

A novel algorithm for calculating Feynman integrals via differential equations

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Antidifferentiation and the Calculation of Feynman Amplitudes
October 7, 2020



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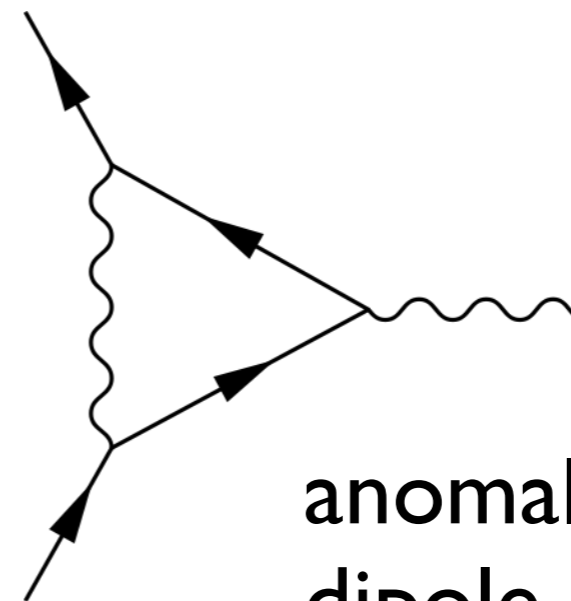
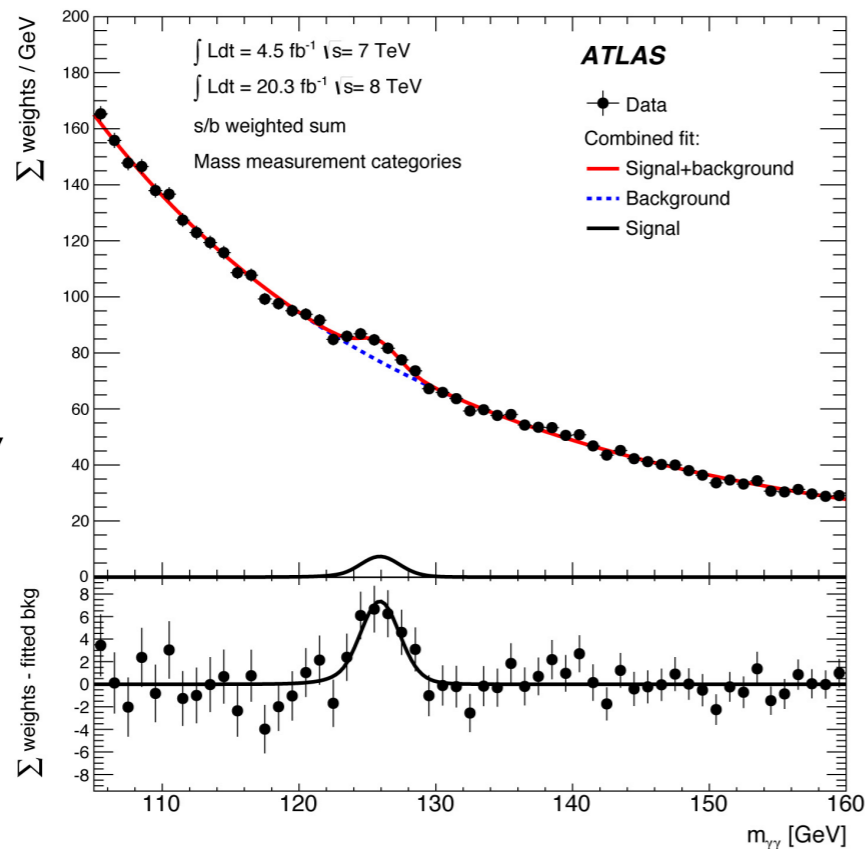


European Research Council
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Feynman integrals are important in QFT

Required to compute physical observables

Higgs
discovery



anomalous magnetic
dipole moment

Fruitful interplay with mathematics

special functions; differential equations; algebraic geometry

Challenge: multi-loop Feynman integrals

$$f(p_i \cdot p_j) = \int d^4 k_1 d^4 k_2 \dots d^4 k_L \mathcal{I}(p_i; k_j)$$

↑
Transcendental
function

↑
rational loop integrand

- What special functions do Feynman integrals evaluate to?
- What singularities do they have?
- How can we determine the functions efficiently?

Usefulness of dimensional regularisation

Integrals in non-integer dimensions:

$$d^4 k \longrightarrow d^{4-2\epsilon} k$$

Physical use: serves to regulate divergences:

$$f(p_i; \epsilon) = \frac{1}{\epsilon^{k_{\max}}} \sum_{k \geq 0} \epsilon^k f^{(k)}(p_i)$$

Mathematical use: organizing principle (grading) $f^{(k)}$
are k-fold iterated integrals (under some assumptions)
We say they have uniform transcendental (UT) weight.

Canonical differential equation (DE) method

n-th order partial differential equations (Picard-Fuchs)

↙ ↘
system of 1st order DE

Typically
complicated



↓
'canonical' DE define special functions

Very
simple

Canonical differential equations

[Henn, 2013]

Generic
form:

$$\frac{d}{dx} f(x, \epsilon) = A(x, \epsilon) f(x, \epsilon)$$

Matrix with complicated rational entries

$$f \longrightarrow T^{-1} f \quad \downarrow \quad A \longrightarrow T^{-1} A T - T^{-1} \partial_x T$$

Canonical
form:

$$\frac{d}{dx} f(x, \epsilon) = \epsilon \left[\sum_k m_k \frac{1}{x - x_k} \right] f(x, \epsilon)$$

Basis of uniform weight
Feynman integrals (UT)

Constant
matrices

Singular
points

How to find UT integrals?

Solution of canonical DE are iterated integrals of uniform transcendental weight (UT)

- Analysis of singularity structure of DE matrix A

$$f \longrightarrow T^{-1} f \qquad A \longrightarrow T^{-1} AT - T^{-1} \partial_x T$$

[Lee '14; JMH, '14; Prausa '17; Meyer '17; Gituliar, Magerya '17]

- Analysis of residues of rational loop integrand

[Arkani-Hamed, Bourjaily, Cachazo, Goncharov, Postnikov, Trnka, 2012]

related to `d-log integrand`:

$$\frac{d^4 \ell (p_1 + p_2)^2 (p_1 + p_3)^2}{\ell^2 (\ell + p_1)^2 (\ell + p_1 + p_2)^2 (\ell - p_4)^2} \\ = d \log \left(\frac{\ell^2}{(\ell - \ell^*)^2} \right) d \log \left(\frac{(\ell + p_1)^2}{(\ell - \ell^*)^2} \right) d \log \left(\frac{(\ell + p_1 + p_2)^2}{(\ell - \ell^*)^2} \right) d \log \left(\frac{(\ell - p_4)^2}{(\ell - \ell^*)^2} \right)$$

Practical need for refined methods

Often easy to find a few UT integrals,
but **hard to find a complete UT basis.**

Some methods **restricted to small matrices,
or to few variables.**

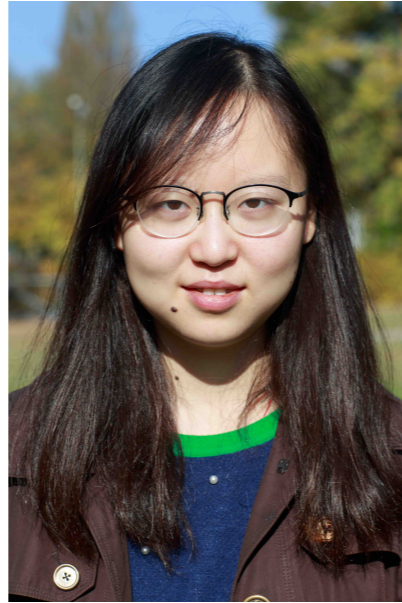
Our new method:

- **needs only one UT integral as starting point**
- **can deal with larger matrices**
- **applies to multiple variables**

The collaboration at the Max Planck Institute for Physics



Christoph Dlapa
(PhD student)



Kai Yan
(Postdoc)



Based on JHEP 05 (2020) 025



We develop further an idea by [Höschele, Hoff, Ueda '14]

- First-order DE only canonical if all integrals are UT
- Picard-Fuchs eq. for a single integral is unique, and contains **valuable information**.
- They applied this knowledge to find the remaining UT integrals, for cases with 2-3 master integrals, and outlined a general procedure.
- We formulate the method in matrix form, and solve the equations systematically



Public algorithm:

[Dlapa, Henn, Yan '19]

<https://github.com/UT-team/INITIAL>

Step I: Picard-Fuchs eq. of a UT integral

Assumption: know one UT integral f_1

- Complete to (any) basis \vec{f} , and compute DE

$$\frac{d}{dx} \vec{f} = A(x, \epsilon) \vec{f}.$$

- Differentiate $(f_1', f_1'', \dots, f_1^{(n)})^T = \Psi(x, \epsilon) \vec{f}$

- Ψ^{-1} yields Picard-Fuchs eq. for f_1 ← given by derivatives of A



Idea: use the infinite amount of information provided by f_1 being UT.

Step 2: Ansatz for canonical system

Assume existence of UT basis \vec{g} with $g_1 = f_1$.

$$\frac{d}{dx} \vec{g} = \epsilon \tilde{A}(x) \vec{g}.$$

$$(g_1', g_1'', \dots, g_1^{(n)})^T = \Phi(x, \epsilon) \vec{g}.$$

Basis transformation: $\vec{f} = T \vec{g}$. where $T \equiv \Psi^{-1} \Phi$.

$$\vec{v}_0 \Psi^{-1} \Phi = \vec{v}_0 \longleftarrow \text{Unit vector}$$

Constraint on Φ from expected canonical form:

$$\tilde{A}(x) = \sum_i \frac{d \ln a_i(x)}{dx} m_i$$

Follow from singularities of $A(x)$

constant matrices: to be found

Step 3: Solve equations algorithmically

$$\vec{v}_0 \Psi^{-1} \Phi = \vec{v}_0$$

Known matrix

x, ϵ

Parametrized by set of unknown

constant matrices m_i

Known polynomial ϵ dependence

Unit vector

Equations valid for any x : can use finite field methods.

Solve at each order in ϵ .

Higher orders in ϵ provide consistency check.

Public algorithm available

All steps implemented in public algorithm

<https://github.com/UT-team/INITIAL>

Presentation was for single-variable case, but method works for multiple variables; key advantage over other methods!

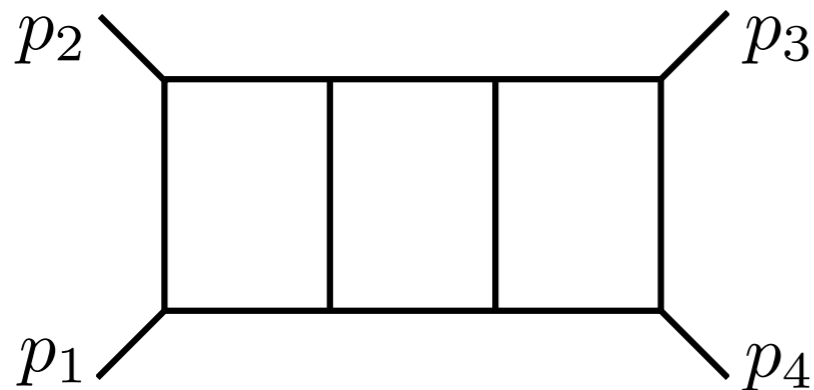
Corollary of our method: test of UT property of a given integral

- Algorithm can be used to test whether candidate integrals can be UT
- Can suggest modifications if an integral is 'almost' UT
- Very useful to algorithmically search for UT integrals

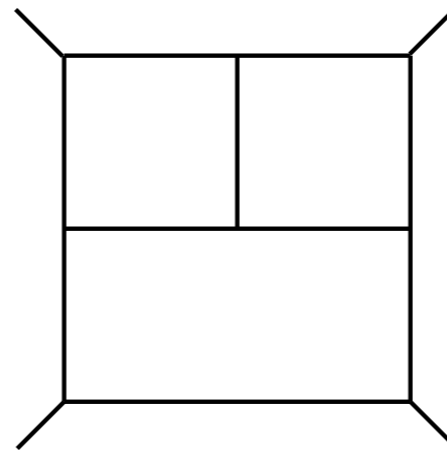
State-of-the art applications

Application I: Planar three-loop on-shell integrals

Four-particle scattering



(a)



(b)

[Henn, Smirnov², '14]

Kinematics: $\sum_{i=1}^4 p_i = 0$, $p_i^2 = 0$. $s = (p_1 + p_2)^2$, $t = (p_2 + p_3)^2$, $x = t/s$.

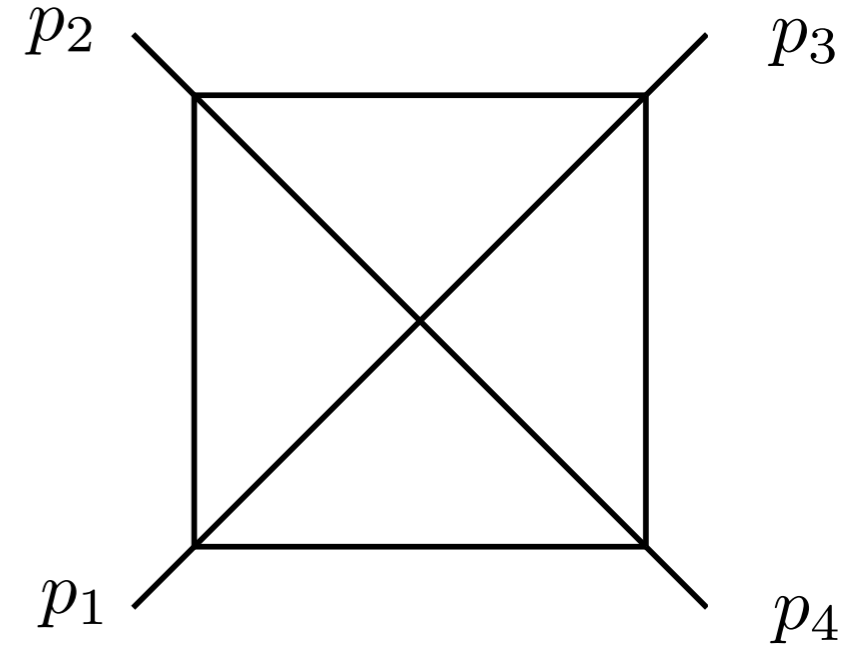
- UT integral in top sector easily found using d-log integrand analysis [Arkani-Hamed et al '11, Henn '13, Wasser, '16]
- Obtained full system of differential equations
Matrix size 26x26 for case (a), 41x41 for case (b)

$$\tilde{A}(x) = m_0 \frac{d}{dx} \ln x + m_1 \frac{d}{dx} \ln(1 + x)$$

Matrix block structure: at most 3 master integrals per sector.

Application 2: Four-loop four-particle scattering

1. Solved the system on the cut (8 MI)
2. Found further UT integrals (off the cut) by testing a large list of candidates
3. Full canonical system (19 MI) of DE
4. Fix boundary constants from analyticity + one trivial integral
5. Analytic result for scalar integral:

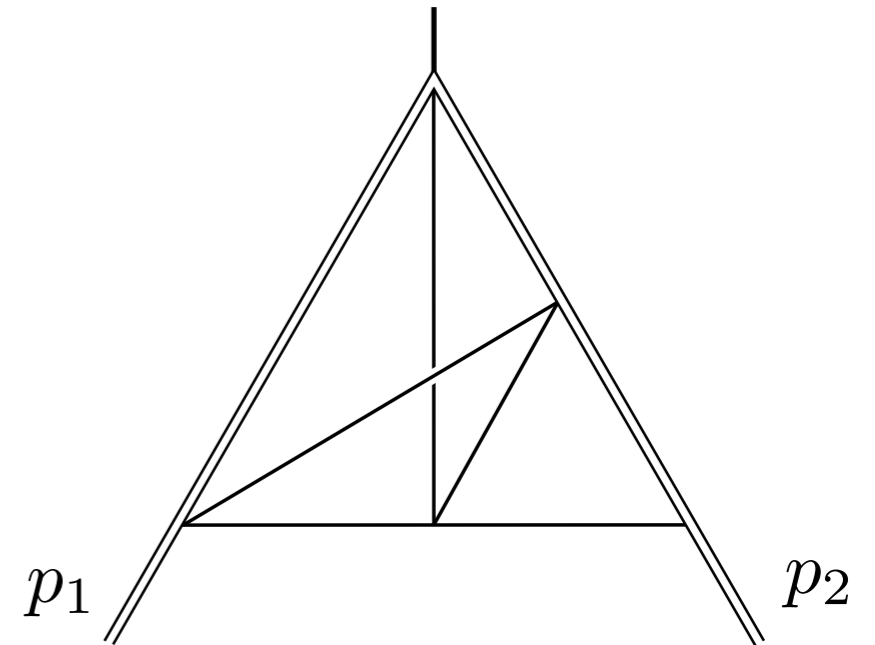


$$\begin{aligned}
 \epsilon^6(1-5\epsilon)(1-6\epsilon) G_{1,1,1,1,1,1,1,1,0,0,0,0,0,0,0,0,0,0,0,0} &= 5\epsilon^5\zeta_5 + \epsilon^6 \left\{ \frac{\pi^6}{30} + 3\zeta_3^2 + \frac{1}{6}\pi^4 H_{-1,-1}(x) + \frac{1}{5}\pi^4 H_{-1,0}(x) \right. \\
 &- \frac{8}{45}\pi^4 H_{0,-1}(x) + \frac{2}{3}\pi^2 H_{-1,-1,0,0}(x) - 2\pi^2 H_{-1,0,0,-1}(x) + \frac{2}{3}\pi^2 H_{-1,0,0,0}(x) - \frac{2}{3}\pi^2 H_{0,-1,-1,0}(x) \\
 &+ 2\pi^2 H_{0,0,-1,-1}(x) - \frac{2}{3}\pi^2 H_{0,0,-1,0}(x) + 4 H_{-1,-1,0,0,0,0}(x) - 4 H_{-1,0,0,-1,0,0}(x) - 4 H_{0,-1,-1,0,0,0}(x) \\
 &\left. + 4 H_{0,0,-1,-1,0,0}(x) - \frac{4}{3}\pi^2\zeta_3 H_{-1}(x) - 4\zeta_3 H_{0,-1,0}(x) + 4\zeta_3 H_{0,0,-1}(x) - 20\zeta_5 H_{-1}(x) \right\}
 \end{aligned}$$

Application 3: many coupled master integrals

- State of the art for heavy quark effective theory (HQET) integrals: 3 loops [Henn, Korchemsky Marquard '15]
- Here: **four-loop** non-planar integrals
- **17 coupled master integrals**

Kinematics $\cos \phi = \frac{p_1 \cdot p_2}{\sqrt{p_1^2 p_2^2}} = \frac{1}{2} \left(x + \frac{1}{x} \right)$



- Form of singularities:

$$\tilde{A}(x) = m_0 \frac{d}{dx} \ln x + m_1 \frac{d}{dx} \ln(1+x) + m_{-1} \frac{d}{dx} \ln(1-x)$$

Solved easily (~ 10 min) using our algorithm.

Recently applied to four-loop cusp anomalous dimension!

[Brüser, Dlapa, Henn, Yan, 2007.04851 [hep-th] (submitted to PRL)]

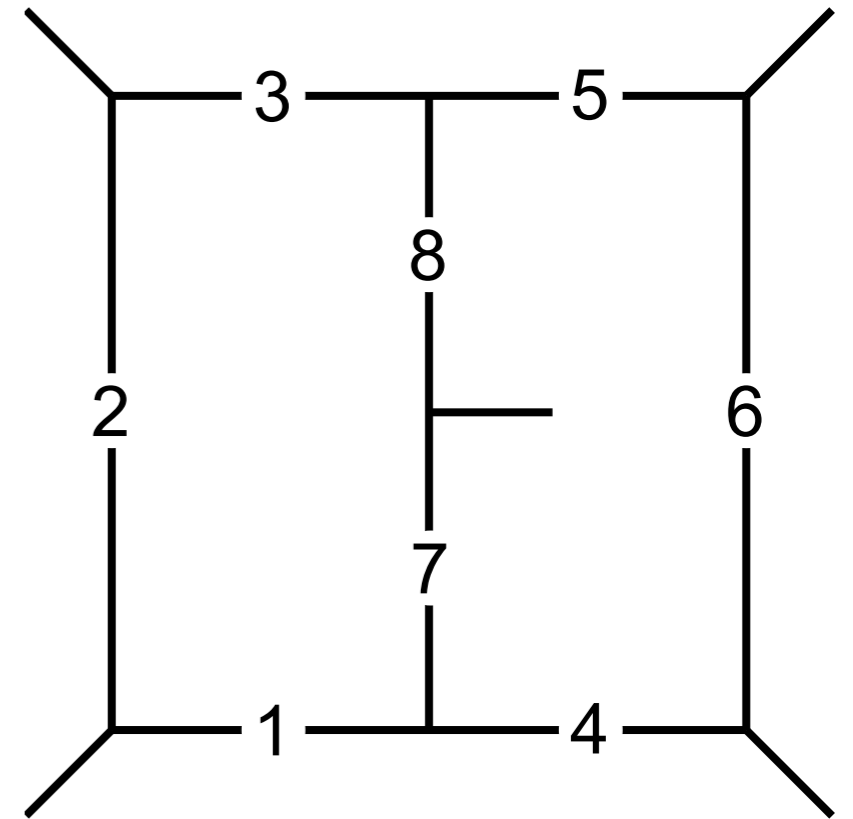
Application 4: multi-variable case

- Two-loop double pentagon integrals
- Computed only 1 year ago using state-of-the-art methods (e.g. D-dimensional leading singularities)

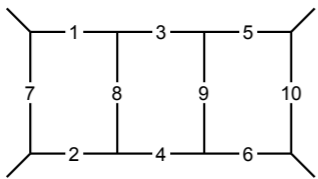
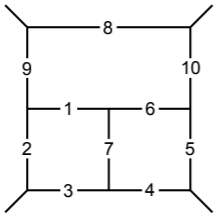
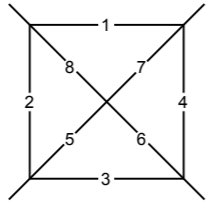
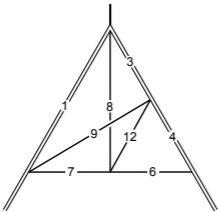
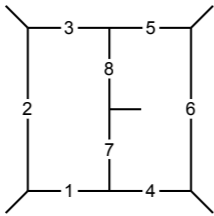
[Abreu, Dixon, Herrmann, Zeng '18; Chicherin, Gehrmann, Henn, Wasser, Zhang, Zoia '18]

- 9 coupled integrals in top sector
- Four different kinematic variables
- 17 relevant alphabet letters

Our algorithm takes 5 minutes to find the UT basis on the cut.



Algorithm is efficient for many coupled integrals, and in multi-variable case

Type of problem		#MI	#vars	#letters	time [min.]	Memory [MB]
Full three-loop DE		26 3	1	2	2	330
		41 3	1	2	34	1710
Full four-loop DE		19 12	1	2	1	240
HQET DE on cut		17 17	1	3	2	390
Five-point integrals DE on cut		9 9	4	17	5	510

Discussion

Our work provides an **automated public tool** for the calculation of canonical differential equations. It removes an important bottleneck in the **calculation of Feynman integrals**.

- Our method is **efficient** for solving large systems of coupled integrals
- Applies to **multi-variables** case
- Corollary: **test of the UT property**

Outlook: more complicated integrals

The idea of **canonical form of differential equations** has also been explicitly applied **for elliptic polylogarithms**. We find it conceivable that our new ideas can be applied here as well.

‘Pre-canonical form’:

$$d g(\mathbf{x}, \epsilon) = [dA_0(\mathbf{x}) + \epsilon dA_1(\mathbf{x})] g(\mathbf{x}, \epsilon)$$

[Henn '14; Mizera, Pokraka '19]

Canonical form

Integrating out A_0 introduces elliptic functions:

$$d f(\mathbf{x}, \epsilon) = \epsilon dA(\mathbf{x}) f(\mathbf{x}, \epsilon)$$

Beyond logarithmic kernels.

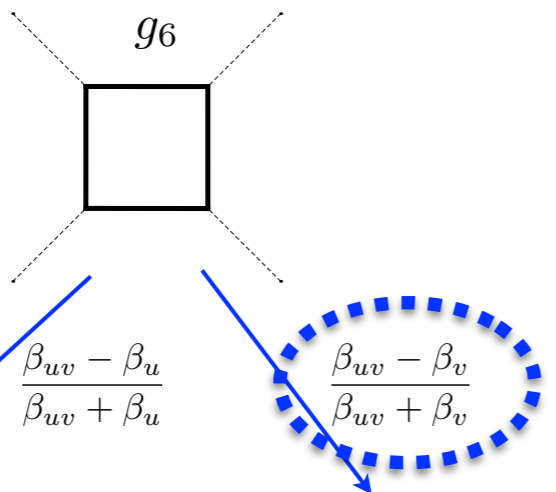
[Broedel, Duhr, Dulat, Penante, Tancredi '18;
Adams and Weinzierl '18]

Outlook: finite integrals (e.g. D=4)

For **finite integrals**, simplifications occur, and the matrices become nilpotent [Caron-Huot, Henn '14].

Transcendental weight

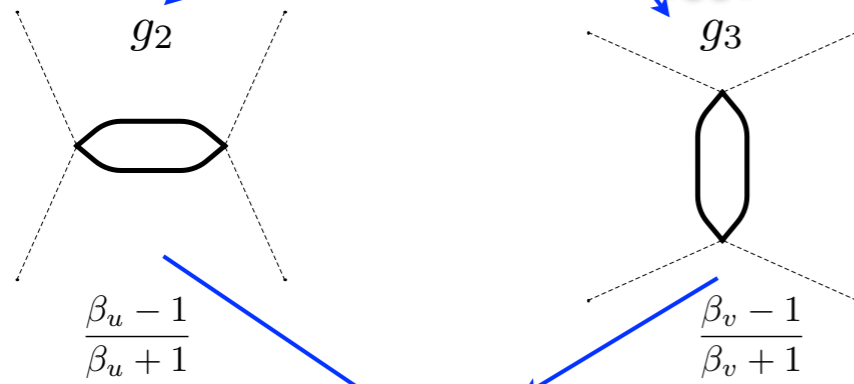
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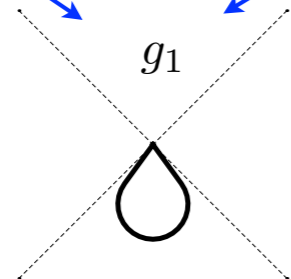
Arrows represent non-zero matrix entries

$$d \begin{pmatrix} g_6 \\ g_3 \\ g_2 \\ g_1 \end{pmatrix} = d \begin{pmatrix} 0 & \log \frac{\beta_{uv} - \beta_v}{\beta_{uv} + \beta_u} & \log \frac{\beta_{uv} - \beta_u}{\beta_{uv} + \beta_v} & 0 \\ 0 & 0 & 0 & \log \frac{\beta_u - 1}{\beta_u + 1} \\ 0 & 0 & 0 & \log \frac{\beta_v - 1}{\beta_v + 1} \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} g_6 \\ g_3 \\ g_2 \\ g_1 \end{pmatrix}$$

1



0



Exploit simpler structure to solve larger systems?

Thank you for your attention!