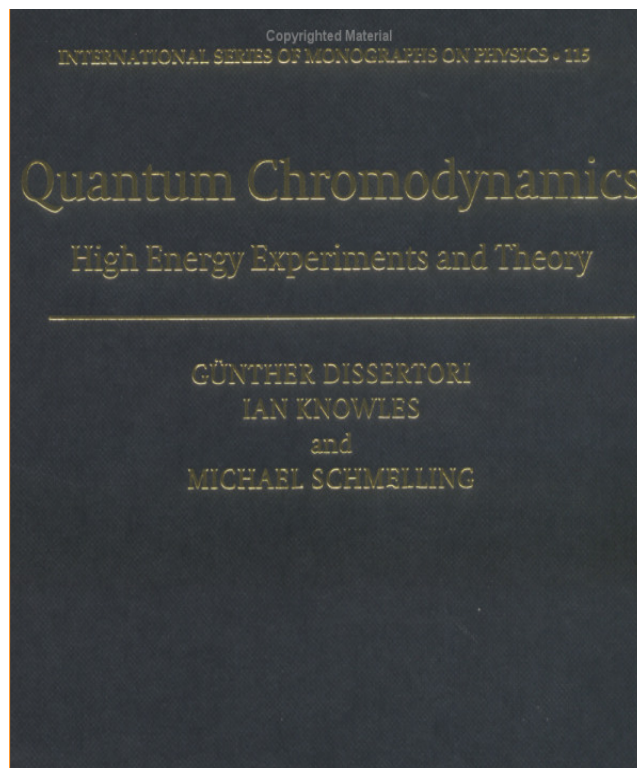
QCD parton dynamics, hadronic final states, MCs

H. Jung (DESY)
hannes.jung@desy.de

Literature



http://www.desy.de/~jung/qcd_and_mc_2015



QCD AND MONTE CARLOS (2015)

**QCD and Monte Carlos
Lecture Course
Winter 2015
Hannes Jung (DESY, U Antwerp)
University Antwerp**

The lecture course is intended for master students, PhD students and postdocs. It covers a basic introduction to QCD and the QCD evolution equations (DGLAP, BFKL and CCFM). A basic introduction to Monte Carlo methods will be given, and these methods will be applied to calculate cross sections and the evolution of parton densities. We will go through the latest QCD results from LHC.

During the course we will learn, how to write a small program to integrate a partonic cross section. We will apply Monte Carlo methods to solving the DGLAP evolution equation and to calculate the transverse momentum spectrum of Higgs production in proton-proton collisions at the LHC.

The course is held at University Antwerp, in blocks of lectures with exercise sessions.

Lectures will be on

Monday 12. Oct	13:45 - 17:00	L203
Tuesday 13. Oct	13:45 - 17:00	L203
Monday 26. Oct	13:45 - 17:00	L203
Tuesday 27. Oct	13:45 - 17:00	L203
Monday 16. Nov	13:45 - 17:00	L203
Thursday 19. Nov	13:45 - 17:00	L203
Wednesday 14. Oct	13:45 - 17:00	L406
Thursday 15. Oct	13:45 - 17:00	L203
Wednesday 28. Oct	13:45 - 17:00	L406
Thursday 29. Oct	13:45 - 17:00	L203
Thursday 17. Nov	13:45 - 17:00	L203
Wednesday 18. Nov	13:45 - 17:00	L406

Literature
MC and statistics:
G. Stasin and G. Zech, *Introduction to statistics and measurement analysis for physicists*, Harburg, Germany: DESY (2005) 391 p.
G. Cowan, *Statistical data analysis*, Oxford, UK: Clarendon (1998) 197 p.
D. E. Knuth, *The Art of Computer Programming Volumes 1-3* (Second Ed.), Addison-Wesley Longman Publishing Co., Inc., Boston, MA, USA, 1996.
QCD:
Y.L. Dokshitzer, V.A. Khoze, A.H. Mueller, and S.I. Troyan, *Basics of perturbative QCD*, GIF sur/Vietri, France: Ed. Frontiers (1991).
R.K. Ellis, W.J. Stirling, and B.R. Webber, *QCD and collider physics*, Camb. Monogr. Part. Phys. Nucl. Phys. Cosmol., 81-435, 1996.
R. D. Field, *Applications of Perturbative QCD*, Front. Phys., 77:1-366, 1989.
F. Halzen and A. D. Martin, *QUARKS AND LEPTONS: AN INTRODUCTORY COURSE IN MODERN PARTICLE PHYSICS*, Wiley (1984) 396p.
Email Me: hannes.jung@desy.de

PHYSICS REPORTS (Review Section of Physica Letters) 81, No. 1 (1982) 1-129. NORTH-HOLLAND PUBLISHING COMPANY

PARTONS IN QUANTUM CHROMODYNAMICS

Guido Altarelli
Istituto di Fisica, Università di Roma,
Istituto Nazionale di Fisica Nucleare, Sezione di Roma, Italy

Received 20 July 1981

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Abstract:

An overall view of the physics of QCD in the perturbative domain is presented in a form that could be of use both as an introduction to the subject with its main lines of current development and as a reference review text for more expert readers as well.

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What is QCD ?

- Quantum Chromo Dynamics = quantum field theory of strong interactions

- Basic elements;

- quarks: u,d,s,c,b,t

- fractional charges $1/3$ or $2/3$

- carry color: $r, g, b = white$

- anti-quarks carry anti-color: $r\bar{r} = white, g\bar{g} = white, b\bar{b} = white$

- exchange bosons are gluons

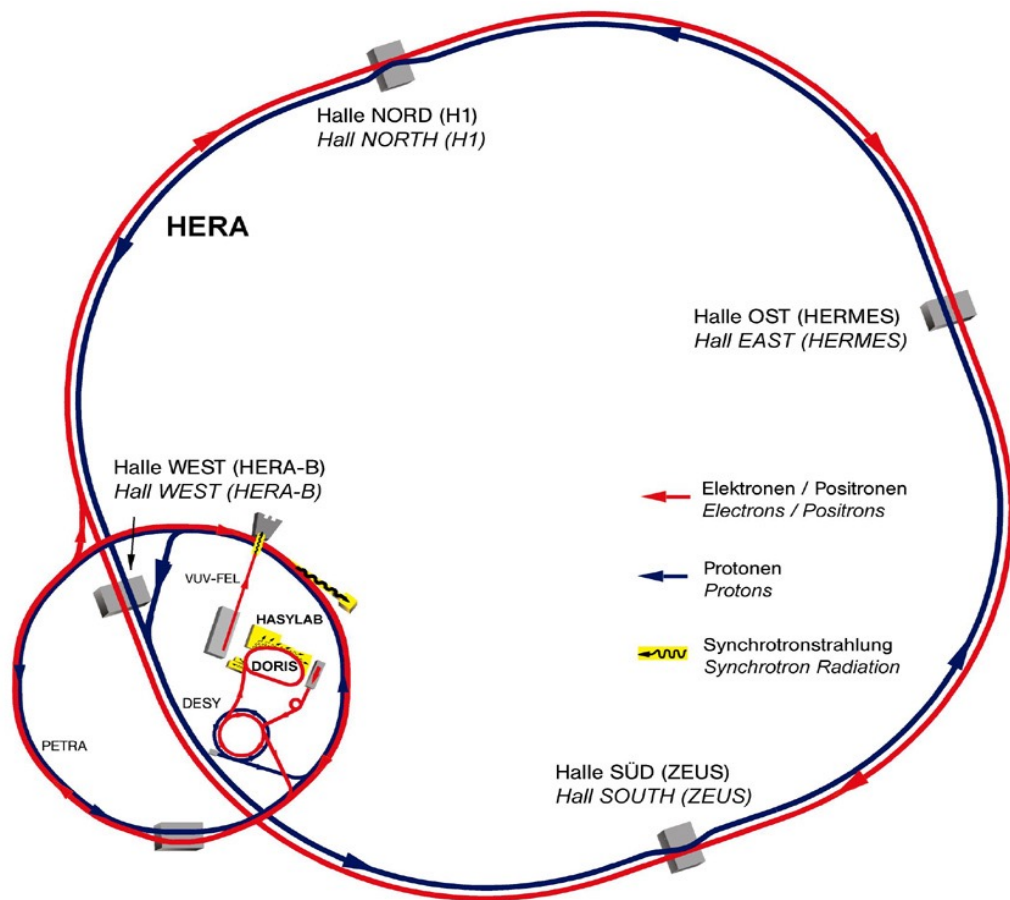
- gluons couple to quarks, anti-quarks and gluons

- gluons carry color+anti-color (but not in a color singlet):

$$r\bar{g}, r\bar{b}, g\bar{r}, g\bar{b}, \sqrt{\frac{1}{2}}(r\bar{r} - g\bar{g}), \sqrt{\frac{1}{2}}(r\bar{r} + g\bar{g} - 2b\bar{b})$$

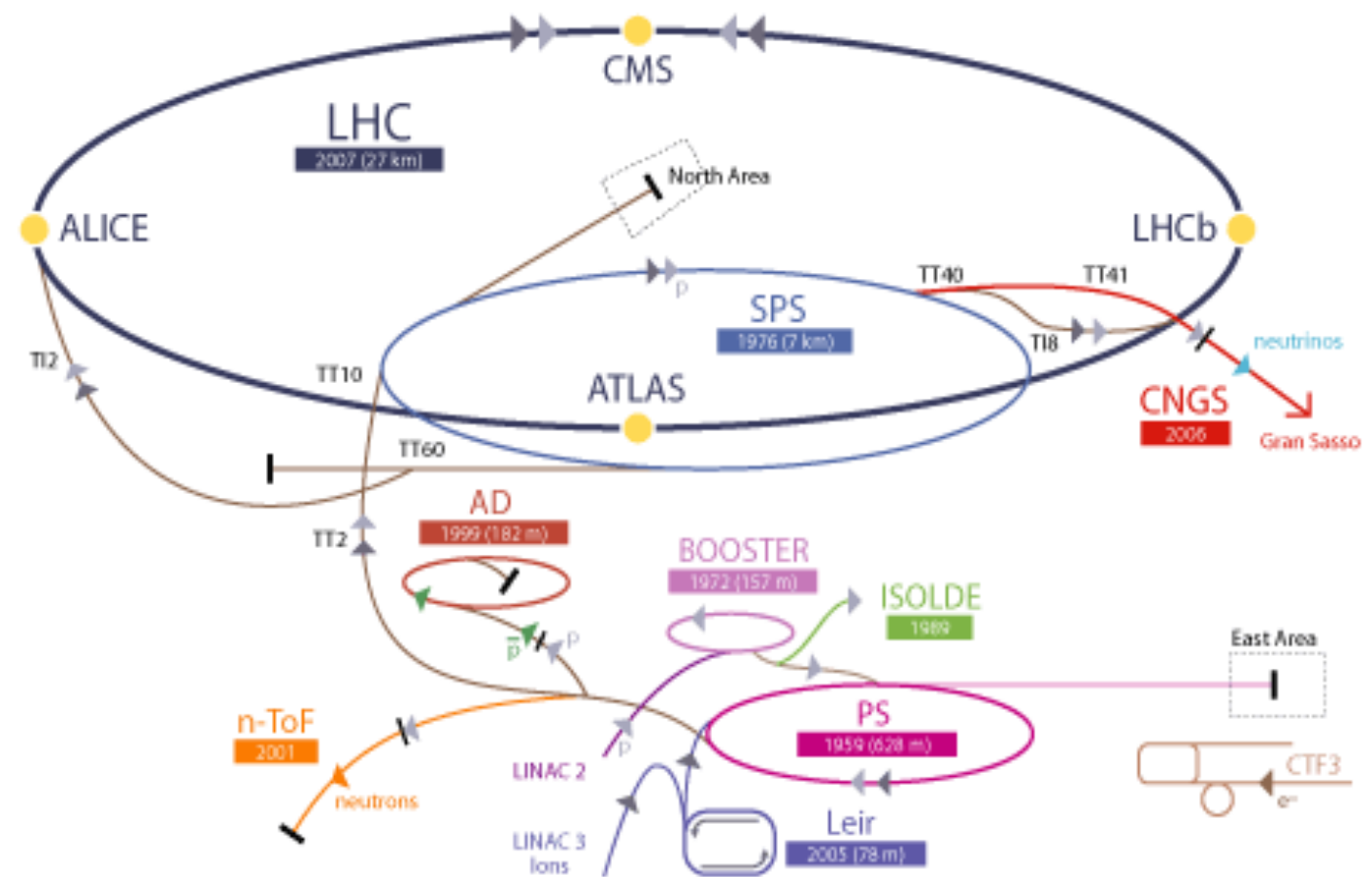
Fixed Target, HERA & LHC

- HERA at DESY



- SPS & LHC at CERN

CERN Accelerator Complex



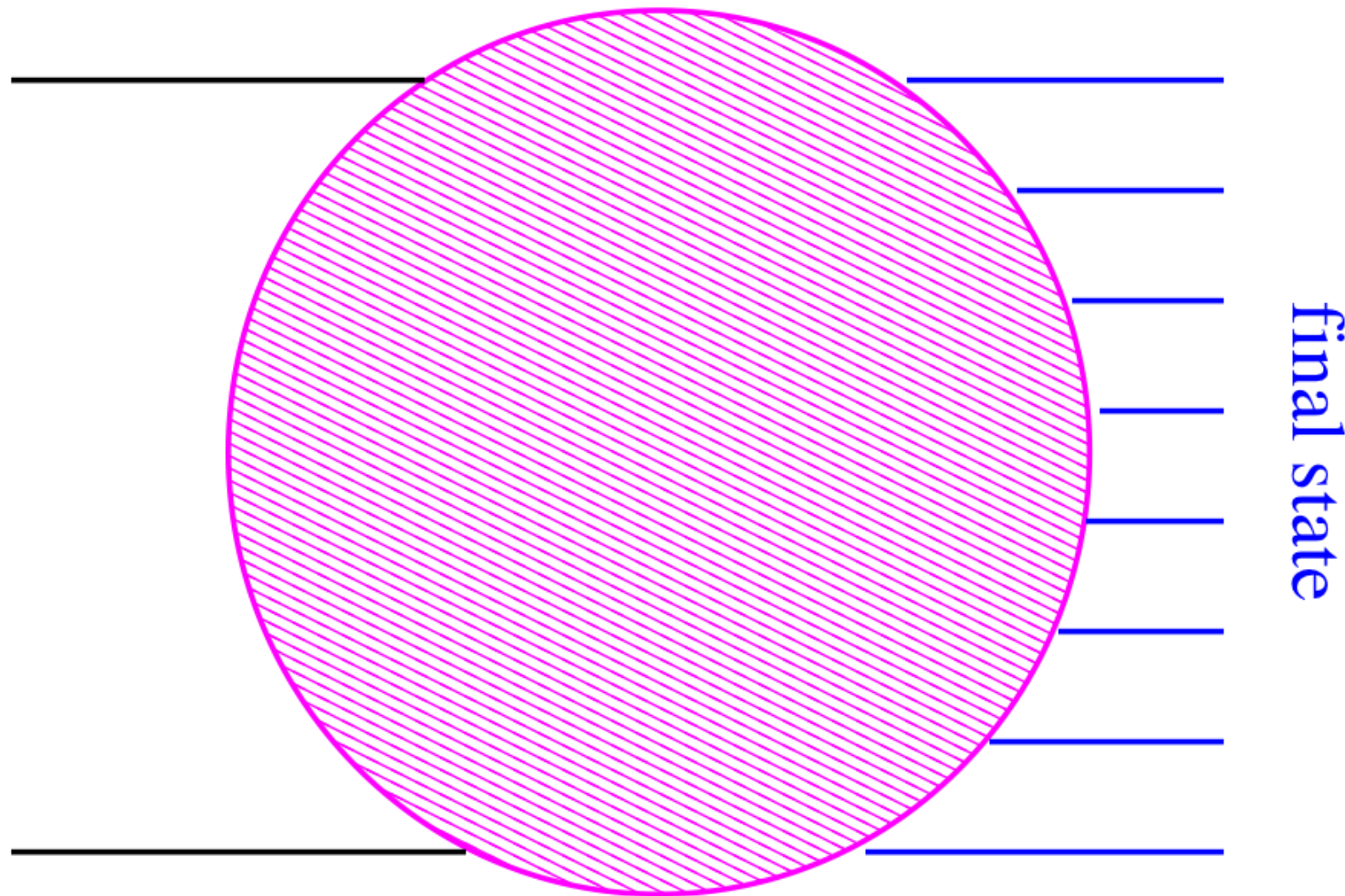
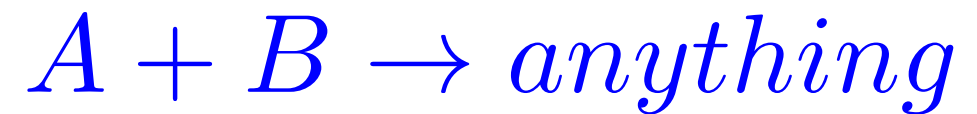
▶ p (proton) ▶ ion ▶ neutrons ▶ \bar{p} (antiproton) ▶ neutrinos ▶ electron
 ↔↔↔ proton/antiproton conversion
 LHC Large Hadron Collider SPS Super Proton Synchrotron PS Proton Synchrotron
 AD Antiproton Decelerator CTF3 Clic Test Facility
 CNGS Cern Neutrinos to Gran Sasso ISOLDE Isotope Separator OnLine DEvice
 LEIR Low Energy Ion Ring LINAC LINEar ACcelerator n-ToF Neutrons Time Of Flight

$$\sqrt{s}_{HERA} = 300 GeV$$

$$\sqrt{s}_{LHC} = 7 - 14 TeV$$

The general case

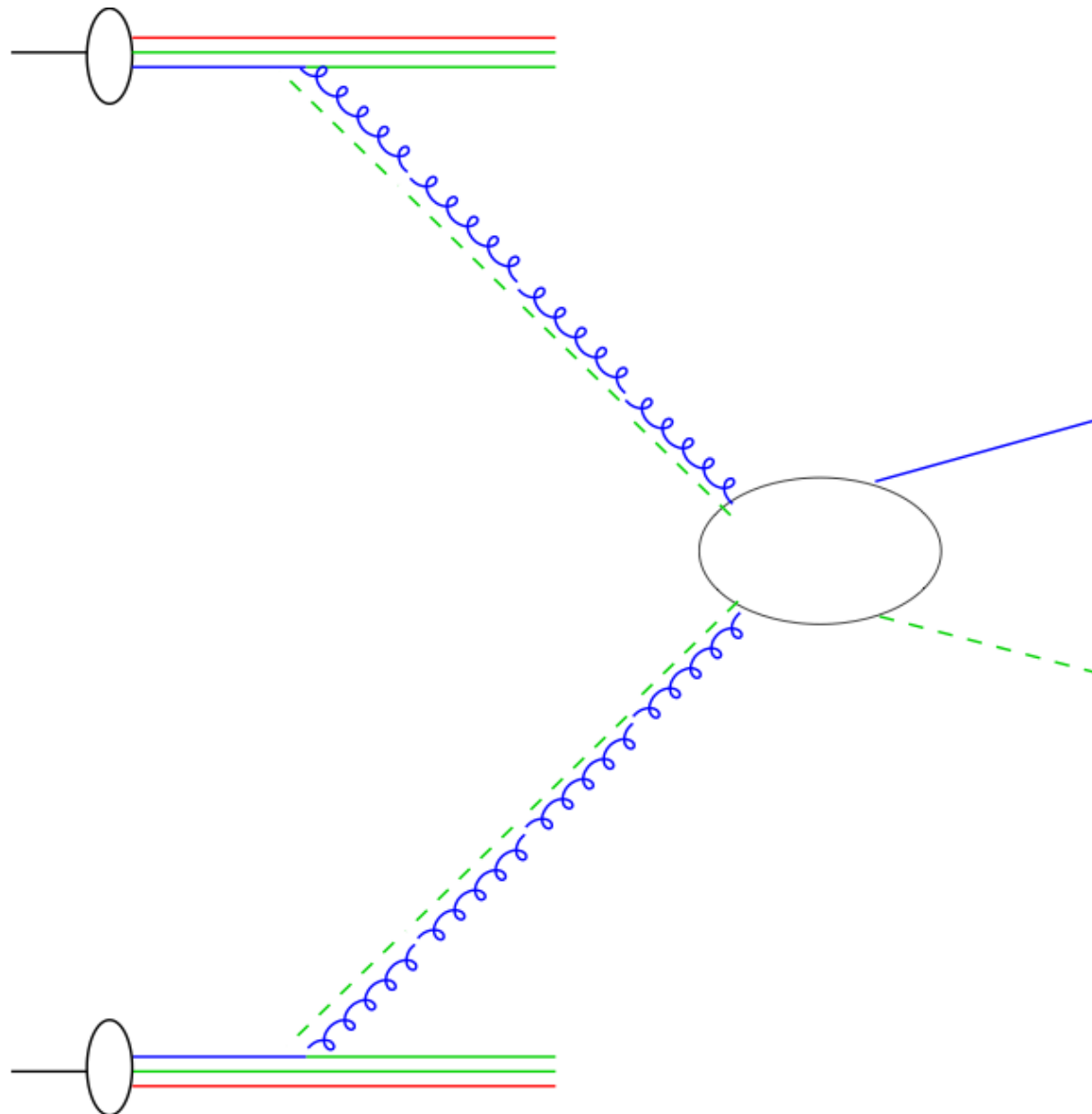
- Calculation of cross section of



→ Start with jet production

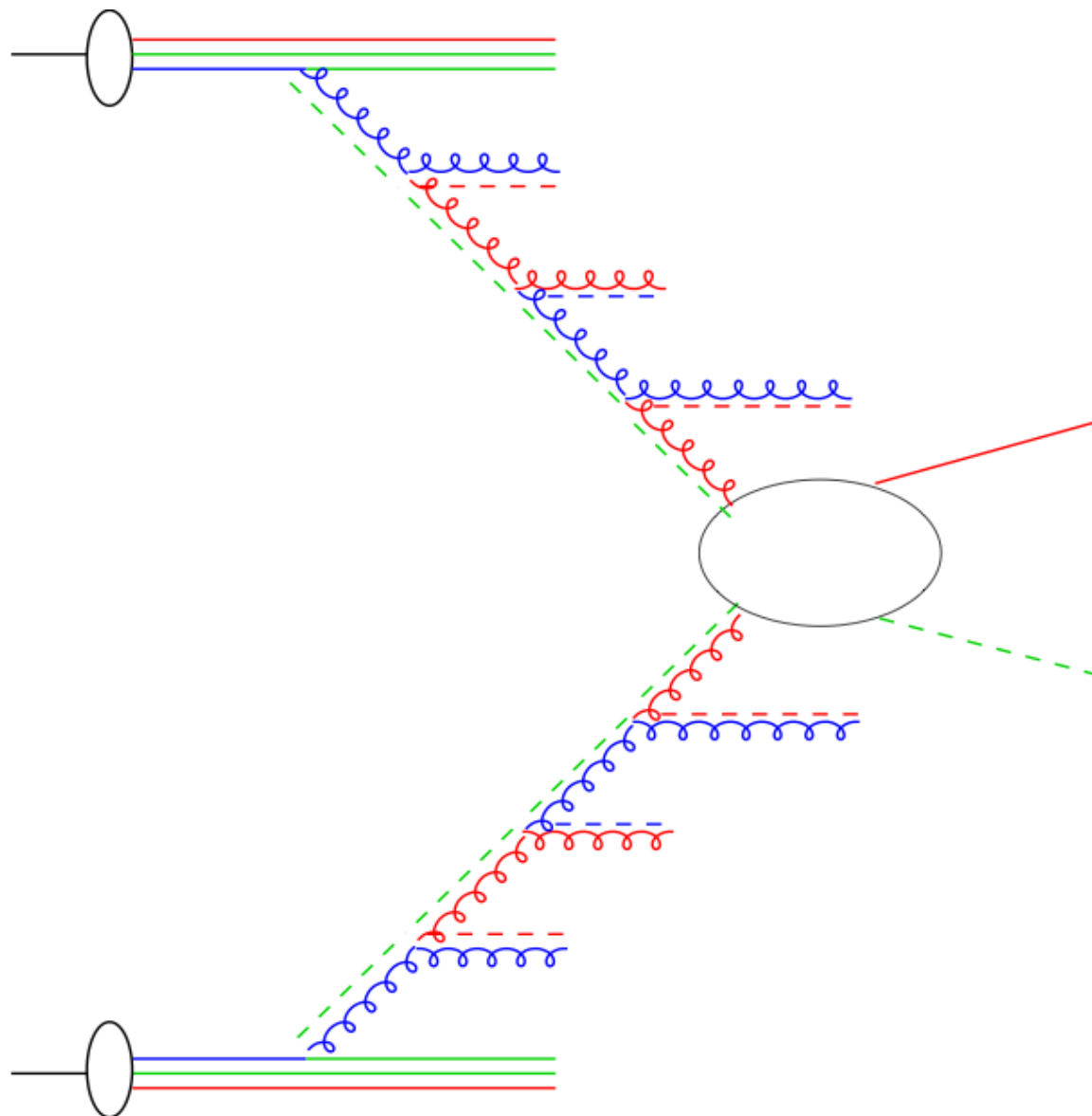
Picture of jet production

- General approach to hard scattering processes



Picture of jet production

- General approach to hard scattering processes
 - including higher order parton radiation



Picture of jet production

- General approach to hard scattering processes
 - including higher order parton radiation
 - adding hadronization and fragmentation
- leads to the concept of factorization:

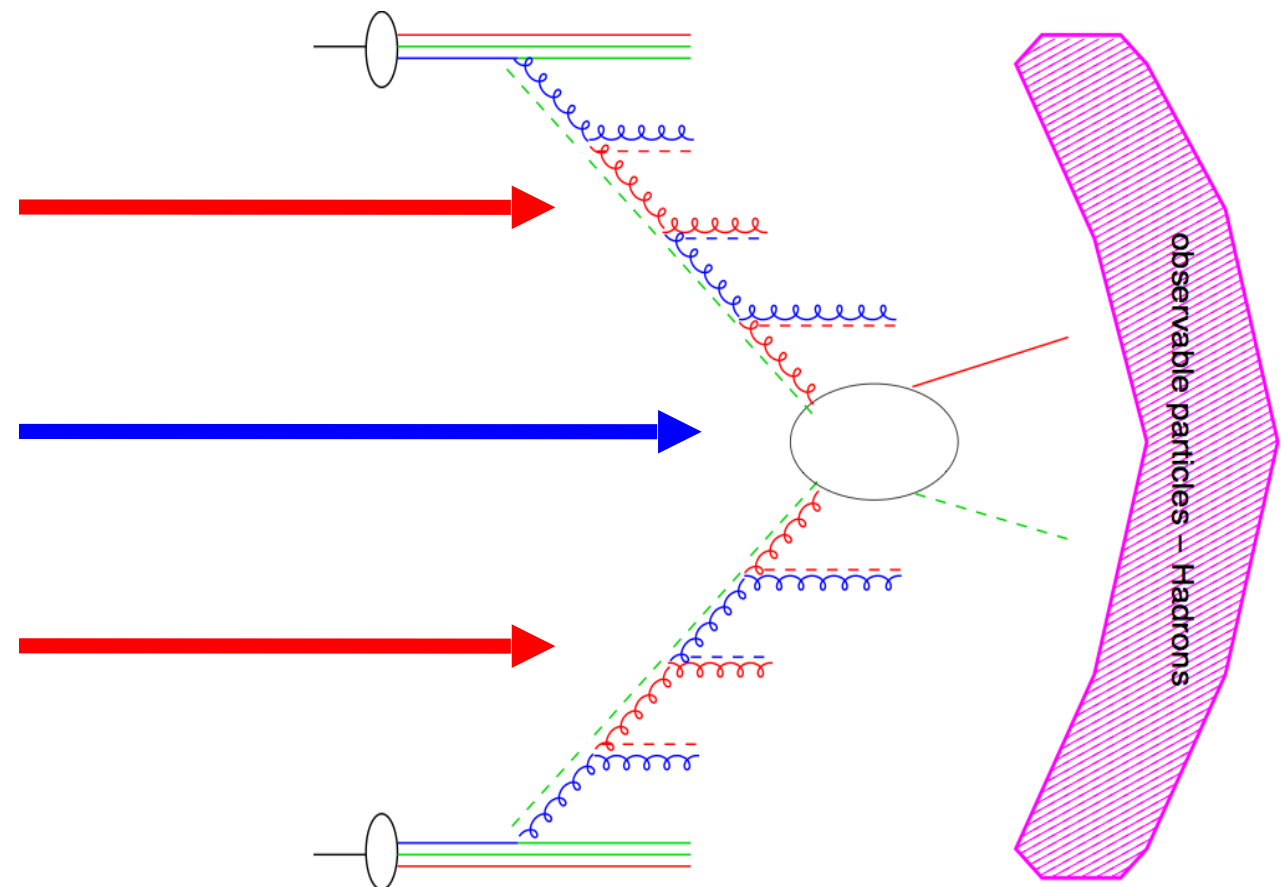
$$\sigma = f(x_1, \mu^2)$$

⊗

$$\hat{\sigma}(x_1, x_2, \mu^2)$$

⊗

$$f(x_2, \mu^2)$$



→ factorization ansatz is used in

any calculation (LO, NLO, MC event generators ...)

How can processes be calculated ?

- Monte Carlo method
 - **refers** to any procedure that makes use of random numbers
 - **uses** probability statistics to solve the problem
- Monte Carlo methods are used in:
 - Simulation of natural phenomena
 - Simulation of experimental apparatus
 - Numerical analysis

Monte Carlo method

- Monte Carlo method
 - **refers** to any procedure that makes use of random numbers
 - **uses** probability statistics to solve the problem
- Monte Carlo methods are used in:
 - Simulation of natural phenomena
 - Simulation of experimental apparatus
 - Numerical analysis
- Random Numbers
 - one of them is **3**
 - **No such thing as a single random number**
 - A sequence of random numbers is a set of numbers that have nothing to do with the other numbers in a sequence

Random Numbers

- In a uniform distribution of random numbers in $[0,1]$ every number has the same chance of showing up
- Note that 0.0000000001 is just as likely as 0.5

To obtain random numbers:

- Use some chaotic system like roulette, lotto, 6-49, ...
- Use a process, inherently random, like radioactive decay
- Tables of a few million truly random numbers exist

(.....until a few years ago.....)

BUT not enough for most applications

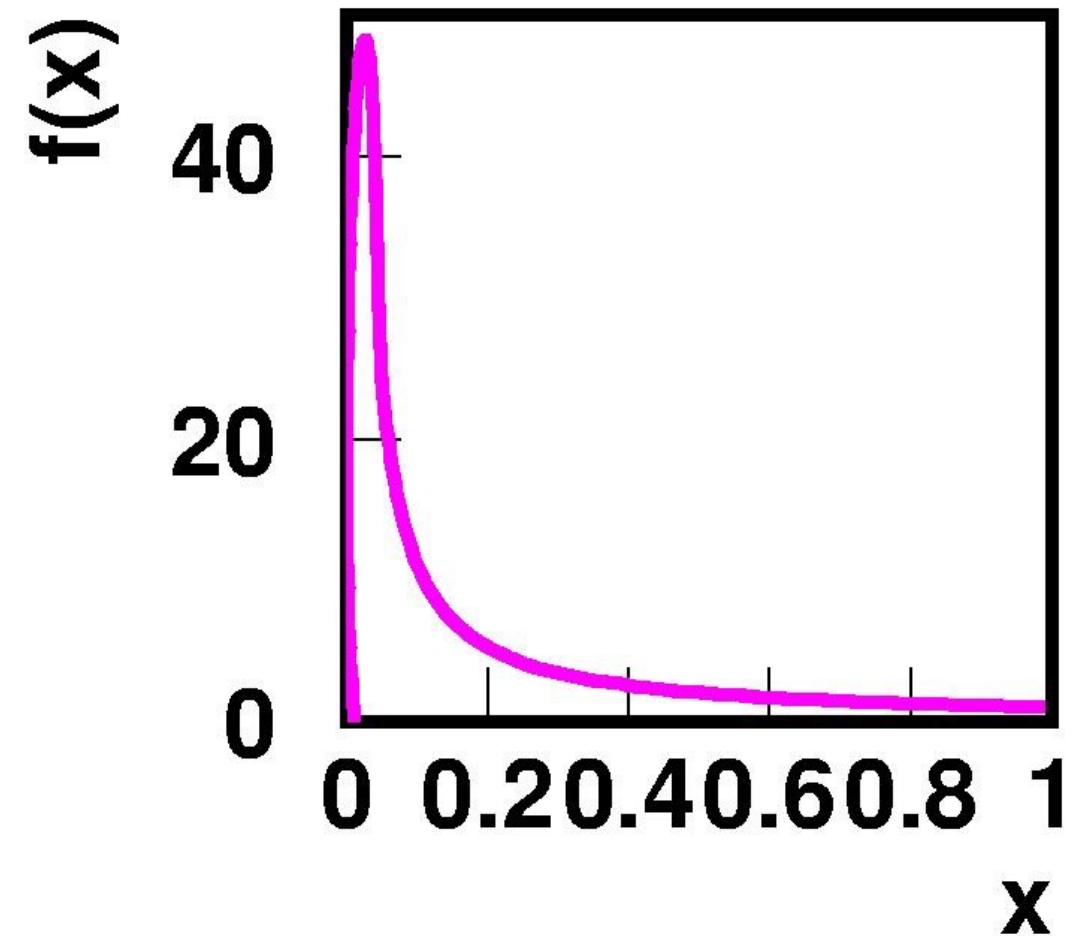
→..... we have true random number generators ...

Generating distributions

- Brute Force or Hit & Miss method
 - use this if there is no easy way to find a analytic integrable function
 - find $c \leq \max f(x)$
 - reject if $f(x_i) \leq u_j \cdot c$
 - accept if $f(x_i) \geq u_j \cdot c$

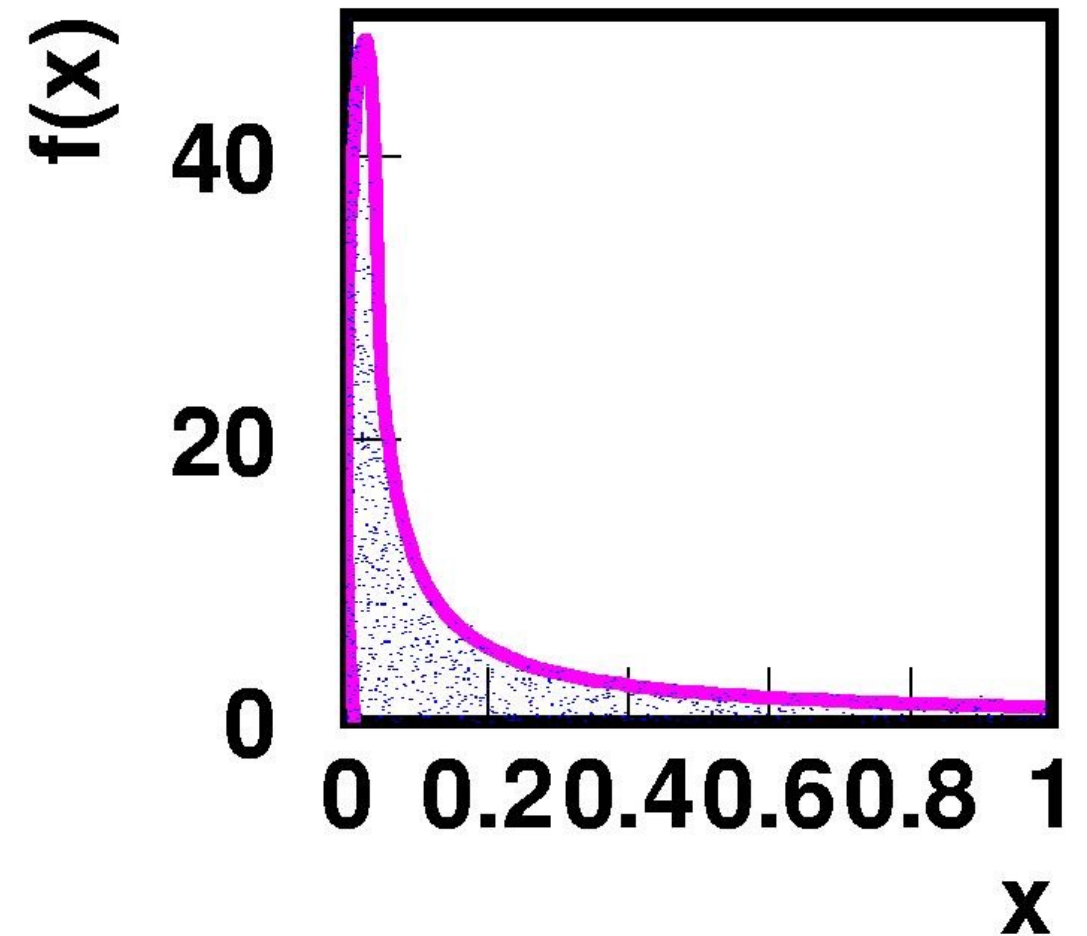
Generating distributions: Hit & Miss

MC for function $f(x)$:
get random number:
R1 in $(0,1)$ and R2 in $(0,1)$
calculate $x = R1$
reject event if: $f_x < f_{\max} R2$



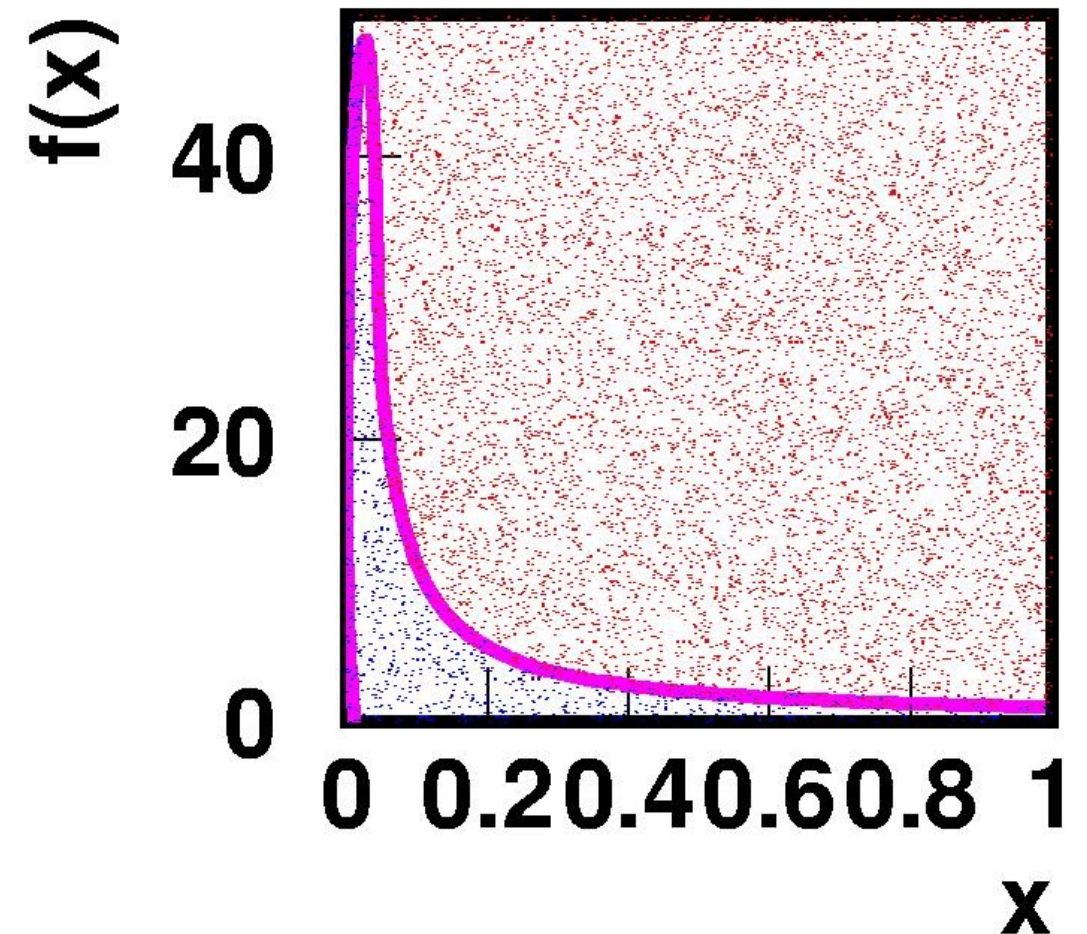
Generating distributions: Hit & Miss

MC for function $f(x)$:
get random number:
R1 in $(0,1)$ and R2 in $(0,1)$
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Generating distributions: Hit & Miss

MC for function $f(x)$:
get random number:
 $R1$ in $(0,1)$ and $R2$ in $(0,1)$
calculate $x = R1$
reject event if: $f_x < f_{\max} R2$

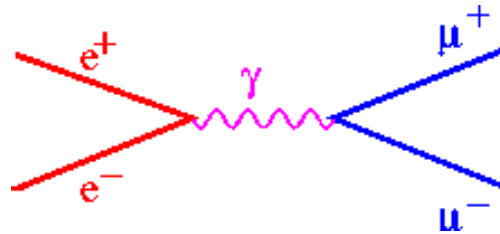


Works always:

→ but can be very inefficient

Constructing a MC for $e^+ e^- \rightarrow X$

- process: $e^+ e^- \rightarrow \mu^+ \mu^-$



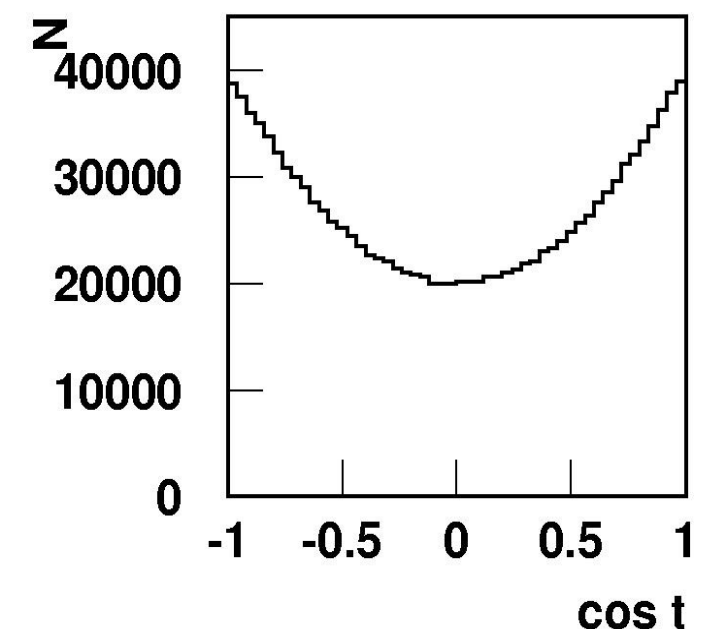
$$\frac{d\sigma}{d\cos\theta d\phi} = \frac{\alpha_{em}^2}{4s} (1 + \cos^2\theta)$$

- goal: generate 4-momenta of μ 's, need cm energy s , $\cos\theta$, ϕ

random number $R1(0,1)$: $\phi = 2\pi R1$
random number $R2(0,1)$: $\cos\theta = -1 + 2R2$

for every $R1, R2$ use weight with
repeat many times

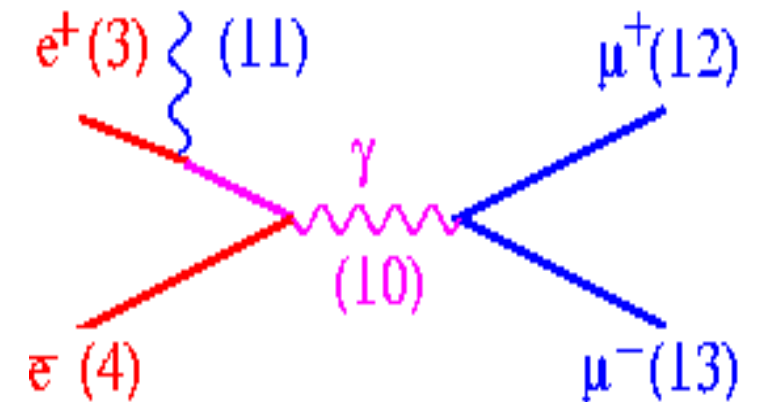
$$\frac{d\sigma}{d\cos\theta d\phi}$$



Example event for $e^+e^- \rightarrow X$

- example from PYTHIA: Event listing

I	particle/jet	KS	KF	orig	p_x	p_y	p_z	E	m
1	!e+	21	-11	0	0.000	0.000	30.000	30.000	0.001
2	!e-	21	11	0	0.000	0.000	-30.000	30.000	0.001
=====									
3	!e+	21	-11	1	0.000	0.000	30.000	30.000	0.000
4	!e-	21	11	2	0.000	0.000	-30.000	30.000	0.000
5	!e+	21	-11	3	0.143	0.040	26.460	26.460	0.000
6	!e-	21	11	4	0.000	0.000	-29.998	29.998	0.000
7	!Z0!	21	23	0	0.143	0.040	-3.539	56.458	56.347
8	!mu-	21	13	7	-9.510	1.741	24.722	26.546	0.106
9	!mu+	21	-13	7	9.653	-1.700	-28.261	29.913	0.106
=====									
10	(Z0)	11	23	7	0.143	0.040	-3.539	56.458	56.347
11	gamma	1	22	3	-0.143	-0.040	3.539	3.542	0.000
12	mu-	1	13	8	-9.510	1.741	24.722	26.546	0.106
13	mu+	1	-13	9	9.653	-1.700	-28.261	29.913	0.106
=====									
	sum:		0.00		0.000	0.000	0.000	60.000	60.000



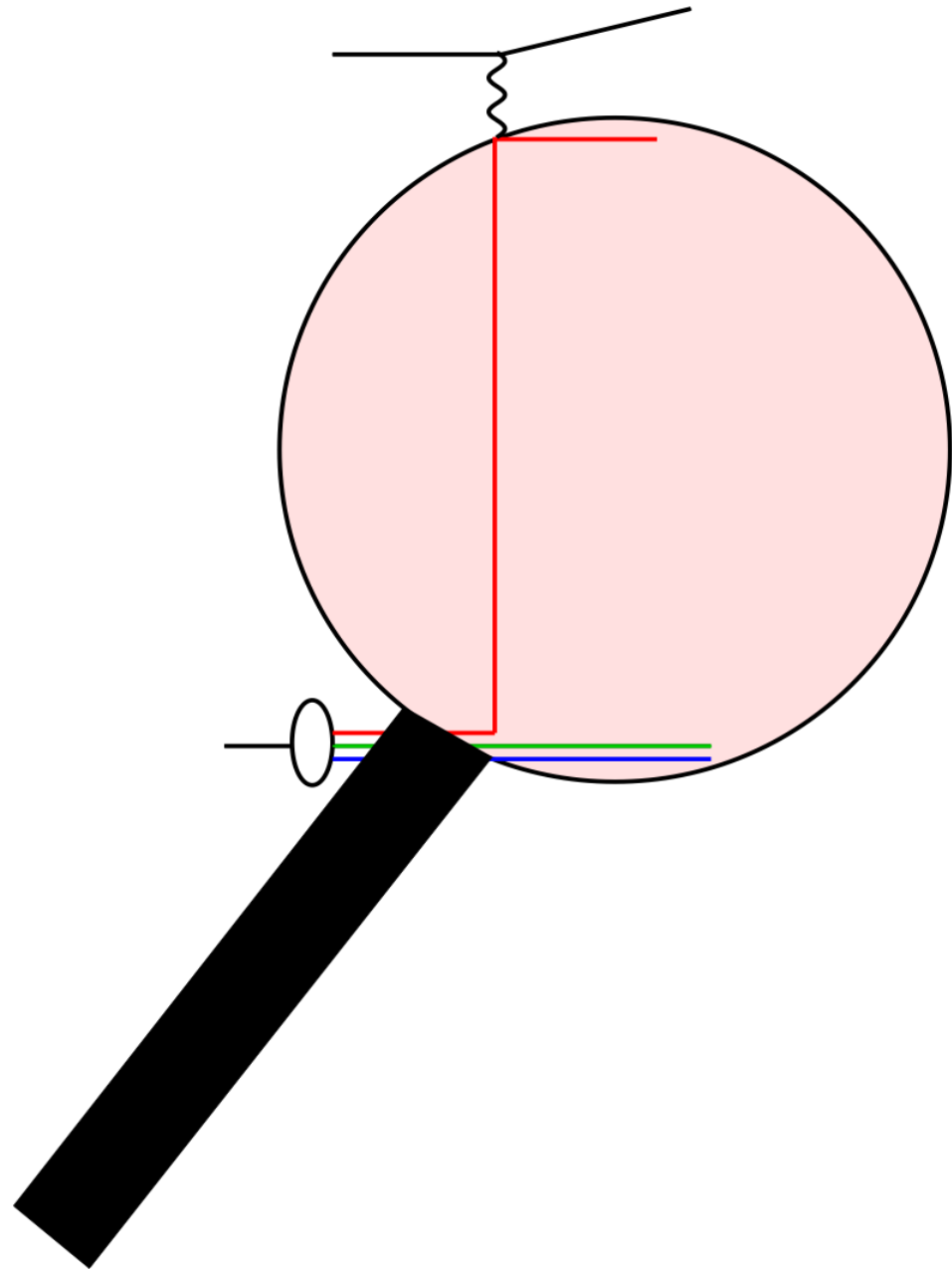
- technicalities/advantages
- can work in any frame
- Lorentz-boost 4-vectors back and forth
- can calculate any kinematic variable
- x-section:
 - fill histogram in any variable, no Jacobians !

From e^+e^- to hadron scattering ...

... add a hadron in the initial state:

ep scattering

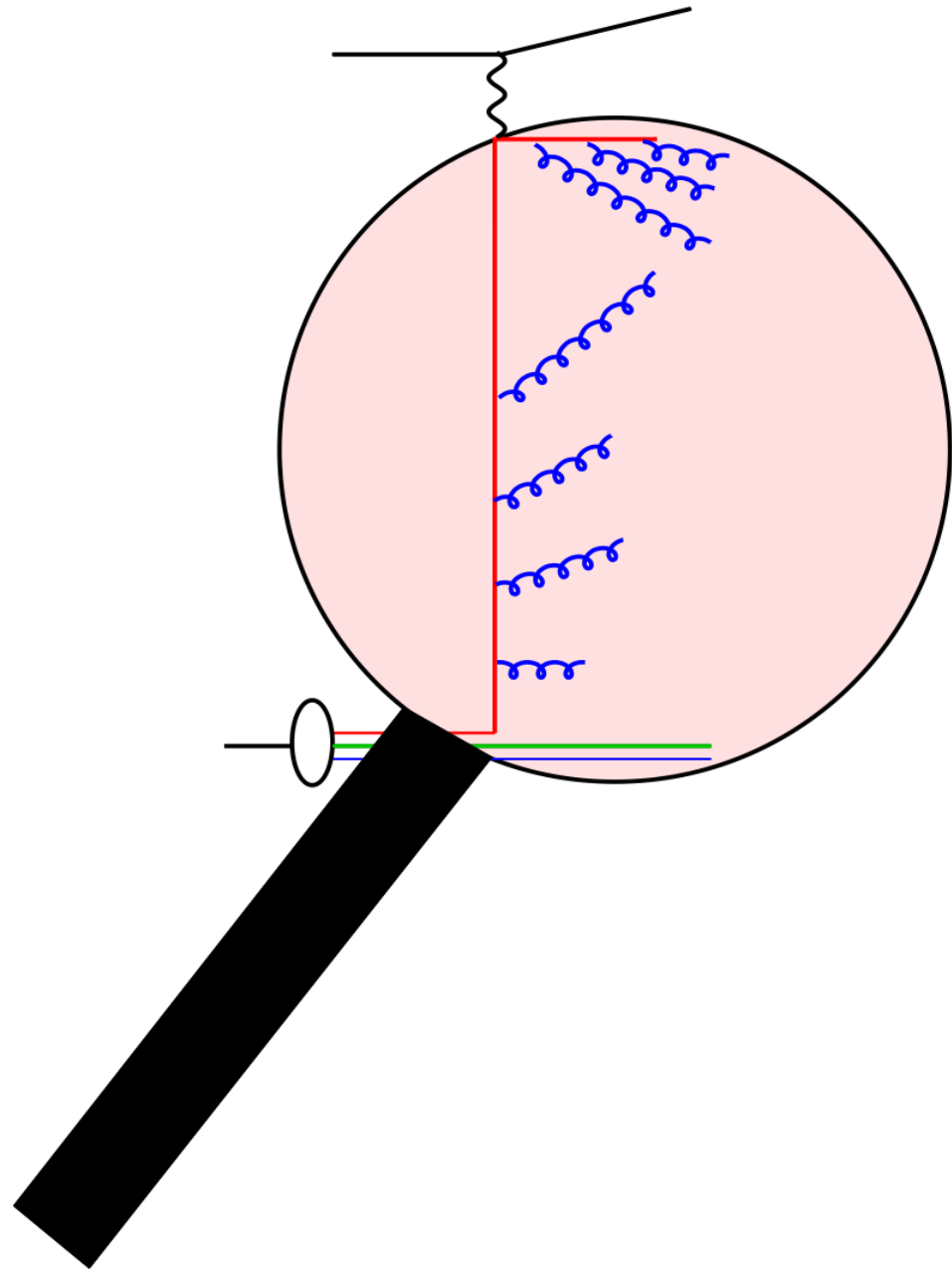
A proton in the initial state



- Deep Inelastic Scattering is a incoherent sum of $e^+ q \rightarrow e + q$
- only 50 % of p momentum carried by quarks
- need a large gluon component
- partonic part convoluted with parton density function $f_i(x)$

$$\sigma(e^+ p \rightarrow e^+ X) = \sum_i f_i(x, \quad) \sigma(e^+ q_i \rightarrow e^+ q_i)$$

A proton in the initial state



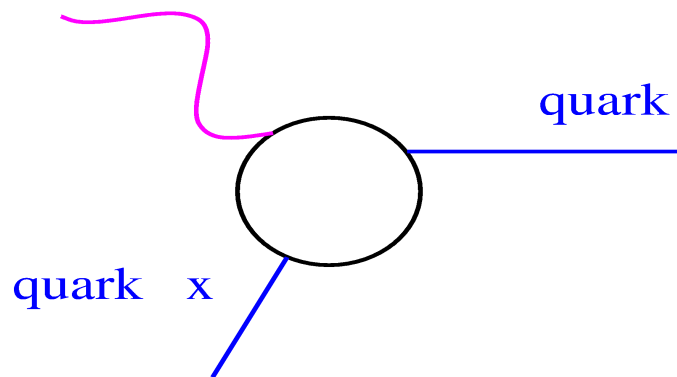
- Deep Inelastic Scattering is a incoherent sum of
$$e^+ q \rightarrow e + q$$
- only 50 % of p momentum carried by quarks
- need a large gluon component
- partonic part convoluted with parton density function $f_i(x)$
- BUT we know, PDF depends on resolution scale Q^2

$$\sigma(e^+ p \rightarrow e^+ X) = \sum_i f_i(x, Q^2) \sigma(e^+ q_i \rightarrow e^+ q_i)$$

$F_2(x, Q^2)$: DGLAP evolution equation

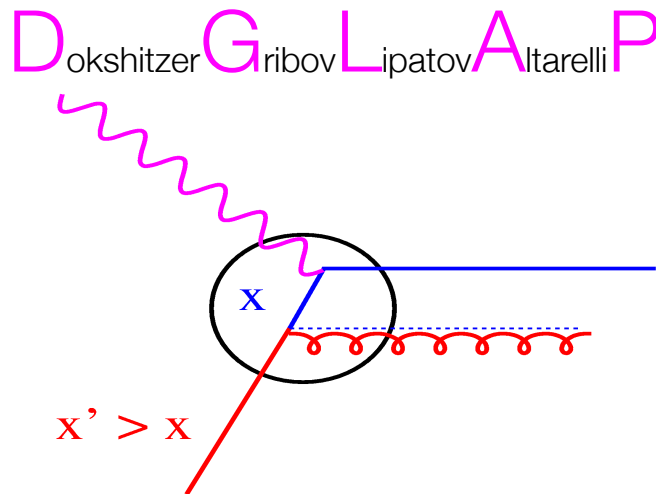
- QPM: F_2 is independent of Q^2

- Q^2 dependence of structure function: D_{okshitzer} G_{ribov} L_{ipatov} A_{ltarelli} P_{arisi}



Q^2 small

small resolution power



Q^2 small

better resolution power

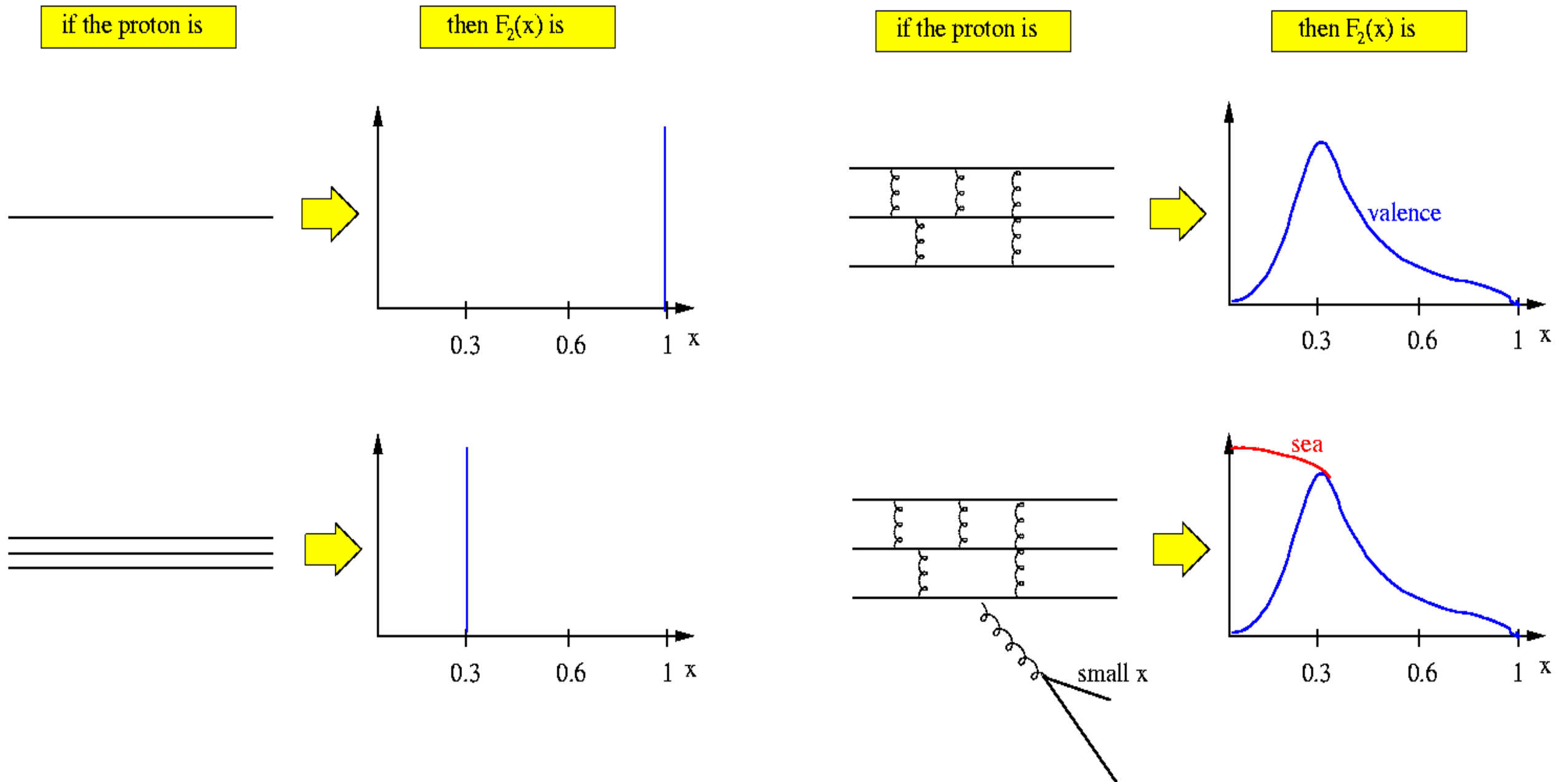
→ Probability to find parton at small x increases with Q^2

$$F_2 = \left| \begin{array}{c} \text{Diagram 1} \\ \text{OPM} \end{array} \right|^2 + \left| \begin{array}{c} \text{Diagram 2} \\ \text{QCDC} \end{array} \right|^2 + \left| \begin{array}{c} \text{Diagram 3} \\ \text{BGF} \end{array} \right|^2$$

→ Test of theory: Q^2 evolution of $F_2(x, Q^2)$!!!!!

From Naïve F_2 picture to QCD ...

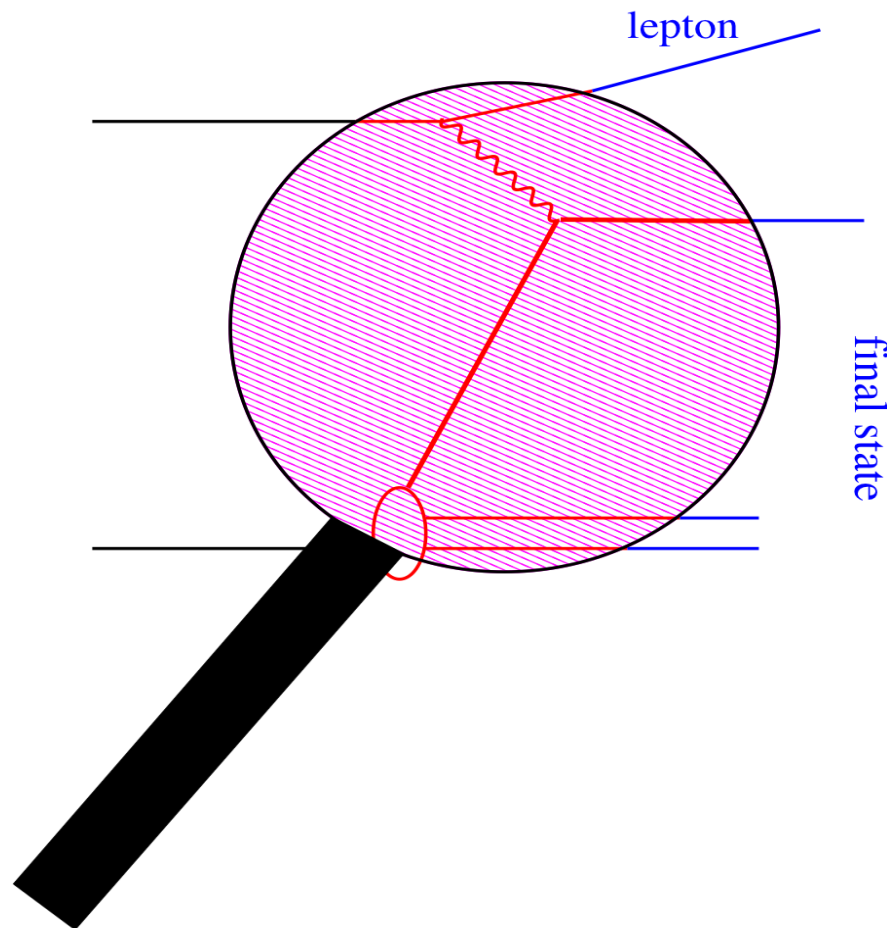
From Halzen & Martin: Quarks & Leptons, p201



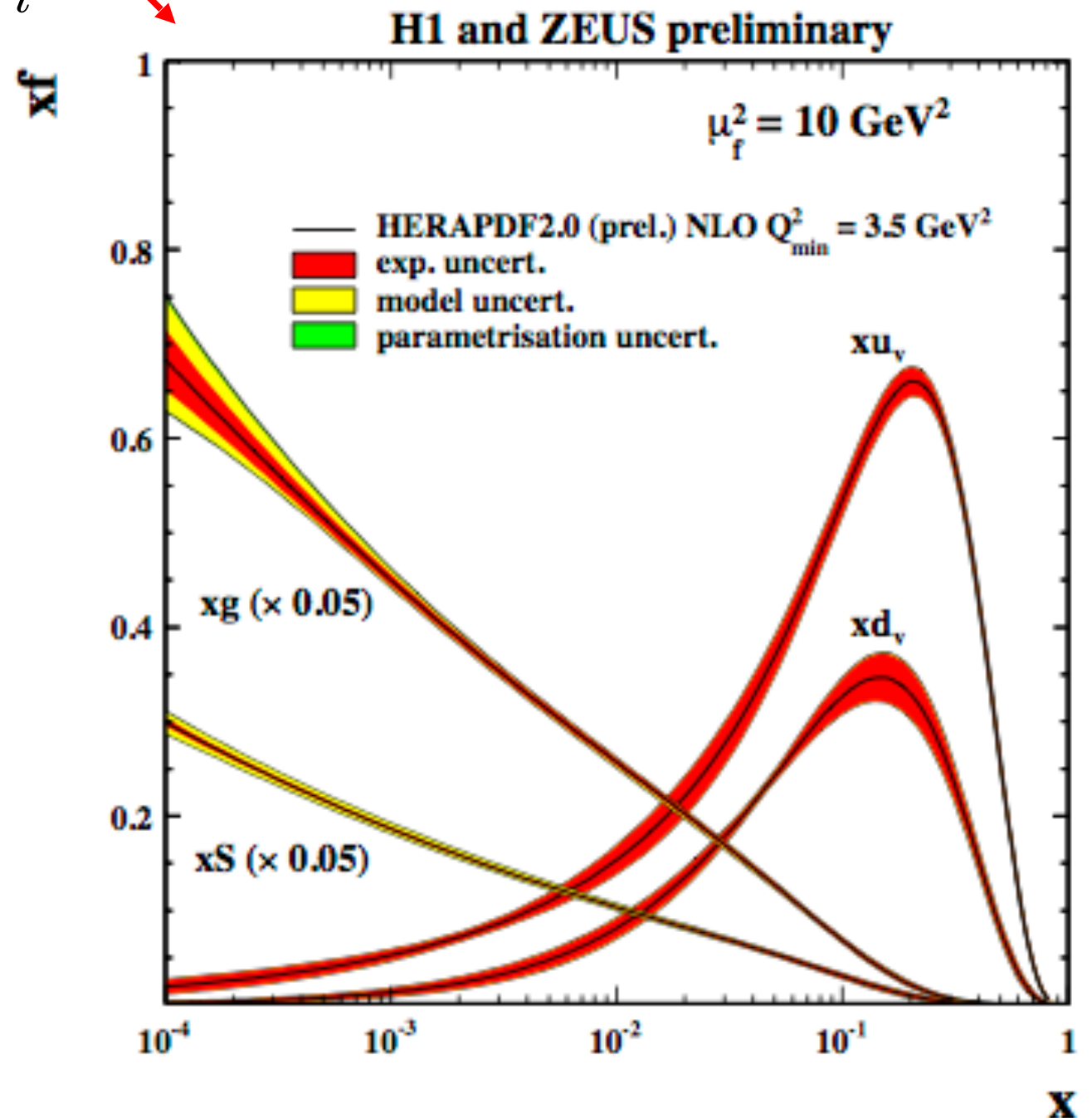
Lepton Hadron scattering

- Deep Inelastic Scattering is an incoherent sum of $e^+ q \rightarrow e + q$

$$\sigma(e^+ p \rightarrow e^+ X) = \sum_i f_i(x, Q^2) \sigma(e^+ q_i \rightarrow e^+ q_i)$$



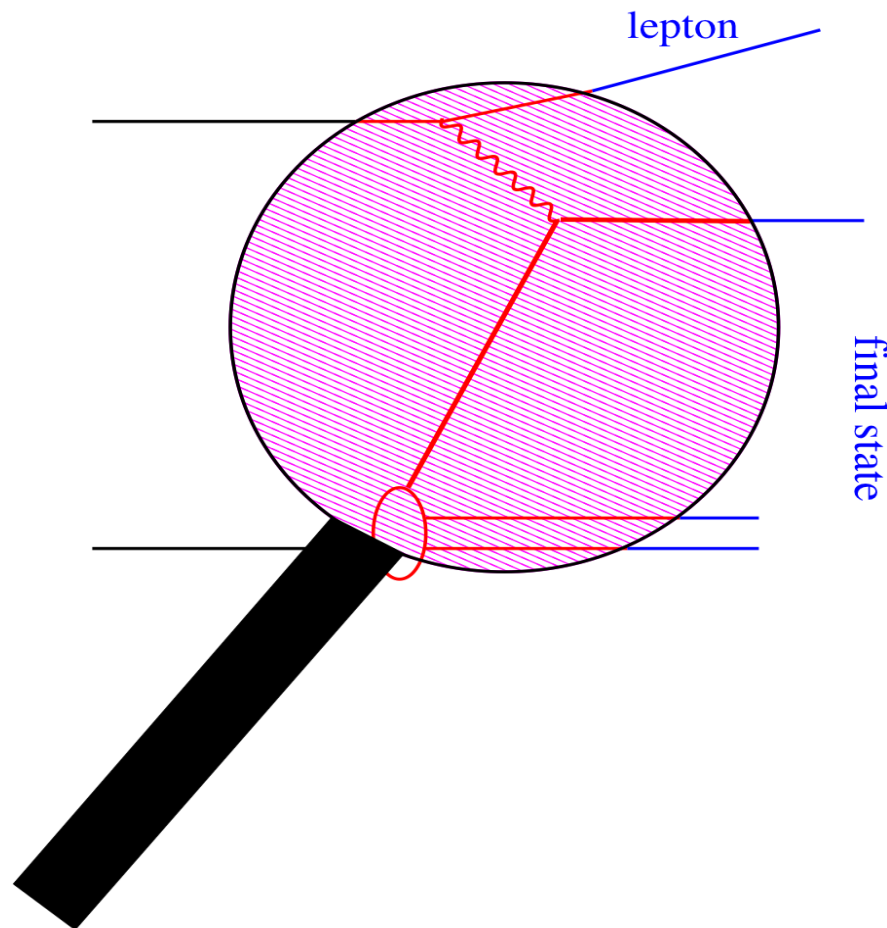
- only 50 % of p momentum carried by quarks
- need a large gluon component



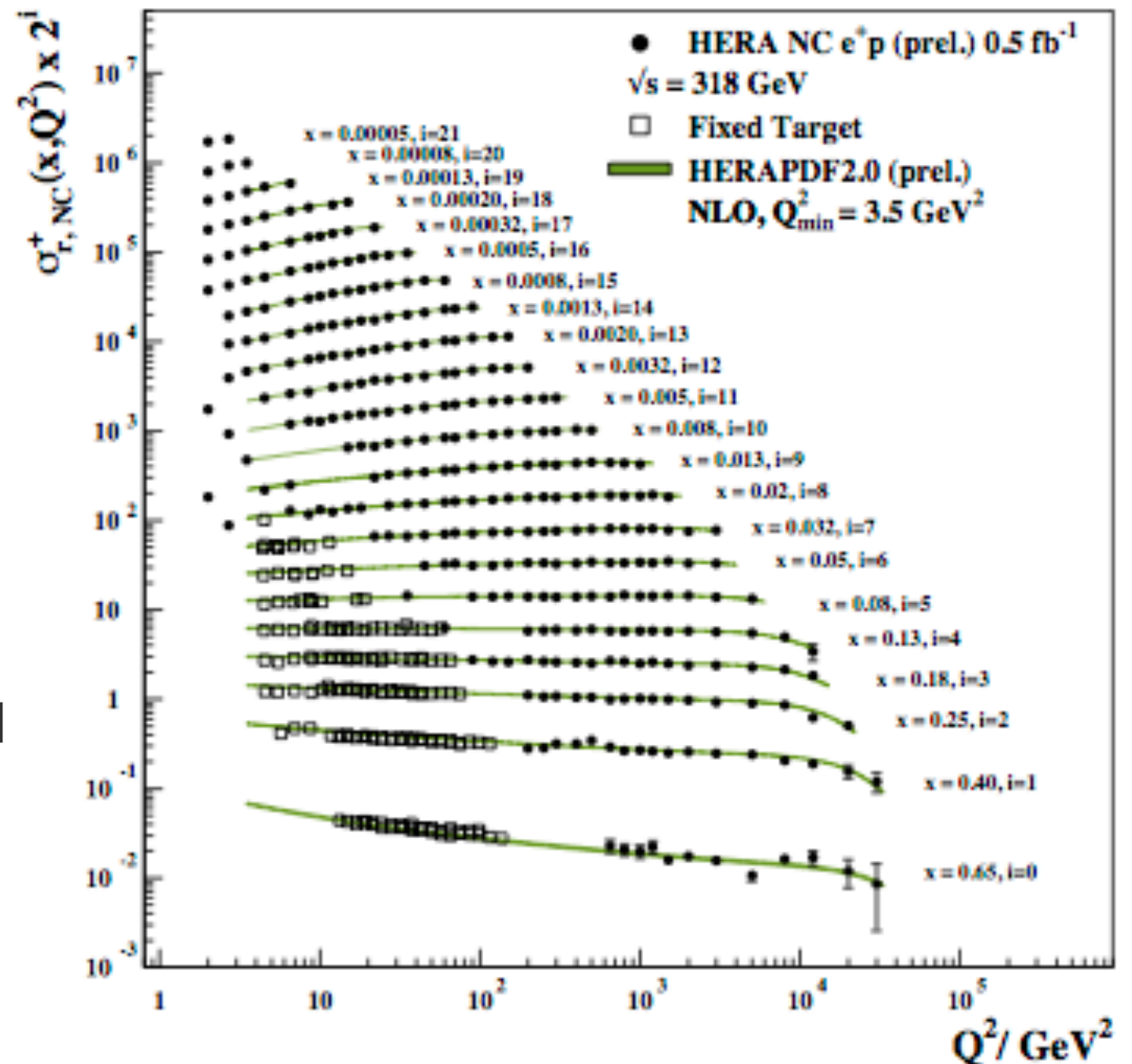
Lepton Hadron scattering

- Deep Inelastic Scattering is a incoherent sum of $e^+ q \rightarrow e + q$

$$\sigma(e^+ p \rightarrow e^+ X) = \sum_i f_i(x, Q^2) \sigma(e^+ q_i \rightarrow e^+ q_i)$$



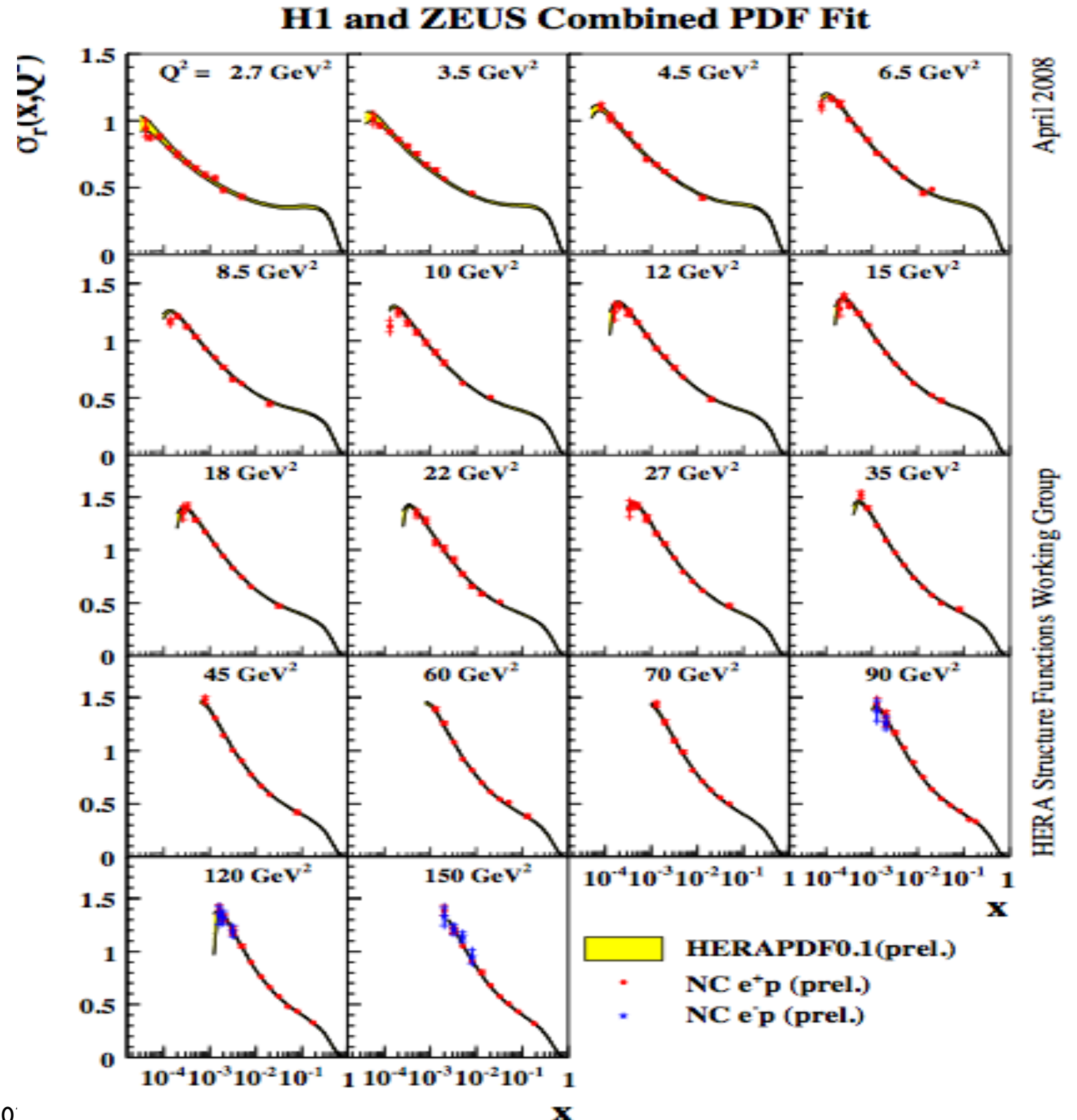
H1 and ZEUS preliminary



- only 50 % of p momentum carried by quarks
- need a large gluon component

Inelastic Scattering: main results

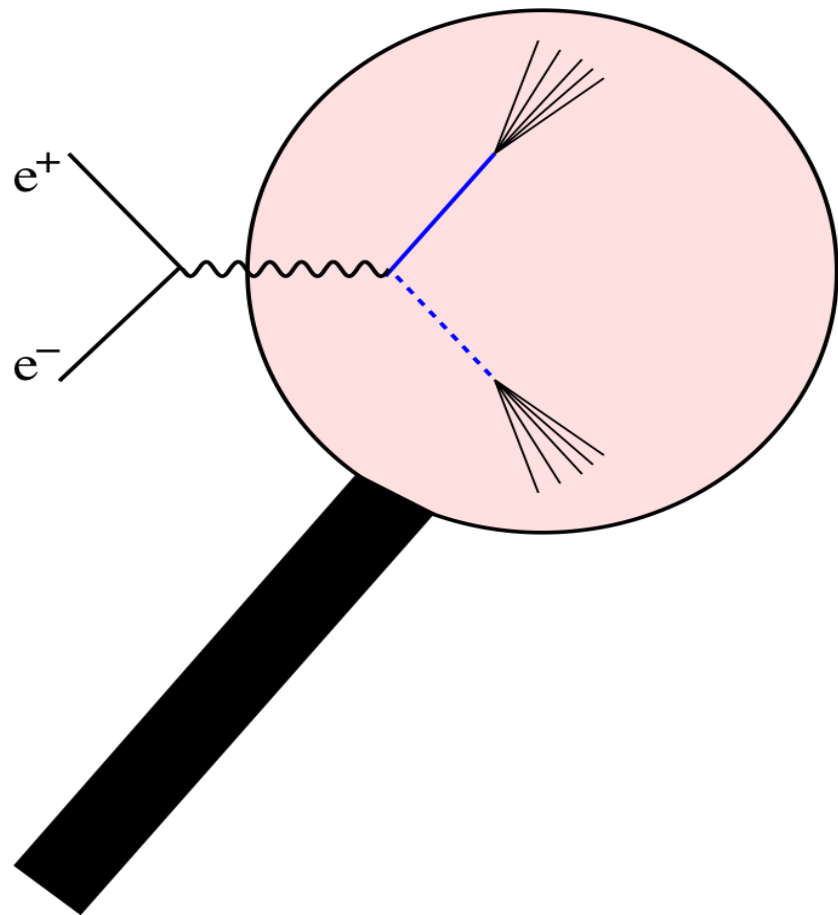
- F_2 scaling at large x
- $\sim 50\%$ gluons
- F_2 rise at small x
 - How can rising F_2 be understood?
 - Does rise continue forever?
 - What limits F_2 ?



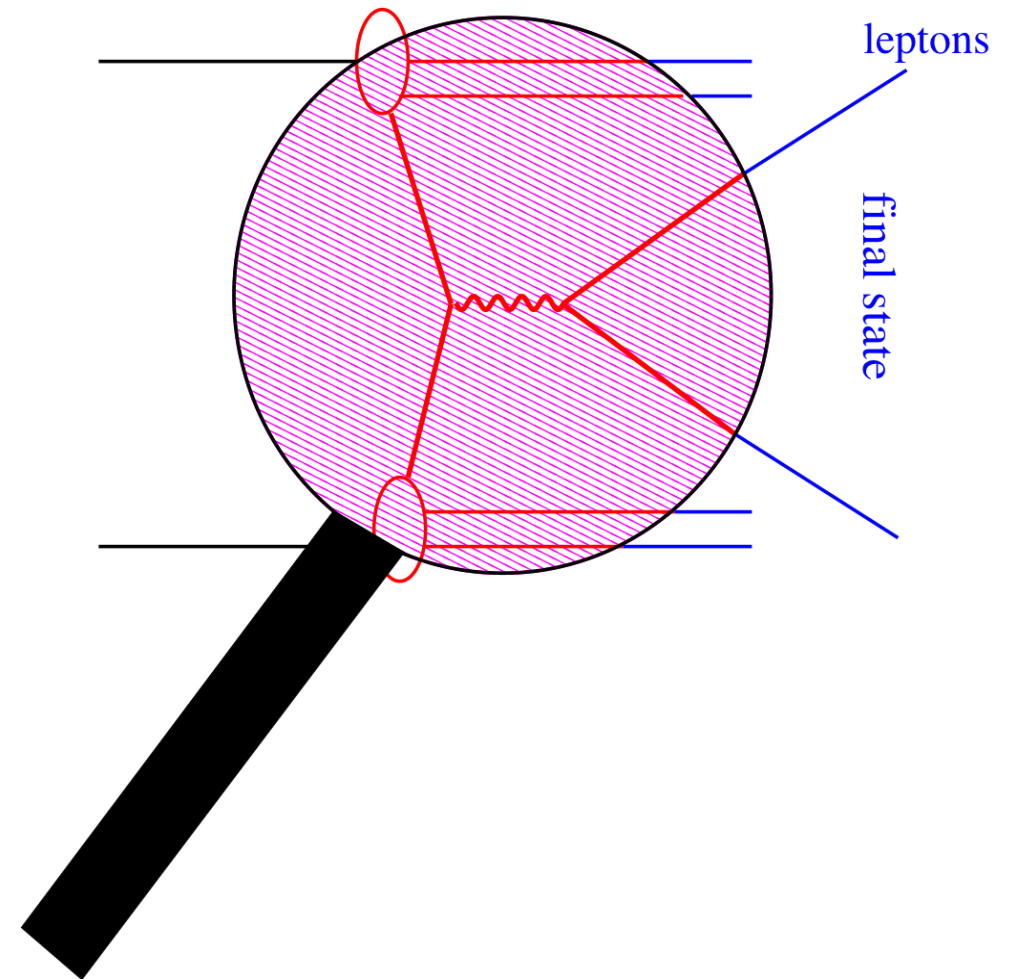
From ep scattering ...

... to pp scattering

Rotating the diagrams



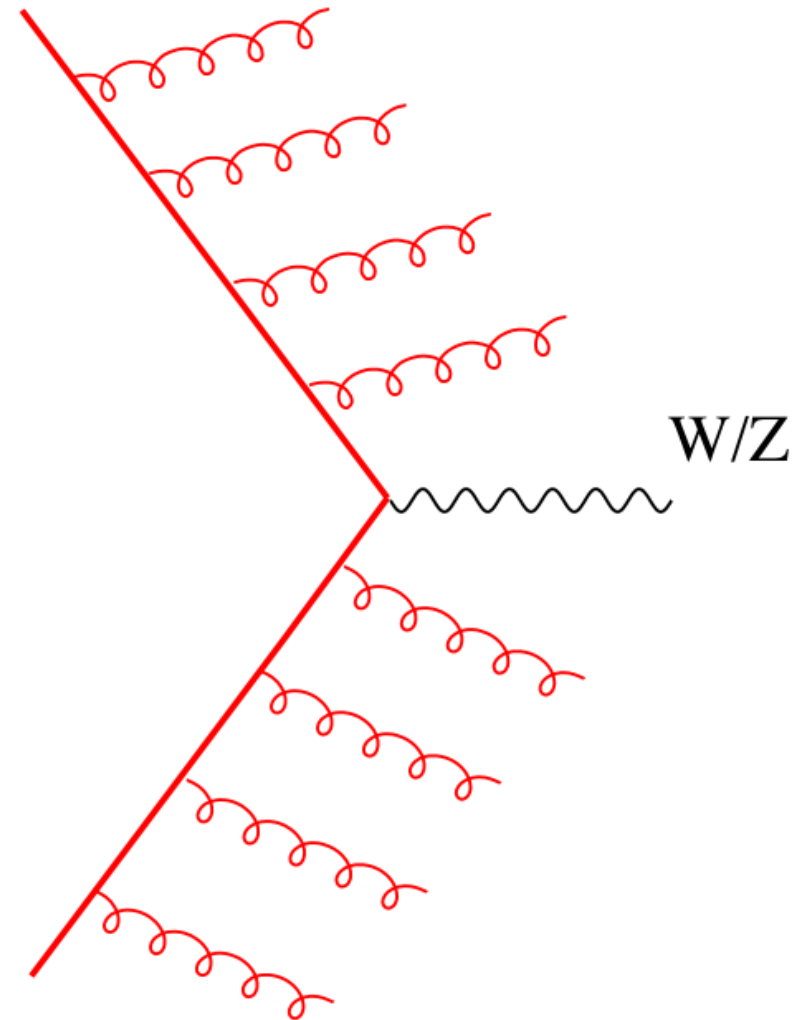
$$\sigma(e^+e^- \rightarrow q\bar{q}) = 3 \frac{4\pi\alpha^2}{3s} e_q^2$$



$$\sigma(q\bar{q} \rightarrow l^+l^-) = \frac{4\pi\alpha^2}{3 \times 3s} e_q^2$$

Monte Carlo approach

- simulate explicitly parton radiation with evolution of parton densities
- advantage to include properly energy momentum conservation in each step
- perform resummation numerically



Jet production in pp

- x-section (i.e. for light and heavy quarks ($t\bar{t}$) production)

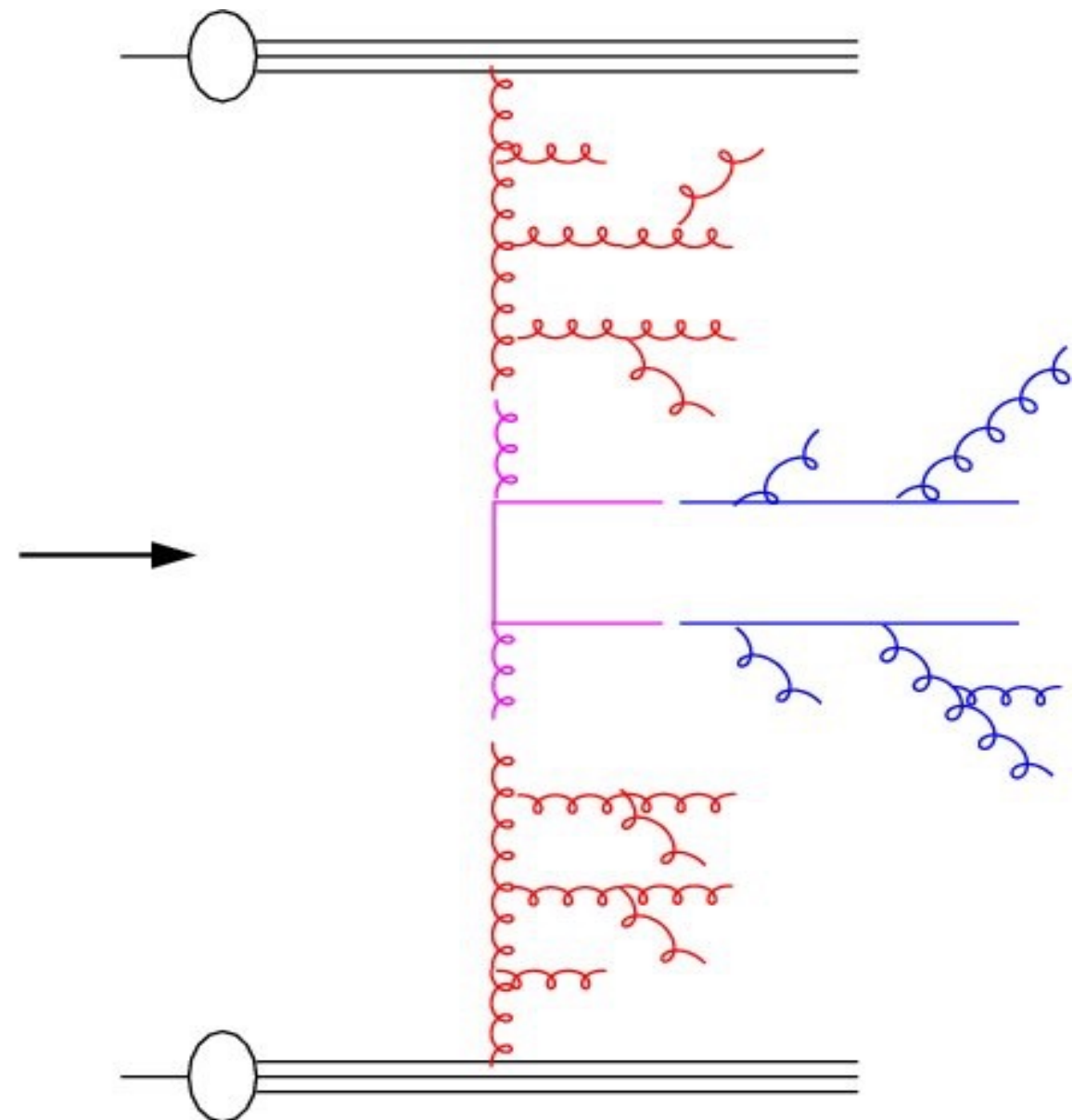
$$\sigma(\text{pp} \rightarrow \text{q}\bar{\text{q}}\text{X}) = \int \frac{dx_1}{x_1} \frac{dx_2}{x_2} x_1 G(x_1, \bar{q}) x_2 G(x_2, \bar{q}) \times \hat{\sigma}(\hat{s}, \bar{q})$$

- with gluon densities

$$xG(x, \bar{q})$$

- hard x-section:

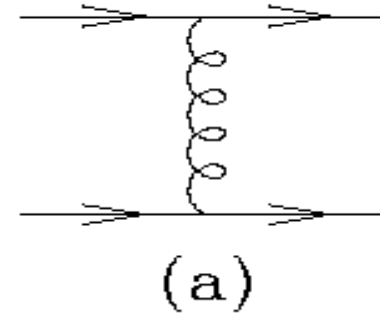
$$\frac{d\sigma}{dt} = \frac{1}{64\hat{s}^2} |M_{ij}|^2$$



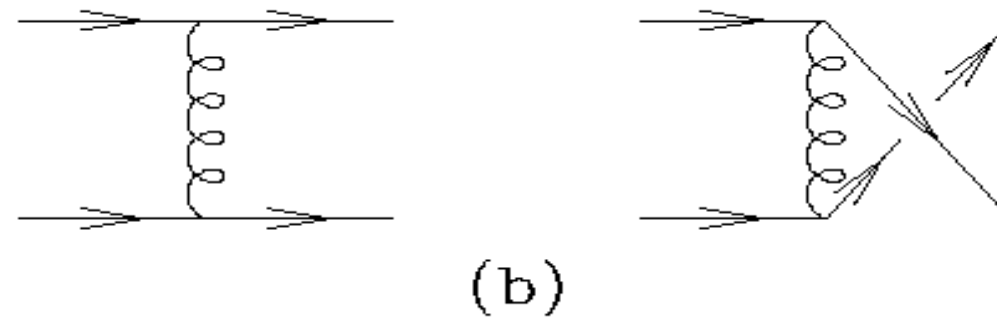
Lowest Order Diagrams

Ellis, Stirling, Webber
QCD & collider physics p248

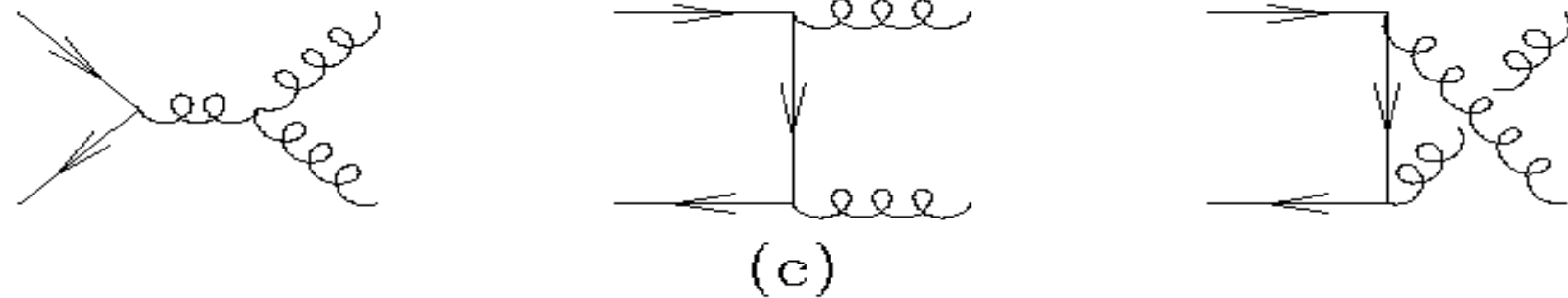
$$qq' \rightarrow qq'$$



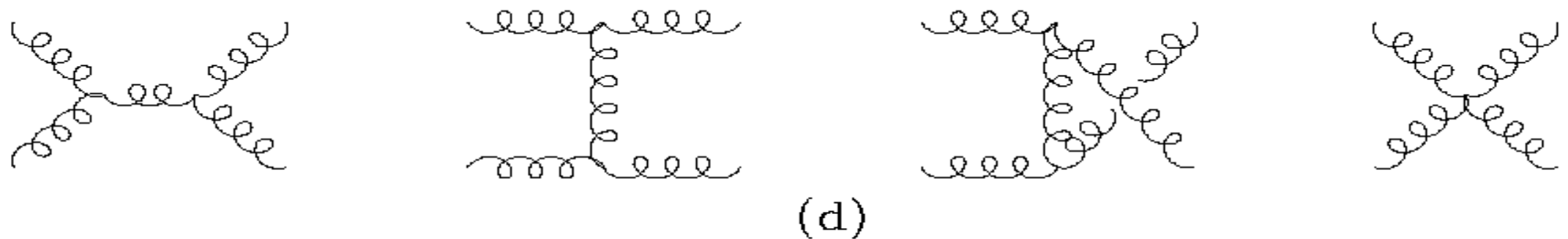
$$qq \rightarrow qq$$



$$q\bar{q} \rightarrow gg$$



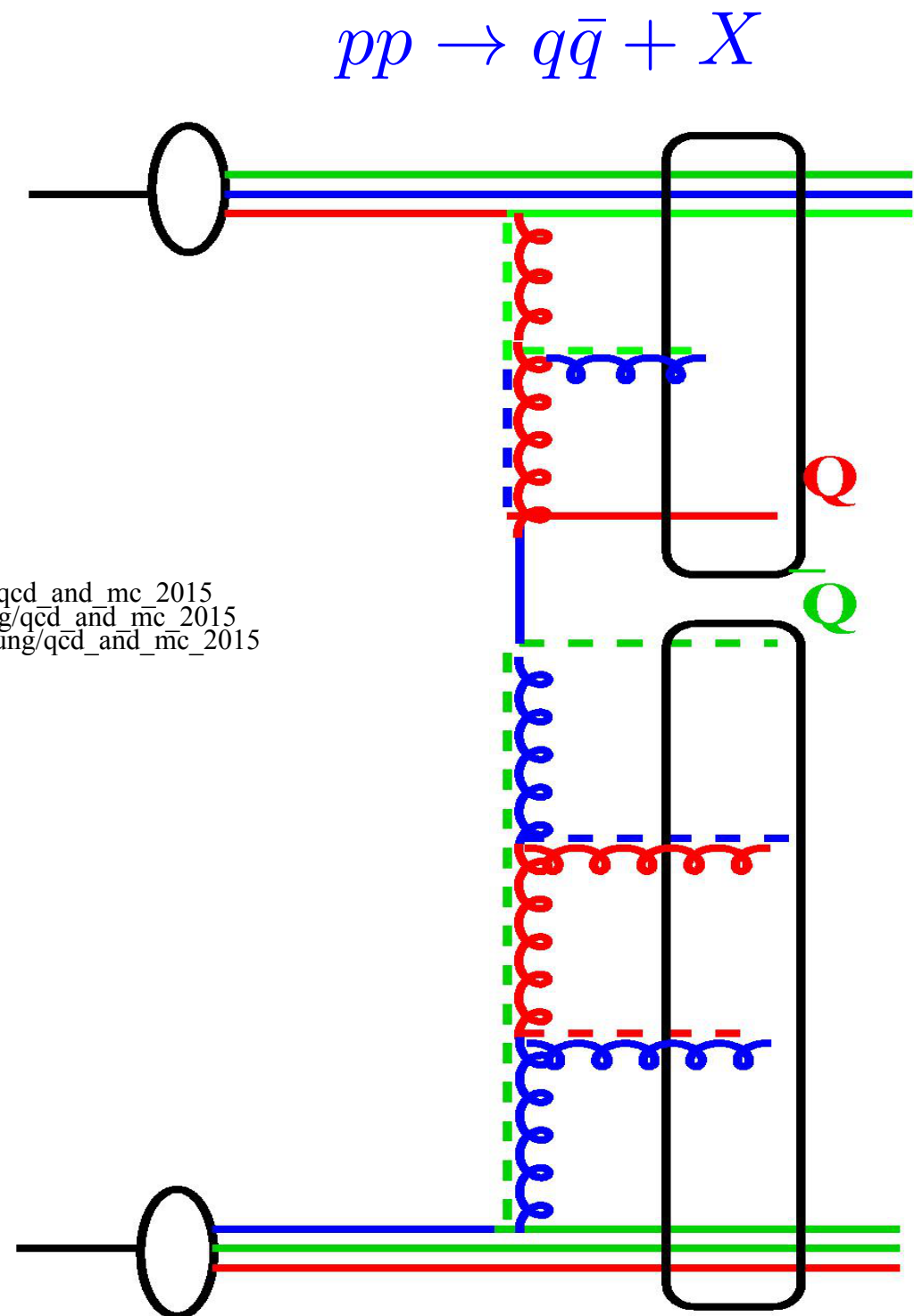
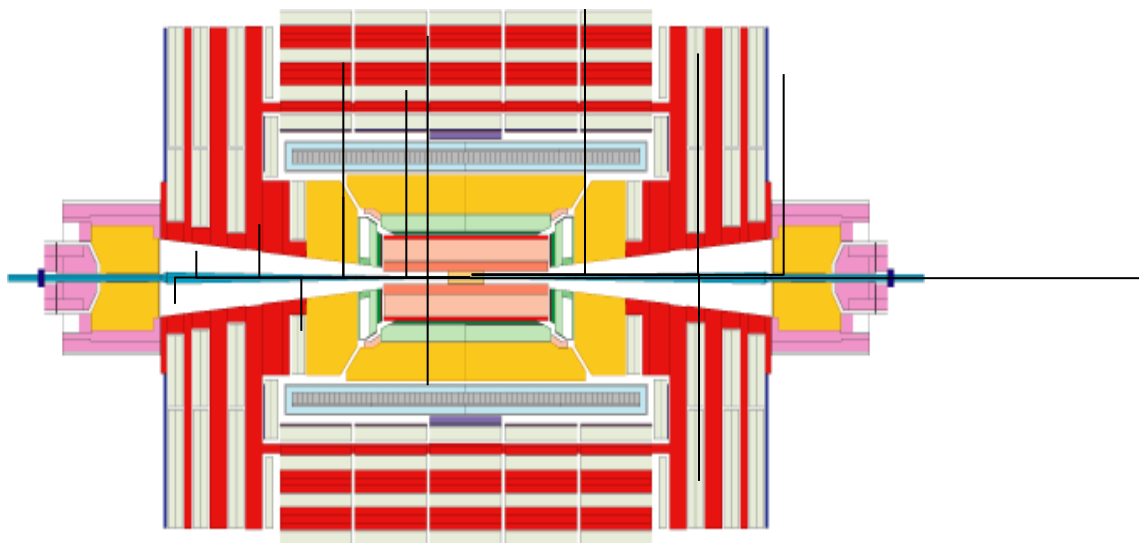
$$gg \rightarrow gg$$



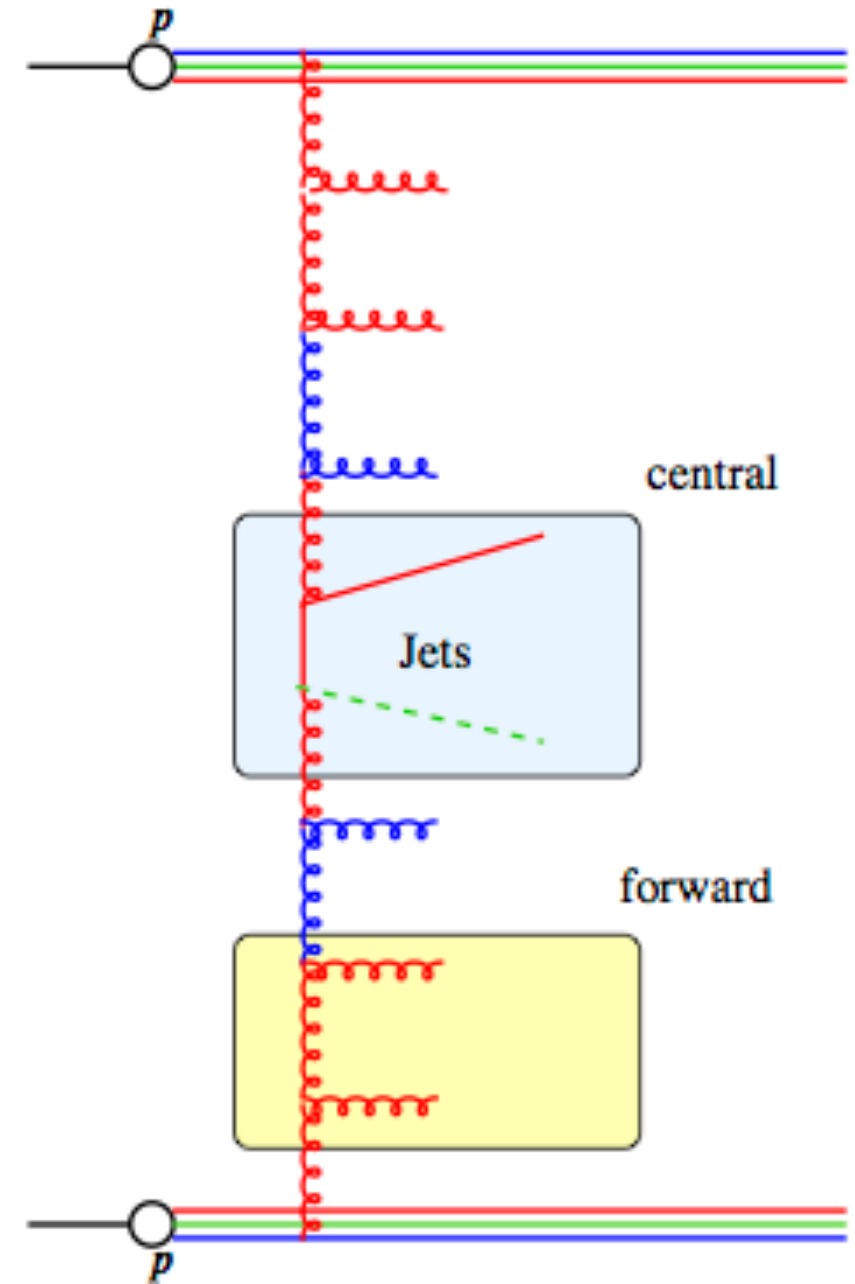
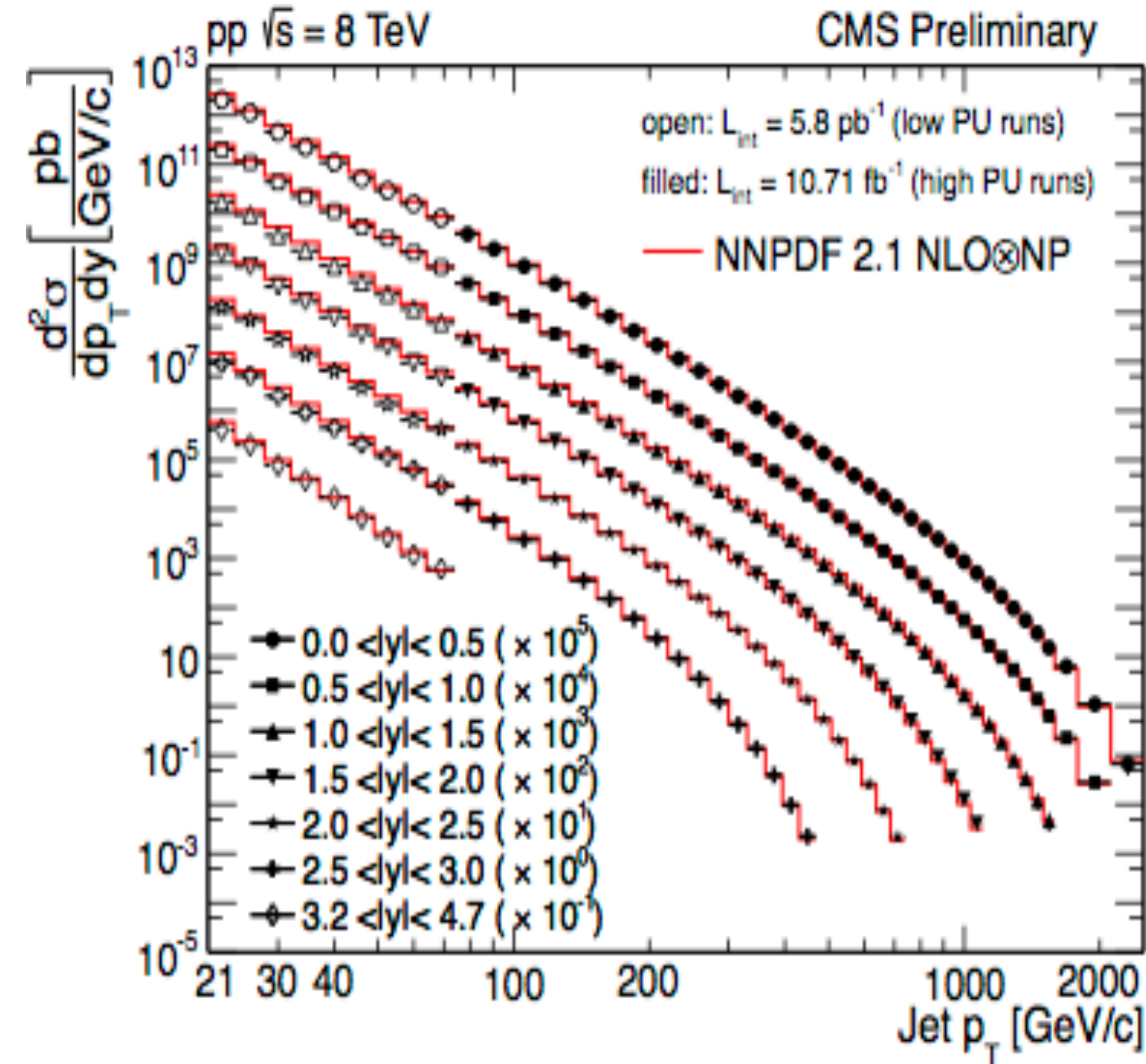
Color Flow in pp

- quarks carry color
- anti-quarks carry anti-color
- gluons carry color – anti-color
 - connect to color singlet systems
 - watch out pp or $p\bar{p}$

http://www.desy.de/~jung/qcd_and_mc_2015
http://www.desy.de/~jung/qcd_and_mc_2015
http://www.desy.de/~jung/qcd_and_mc_2015



Jet production at the LHC

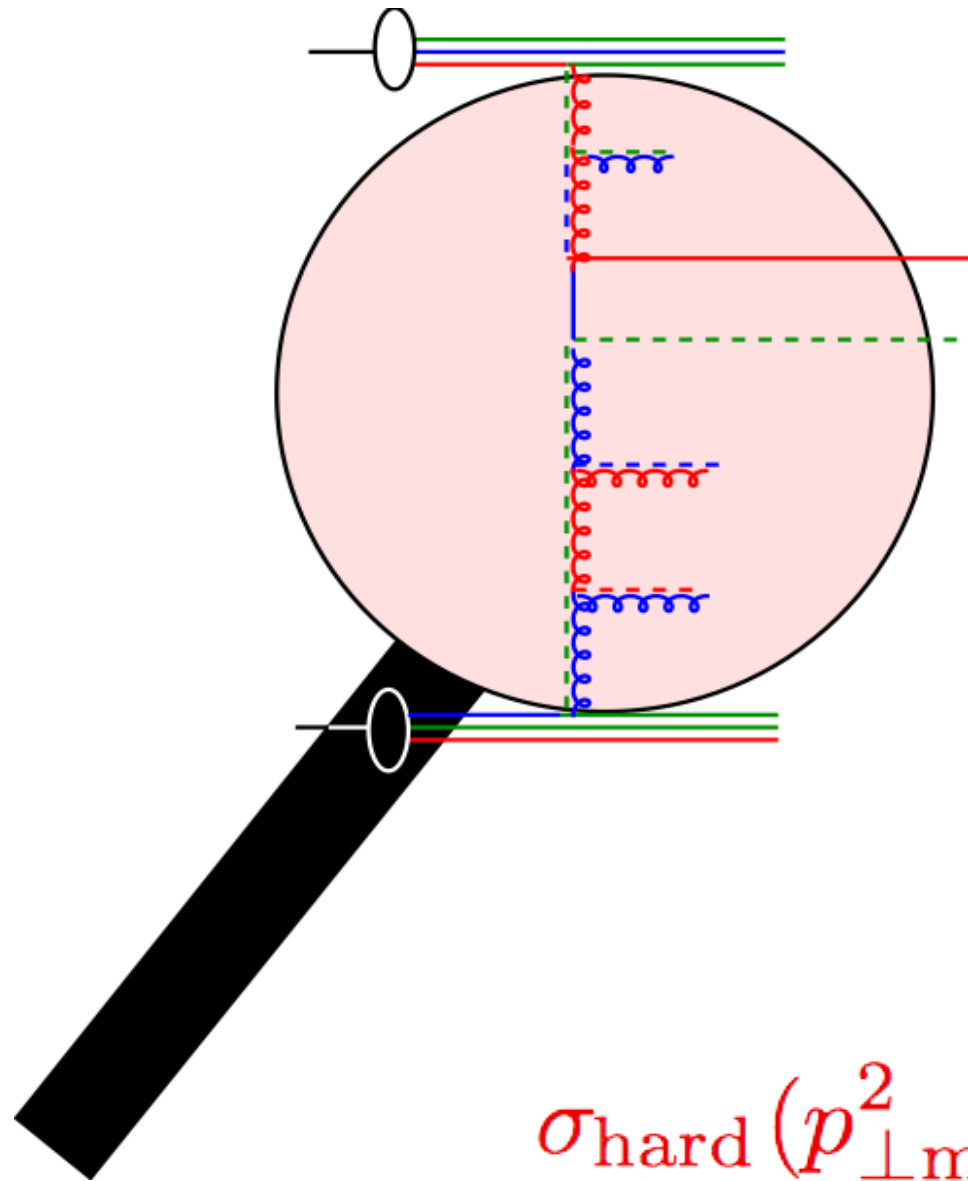


What happens at highest energies ?

High energy behavior of xsections

Partonic Cross sections

- Cross section $\sigma(p_1 + p_2 \rightarrow j_1 + j_2 + X) = f(x_1, \mu^2) \otimes \hat{\sigma}(x_1 p_1 + x_2 p_2 \rightarrow j_1 + j_2) \otimes f(x_2, \mu^2)$



- partonic cross section diverges with p_{\perp}
- calculate x-section as function of $p_{\perp, \min}$

$$\sigma_{\text{hard}}(p_{\perp, \min}^2) = \int_{p_{\perp, \min}^2} \frac{d\sigma_{\text{hard}}(p_{\perp}^2)}{dp_{\perp}^2} dp_{\perp}^2$$

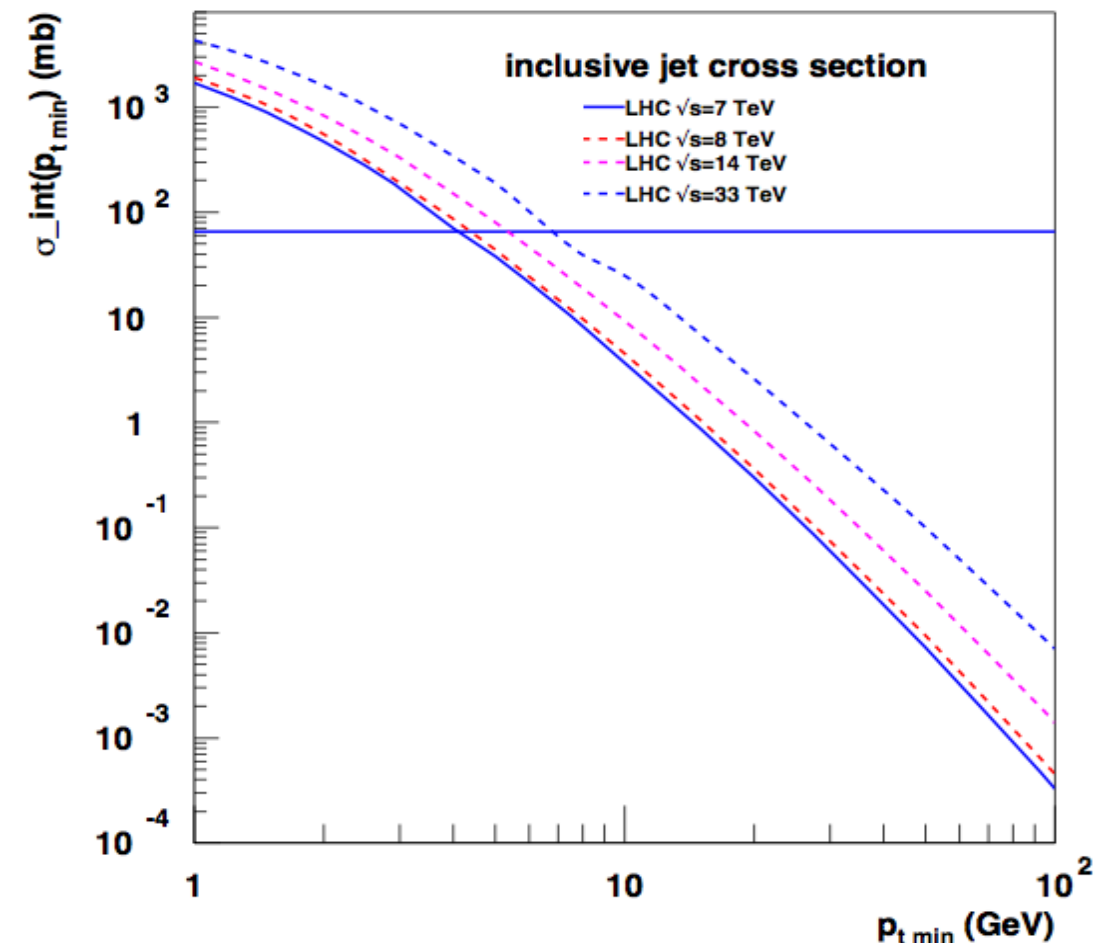
Partonic Cross Section

- Basic partonic perturbative cross section

$$\sigma_{\text{hard}}(p_{\perp \text{min}}^2) = \int_{p_{\perp \text{min}}^2} \frac{d\sigma_{\text{hard}}(p_{\perp}^2)}{dp_{\perp}^2} dp_{\perp}^2$$

diverges faster than $1/p_{\perp, \text{min}}^2$ as $p_{\perp, \text{min}}$ and exceeds eventually total inelastic (non-diffractive) cross section

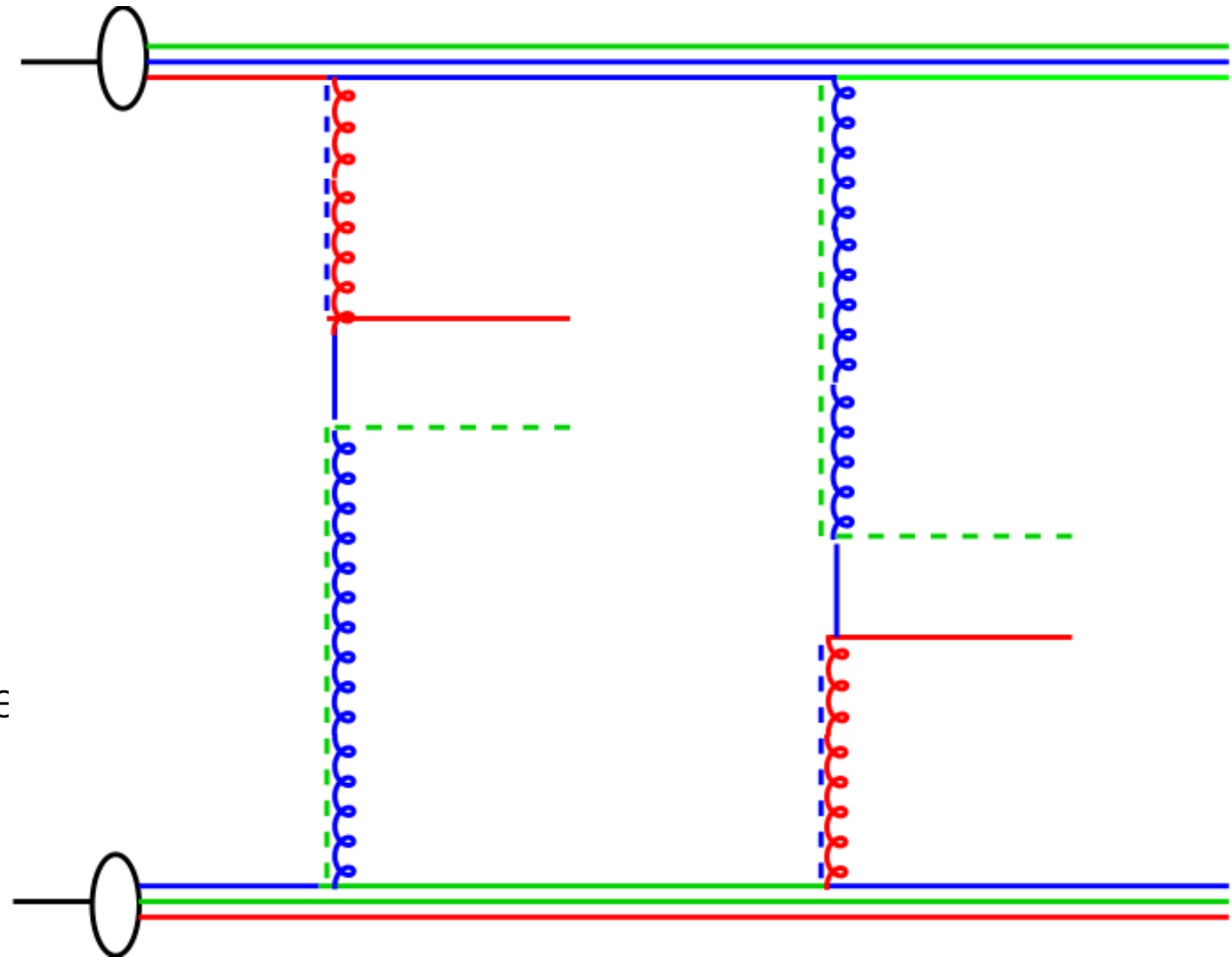
- Interaction x-section exceeds total xsection
- happens well above λ_{QCD}
- in perturbative region



Models for Multi-Parton Interaction

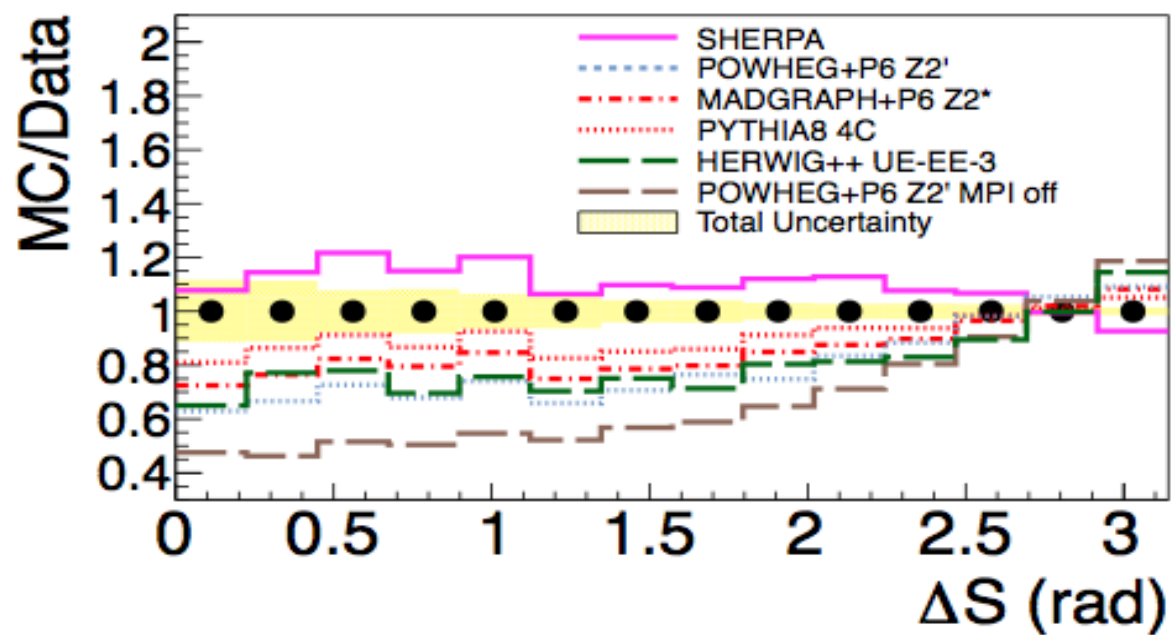
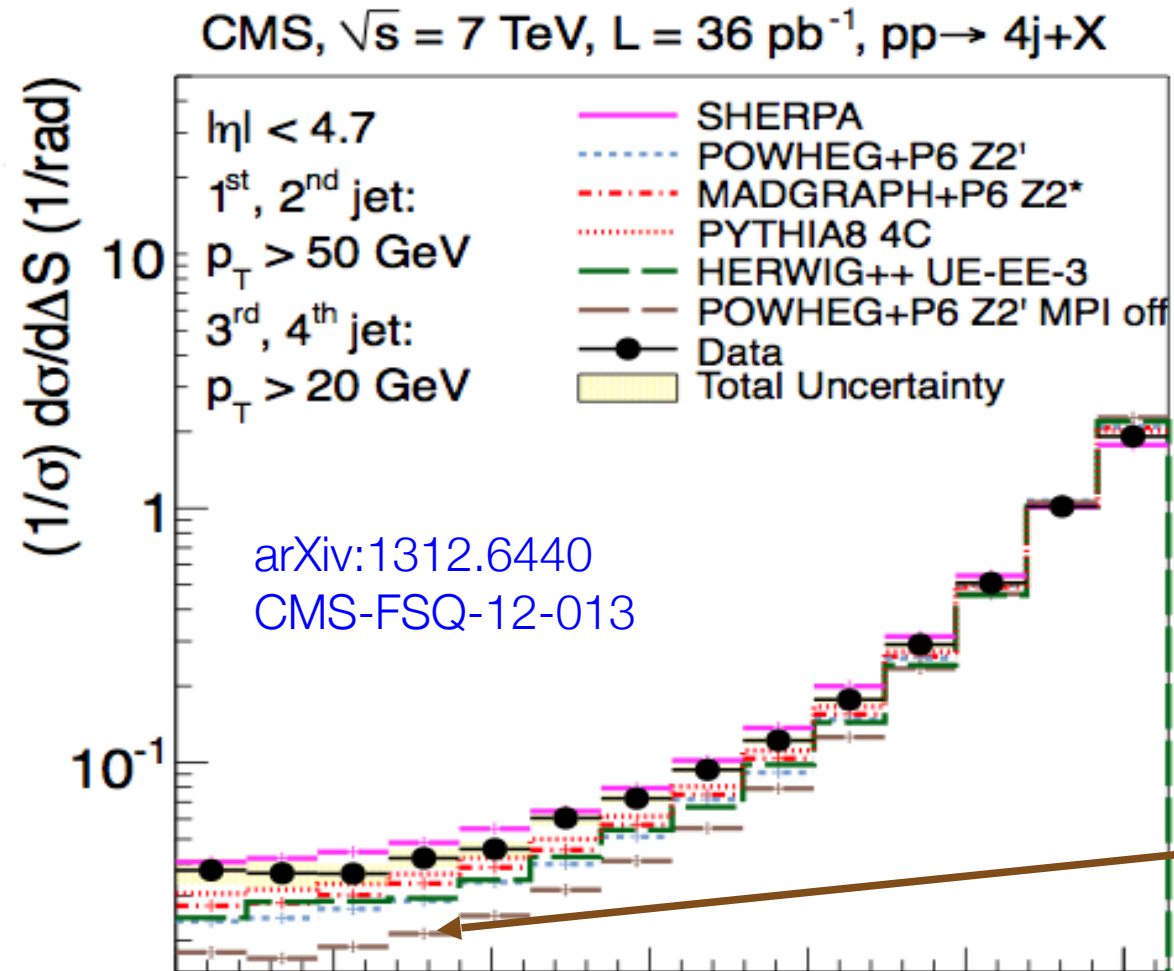
- The very simple model
 - add secondary interactions
 - first model by: T. Sjostrand, M. Ziji PRD (1987) 2019
- order scatterings in p_t
 - use of sudakov form factor
 - result in Poisson distribution of number of scatterings:

with
$$p_r = \frac{\mu^r}{r!} \exp(-\mu)$$

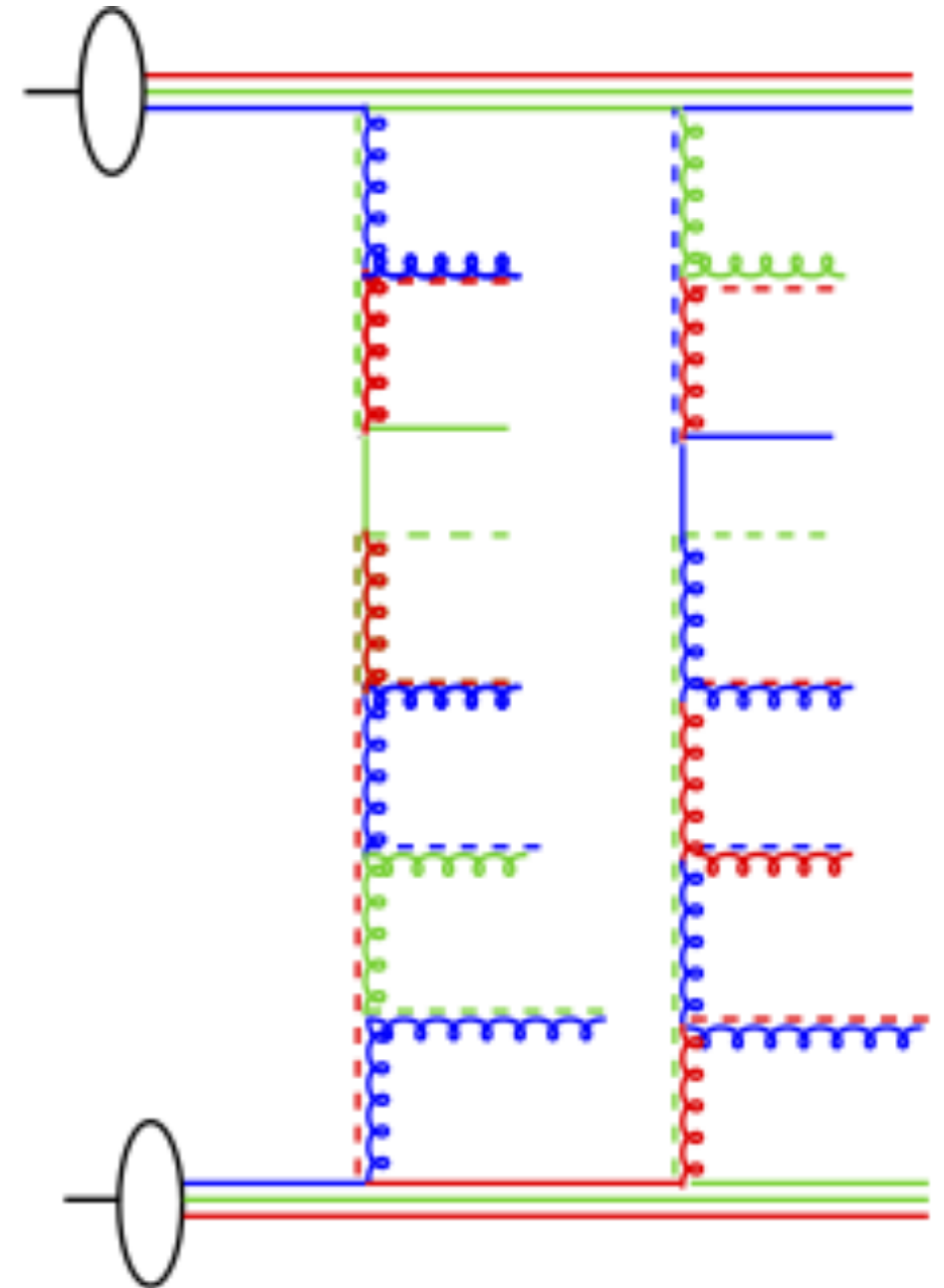


$$\mu = \langle n \rangle = \frac{1}{\sigma_{nd}} \int_{p_{\perp \min}}^{p_{\perp \max}} \frac{d\sigma_{\text{hard}}}{dp'_{\perp}} dp'_{\perp}$$

Multiparton interaction at the LHC



no MPI



Partonic Cross Section

- Basic partonic perturbative cross section

$$\sigma_{\text{hard}}(p_{\perp,\text{min}}^2) = \int_{p_{\perp,\text{min}}^2} \frac{d\sigma_{\text{hard}}(p_{\perp}^2)}{dp_{\perp}^2} dp_{\perp}^2$$

diverges faster than $1/p_{\perp,\text{min}}^2$ as $p_{\perp,\text{min}}$ and exceeds eventually total inelastic (non-diffractive) cross section

- Interaction x-section exceeds total xsection

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